## STUDY GUIDE

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VOLUME 1: CHAPTERS 1-20
LAIRD KRAMER

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Used successfully by a broad range of students, the Study Guide for University Physics, Twelfth Edition, by Laird Kramer, highlights how you can take advantage of the learning features of the text and make the most of your valuable study time. In addition, it will help you develop the intuition and strong problem-solving skills you'll need for success in your physics course.

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- Problem-Solving Strategies are demonstrated in numerous examples from the text to help you tackle problems more successfully. Examples cover all the key strategies you will need, including free-body diagrams, coordinate systems, and sign conventions.
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## STUDY GUIDE <br> VOLUME 1: CHAPTERS 1-20

## SEARS \& ZEMANSKY'S

# UNIVERSITY <br>  <br> 12TH EDITION 

YOUNG AND FREEDMAN

## LAIRD KRAMER

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SEARS \& ZEMANSKY'S

# UNIVERSITY <br>  <br> 12TH EDITION 

## YOUNG AND FREEDMAN

## LAIRD KRAMER

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## Preface

What do an Olympic athlete, your favorite music artist, and Albert Einstein have in common? They all became experts in their fields through practice. To understand physics and to do well in your course, you must practice. When you learned to walk, ride a bike, and drive a car; you had to practice to master those skills. It would be silly to think you can learn physics by listening to lectures and skimming the book. This study guide is designed to help you practice and to build a deep understanding of physics.

Expert problem solvers in physics follow a systematic approach in their problem solving. Elite athletes also follow a systematic approach in their training to reach the upper level of their sport. You should also follow a systematic approach in your physics course to fully develop your skills. To encourage you in building good problem-solving skills, this study guide follows a systematic problem-solving procedure throughout-the Identify, Set Up, Execute, and Evaluate procedure developed in the textbook.

In the Identify phase of the problem, you should identify the relevant concepts. Decide which physics concepts can be used to solve the problem. Identify the target variable in the problem, and keep this target variable in mind as you solve the problem. Don't think you can save time by skipping this step and jumping right into an equation search. You need to plan a strategy for solving the problem: Decide what you know, where you are going, and how to proceed to the solution.

In the Set Up phase of the problem, you should select the equations you will use to solve the problem and how to use them to determine the solution. Make sure you select equations that are appropriate for the physics of the problem, and don't select equations based solely on the variables in the equation. You should sketch each problem to help you visualize the physical situation and guide you to the solution. Rarely do physicists discuss cutting-edge research problems without first sketching their ideas.

When you proceed to Execute the solution, work through the solution step-by-step. Identify all of the known and unknown quantities in the equations, making a note of the target variable. Then do the calculations to find the solution, writing down all of your work so you may return and check it later. If you run into a dead end, don't erase your work as you may find it useful in a later phase of the problem. Try another avenue when you get stuck and you will eventually find the solution.

After completing the problem, Evaluate your result. Your goal is to learn from the problem, and build your physics intuition. Does the answer make sense? If you were estimating how high an elephant can jump, you'd expect it ought to be less than a meter or two. Consider how this problem compares to the last problem you completed, the example in the text, and the example shown in class. Physicists constantly compare and contrast their new results to previous work as they observe natural phenomena, find patterns, and build principles to connect various phenomena.

Questions and problems chosen for this study guide cover the most critical topics you'll encounter. Working through the guide will better prepare you for homework (including MasteringPhysics) and exams as well as
assist in developing a deeper understanding of physics. One way to build confidence is to try working through the questions and problems in the guide for practice, referring to the solutions only when you get stuck. Building confidence before an exam reduces stress during the exam, improving performance. Several 'Try It Yourself' problems are included at the end of each chapter to help build your confidence. Solution checkpoints are included for each of these problems to help you if you get stuck. Summaries, Objectives, Concepts and Equations, and Problem Summaries are ancillary materials that help bring the physics topics of each chapter into coherence. Taking advantage of all of the components in this study guide will help build your problem-solving repertoire.

This study guide is but one of many resources at your disposal when learning physics. Your instructor, class, and textbook are also important resources. But you should also consider who approaches the material from the same level and perspective as yourself-your fellow students. The best untapped resource in a physics class is often other students learning physics for the first time. Discuss physics as a group and confront your questions together, just as many professionals collaborate in the workplace.

We know physics has a reputation for being challenging. While it can be challenging, many students have succeeded in learning physics. Their success was built on a series of small steps, regular practice, and following a systematic approach. Follow their footsteps and you will master physics as they did. You'll also find physics to be a rich and beautiful subject.

Good luck and enjoy learning physics!
Dedicated to Marley, a 45-pound lab mutt that brought a petaton of happiness to our lives.
Laird Kramer
Miami, Florida, 2007

## Units, Physical Quantities, and Vectors

## Summary

Physics is the study of natural phenomena. In physics, we build theories based on observations of nature, and those theories evolve into physical laws. We often seek simplicity: Models are simplified versions of physical phenomena that allow us to gain insight into a physical process. This chapter covers foundational material that we will use throughout our study. We begin with measurements that include units, conversions, precision, significant figures, estimates, orders of magnitude, and scientific notation. We will also examine physical quantities. Scalar quantities, such as temperature, are described by a single number. Vector quantities, such as velocity, require both a magnitude and a direction for a complete description. We will delve deeper into vectors, as they are used throughout physics. We'll also include a summary of critical mathematics skills you will apply throughout your physics career. Techniques developed in this chapter will be used throughout our investigation of physics.

## Objectives

After studying this chapter, you wil! understand

- The process of experimentation and its relation to theory and laws.
- The SI units for length, mass, and time and common metric prefixes.
- How to express results in proper units and how to convert between different sets of units.
- Measurement uncertainties and how significant figures express precision.
- The use and meaning of scalar and vector quantities.
- Various ways to represent vectors, including graphical, component, and unit-vector representations.
- How to add and subtract vectors, both graphically and componentwise.
- How to multiply vectors (the dot product and the cross product).
- The six most critical mathematics techniques you will encounter in physics.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Physical Law | A physical law is a well-established description of a physical phenomenon. |
| Model | A model is a simplified version of a physical system that focuses on its most <br> important features. |
| Système International (SI) | The Systeme International (SI) is the system of units based on metric mea- <br> sures. It established refined definitions of units, including definitions of the <br> second, meter, and kilogram. |
| Significant Figures | The accuracy of a measurement is indicated by the number of significant fig- <br> ures, or the number of meaningful digits, in a value. In multiplying or divid- <br> ing, the number of significant figures in the result is no greater than in the <br> factor with the fewest significant figures. In adding or subtracting, the result <br> can have no more decimal places than the term with the fewest decimal places. |
| Scalar Quantity | A scalar quantity is expressed by a single number. Examples include tempera- <br> ture, mass, length, and time. |
| Vector Quantity | A vector quantity is expressed by both a magnitude and a direction and is <br> often shown as an arrow in sketches. Vectors are frequently represented as <br> single letters with arrows above them or in boldface type. Common examples <br> include velocity, displacement, and force. |
| Component of a Vector | The vector $\vec{A}$ lying in the $x y$ plane has components $A_{x}$ parallel to the $x$-axis and |
| $A_{y}$ parallel to the $y$-axis; $A_{x}$ and $A_{y}$ are the $x$ and $y$ component vectors of $\vec{A}$. |  |
| Vector $\vec{A}$ can be described by unit vectors-vectors that have unity magnitude |  |
| and that align along a particular axis. The unit vectors $\hat{i}, \hat{j}$, and $\hat{k}$ respectively |  |
| align along the $x-, y$-, and $z$-axes of the rectangular coordinate system. Here, |  |



| Magnitude of a Vector | The magnitude of a vector is the length of the vector. Magnitude is a scalar <br> quantity that is always positive. It has several representations, including |
| :--- | :--- |
| Magnitude of $\vec{A}=A=\|\vec{A}\|$. |  |


|  | $A=\sqrt{A_{x}^{2}+A_{y}^{2}} .$ |
| :---: | :---: |
| Vector Addition and Subtraction | Two vectors, $\vec{A}$ and $\vec{B}$, are added graphically by placing the tail of $\vec{A}$ at the tip of $\vec{B}$ : |

Vector $\vec{B}$ is subtracted from vector $\vec{A}$ by reversing the direction of $\vec{B}$ and then adding it to $\vec{A}$ :


| Scalar Product | Vector addition can also be done with component vector <br> and $A_{y}$ of the vector $\vec{A}$ and components $B_{x}$ and $B_{y}$ of the <br> nents $R_{x}$ and $R_{y}$ of the resultant vector $\vec{R}$ are given by |
| :--- | :--- |
| $\qquad$$R_{x}=A_{x}+B_{x}$ and $R_{y}=A_{y}+B_{y}$. |  |
| $\qquad$The scalar, or dot, product $C=\vec{A} \cdot \vec{B}$ of two vectors $\vec{A}$ and <br> tity. It can be expressed in terms of the magnitudes of $\vec{A}$ <br> $\phi$ between the two vectors—that is, |  |
| or in terms of the components of $\vec{A}$ and $\vec{B}-$ that is, |  |
| $\vec{A} \cdot \vec{B}=A_{x} B_{x}+A_{y} B_{y}+A_{z} B_{z}$. |  |

## Vector Product

The scalar product of two perpendicular vectors is zero.
The vector, or cross, product $\vec{C}=\vec{A} \times \vec{B}$ of two vectors $\vec{A}$ and $\vec{B}$ is a vector quantity. The magnitude of $\vec{A} \times \vec{B}$ depends on the magnitudes of $\vec{A}$ and $\vec{B}$ and the angle $\phi$ between the two vectors. The direction of $\vec{A} \times \vec{B}$ is perpendicular to the plane in which vectors $\vec{A}$ and $\vec{B}$ lie and is given by the right-hand rule. The magnitude of $\vec{C}=\vec{A} \times \vec{B}$ is

$$
C=A B \sin \phi
$$

and the components are

$$
\begin{aligned}
& C_{x}=A_{y} B_{z}-A_{z} B_{y} \\
& C_{y}=A_{z} B_{x}-A_{x} B_{z} \\
& C_{z}=A_{x} B_{y}-A_{y} B_{x} .
\end{aligned}
$$

$$
\text { (Magnitude of } \vec{A} \times \vec{B})=A B \sin \phi
$$

The vector product of two parallel or antiparallel vectors is zero.

## Mathematics Review: Top Six Math Skills You Will Need in Introductory Physics

Mathematics is the main language of physics. You will rely on mathematics throughout your study of physics and therefore must become comfortable with mathematical techniques. Here, we present the six most important mathematical techniques you will use throughout your physics career. We strongly encourage you to review these materials thoroughly. When your knowledge of mathematics becomes second nature, your understanding of physics will blossom.

## Math 1: Trigonometry

We will use trigonometry throughout physics; problems involving objects tossed into the air, ramps, velocities, and forces all require trigonometry. The basic trigonometric functions relate the lengths of the sides of a right triangle to the inside angle. We define $\sin \theta, \cos \theta$, and $\tan \theta$ for the right triangle shown:


$$
\sin \theta=\frac{\text { opposite side }}{\text { hypotenuse }}, \cos \theta=\frac{\text { adjacent side }}{\text { hypotenuse }}, \tan \theta=\frac{\text { opposite side }}{\text { adjacent side }} .
$$

Often, the triangle is formed in an $x y$ coordinate system, as is seen in Figure 1.1. Note that the two inside angles complement each other (add to $90^{\circ}$ ), so two sets of relations can be used:

$$
\begin{aligned}
\cos \theta & =\frac{x}{r}, \sin \theta=\frac{y}{r}, \text { and } \tan \theta=\frac{y}{x} \\
\cos \phi & =\frac{y}{r}, \sin \phi=\frac{x}{r}, \text { and } \tan \theta=\frac{x}{y} .
\end{aligned}
$$



Figure 1.1 xy coordinate system.
CAUTION Watch sines and cosines! One common mistake is to automatically associate the $x$ component with cosine and the $y$-component with sine. As you can see from Figure 1.1, this association does not always hold. One of the most common mistakes encountered in physics is confusing components of sines and cosines. By checking components every time, you avoid this mistake.

We will also need to manipulate trigonometric relations, so we will use common trigonometric identities, including the following:

$$
\begin{aligned}
\tan \theta & =\frac{\sin \theta}{\cos \theta} \\
\sin ^{2} \theta+\cos ^{2} \theta & =1 \\
\sin 2 a & =2 \sin a \cos a \\
\cos 2 a & =\cos ^{2} a-\sin ^{2} a
\end{aligned}
$$

Other trigonometric identities are in Appendix B.

## Math 2: Derivatives

Physics often investigates changes and rates of changes of various quantities. Derivatives provide the instantaneous rate of change of a quantity. Derivatives are therefore the natural choice for finding rates of change in physics. They are especially useful when we have functional definitions of quantities. Speed is the rate of change of position. If we are given speed in terms of a position function, we can easily find the speed by taking the derivative of position. The most common derivatives you will encounter in physics are given in Table 1.

$$
\begin{aligned}
& \text { TABLE } 1 \text { : Common } \\
& \text { derivatives. } \\
& \begin{array}{|c|}
\hline \frac{d}{d x} x^{n}=n x^{n-1} \\
\hline \frac{d}{d x} \sin a x=a \cos a x \\
\hline \frac{d}{d x} \cos a x=-a \sin a x \\
\hline \frac{d}{d x} e^{a x}=a e^{a x} \\
\hline \frac{d}{d x} \ln a x=\frac{1}{x} \\
\hline
\end{array}
\end{aligned}
$$

## Math 3: Integrals

In physics, we also need to sum various quantities-quantities that are often given in terms of functions. Integrals sum functions and therefore are used to sum physical quantities. For example, one may find the total mass of an object by integrating its density function. The most common integrals you will
encounter in physics are given in Table 2. In addition, you may want to review integration by parts and trigonometric substitution for integrals in order to solve the more complicated ones.

TABLE 2: Common integrals.

| $\int x^{n} d x=\frac{x^{n+1}}{n+1} \quad(n \neq-1)$ |
| :--- |
| $\int \frac{d x}{x}=\ln x$ |
| $\int \sin a x d x=-\frac{1}{a} \cos a x$ |
| $\int \cos a x d x=\frac{1}{a} \sin a x$ |
| $\int e^{a x} d x=\frac{1}{a} e^{a x}$ |
| $\int \frac{d x}{\sqrt{a^{2}-x^{2}}}=\arcsin \frac{x}{a}$ |
| $\int \frac{d x}{\sqrt{x^{2}+a^{2}}}=\ln \left(x+\sqrt{x^{2}+a^{2}}\right)$ |
| $\int \frac{d x}{x^{2}+a^{2}}=\frac{1}{a} \arctan \frac{x}{a}$ |
| $\int \frac{d x}{\left(x^{2}+a^{2}\right)^{3 / 2}}=\frac{1}{a^{2}} \frac{x}{\sqrt{x^{2}+a^{2}}}$ |
| $\int \frac{x d x}{\left(x^{2}+a^{2}\right)^{3 / 2}}=\frac{1}{\sqrt{x^{2}+a^{2}}}$ |

## Math 4: Graphs

Graphs are common in many fields, including physics, in which data or information is plotted as a function of time. However, many students do not gain a full appreciation of graphs, and instructors often take graph interpretation for granted. In physics, graphs provide added insight into complex phenomena. We will review important features of graphs to help your physics interpretations, as well as to help your interpretation of graphs wherever you encounter them.

Let's begin by examining the graph in Figure 1.2. Here, position is plotted as a function of time for three different objects. As time increases, the positions of all three objects increase, so the object is moving away from the origin. The slope gives the rate of change of the position-the speed-of the object:

$$
\text { slope }=\text { rate of change of position }=\text { speed }=\frac{\Delta(\text { position })}{\Delta(\text { time })} .
$$

All three lines are straight lines, indicating that the speed is constant for each object. Both objects $A$ and $B$ start at the same initial position. Object $A$ 's line has the greatest slope, so object $A$ moves the fastest or has the greatest speed. Object $C$ starts away from objects $A$ and $B$ and moves away at a
slower rate than the other two. Where the lines intersect, objects $A$ and $C$ are at the same position at the same time. After the intersection, object $A$ moves away from object $C$ and so passes object $C$.


Figure 1.2 Position-versus-time graph.
Moving on to a more interesting case, we see that the slope in Figure 1.3 is not constant; calculus will be necessary to interpret this graph. The rate of change of the position varies, so we will need to consider the instantaneous slope, or the derivative of the position:

$$
\text { instantaneous slope }=\lim _{\Delta t \rightarrow 0} \frac{\Delta(\text { position })}{\Delta(\text { time })}=\frac{d x}{d t}
$$

Graphically, the instantaneous slope is the tangent to the line in the position-versus-time graph.
If we look at the figure, we see that the object moves away from the initial position, remains at a constant position for a period of time, and then moves toward the initial position. How it moves is found by looking at the slope. The slope of the line increases between point $A$ and point $B$, indicating that the object begins by moving slowly and then speeding up as it moves away. After point $B$, the slope decreases until point $C$, where the slope becomes zero. Thus, the object begins slowing down at point $B$ and stops at point $C$. At point $D$, the slope decreases and then becomes constant, indicating that the object moves back toward its initial position. The slope past point $D$ is negative, indicating that the object moves in the direction opposite that of its initial movement. Points $E$ and $F$ show the instantaneous slope at two points on the curve. The slope at point $E$ is greater than the slope at point $F$, indicating that the object is moving faster at point $E$.


Figure 1.3 Position-versus-time graph.
The speed of another object is plotted in Figure 1.4. From this graph, we can use calculus to determine the distance the object travels in a given time interval. The distance traveled by an object between times $t_{a}$ and $t_{b}$ is the area under the curve, or the shaded area shown in the figure. The area under the
curve is the speed multiplied by the time, which is the distance traveled, as we will learn in Chapter 2. To find the area, we will need to sum up, or integrate, the speed over time:

$$
\text { distance }=\int_{t_{a}}^{t_{b}}(\text { speed }) d t
$$



Figure 1.4 Speed-versus-time graph.
We have now seen somewhat how calculus and certain graphs are intertwined. Derivatives and integrals, respectively, give us the slope and the areas under a curve. We will see that switching back and forth between graphs and the mathematics of calculus will help develop our physics intuition and skills. When you're confused about integrals, consider using a graph to clarify your understanding.

## Math 5: Solving quadratic equations

We will encounter quadratic equations throughout our physics investigations. We can either factor the equation to obtain the solutions or use the quadratic equation. It is often easier to apply the quadratic equation and not even attempt to factor, as the quadratic equation leads directly to the solutions. For a quadratic equation of the form $a x^{2}+b x+c=0$ with real numbers $a, b$, and $c$, the solutions are given by the quadratic equation:

$$
x=\frac{-b \pm \sqrt{b^{2}-4 a c}}{2 a}
$$

## Math 6: Solving simultaneous equations

Our goal in solving problems is to determine unknown quantities in equations. When we have multiple unknowns, we will need multiple equations to find solutions. We will need at least as many equations as we have unknowns to solve any problem.

When we are presented with multiple equations, one option is to rewrite one of the equations, solving for one unknown in terms of the other unknown(s), and substitute the result into the other equation(s) to eliminate one variable. For example, if we need to solve for both $x$ and $y$, given two equations, we first rewrite one equation to solve for $x$ in terms of $y$. Then we replace the $x$ terms in the second equation with the solution of the first equation, leaving an equation having only $y$ terms.

A second technique comes from linear algebra. We multiply each equation by a factor and then add or subtract the two equations. By choosing the proper factor, we eliminate one variable in the process. For example, with the equations

$$
\begin{aligned}
& 3 x+2 y=3 \\
& 4 x-3 y=7
\end{aligned}
$$

we multiply the top equation by 4 and the bottom equation by 3 , leaving

$$
\begin{aligned}
& 12 x+8 y=12 \\
& 12 x-9 y=21
\end{aligned}
$$

If we now subtract the two equations, the $x$ variable is eliminated. We can also multiply the top equation by 3 and the bottom equation by 2 and add the two equations to eliminate $y$. The key in this process is to multiply all terms of each equation by the same factor.

## Other mathematical topics

Other mathematical relations that we will encounter as we cover the material of this course include the following:

- Circumference, area, surface area, and volume of spheres and cylinders
- Exponentials, logarithms, and their identities
- The binomial theorem
- Power series expansions of algebraic, trigonometric, and exponential functions
- Multivariable calculus, including derivatives and integrals in two and three dimensions.

It is best to review the preceding topics as you encounter them in the course. Appendix B includes a summary of these topics. You may also want to consult your mathematics textbooks or the Internet for more information.

CAUTION Don't be afraid to review math! Knowing math will let you focus on physics and save time when you solve problems.

## Conceptual Questions

## 1: Sketch the situation

A man uses a cable to drag a trunk up the loading ramp of a mover's truck. The ramp has a slope angle of $20.0^{\circ}$, and the cable makes an angle of $30.0^{\circ}$ with the ramp. Make a sketch of this situation.

## Solution



Figure 1.5 Sketch of trunk being dragged up a loading ramp.

IDENTIFY, SET UP, AND EXECUTE: The sketch is shown in Figure 1.5. The ramp makes an angle of $20^{\circ}$ with the ground. The trunk is on the ramp and the cable is attached to the trunk. The cable makes an angle of $30^{\circ}$ with respect to the ramp, clearly marked. The mover is shown pulling the trunk up the ramp.

EVALUATE: Understanding the physical situation in physics problems is critical for a correction interpretation. You should always draw a diagram (or diagrams) of the physical system you are investigating. Even when a figure is provided, it is often useful to sketch the important aspects. Only after creating a diagram should you proceed to interpret the physics and determine the proper equations to apply.

## 2: Dimensional analysis practice

Based only on consistency of units, which of the following formulas could not be correct? In each case, $x$ is distance, $v$ is speed, and $t$ is time.
(a) $t=\sqrt{\frac{2 x}{9.8 \mathrm{~m} / \mathrm{s}^{2}}}$
(b) $\quad x=v t+\left(4.9 \mathrm{~m} / \mathrm{s}^{2}\right) t$
(c) $\quad v=v_{0} \sin \theta+\frac{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) x}{v_{0} \cos \theta}$
(d) $\quad x^{2}-\frac{2 v_{0}^{2} \sin \theta}{9.8 \mathrm{~m} / \mathrm{s}^{2}}-\frac{2 v^{2}}{9.8 \mathrm{~m} / \mathrm{s}^{2}}=0$
(e) $\quad v^{2}=v_{0}^{2}-2\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)\left(v_{0} \tan \theta-\frac{1}{2}\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) t^{2}\right)$
(f) $t=\frac{v^{2}+\left(4.9 \mathrm{~m} / \mathrm{s}^{2}\right) x}{\left(3.0 \mathrm{~m} / \mathrm{s}^{2}\right) x}$

## Solution

IDENTIFY, SET UP, AND EXECUTE: For each of the six equations, carefully examine the units of each term in the equation. Equations $(a),(c)$, and $(d)$ are dimensionally correct; however, $(b),(e)$, and $(f)$ are incorrect. The far-right term in equation $(b)$ has units of $(\mathrm{m} / \mathrm{s})$, while the other two terms have units of $(\mathrm{m})$. In equation $(e)$, the left term inside the rightmost set of parentheses (the term $v_{0} \tan \theta$ ) has units of ( $\mathrm{m} / \mathrm{s}$ ), which, when combined with the $\left(\mathrm{m} / \mathrm{s}^{2}\right)$ outside of the left parenthesis, would result in units of $\left(\mathrm{m}^{2} / \mathrm{s}^{3}\right)$. The other three terms in the equation have units of $\left(\mathrm{m}^{2} / \mathrm{s}^{2}\right)$. The fraction in equation $(f)$ has no units, while the left-hand side has units of $(s)$.

Thus, an error exists in each of the three equations $((b),(e)$, and $(f))$, since the units on the two sides of the equation do not agree. The next step would be to recheck our derivation to locate the source of the mistake.

EVALUATE: Dimensional analysis is a powerful technique to help keep you from making errors. Catching the three errors in this problem would save time while reducing confusion. Always check your units!

## 3: Maximum and minimum magnitudes of a vector

Given vector $\vec{A}$ with magnitude 1.3 N and vector $\vec{B}$ with magnitude 3.4 N , what are the minimum and maximum magnitudes of $\vec{A}+\vec{B}$ ?

## Solution

IDENTIFY, SET UP, AND EXECUTE: The maximum magnitude is achieved when the two vectors are parallel and point in the same direction. The minimum magnitude is achieved when the two vectors are parallel and point in opposite directions. (The vectors are then called antiparallel.)

For parallel vectors, the magnitude is the sum of their magnitudes, 4.7 N in this case. For antiparallel vectors, the magnitude is the difference of their magnitudes, 2.1 N here.
EVALUATE: This example helps illustrate the fact that vectors do not add like ordinary scalar numbers. They do not subtract like scalar numbers either. The magnitude of $\vec{A}+\vec{B}$ for any arbitrary alignment of the two vectors must lie between 2.1 N and 4.7 N .

## 4: Finding parallel and perpendicular vectors

If you are given two vectors, how can you determine whether the vectors are parallel or perpendicular?

## Solution

IDENTIFY, SET UP, AND EXECUTE: The cross product of two parallel vectors is zero. The dot product of two perpendicular vectors is zero. By taking the cross and dot products of the two vectors, you will determine whether they are parallel or perpendicular.

EVALUATE: There are circumstances in which you cannot easily identify parallel and perpendicular vectors, such as vectors lying in the $x y z$ plane and that are given in terms of their components. Here, the best way to identify their orientation is to take dot and cross products.

## Problems

## 1: Convert knots to m/s

A yacht is traveling at 18.0 knots. (One knot is 1 nautical mile per hour.) Find the speed of the yacht in $\mathrm{m} / \mathrm{s}$.

## Solution

IDENTIFY AND SET UP: We'll use a series of conversion factors to solve this problem. Appendix E has 1 nautical mile $=6080 \mathrm{ft}$ and $1 \mathrm{mi}=5280 \mathrm{ft}=1.609 \mathrm{~km}$. We know that $1 \mathrm{~km}=1000 \mathrm{~m}$ and that 1 hour $=60 \mathrm{~min}=60 \times(60 \mathrm{~s})=3600 \mathrm{~s}$.
EXECUTE: We apply the conversion factors to the speed in knots to solve:

$$
18.0 \text { knots }=\left(\frac{18.0 \text { nautieatmites }}{1 \mathrm{k}}\right)\left(\frac{6080 \mathrm{ft}}{1 \text { nautieatmite }}\right)\left(\frac{1.609 \mathrm{~km}}{5280 \mathrm{ft}}\right)\left(\frac{1000 \mathrm{~m}}{1 \mathrm{~km}}\right)\left(\frac{1 \mathrm{k}}{3600 \mathrm{~s}}\right)=9.26 \mathrm{~m} / \mathrm{s}
$$

EVALUATE: Using a combination of several conversion factors, we have found that 18.0 knots is equal to $9.26 \mathrm{~m} / \mathrm{s}$. The last four quantities in parentheses are each equal to unity; hence, multiplying 18.0 knots by several factors of unity doesn't change the magnitude of the quantity. Crossing out the units helps prevent mistakes.

## 2: Finding components of vectors

Find the $x$ and $y$ components of the vector $\vec{A}$ in Figure 1.6. The magnitude of vector $\vec{A}$ is 26.2 cm .


Figure 1.6 Problem 2.

## Solution



Figure 1.7 Problem 2 with components.

IDENTIFY AND SET UP: We will find the components of a vector by examining the triangle formed by the vector and the coordinate axes. Figure 1.7 shows Figure 1.6 redrawn to include the component vectors.
EXECUTE: The $x$ component of $\vec{A}$ is located opposite the $67.1^{\circ}$ angle; hence, we'll use the sine function:

$$
A_{x}=A \sin 67.1^{\circ}=(26.2 \mathrm{~cm}) \sin 67.1^{\circ}=24.1 \mathrm{~cm}
$$

The $y$ component of $\vec{A}$ is located adjacent to the $67.1^{\circ}$ angle; thus, we'll use the cosine function:

$$
A_{y}=A \cos 67.1^{\circ}=(26.2 \mathrm{~cm}) \cos 67.1^{\circ}=10.2 \mathrm{~cm}
$$

The vector has an $x$ component of 24.1 cm and a $y$ component of 10.2 cm .
EVALUATE: Finding the components of the vector required applying the sine and cosine functions. Often, but not always, the horizontal components will use cosine and the vertical components will use sine. This example illustrates an exception to that general assertion. It is important to examine a problem carefully in order to identify the proper trigonometric function for each component.

## 3: Vector addition

Find the vector sum $\vec{A}+\vec{B}$ of the two vectors in Figure 1.8. Express the results in terms of components.


Figure 1.8 Problem 3.

## Solution



Figure 1.9 Sketch of Problem 3.
IDENTIFY AND SET UP: Figure 1.9 shows a sketch of the two vectors added together, head to tail. The sketch indicates that we should expect a resultant in the first quadrant, with positive $x$ and $y$ components. We will add the vectors by adding their $x$ and $y$ components, using the Cartesian coordinate system provided.

EXECUTE: We find the components of the vectors by examining the triangles made by the vectors and their components. For $\vec{A}$,

$$
\begin{aligned}
& A_{x}=A \cos 60.0^{\circ}=(15.0 \mathrm{~N}) \cos 60.0^{\circ}=7.50 \mathrm{~N} \\
& A_{y}=A \sin 60.0^{\circ}=(15.0 \mathrm{~N}) \sin 60.0^{\circ}=13.0 \mathrm{~N}
\end{aligned}
$$

For $\vec{B}$,

$$
\begin{aligned}
& B_{x}=B \sin 40.0^{\circ}=(10.0 \mathrm{~N}) \sin 40.0^{\circ}=6.43 \mathrm{~N} \\
& B_{y}=-B \cos 40.0^{\circ}=-(10.0 \mathrm{~N}) \cos 40.0^{\circ}=-7.66 \mathrm{~N}
\end{aligned}
$$

We can now sum the components:

$$
\begin{aligned}
& R_{x}=A_{x}+B_{x}=7.50 \mathrm{~N}+6.43 \mathrm{~N}=13.9 \mathrm{~N} \\
& R_{y}=A_{y}+B_{y}=13.0 \mathrm{~N}-7.66 \mathrm{~N}=5.34 \mathrm{~N}
\end{aligned}
$$

The resultant vector has an $x$ component of 13.9 N and a $y$ component of 5.34 N .
EVALUATE: The resultant vector has positive components and resides in the first quadrant, as expected. Note how the components of the two vectors include both sine and cosine terms (i.e., vector A's $x$ component includes the cosine component, and vector $B$ 's $x$ component includes the sine component). This results from how the vectors' angles were given: vector $A$ 's angle was with respect to the horizontal axis and vector $B$ 's angle was with respect to the vertical axis. It is critical not to automatically associate all horizontal components with the cosine and all vertical components with the sine.
Practice Problem: Find the magnitude and direction of the resultant vector. Answer: The magnitude is 14.9 N , and its direction is $21.0^{\circ}$ above the positive $x$-axis.

## 4: Determine displacement on a lake

Marie paddles her canoe around a lake. She first paddles 0.75 km to the east, then paddles 0.50 km $30^{\circ}$ north of east, and finally paddles $1.0 \mathrm{~km} 50^{\circ}$ north of west. Find the resulting displacement from her origin.

## Solution

IDENTIFY: Displacement is a vector indicating change in position. The displacement vector points from the starting point of a journey to the endpoint. If we represent each of the three segments of the journey as a vector, the displacement vector is the sum of the three vectors. The goal is to find the sum of the three displacement vectors.


Figure 1.10 Problem 4.

SET UP: Figure 1.10 shows a sketch of the three displacement segments (labeled $\vec{A}, \vec{B}$, and $\vec{C}$ ) and the resultant displacement vector $(\vec{R})$. We will add the three vectors, using the Cartesian coordinate system in the figure.

EXECUTE: We find the components of the vectors by examining the triangles made by the vectors and their components. For $\vec{A}$, there is only a horizontal component:

$$
\begin{aligned}
& A_{x}=A=0.75 \mathrm{~km} \\
& A_{y}=0
\end{aligned}
$$

For $\vec{B}$,

$$
\begin{aligned}
& B_{x}=B \cos 30^{\circ}=(0.50 \mathrm{~km}) \cos 30^{\circ}=0.443 \mathrm{~km} \\
& B_{y}=B \sin 30^{\circ}=(0.50 \mathrm{~km}) \sin 30^{\circ}=0.250 \mathrm{~km}
\end{aligned}
$$

For $\vec{C}$,

$$
\begin{aligned}
& C_{x}=-C \cos 50^{\circ}=-(1.0 \mathrm{~km}) \cos 50^{\circ}=-0.643 \mathrm{~km}, \\
& C_{y}=C \sin 50^{\circ}=(1.0 \mathrm{~km}) \sin 50^{\circ}=0.766 \mathrm{~km} .
\end{aligned}
$$

The $x$ component is negative, as it points to the west. We can now sum the components:

$$
\begin{aligned}
& R_{x}=A_{x}+B_{x}+C_{x}=0.75 \mathrm{~km}+0.443 \mathrm{~km}-0.643 \mathrm{~km}=0.550 \mathrm{~km} \\
& R_{y}=A_{y}+B_{y}+C_{y}=0 \mathrm{~km}+0.250 \mathrm{~km}+0.766 \mathrm{~km}=1.016 \mathrm{~km}
\end{aligned}
$$

The resultant displacement vector has an $x$ component of 0.55 km and a $y$ component of 1.02 km . We can express the displacement vector in terms of magnitude and direction. To find the magnitude, we use the Pythagorean theorem:

$$
R=\sqrt{R_{x}^{2}+R_{y}^{2}}=\sqrt{(0.550 \mathrm{~km})^{2}+(1.016 \mathrm{~km})^{2}}=1.16 \mathrm{~km}
$$

The inverse tangent gives us the angle:

$$
\theta=\tan ^{-1} \frac{R_{y}}{R_{x}}=\tan ^{-1} \frac{1.016 \mathrm{~km}}{0.550 \mathrm{~km}}=61.6^{\circ}
$$

The resultant displacement vector has a magnitude of 1.16 km and points $61.6^{\circ}$ above the positive $x$-axis.
EVALUATE: Marie paddled a total of 2.25 km , only to end up 1.16 km away from her starting point. This shows how the magnitude of a vector sum can be smaller than the sum of the magnitudes of the individual vectors. Note that we carried an extra significant figure through the calculations and rounded off only in the final step.

## 5: Finding the dot product

Find the dot product of vectors $\vec{A}$ and $\vec{B}$ if $\vec{A}=5.0 \hat{i}+2.3 \hat{j}-6.4 \hat{k}$ and $\vec{B}=12.0 \hat{i}-4.7 \hat{j}+9.3 \hat{k}$.

## Solution

IDENTIFY AND SET UP: We are given the $x, y$, and $z$ components of the two vectors. We will find the scalar product by using the component form of the dot-product relation.
EXECUTE: The dot product is the sum of the products of the components of the vectors:

$$
\vec{A} \cdot \vec{B}=A_{x} B_{x}+A_{y} B_{y}+A_{z} B_{z} .
$$

For our two vectors,

$$
\begin{aligned}
\vec{A} \cdot \vec{B} & =(5.0)(12.0)+(2.3)(-4.7)+(-6.4)(9.3) \\
& =60-10.8-59.5 \\
& =-10.3
\end{aligned}
$$

The dot product is -10.3 .
EVALUATE: The result is a negative number, indicating that the projection of one vector onto the other points in the direction opposite that of the other. The vectors are not perpendicular, since the dot product is not zero.

We could have attempted to sketch these vectors, but since they are in three-dimensional $x y z$ space, it is difficult to represent them accurately on a two-dimensional page. Building intuition in twodimensional space helps us when we work in three-dimensional space.

## 6: Finding the cross product

Find the cross product $\vec{A} \times \vec{B}$, given $\vec{A}=5.0 \hat{i}+2.3 \hat{j}-6.4 \hat{k}$ and $\vec{B}=12.0 \hat{i}-4.7 \hat{j}+9.3 \hat{k}$.

## Solution

IDENTIFY AND SET UP: We are given the $x, y$, and $z$ components of the two vectors. We will find the cross product by using the component form of the cross-product relation.

EXECUTE: The three components of the cross product are given by various products of the components of the two vectors:

$$
\begin{aligned}
C_{x} & =A_{y} B_{z}-A_{z} B_{y} \\
C_{y} & =A_{z} B_{x}-A_{x} B_{z} \\
C_{z} & =A_{x} B_{y}-A_{y} B_{x} .
\end{aligned}
$$

For this problem, the components are

$$
\begin{aligned}
& C_{x}=(2.3)(9.3)-(-6.4)(-4.7)=(21.4)-(30.1)=-8.7 \\
& C_{y}=(-6.4)(12.0)-(5.0)(9.3)=(-76.8)-(46.5)=-123.3 \\
& C_{z}=(5.0)(-4.7)-(2.3)(12.0)=(-23.5)-(27.6)=-51.1
\end{aligned}
$$

The cross product is $\vec{C}=-8.7 \hat{i}-123.3 \hat{j}-51.1 \hat{k}$.
EVALUATE: The result is a vector, as is expected for the cross product. The vectors are not parallel, since the cross product is not zero.
Practice Problem: Find the magnitude of the resultant vector. Answer: The magnitude is 133.8.

## 7: Review of simultaneous equations

Solve the following expressions for $T_{A}$ and $T_{B}$.

$$
\begin{aligned}
& 27 T_{A}+13 T_{B}=0 \\
& 32 T_{A}+52 T_{B}=22 .
\end{aligned}
$$

## Solution

IDENTIFY AND SET UP: Both of the expressions involve two unknowns, so we cannot find a solution by using only one equation. We will multiply the first equation by 32 , multiply the second by 27, and then subtract the second equation from the first.

EXECUTE: Multiplying the first expression by 32 and the second expression by 27 gives

$$
\begin{aligned}
& 864 T_{A}+416 T_{B}=0 \\
& 864 T_{A}+1404 T_{B}=594
\end{aligned}
$$

Subtracting the second equation from the first gives

$$
\begin{aligned}
864 T_{A}+416 T_{B}-864 T_{A}-1404 T_{B} & =0-594 \\
988 T_{B} & =594 \\
T_{B} & =\frac{594}{988}=0.601
\end{aligned}
$$

Substituting the value for $T_{B}$ back into either expression to find $T_{A}$ yields

$$
\begin{aligned}
32 T_{A}+52(0.601) & =22, \\
T_{A} & =-0.289
\end{aligned}
$$

The two equations together result in $T_{A}=-0.29$ and $T_{B}=0.60$.
EVALUATE: An alternative solution would be to write $T_{B}$ in terms of $T_{A}$, using the first expression, and then substitute for $T_{B}$ into the second expression. This gives the same result. You may choose the method you prefer and may end up applying both to particular classes of problems.

If we encounter three unknowns in an expression, how many equations will we need to solve for each unknown simultaneously? Three equations will be needed to solve for the three unknown quantities.

## 8: Review of the quadratic formula

The position of a ball tossed in the air depends on the initial speed of the ball and the time elapsed and is given by

$$
y=v_{0} t-\left(4.9 \mathrm{~m} / \mathrm{s}^{2}\right) t^{2}
$$

where $v_{0}$ is the initial speed and $t$ is the time elapsed. For a ball tossed with an initial speed of $30.0 \mathrm{~m} / \mathrm{s}$, find the time(s) when the ball is at a height of 12.5 m .

## Solution

IDENTIFY AND SET UP: We recognize that the equation is quadratic, since it has a $t^{2}$ term, a $t$ term, and a constant term. If we try to rewrite the equation in terms of $t$ alone, we find that we cannot easily isolate the $t$ term. We will employ the quadratic formula to solve the problem.
EXECUTE: We rewrite the equation, substituting the given values:

$$
12.5 \mathrm{~m}=(30.0 \mathrm{~m} / \mathrm{s}) t-\left(4.9 \mathrm{~m} / \mathrm{s}^{2}\right) t^{2}
$$

The quadratic formula requires that the equation be written as $a x^{2}+b x+c=0$, so we rearrange terms to yield

$$
\left(-4.9 \mathrm{~m} / \mathrm{s}^{2}\right) t^{2}+(30.0 \mathrm{~m} / \mathrm{s}) t+(-12.5 \mathrm{~m})=0
$$

From this rearrangement, we see that $a=-4.9 \mathrm{~m} / \mathrm{s}^{2}, b=30.0 \mathrm{~m} / \mathrm{s}$, and $c=-12.5 \mathrm{~m}$. The solutions of the quadratic equation are

$$
x=\frac{-b \pm \sqrt{b^{2}-4 a c}}{2 a}
$$

Substituting our values into the quadratic equation gives

$$
t=\frac{-(30.0 \mathrm{~m} / \mathrm{s}) \pm \sqrt{(30.0 \mathrm{~m} / \mathrm{s})^{2}-4\left(-4.9 \mathrm{~m} / \mathrm{s}^{2}\right)(-1.25 \mathrm{~m})}}{2\left(-4.9 \mathrm{~m} / \mathrm{s}^{2}\right)}
$$

Multiplying out the terms and canceling the units produces

$$
t=\frac{-(30.0 \mathrm{mI} / \mathrm{s}) \pm \sqrt{(900.0-245.0)\left(\mathrm{m}^{2} / \mathrm{s}^{2}\right)}}{-9.8\left(\mathrm{~m} / \mathrm{s}^{\not x}\right)}=\frac{-30.0 \pm 25.59}{-9.8} \mathrm{~s}=0.445 \mathrm{~s}, 5.67 \mathrm{~s}
$$

There are two times when the ball is at a height of $12.5 \mathrm{~m}: 0.445 \mathrm{~s}$ and 5.67 s . These are, respectively, when the ball is rising to its maximum height and when it is falling from its maximum height.

EVALUATE: You must learn to recognize quadratic equations. Once you identify a quadratic equation, the solution is straightforward (although it requires careful algebra). Quadratic equations result in two solutions, and you must be able to interpret their meanings. In this case, the two solutions corresponded to the upward and downward motion of the ball. You may need only one of the solutions for your situation. If neither solution seems reasonable, then you should check your work. We'll encounter quadratic equation problems again in Chapter 2.

CAUTION Watch units! It is important to check units every time you write equations. If we found incorrect units when we solved for time, we would have discovered a mistake that would have been quickly corrected.

## Problem Summary

The problems in this chapter represent a foundation that you will use throughout your physics course. Common elements make up good problem-solving techniques, including

- Identifying a procedure to find the solution.
- Making a sketch when no figure is provided.
- Adding appropriate coordinate systems to the sketch.
- Identifying the known and unknown quantities in the problem.
- Finding appropriate equations to solve for the unknown quantities.
- Checking for consistency of units in derived equations.
- Evaluating results to check for inconsistencies.

We will see how these techniques apply to a wide variety of problems as we progress. Although they may seem cumbersome right now, they will help you solve the problems you encounter.

# Motion along a Straight Line 

## Summary

We will introduce kinematics, the study of an object's motion, or change of position with time, in this chapter. Motion includes displacement, the change in position of an object; velocity, the rate of change of position with respect to time; and acceleration, the rate of change of velocity with respect to time. We introduce average velocity and average acceleration as changes over a time interval and instantaneous velocity and instantaneous acceleration as changes over an infinitely short time interval. We'll learn relationships between displacement, velocity, and acceleration and see how they are modified for freely falling objects. We'll restrict ourselves to motion along a straight line, or one-dimensional motion, in this chapter and expand our examination to motion in two or three dimensions in the next chapter. This is our first step into understanding mechanics, the study of the relationships among force, matter, and motion that we'll cover in the upcoming chapters.

## Objectives

After studying this chapter, you will understand

- The definitions of kinematic variables for position, velocity, and acceleration.
- How to calculate and interpret average and instantaneous velocities.
- How to calculate and interpret average and instantaneous accelerations.
- How to apply the equations of motion for constant acceleration.
- How to apply equations of constant acceleration to freely falling objects.
- How to analyze motion when acceleration is not constant.

| Term Description |  |
| :---: | :---: |
| Average Velocity | A particle's average $x$-velocity $v_{\mathrm{a} v-x}$ over a time interval $\Delta t$ is its displacement $\Delta x$ divided by the time interval $\Delta t$ : $v_{a v-x}=\frac{x_{2}-x_{1}}{t_{2}-t_{1}}=\frac{\Delta x}{\Delta t} .$ <br> The SI unit of velocity is meters per second ( $\mathrm{m} / \mathrm{s}$ ). |
| Instantaneous Velocity | A particle's instantaneous velocity is the limit of the average velocity as $\Delta t$ goes to zero, or the derivative of position with respect to time. The $x$ component is defined as $v_{x}=\lim _{t \rightarrow \infty} \frac{\Delta x}{\Delta t}=\frac{d x}{d t} .$ <br> The term velocity refers to the instantaneous velocity. |
| Average Acceleration | The average $x$ acceleration of a particle over a time interval $\Delta t$ is the change in the $x$ component of velocity, $\Delta v_{x}=v_{2 x}-v_{1 x}$, divided by the time interval $\Delta t$ : $a_{\mathrm{av}-x}=\frac{v_{2 x}-v_{1 x}}{t_{2}-t_{1}}=\frac{\Delta v_{x}}{\Delta t} .$ <br> The SI unit of acceleration is meters per second per second $\left(\mathrm{m} / \mathrm{s}^{2}\right)$. |
| Instantaneous Acceleration | A particle's instantaneous acceleration is the limit of the average acceleration as $\Delta t$ goes to zero, or the derivative of the velocity with respect to time. The $x$ component is defined as $a_{x}=\lim _{t \rightarrow \infty} \frac{\Delta v_{x}}{\Delta t}=\frac{d v_{x}}{d t} .$ <br> The term acceleration refers to the instantaneous acceleration. |
| Motion with Constant Acceleration | When the $x$ acceleration is constant, position, $x$ velocity, acceleration, and time are related by $\begin{aligned} x & =x_{0}+v_{0 x} t+\frac{1}{2} a_{x} t^{2} \\ v_{x} & =v_{0 x}+a_{x} t \\ v_{x}^{2} & =v_{0 x}^{2}+2 a_{x}\left(x-x_{0}\right) \\ x-x_{0} & =\left(\frac{v_{0 x}+v_{x}}{2}\right) t . \end{aligned}$ |
| Freely Falling Body | A freely falling body is a body that moves under the influence of the gravity. The acceleration due to gravity is denoted by $g$, is directed downwards, and has a value of $9.8 \mathrm{~m} / \mathrm{s}^{2}$ near the surface of the earth. |
| Motion with Varying Acceleration | When the acceleration is not constant, we can find the position and velocity as a function of time by integrating the acceleration function: $\begin{aligned} x & =x_{0}+\int_{0}^{t} v_{x} d t \\ v_{x} & =v_{0 x}+\int_{0}^{t} a_{x} d t . \end{aligned}$ |

## Conceptual Questions

## 1: Velocity and acceleration at the top of a ball's path

A ball is tossed vertically upward. (a) Describe the velocity and acceleration of the ball just before it reaches the top of its flight. (b) Describe the velocity and acceleration of the ball at the instant it reaches the top of its flight. (c) Describe the velocity and acceleration of the ball just after it reaches the top of its flight.


Figure 2.1 Question 1.

## Solution

IDENTIFY, SET UP, AND EXECUTE: Figure 2.1 shows the three time frames we will examine. During its flight, the ball undergoes acceleration due to gravity. The initial velocity is directed upward, slowing to zero at the top of the flight. Then the velocity increases downward.

PART (A): The velocity is directed upward and is very small just before the top of the flight. The acceleration due to gravity is directed downward.

PART (B): The velocity is zero at the top of the flight. The acceleration due to gravity remains constant and is directed downward. The acceleration has caused the velocity to decrease from its small positive value in part (a) to zero.

PART (C): The velocity is directed downward and is very small just after the top of the flight. The acceleration due to gravity remains constant and directed downward. The acceleration has caused the velocity to increase downward from zero in part (b).

EVALUATE: The acceleration due to gravity causes a change in the velocity of the ball during its flight. The ball starts with an upward velocity, which slows, drops to zero, and then increases downward. The acceleration due to gravity is constant throughout the motion. The velocity is zero for an instant at the top, changing from slightly upward to slightly downward around this instant.

CAUTION Gravity doesn't turn off! A common misconception is that there is no acceleration due to gravity at the top of an object's trajectory, but how would gravity know to turn off at that instant?

## 2: Comparing two cyclists

The position-versus-time graph depicting the paths of two cyclists is shown in Figure 2.2. (a) Do the cyclists start from the same position? (b) Are there any times that they have the same velocity? (c) What is happening at the intersection of lines $A$ and $B$ ?


Figure 2.2 Question 2.

## Solution

IDENTIFY, SET UP, AND EXECUTE PART (A): We find the starting location by examining the position when time is zero (i.e., by looking at the $x$-intercept). At $t=0$, the two cyclists are at different locations.

PART (B): The velocity is found by examining the slope of the position-versus-time graph. The slopes of the two lines are different; hence, the cyclists never have the same velocity.

PART (C): At the intersection of lines $A$ and $B$, both cyclists are at the same position at the same time. At this point, cyclist $A$ is passing cyclist $B$, since cyclist $A$ started closer to the origin and has a higher velocity.

EVALUATE: These three questions show only a small part of what can be learned from graphs, which offer a parallel representation of physical phenomena. The interpretation of graphs is an important tool in physics and, indeed, science in general.

## 3: Interpreting a position-versus-time graph

Figure 2.3 shows a position-versus-time graph of the motion of a car. Describe the velocity and acceleration during segments $O A, A B$, and $B C$.


Figure 2.3 Question 3.

## Solution

IDENTIFY, SET UP, AND EXECUTE: Velocity is the change in position with respect to time and acceleration is the change in velocity with respect to time. We can describe the velocity by examining the slope of the position-versus-time graph, and we can describe the acceleration by noting how the velocity changes.

In segment $O A$, the slope is positive and constant, indicating that the velocity is positive and constant. With constant velocity, there is no acceleration.

In segment $A B$, the slope is increasing smoothly, indicating that the velocity is increasing. There must be acceleration in order for the velocity to increase.

In segment $B C$, the slope is again constant, indicating that the velocity is constant. This velocity is greater in magnitude than the velocity in segment $O A$, since the slope is larger. With constant velocity, there is no acceleration.

EVALUATE: This question illustrates how we can describe the velocity and acceleration from the position-versus-time graph.

## 4: A falling ball

A ball falls from the top of a building. If the ball takes time $t_{A}$ to fall halfway from the top of the building to the ground, is the time it takes to fall the remaining distance to the ground equal to, greater than, or smaller than $t_{A}$ ?

## Solution

IDENTIFY, SET UP, AND EXECUTE: We can break the problem up into two segments: the first half and the second half. In the first segment, the falling ball starts with an initial velocity of zero. In the second segment, the ball has acquired velocity, so it has an initial velocity. The time to complete the second segment must be shorter than $t_{A}$.

EVALUATE: If you watch a ball fall, you should be able to see that it spends more time in the first half of the motion than in the second half. We can also look at the equation for the position of a falling body:

$$
y=y_{0}+v_{0 y} t+\frac{1}{2} a_{y} t^{2}
$$

For the first half of the motion, the velocity term is zero; for the second half, it is not zero. Given equal time and equal acceleration, a segment with an initial velocity will cover a larger distance, or cover the same distance in a shorter time.

## Problems

## 1: Throwing a ball upward

Robert throws a ball vertically upward from the edge of a $150-\mathrm{m}$-tall building. The ball falls to the ground 9.5 s after leaving Robert's hand. Assume that the ball leaves Robert's hand when it is 2.0 m above the roof of the building. Find the initial velocity of the ball and the time the ball reaches its maximum height.


Figure 2.4 Problem 1.

## Solution

IDENTIFY: The ball undergoes constant acceleration due to gravity, so we will use the constantacceleration kinematics equation to solve the problem.

SET UP: Figure 2.4 shows a sketch of the problem. Once thrown, the ball has an initial upward velocity and will undergo downward gravitational acceleration.

We ignore effects due to the air. A vertical coordinate system is shown in the diagram, with the origin located at the edge of the building and positive values directed upward.

EXECUTE: We first determine the initial velocity of the ball. We know the initial and final positions, times, and accelerations of the ball; therefore, we use the equation for position as a function of time:

$$
y=y_{0}+v_{0 y} t+\frac{1}{2} a_{y} t^{2}
$$

The initial position of the ball $\left(y_{0}\right)$ is +2.0 m , the final position $(y)$ is -150 m (the ground is below the edge of the building), the acceleration is $-g$, and the time is 9.5 s . Solving for the initial velocity $v_{0 y}$ gives

$$
v_{0 y}=\frac{y-y_{0}-\frac{1}{2} a_{y} t^{2}}{t}
$$

Substituting the given values yields

$$
v_{0 y}=\frac{(-150 \mathrm{~m})-(2.0 \mathrm{~m})-\frac{1}{2}\left(-9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(6.5 \mathrm{~s})^{2}}{(6.5 \mathrm{~s})}=8.5 \mathrm{~m} / \mathrm{s}
$$

The initial velocity of the ball is $8.5 \mathrm{~m} / \mathrm{s}$. The value is positive, indicating that the initial velocity is directed upward. To find the time taken to reach the maximum height, we know that the velocity at that height is momentarily zero, so we can use the equation for velocity as a function of time:

$$
v_{y}=v_{0 y}+a_{y} t .
$$

We now solve for the time $t$ when the velocity $v_{y}$ is zero

$$
t=\frac{v_{y}-v_{0 y}}{a_{y}}=\frac{(0-8.5 \mathrm{~m} / \mathrm{s})}{\left(-9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}=0.87 \mathrm{~s}
$$

The ball reaches its maximum height 0.87 s after leaving Robert's hand.
EVALUATE: This is a straightforward application of constant-acceleration kinematics. We identified the known and unknown quantities and substituted into appropriate equations to find the unknown quantities.

PRACTICE PROBLEM: Find the maximum height $y_{\max }$ of the ball. Answer: $y_{\max }=5.7 \mathrm{~m}$ above the top of the building.

## 2: Dropping a stone from a moving helicopter

A helicopter is ascending at a constant rate of $18 \mathrm{~m} / \mathrm{s}$. A stone falls from the helicopter 12 s after it leaves the ground. How long does it take for the stone to reach the ground?


Figure 2.5 Problem 2.

## Solution

IDENTIFY: There are two segments of the stone's motion: (1) moving upward with the helicopter at constant velocity and (2) free fall after the stone breaks free of the helicopter. To solve the problem, we
will apply the constant-acceleration kinematics equations to the two segments, using the final quantities from the first segment as the initial quantities in the second.

SET UP: The two segments of the stone's motion are sketched in Figure 2.5. As it moves upward with the helicopter, the stone has a constant velocity $v_{y 0}$. When it breaks loose and begins to fall freely, the stone has an initial velocity that is the same as the helicopter's and undergoes acceleration due to gravity. We need to know the position, velocity, and time at the end of the first segment to solve for the second segment. The velocity and time are given in the statement of the problem.

We ignore effects due to the air. A vertical coordinate system is shown in the diagram, with the origin located on the ground and positive values directed upward.

EXECUTE: The position is found from the equation for position as a function of time with zero acceleration:

$$
y=y_{0}+v_{0 y} t
$$

The initial position is zero (the helicopter starts at the ground) and the helicopter is ascending, so $v_{y 0}$ is $+18 \mathrm{~m} / \mathrm{s}$ and the time is 12 s . Substituting yields

$$
y=y_{0}+v_{0 y} t=0+(18 \mathrm{~m} / \mathrm{s})(12 \mathrm{~s})=216 \mathrm{~m}
$$

For the second segment, the initial position is 216 m , the initial velocity is $+18 \mathrm{~m} / \mathrm{s}$, the final position is zero, and the acceleration is $-g$. The equation for position as a function of time with constant acceleration can be used to find the time:

$$
y=y_{0}+v_{0 y} t+\frac{1}{2} a_{y} t^{2} .
$$

Substituting values gives

$$
0=(216 \mathrm{~m})+(18 \mathrm{~m} / \mathrm{s}) t+\frac{1}{2}\left(-9.8 \mathrm{~m} / \mathrm{s}^{2}\right) t^{2}
$$

We cannot solve this equation directly for $t$, so we resort to the quadratic equation. In this case, $a=-4.9 \mathrm{~m} / \mathrm{s}^{2}, b=18 \mathrm{~m} / \mathrm{s}$, and $c=216 \mathrm{~m}$. The result is given by

$$
t=\frac{-b \pm \sqrt{b^{2}-4 a c}}{2 a}
$$

Substituting and solving yields

$$
t=\frac{-(18 \mathrm{~m} / \mathrm{s}) \pm \sqrt{(18 \mathrm{~m} / \mathrm{s})^{2}-4\left(-4.9 \mathrm{~m} / \mathrm{s}^{2}\right)(216 \mathrm{~m})}}{2\left(-4.9 \mathrm{~m} / \mathrm{s}^{2}\right)}=-5.1 \mathrm{~s},+8.7 \mathrm{~s}
$$

The positive solution, 8.7 s , corresponds to the time the stone hits the ground. The stone hits the ground 8.7 s after falling from the helicopter, or 20.7 s after the helicopter originally left the ground.

EVALUATE: We applied the equations for motion with constant acceleration to each of the two segments in this problem, using the results from the first part as input into the second part. The negative solution of the quadratic equation corresponds to the time the stone would have left the ground, assuming that it was thrown from the ground. Because this aspect of the motion doesn't apply to our problem, we ignore that solution.

## 3: Avoiding a ticket

A speed trap is made by placing two pressure-sensitive tracks across a highway, 150 m apart. Suppose you are driving and you notice the speed trap and begin slowing down the instant you cross the first track. If you are moving at a rate of $42 \mathrm{~m} / \mathrm{s}$ and the speed limit is $35 \mathrm{~m} / \mathrm{s}$, what must your minimum acceleration be in order for you to have an average speed within the speed limit by the time your car crosses the second track?

## Solution

IDENTIFY: For the average speed over the interval to be under the speed limit, the final speed at the second track must be less than the speed limit. We will use the kinematic equations to find the acceleration necessary to avoid a speeding ticket.

SET UP: A sketch of the problem is shown in Figure 2.6. We determine the final speed by writing the average speed in terms of the initial and final speeds, setting the average speed to the speed limit, and solving for the final speed. Once we determine the final speed, we find the acceleration from the kinematics equations.


Figure 2.6 Problem 3.
EXECUTE: The average speed (for constant acceleration) is

$$
v_{\mathrm{av}, x}=\frac{v_{0 x}+v_{x}}{2}
$$

Substituting and solving for the final speed gives

$$
v_{x}=2_{\mathrm{av}, x}-v_{0 x}=2(35 \mathrm{~m} / \mathrm{s})-(42 \mathrm{~m} / \mathrm{s})=28 \mathrm{~m} / \mathrm{s} .
$$

The final speed must be $28 \mathrm{~m} / \mathrm{s}$ in order for the average speed to be $35 \mathrm{~m} / \mathrm{s}$. We use the equation for velocity as a function of position with constant acceleration, or

$$
v_{x}^{2}=v_{0 x}^{2}+2 a_{x}\left(x-x_{0}\right)
$$

In our coordinate system, the difference between the final and initial positions is 150 m . Substituting and solving for the acceleration gives

$$
a_{x}=\frac{v_{x}^{2}-v_{0 x}^{2}}{2\left(x-x_{0}\right)}=\frac{(28 \mathrm{~m} / \mathrm{s})^{2}-(42 \mathrm{~m} / \mathrm{s})^{2}}{2(150 \mathrm{~m})}=-3.3 \mathrm{~m} / \mathrm{s}^{2}
$$

You will need to accelerate at a rate of $-3.3 \mathrm{~m} / \mathrm{s}^{2}$ to avoid a ticket, with the minus indicating that you will need to slow down.

EVALUATE: The challenge in this problem was to recognize that we needed a final velocity that would result in the correct average velocity. A common mistake is to take the desired average velocity as the final velocity. Understanding the difference can help you avoid errors (and a ticket!).

## 4: Graphical solution to an accelerating car

A car undergoing constant acceleration moves 250 m in 8.5 s . If the speed at the end of the 250 m segment is $33 \mathrm{~m} / \mathrm{s}$, what was the car's speed at the beginning of the segment?

## Solution

IDENTIFY: We can approach this problem in two ways. First, we can use the kinematics equations to solve the problem, but doing so will require several equations. Second, we can use the velocity-versustime plot and solve the problem graphically. We will choose the graphical method in this case.

SET UP: A sketch of the problem is shown in Figure 2.7. We realize that there is no single kinematics equation that ties these quantities together, so we construct the velocity-versus-time graph. The graph must start with initial velocity $v_{0}$ and result in final velocity $v_{1}$. The slope of the line between the two velocities must be constant, since the car exhibits constant acceleration. The time interval between the two velocities must be the given 8.5 s , so we construct the graph shown in Figure 2.8.


Figure 2.7 Problem 4 sketch.


Figure 2.8 Problem 4 velocity-
vs.-time graph.
Examining the graph, we realize that the area under the velocity line is the distance the car travels. We can therefore find the initial velocity by calculating the area under the curve.

EXECUTE: The area under the curve is the sum of the area of the rectangle and the area of the triangle shown in the figure. We find these areas by multiplying the time by the velocities:

$$
\text { Area }=v_{0} t+\frac{1}{2}\left(v_{1}-v_{0}\right) t
$$

The area under the curve is just the distance the car travels, 250 m . We rewrite the equation to solve for the initial velocity $v_{0}$ :

$$
v_{0}=\frac{2 \mathrm{Area}-v_{1} t}{t}=\frac{2((250 \mathrm{~m})-(33 \mathrm{~m} / \mathrm{s})(8.5 \mathrm{~s}))}{(8.5 \mathrm{~s})}=25.8 \mathrm{~m} / \mathrm{s}
$$

The initial velocity is $25.8 \mathrm{~m} / \mathrm{s}$.

EVALUATE: The example illustrates how graphical analysis can lead to a straightforward solution. The key was to realize that the area under the curve is the distance the car traveled, or its displacement. If we were to solve the problem with kinematics equations, we would have had two unknowns (initial velocity and acceleration), requiring us to utilize two equations in the solution.

## 5: Two objects falling from a building

A ball is dropped from the top of a tall building. One second later, another ball is thrown from the top of the building with a velocity of $30 \mathrm{~m} / \mathrm{s}$ directed vertically downwards. Will the balls ever meet? If so, when and where?

## Solution

IDENTIFY: Both balls undergo acceleration due to gravity after being dropped or thrown, so we will apply constant-acceleration kinematics. We will use separate sets of kinematic equations for the two balls; ball 1 will be the dropped ball and ball 2 will be the thrown ball.

SET UP: A sketch of the problem is shown in Figure 2.9. We are interested in where the balls meet, so we will use the position equations, setting their positions equal to each other to find out whether they meet at any point in time. The coordinate system is shown in the sketch, with the origin at the top of the building and positive values directed downward.


Figure 2.9 Problem 5.

EXECUTE: The equation for the position as a function of time for ball 1 , the dropped ball, is

$$
x_{1}=x_{01}+v_{01} t+\frac{1}{2} a_{x} t^{2}=\frac{1}{2} g t^{2} .
$$

The equation for the position as a function of time for ball 2 , the thrown ball, is

$$
x_{2}=x_{02}+v_{02} t^{\prime}+\frac{1}{2} a_{x} t^{\prime 2}=v_{02}(t-1 \mathrm{~s})+\frac{1}{2} a_{x}(t-1 \mathrm{~s})^{2}
$$

where we have included the initial velocity $v_{02}$ and the replaced the time $t^{\prime}$ with $(t-1 \mathrm{~s})$. For the two balls to meet, the two positions must be the same. We set the two equations equal to each other and solve for the time they meet:

$$
\begin{aligned}
\frac{1}{2} a_{x} t^{2} & =v_{02}(t-1 \mathrm{~s})+\frac{1}{2} a_{x}(t-1 \mathrm{~s})^{2} \\
\frac{1}{2} g t^{2} & =v_{02} t-v_{02} 1 \mathrm{~s}+\frac{1}{2} g t^{2}+\frac{1}{2} g(-2 t \mathrm{~s})+\frac{1}{2} g(1 \mathrm{~s})^{2}
\end{aligned}
$$

The $t^{2}$ term cancels, leaving

$$
0=v_{02} t-v_{02} 1 \mathrm{~s}+\frac{1}{2} g(-2 t \mathrm{~s})+\frac{1}{2} g(1 \mathrm{~s})^{2}
$$

Solving for $t$ gives

$$
t=\frac{\frac{1}{2} \mathrm{~g}(\mathrm{~s})-v_{02}}{g(\mathrm{~s})-v_{02}}=\frac{\frac{1}{2}\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(\mathrm{s})-(30 \mathrm{~m} / \mathrm{s})}{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(\mathrm{s})-(30 \mathrm{~m} / \mathrm{s})}=1.24 \mathrm{~s}
$$

The balls meet 1.24 s after the first ball is dropped. The position of the balls at this time is

$$
x_{1}=\frac{1}{2} g t^{2}=\frac{1}{2}\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(1.24 \mathrm{~s})^{2}=7.57 \mathrm{~m}
$$

The balls meet 7.57 m below the top of the building.
EVALUATE: We check our results and see that the balls meet 1.24 s after the first ball is dropped, or 0.24 s after the second ball is thrown. Since the balls meet after the second ball is thrown, we conclude these times represent a reasonable result. They meet 7.57 m below the top of the building. (A positive value indicates that they meet below the top of the building.) We also see that the building must be at least 7.57 m tall.

## 6: Nonconstant acceleration

A particle has an initial velocity of $12.0 \mathrm{~m} / \mathrm{s}$ and starts at $x=14.2 \mathrm{~m}$. It moves along the $x$ axis and possesses an acceleration given by

$$
a_{x}=b t^{2}
$$

where $b$ is a constant equal to $3.5 \mathrm{~m} / \mathrm{s}^{4}$. Find the particle's velocity and position as a function of time.

## Solution

IDENTIFY AND SET UP: The acceleration varies with time, so we cannot use the constantacceleration kinematic equations. Instead we integrate the acceleration to find the velocity as a function of time and we integrate the velocity to find the position as a function of time. Both the initial position and initial velocity are zero.

EXECUTE: We begin by integrating the acceleration to find the velocity:

$$
v_{x}=v_{0 x}+\int_{0}^{t} a_{x} d t
$$

For our problem, $v_{0 x}$ is $12.0 \mathrm{~m} / \mathrm{s}$ and $a_{x}$ is given. Substituting and solving, we obtain

$$
\begin{aligned}
v_{x} & =(12.0 \mathrm{~m} / \mathrm{s})+\int_{0}^{t} b t^{2} d t \\
& =(12.0 \mathrm{~m} / \mathrm{s})+\left.\frac{1}{3} b t^{3}\right|_{0} ^{t} \\
& =(12.0 \mathrm{~m} / \mathrm{s})+\frac{1}{3} b t^{3}
\end{aligned}
$$

To find the position as a function of time, we integrate the velocity:

$$
x=x_{0}+\int_{0}^{t} v_{x} d t
$$

Substituting $x_{0}=14.2 \mathrm{~m}$ as given and the value of $v_{x}$ that we found and solving yields

$$
\begin{aligned}
v_{x} & =\int_{0}^{t}\left[(12.0 \mathrm{~m} / \mathrm{s})+\frac{1}{3} b t^{3}\right] d t \\
& \left.=(14.2 \mathrm{~m})+\left[(12.0 \mathrm{~m} / \mathrm{s}) \mathrm{t}+\frac{1}{12} b t^{4}\right]\right]_{0}^{t} \\
& =(14.2 \mathrm{~m})+(12.0 \mathrm{~m} / \mathrm{s}) \mathrm{t}+\frac{1}{12} b t^{4} .
\end{aligned}
$$

We have found both the position and velocity as a function of time.
EVALUATE: We can check our integration by taking derivatives of our results. When we do, we find the original acceleration.

## Try It Yourself!

Learning physics requires that you practice problems without having solutions next to the problem. To help you prepare for homework problems and exams, we have included sample problems with checkpoints to help you through them. We encourage you to try these problems on your own and refer to the checkpoints only when you get stuck. So go ahead and Try It Yourself!

## 1: Police chase

A speeder traveling at a constant speed of $100 \mathrm{~km} / \mathrm{hr}$ passes a waiting police car that immediately starts from rest and accelerates at a constant $2.5 \mathrm{~m} / \mathrm{s}^{2}$. (a) How long will it take for the police car to catch the speeder? (b) How fast will the police car be traveling when it catches the speeder? (c) How far will the police car have traveled when it catches the speeder?

## Solution Checkpoints

IDENTIFY AND SET UP: Constant-acceleration kinematics are appropriate in this problem. Two separate sets of kinematics equations should be used to represent the police car and the speeder. We set appropriate values equal to each other to solve the problem.

EXECUTE: The positions of the police car and speeder must be the same when the police car catches the speeder:

$$
v_{0, \text { speeder }} t=\frac{1}{2} a t^{2} .
$$

This leads to the conclusion that the speeder is caught 22.2 s after passing the police car. Kinematics then indicates that the police car had a velocity of $200 \mathrm{~km} / \mathrm{hr}$ and a position of 616 m .

EVALUATE: Would you expect the police car to accelerate at a constant rate or the speeder not to slow down? How would these changes affect the result? $t=0 \mathrm{~s}$ is also a solution of the equation. To what event does $t=0 \mathrm{~s}$ correspond?

## 2: Don't hit the truck

A car traveling $100 \mathrm{~km} / \mathrm{hr}$ is 200 m away from a truck traveling $50 \mathrm{~km} / \mathrm{hr}$ in the same direction. What minimum acceleration must the car have in order to avoid hitting the truck? Assume constant acceleration during braking.

## Solution Checkpoints

IDENTIFY AND SET UP: Constant-acceleration kinematics are valid in this problem. Two sets of kinematics equations should be used to represent the car and truck separately. Choose an appropriate coordinate system.

EXECUTE: To avoid the collision with the minimum acceleration, the car and truck will meet at the same point at the same time and with the same velocity. Setting the car's and truck's position equations equal to each other gives

$$
v_{0 \text { car }} t+\frac{1}{2} a t^{2}=x_{0 \text { truck }}+v_{0 \text { truck }} t
$$

This equation has two unknowns, so you will need to set the velocities equal to each other and solve by using both relations. You should find that the magnitude of the acceleration is $0.48 \mathrm{~m} / \mathrm{s}^{2}$.

EVALUATE: What is the sign of the acceleration you found? Is it what you would expect for a car slowing down?

## 3: Throwing a ball upward

A ball is thrown vertically upward from a $125-\mathrm{m}$-high building with an initial velocity of $45 \mathrm{~m} / \mathrm{s}$. (a) What is the ball's maximum height? (b) What is its velocity as it passes the top of the building on its way down? (c) How long does it take the ball to reach the ground?

## Solution Checkpoints

IDENTIFY AND SET UP: Constant-acceleration kinematics are valid in this problem. You will need equations for position as a function of time and velocity as a function of time. Choose an appropriate coordinate system.

EXECUTE (A): At the maximum height, the velocity is zero. Solving will give a height of 103 m above the building.
(b) On the way down, the velocity of the ball when it passes the top of the building will have the same magnitude as the initial velocity, but will be opposite in direction.
(c) The equation for position as a function of time can be used to find the time taken for the ball to hit the ground. The position equation will lead to a quadratic equation and a result of 11.4 s .

EVALUATE: The quadratic equation of part (c) had two roots. Why did you omit one root? What is the physical interpretation of the omitted root?

## Motion in Two or Three Dimensions

## Summary

In this chapter we expand our kinematics to motion of bodies in two or three dimensions. In doing so, we will find that we can simultaneously apply our one-dimensional kinematics equations to multiple axes independently. Displacement, velocity, and acceleration take on their vector qualities as we expand to more dimensions, requiring us to work with components of each quantity. Our new skills will allow us to investigate projectile motion and the interesting case of uniform circular motion. We will also learn to analyze motion viewed from different moving reference frames. By the end of this chapter, we will have laid a strong foundation in kinematics and will be ready to investigate the causes of motion.

## Objectives

After studying this chapter, you will understand

- How to describe a body's position, velocity, and acceleration in terms of vector quantities.
- How to apply equations of motion to bodies moving in a plane.
- How to describe and analyze the motion of projectiles.
- How to analyze an object in uniform circular motion.
- How to combine components of acceleration that are parallel and perpendicular to a body's path.
- How to relate the velocities of objects to different reference frames.

Concepts and Equations

| Term | Description |
| :--- | :--- |
| Position Vector | The position vector $\vec{r}$ of a point $P$ in space is the displacement vector from <br> the origin to $P$. It has components $x, y$, and $z$. |
| Average Velocity | The average velocity $\vec{v}_{\mathrm{av}}$ during a time interval $\Delta t$ is the displacement $\Delta \vec{r}$ <br> divided by $\Delta t:$ |

$\xrightarrow[\text { Instantaneous Velocity }]{ }$

## Average Acceleration

A body's instantaneous velocity is the derivative of $\vec{r}$ with respect to time:

$$
\vec{v}=\lim _{t \rightarrow \infty} \frac{\Delta \vec{r}}{\Delta t}=\frac{d \vec{r}}{d t} .
$$

The instantaneous velocity has components

$$
v_{x}=\frac{d x}{d t}, \quad v_{y}=\frac{d y}{d t}, \quad v_{z}=\frac{d z}{d t} .
$$

The average acceleration $\vec{a}_{\mathrm{av}}$ during a time interval $\Delta t$ is the change in velocity $\Delta \vec{v}$ divided by $\Delta t$ :

$$
\vec{a}_{\mathrm{av}}=\frac{\vec{v}_{2}-\vec{v}_{1}}{t_{2}-t_{1}}=\frac{\Delta \vec{v}}{\Delta t}
$$

## Instantaneous Acceleration

A body's instantaneous acceleration is the derivative of $\vec{v}$ with respect to time:

$$
\vec{a}=\lim _{i \rightarrow \infty} \frac{\Delta \vec{v}}{\Delta t}=\frac{d \vec{v}}{d t} .
$$

The instantaneous velocity has components

$$
a_{x}=\frac{d v_{x}}{d t}, \quad a_{y}=\frac{d v_{y}}{d t}, \quad a_{z}=\frac{d v_{z}}{d t} .
$$

The component of acceleration parallel to the velocity affects the speed of the body. The component of acceleration perpendicular to the velocity affects the body's direction of motion. A body has acceleration if either its speed or direction changes.

## Projectile Motion

A body undergoes projectile motion when it is given an initial velocity and then follows a path determined entirely by the effect of a constant gravitational force. The path, or trajectory, is a parabola. The projectile's vertical motion is independent of its horizontal motion. The horizontal acceleration is zero and the vertical acceleration is $-g$. The coordinates and velocities of a projectile with an initial velocity of magnitude $v_{0}$ and direction $\alpha_{0}$ (measured with respect to the ground) are given as a function of time by

$$
\begin{aligned}
x & =\left(v_{0} \cos \alpha_{0}\right) t, \\
y & =\left(v_{0} \sin \alpha_{0}\right) t-\frac{1}{2} g t^{2}, \\
v_{x} & =v_{0} \cos \alpha_{0}, \\
v_{y} & =v_{0} \sin \alpha_{0}-g t .
\end{aligned}
$$

| Uniform Circular Motion | A particle moving in a circular path of radius $R$ and constant speed $v$ is said to move in uniform circular motion. The particle possesses an acceleration of magnitude $a_{\mathrm{rad}}=\frac{v^{2}}{R}$ <br> directed toward the center of the circle. If the particle's speed is not constant in circular motion, then the radial component of acceleration remains as just given and there is also a component parallel to the path of the particle. |
| :---: | :---: |
| Relative Velocity | When a body $P$ moves relative to a reference frame $B$, and $B$ moves relative to a second reference frame $A$, the velocity of $P$ relative to $B$ is denoted by $\vec{v}_{P \mid B}$, the velocity of $P$ relative to $A$ is denoted by $\vec{v}_{P / A}$, and the velocity of $B$ relative to $A$ is denoted by $\vec{v}_{B A A}$. These velocities are related by $\vec{v}_{P / A}=\vec{v}_{P / B}+\vec{v}_{B / A} .$ |

## Conceptual Questions

## 1: Velocity and acceleration at the top of a projectile's path

A projectile is launched with initial nonzero $x$ and $y$ velocities. (a) Describe the velocity and acceleration just before the projectile reaches the top of its trajectory. (b) Describe the velocity and acceleration at the instant the projectile reaches the top of its trajectory. (c) Describe the velocity and acceleration just after the projectile reaches the top of its trajectory.


Figure 3.1 Question 1.

## Solution

IDENTIFY AND SET UP: Figure 3.1 shows the three time frames we will examine. During the flight, the projectile undergoes acceleration due to gravity. The projectile's initial velocity has both $x$ and $y$ components; the $x$ component remains constant while the $y$ component is accelerated by gravity. We have to consider each component of velocity separately.

EXECUTE PART (A): The $x$ component of velocity is constant and directed to the right, and the $y$ component of velocity is directed upward and is very small just before the top of the flight. The acceleration due to gravity is directed downward.

PART (B): The $x$ component of velocity is constant and directed to the right, and the $y$ component of velocity is zero at the top of the flight. The acceleration due to gravity remains constant and directed downward.

PART (C): The $x$ component of velocity is constant and directed to the right, and the $y$ component of velocity is directed downward and is very small just after the top of the flight. The acceleration due to gravity remains constant and directed downward.

EVALUATE: The acceleration due to gravity causes a change in velocity during the flight, but affects only the vertical component of velocity. The ball starts with a nonzero velocity, which decreases, drops to a minimum value at the top, and then increases downward. The acceleration due to gravity is constant throughout the motion. The velocity is changing throughout the motion. Question 1 from Chapter 2 is similar to this problem.

## 2: Launching a marble off the edge of a table

A marble is launched off the edge of a horizontal table and lands on the floor. Draw the trajectory of the ball from the table to the floor. Draw a second line showing the trajectory of the marble if it were given a smaller initial velocity. Draw a third line showing the trajectory if the marble were given a larger initial velocity than the original initial velocity.


Figure 3.2 Question 2.

## Solution

IDENTIFY, SET UP, AND EXECUTE: The table and marble are sketched in Figure 3.2. The initial trajectory is shown and labeled " 1 ." The marble follows a parabolic path, starting with an initial nonzero horizontal velocity. For the smaller initial velocity, the marble also follows a parabolic path, but with a termination point closer to the edge of the table. This path is shown in the figure and is labeled "2." For the larger initial velocity, the marble again follows a parabolic path, but with a termination point farther from the edge of the table. This path is shown in the figure and is labeled " 3 ."

EVALUATE: The paths are similar; their differences owe to the different initial velocities. How does the time the marble spends in the air compare for the three paths? All three take the same amount of time to reach the ground, as they all start with zero initial vertical velocity and fall the same distance. Since they spend the same time in the air, those with larger initial velocities reach greater horizontal distances.

## 3: Comparing projectiles

Figure 3.3 graphs the paths of two projectiles in the $x y$ plane. If we ignore air resistance, how do the initial velocities compare (magnitude and direction)?


Figure 3.3 Question 3.

## Solution

IDENTIFY, SET UP, AND EXECUTE: At the origin, we see that the two paths begin identically. This indicates that the initial directions of the velocities of both projectiles are the same.

The graph does not include a time axis, so we need to look to other clues to compare the magnitudes of the initial velocities. Trajectory $B$ reaches a greater height, indicating that its initial vertical component of velocity was larger than trajectory $A$ 's initial vertical component. Therefore, trajectory $B$ has a greater magnitude of initial velocity.

EVALUATE: Without a time axis, examine the $x$ motion and we cannot assume that trajectory $B$ has a greater magnitude of initial velocity.

How do the horizontal components of the initial velocities compare? Both projectiles have the same initial launch angle; therefore, the ratio of their velocity components must be the same. If the vertical component of $B$ 's velocity is larger, so must $B$ 's horizontal velocity component be larger.

## 4: Comparing projectiles again

Figure 3.4 shows the graph of the paths of two projectiles in the $x y$ plane. If we ignore air resistance, how do the initial velocities compare (in magnitude and direction)? Which projectile lands first?


Figure 3.4 Question 4.

## Solution

IDENTIFY, SET UP, AND EXECUTE: We see that the paths do not coincide at the origin: Projectile $A$ has a larger launch angle.

Again, the graph does not include a time axis, so we need to look to other clues to compare the magnitudes of the initial velocities. Both trajectories reach the same maximum height, indicating that both have the same vertical velocity components. However, their initial directions were different, requiring their initial horizontal components to be different. Trajectory $B$ reaches a greater horizontal distance, so must have a larger initial horizontal velocity. Therefore, trajectory $B$ has a greater magnitude of initial velocity.

Since both trajectories reach the same maximum height and have the same initial vertical velocity, they must end at the same time.

EVALUATE: As an alternative analysis, we could have considered the time first and the velocity second. In that case, it might have been easier to see that the initial horizontal velocity of projectile $B$ was larger because it covered more distance in the same time.

## 5: Falling luggage

A piece of luggage falls out of the cargo door of a airplane flying horizontally at a constant speed. In what direction should the pilot look to follow the luggage to the ground so that it can be recovered?

## Solution

IDENTIFY, SET UP, AND EXECUTE: When the piece of luggage falls from the airplane, its initial velocity is the same as the plane's velocity. As it falls, the luggage accelerates in the vertical direction and its horizontal velocity remains constant (assuming no air resistance). Since the plane and the piece of luggage are moving at the same horizontal velocity, the luggage falls directly below the plane. The pilot should look straight down to see where the luggage will land.

EVALUATE: You might expect that the piece of luggage would fall behind the airplane. Now, what would cause it to fall behind the plane? For it to fall behind the plane, the luggage would have to slow down, or accelerate in a direction opposite that of its horizontal motion. Air resistance could slow down the bag, because air resistance opposes the motion of a body.

## Problems

## 1: Water Balloon Launch

Your physics professor is walking past the physics building at a constant $3.5 \mathrm{~m} / \mathrm{s}$. You're on the thirdfloor balcony ( 25 m above the ground) of the building with your new water balloon launcher. The launcher allows you to adjust the speed of the water balloon, but you can launch the balloon only horizontally. What launch speed should be set for the balloon to land on your professor if you launch it just
as she passes below? What will be your professor's horizontal distance from the building when the balloon hits her, as measured from a point on the ground directly below you?


Figure 3.5 Problem 1.

## Solution

IDENTIFY: Once launched, the water balloon will undergo gravitational acceleration in the vertical direction and continue with constant velocity in the horizontal direction. We will apply the constantacceleration kinematics equations separately to the horizontal and vertical components to solve the problem.

SET UP: Figure 3.5 shows a sketch of the situation. We ignore effects due to the air. Your professor is roughly 1.7 m tall, but we'll ignore her height and determine the position where the balloon hits the ground. An $x y$ coordinate system is shown in the diagram.

We first determine the launch speed. Since there is no acceleration in the horizontal direction, the water balloon must be launched at the same speed as your professor is walking, $3.5 \mathrm{~m} / \mathrm{s}$.

EXECUTE: To find where the balloon hits her, we find the time from the start of the vertical motion and use that to find the horizontal distance the water balloon travels as it falls. The vertical position for constant acceleration is given by

$$
y=y_{0}+v_{0 y} t+\frac{1}{2} a t^{2} .
$$

We've set the origin at the ground; therefore, the initial position becomes 25 m and the final position becomes 0 . The launcher imparts only a horizontal velocity, so the initial vertical velocity is zero. The acceleration is $-g$, since the positive vertical axis is directed upward. Combining these parameters gives

$$
0=25 \mathrm{~m}-\frac{1}{2} g t^{2} .
$$

Solving for $t$ produces

$$
t=\sqrt{\frac{2(25 \mathrm{~m})}{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}}=2.26 \mathrm{~s}
$$

It takes 2.26 s for the balloon to fall to the ground. During this time, it is traveling with constant horizontal velocity. We find the horizontal distance it travels from the formula

$$
x=x_{0}+v_{0 x} t
$$

Your origin is directly below your position on the balcony $\left(x_{0}=0\right)$. Substituting the horizontal velocity and time we calculated, we find the horizontal distance:

$$
x=v_{0 x} t=(3.5 \mathrm{~m} / \mathrm{s})(2.26 \mathrm{~s})=7.9 \mathrm{~m}
$$

The water balloon will hit your professor a horizontal distance 7.9 m away from your location.
EVALUATE: This is a straightforward application of two-dimensional kinematics. We solved for one component of the motion and substituted the result into the equation for the other component to arrive at the solution. Note that we solved for the vertical motion and substituted the result into the equation for the horizontal motion, the opposite order of the previous problem. Practicing solving a variety of problems will build proficiency in solving problems involving motion in a plane.

## 2: Hitting a Baseball in Fenway Park

You win a chance to try hitting a baseball over the "Green Monster" in Fenway Park. The Green Monster is a $37.2-\mathrm{ft}$ ( $11.3-\mathrm{m}$ )-high wall in left field of the ballpark. The left end is closest to home plate, $310 \mathrm{ft}(94.5 \mathrm{~m})$ away. If you give the ball an initial speed of $33 \mathrm{~m} / \mathrm{s}$ at an initial angle of $47^{\circ}$, by how much does the baseball clear (or miss) the top of the wall?


Figure 3.6 Problem 2 sketch.

## Solution

IDENTIFY: The baseball has a nonzero initial velocity, undergoes acceleration due to gravity in the vertical direction, and has no acceleration in the horizontal direction. Constant-acceleration kinematics equations will be applied separately to the horizontal and vertical components to find the solution.

SET UP: We sketch the problem in Figure 3.6. We ignore effects due to the air. The ball is hit roughly 1 m or so above the ground, but we'll neglect this small distance and set the origin at ground level. An $x y$ coordinate system that coincides with this choice is shown in the diagram.

We will solve for the time the baseball arrives at the wall by using the horizontal-position equation. Then we will substitute into the vertical-position equation to find the vertical position of the ball at the wall.

EXECUTE: The horizontal position is given by

$$
x=x_{0}+v_{0 x} t
$$

In this case, we start at the origin $\left(x_{0}=0\right)$ and the $x$ component of velocity includes a cosine term:

$$
x=v_{0 x} t=v_{0} \cos \theta t
$$

We wish to find the time $t$ when the baseball is located at the wall $(x=94.5 \mathrm{~m})$ :

$$
t=\frac{x}{v_{0} \cos \theta}=\frac{(94.5 \mathrm{~m})}{(33 \mathrm{~m} / \mathrm{s}) \cos 47^{\circ}}=4.20 \mathrm{~s}
$$

After 4.20 s , the baseball's horizontal position is 94.5 m . We now find the vertical position at that time. The vertical position is given by

$$
y=y_{0}+v_{0 y} t+\frac{1}{2} a t^{2}
$$

Again, we start at the origin $\left(y_{0}=0\right)$, the acceleration is directed downward (negative) and is of magnitude $g$, and the $y$-component of velocity includes a sine term:

$$
y=v_{0} \sin \theta t+\frac{1}{2}(-g) t^{2}
$$

We can now substitute our values into the equation to find the height of the ball:

$$
y=(33 \mathrm{~m} / \mathrm{s})\left(\sin 47^{\circ}\right)(4.20 \mathrm{~s})+\frac{1}{2}\left(-9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(4.20 \mathrm{~s})^{2}=14.9 \mathrm{~m}
$$

At the wall, the ball's height is 14.9 m , or 3.6 m above the $11.3-\mathrm{m}$-high wall. The ball clears the Green Monster by 3.6 m !

EVALUATE: This is another straightforward application of two-dimensional kinematics. We solved for one component of the motion and substituted the result into the equation for the other component to arrive at the solution. We will follow this procedure often to solve problems involving motion in a plane.

Practice Problem: Find the $x$ and $y$ components of the baseball's velocity at the wall. Answer: $v_{x}=22.5 \mathrm{~m} / \mathrm{s}, v_{y}=-17.0 \mathrm{~m} / \mathrm{s}$.

CAUTION Don't Mix $\boldsymbol{x}$ and $\boldsymbol{y}$ ! It is easy to mix up $x$ and $y$ components for position, velocity, and acceleration. You must label each of these carefully to ensure that you don't make mistakes. Only time is common to both the $x$ and $y$ components.

## 3: Acceleration of a Propeller Tip

The Wright Brothers' plane had a 2.4 -m-long propeller that operated at a constant 350 rpm . Find the acceleration of a particle at the tip of the propeller.

## Solution

IDENTIFY: This is a uniform circular motion problem; the acceleration is determined by the centripetal-acceleration formula.

SET UP: We will need to find the velocity and radius from the information provided. A diagram of the problem is shown in Figure 3.7.

To find the centripetal acceleration, we need the radius and speed of a particle on the tip of the propeller. We are given the diameter of the propeller, and dividing that in half gives the radius.


Figure 3.7 Problem 3.
EXECUTE: The speed of a particle at the end of the propeller is found by dividing the circumference at the tip of the propeller $(2 \pi r)$ by the time it takes the propeller to make one revolution: $(T)$

$$
v=\frac{2 \pi r}{T}
$$

The propeller makes 350 revolutions per minute, so we find the time it takes to make 1 revolution by dividing 1 minute by 350 revolutions:

$$
T=\frac{1 \mathrm{~min}}{350 \mathrm{rev}}=\frac{60 \mathrm{~s}}{350 \mathrm{rev}}=0.171 \mathrm{~s} / \mathrm{rev}
$$

The propeller takes 0.171 s to make one revolution. We can now find the velocity:

$$
v=\frac{2 \pi r}{T}=\frac{2 \pi(1.2 \mathrm{~m})}{0.171 \mathrm{~s}}=44.1 \mathrm{~m} / \mathrm{s}
$$

The centripetal acceleration is then

$$
a_{\text {rad }}=\frac{v^{2}}{r}=\frac{(44.1 \mathrm{~m} / \mathrm{s})^{2}}{(1.2 \mathrm{~m})}=1620 \mathrm{~m} / \mathrm{s}^{2}
$$

The centripetal acceleration of a particle on the tip of the propeller is $1620 \mathrm{~m} / \mathrm{s}^{2}$. This is equivalent to 165 times the acceleration due to gravity!

EVALUATE: We have found the magnitude of the acceleration in this problem. Acceleration is a vector, so where does it point? The acceleration is directed toward the center of the propeller, perpendicular to the velocity.

We did not include gravity in this problem. The particle is affected by attraction toward the ground throughout its motion; however, the resulting acceleration is very small compared with the centripetal acceleration.

## 4: Paddling across a River

You wish to paddle north across a 350 -m-wide river. The river has a $1.2 \mathrm{~m} / \mathrm{s}$ current from east to west, and you can paddle at a steady $1.5 \mathrm{~m} / \mathrm{s}$ pace. In what direction should you paddle, and how long will it take you to cross the river?


Figure 3.8 Problem 4.

## Solution

IDENTIFY: You will need to paddle into the river current to compensate for the river's moving your boat downstream as you cross.

SET UP: Figure 3.8 shows a sketch of the situation. We use relative velocities to solve the problem. The direction in which you must paddle is determined by setting north to be the direction of your resulting relative velocity with respect to the earth.


Figure 3.8 Problem 4.
EXECUTE: Figure 3.9 combines your velocity with respect to the river $\left(v_{Y \mid R}\right)$ with the river current's velocity with respect to the earth $\left(v_{R / E}\right)$ to form your relative velocity with respect to the earth $\left(v_{Y / E}\right)$ :

$$
\vec{v}_{Y \mid E}=\vec{v}_{Y \mid R}+\vec{v}_{R / E}
$$

For you to land directly across from your starting point, the direction of $v_{Y \mid E}$ must be northward. Therefore, the $x$ component of $v_{Y / R}$ must be equal and opposite to $v_{R / E}$. We find the direction in which you should paddle by equating those two magnitudes:

$$
\begin{aligned}
\left(\vec{v}_{Y \mid R}\right)_{x} & =v_{R / E} \\
\left(\vec{v}_{Y \mid R}\right)_{x} & =v_{Y / R} \sin \theta=v_{R / E} \\
\theta=\sin ^{-1}\left(\frac{v_{R / E}}{v_{Y / R}}\right) & =\sin ^{-1}\left(\frac{1.2 \mathrm{~m} / \mathrm{s}}{1.5 \mathrm{~m} / \mathrm{s}}\right)=53^{\circ} .
\end{aligned}
$$

You will need to paddle $53^{\circ}$ east of north to follow a northward path. The time it will take is found from the $y$ component of the displacement. You are traveling at constant velocity, so the vertical component of the displacement is

$$
y-y_{0}=\left(v_{Y \mid E}\right)_{y} t .
$$

$\left(v_{Y \mid E}\right)_{y}$ is the magnitude of $v_{Y \mid E}$, because $v_{Y \mid E}$ has only a $y$ component. $\left(v_{Y \mid E}\right)_{y}$ must also be equal to the $y$ component of $v_{Y \mid R}$. Solving for time gives

$$
t=\frac{y-y_{0}}{\left(v_{Y \mid E}\right)_{y}}=\frac{y-y_{0}}{v_{Y \mid E}}=\frac{y-y_{0}}{\left(v_{Y \mid R}\right)_{y}}=\frac{y-y_{0}}{v_{Y \mid E} \cos \theta}=\frac{350 \mathrm{~m}}{(1.5 \mathrm{~m} / \mathrm{s}) \cos \left(53^{\circ}\right)}=390 \mathrm{~s}
$$

It will take you 390 s to paddle across the river.
EVALUATE: When you paddle across a river perpendicular to its flow, your relative velocity with respect to the earth is always less than your velocity with respect to the river. It also takes longer to cross a river that has a current compared with a calm river when your path is perpendicular to the river. The next practice problem lets you compare the time required to cross a calm river with the time you just found in dealing with a river that has a current.
Practice Problem: How long would it take to paddle across the same river with no current? Answer: 230 s.

## Try It Yourself!

## 1: Ball thrown from a cliff

A boy throws a ball from a cliff at an angle of $30.0^{\circ}$ above the horizontal with an initial velocity of $10.0 \mathrm{~m} / \mathrm{s}$. The ball lands 100.0 m from the base of the cliff. (a) How high is the cliff? (b) What is the time of flight of the ball? (c) What is the velocity of the ball just before impact?

## Solution Checkpoints

IDENTIFY AND SET UP: Constant-acceleration kinematics equations should be applied separately to the horizontal and vertical components of the ball's motion. The initial velocity should be broken down into $x$ and $y$ components. There is acceleration only in the vertical direction.

EXECUTE: Equations for the $x$ and $y$ components of position and velocity can be found. The $x$ position equation can be solved for time and substituted into the $y$ position equation to find the height of the cliff. Doing this leads to

$$
y=x \tan \theta-\frac{1}{2} g\left[\frac{x}{v_{0} \cos \theta}\right]^{2}
$$

from which you will find that the cliff is 595 m high. You can then substitute that number into the $x$ position equation to find that the time is 11.5 s .

To find the velocity just before impact, you can find the velocity components from the kinematic equations. Knowing the components, you can find the magnitude and direction of the velocity. The velocity is $108 \mathrm{~m} / \mathrm{s}$, directed $85^{\circ}$ below the $x$-axis.

EVALUATE: We see that the velocity just before impact is almost straight down. Can it ever be exactly straight down?

## 2: Kicking a soccer ball

A soccer ball is kicked 25 m in the horizontal direction. What is its initial velocity if it reaches a maximum height of 6.0 m ?

## Solution Checkpoints

IDENTIFY AND SET UP: Constant-acceleration kinematics equations should be applied separately to the horizontal and vertical components of the ball's motion. Draw the ball's path and choose an appropriate coordinate system.

EXECUTE: The initial $y$ component of velocity can be found from the formula

$$
v_{v}^{2}=v_{0 y}^{2}+2 a_{y} \Delta y
$$

This gives an initial vertical velocity of $10.8 \mathrm{~m} / \mathrm{s}$. To find the initial horizontal component of velocity, we determine the flight time from the vertical motion and combine that with the horizontal distance traveled. The result is an initial velocity of $15.6 \mathrm{~m} / \mathrm{s}$ at an angle of $44^{\circ}$ above the horizontal.

EVALUATE: By this point, you have seen many kinematics problems. Can you summarize your problem-solving techniques?

## 3: Archer and arrow

An archer shoots an arrow into the air at an angle of $30^{\circ}$ above the horizontal. It lands on a building 100.0 m away at a height of 20.0 m . What was the initial speed of the arrow?

## Solution Checkpoints

IDENTIFY AND SET UP: Can constant-acceleration kinematics be used in this case? What assumptions do you need to make?

EXECUTE: The $x$ and $y$ position equations can be rearranged to yield

$$
y=x \tan \theta-\frac{1}{2} g\left[\frac{x}{v_{0} \cos \theta}\right]^{2}
$$

Rewriting the equation and the $x$ and $y$ positions at the building gives $v_{0}=41 \mathrm{~m} / \mathrm{s}$.
EVALUATE: How can you check your result? How do you know that it is reasonable?

## Newton's Laws of Motion

## Summary

We will define dynamics-the study of the relationship of motion to forces-in this chapter. Newton's laws of motion will lay the foundation for our studies and link forces to acceleration, which we investigated in the previous chapters. We will define force, mass, and weight and apply them to problems. We will use our knowledge of vectors to better understand forces and construct free-body diagrams. By the end of this chapter, we will have built a problem-solving framework that we will apply in the next chapter.

## Objectives

After studying this chapter, you will understand

- The concept of force and why it is a vector quantity.
- How to identify forces acting on a body.
- How to find the resultant force acting on an object by summing multiple forces.
- Newton's three laws of motion.
- The relation between net force, mass, and acceleration.
- How to recognize an inertial frame of reference, in which Newton's laws are valid.
- How to use a free-body diagram to represent forces acting on an object.
- How to use the free-body diagram as a guide in writing force equations for Newton's laws.


## Concepts and Equations

| Term | Description |
| :---: | :---: |
| Force | A force is a quantitative measure of the interaction between two objects, represented by a vector. The SI unit of force is the newton ( N ). One newton equals 1 kilogram-meter per second squared. |
| Combining Forces | The vector sum of forces acting on a body is the resultant, denoted $\vec{R}$ : $\vec{R}=\vec{F}_{1}+\vec{F}_{2}+\vec{F}_{3}+\cdots=\sum \vec{F} .$ <br> The effect of many forces acting on a body can be captured by the resultant force. This principle is called superposition of forces. |
| Contact Force | A contact force is a force between two objects touching at a surface. A contact force has two components: a component perpendicular to the surface (the normal force) and a component parallel to the surface (the frictional force). |
| Normal Force | The normal force is the component of a contact force between two objects that is perpendicular to their common surface. The normal force is denoted by $\vec{n}$. |
| Friction Force | The friction force is the component of a contact force between two objects that is parallel to their common surface. The friction force is denoted by $\vec{f}$. Friction forces often act to resist the sliding of an object. |
| Tension Force | A tension force is conveyed by the pull of a rope or cord and is denoted by $\vec{T}$. |
| Newton's First Law | Newton's first law states that when the vector sum of forces acting on an object is zero, the object is in equilibrium and has zero acceleration. The object will remain at rest or move with constant velocity when no net force acts upon it. The law is valid only in inertial reference frames. |
| Inertial Reference Frame | An inertial reference frame is a reference frame in which Newton's laws are valid. A common example of a noninertial reference frame is that of an accelerating airplane. |
| Newton's Second Law | Newton's second law of motion states that an object which is not in equilibrium is acted upon by a net force and accelerates. The acceleration is given in vector form by $\sum \vec{F}=m \vec{a},$ |

where $m$ is the object's mass, which characterizes the inertia of the object.
Newton's second law can also be written in component form as

$$
\sum F_{x}=m a_{x} \quad \sum F_{y}=m a_{y} \quad \sum F_{z}=m a_{z} .
$$

An object's weight is the gravitational force exerted on the object by the earth or another astronomical body and is denoted by $\vec{w}$. The magnitude of an object's weight is equal to the product of its mass and the magnitude of the acceleration due to gravity:

$$
w=m g .
$$

## Newton's Third Law

Newton's third law states that, for two interacting bodies $A$ and $B$, each exerts a force on the other of equal magnitude and opposite in direction, or

$$
\vec{F}_{A \text { on } B}=-\vec{F}_{B \text { on } A} .
$$

Free-Body Diagram
A diagram showing all forces acting on an object. The object is represented by a point; the forces are indicated by vectors. A free-body diagram is useful in solving problems involving forces.

## Conceptual Questions

## 1: Winning a Tug-of-War

In a tug-of-war shown in Figure 4.1, how does the force applied to the rope by the losing team compare with the force applied to the rope by the winning team? Is the magnitude of the force applied by the losing team less than, greater than, or equal to the magnitude of the force applied by the winning team? How does the winning team win?


Figure 4.1 Question 1.

## Solution

IDENTIFY, SET UP, AND EXECUTE The free-body diagram in Figure 4.2 will guide our analysis. Each team experiences four forces: the tension due to the tug-of-war rope ( $T$ ), a frictional force with the ground $(f)$, the weight of the team $(w)$, and the normal force exerted upward from the ground $(N)$. The subscripts indicate the winning $(w)$ and losing $(l)$ teams. Examining the diagrams, we see that the tension forces must be an action-reaction pair-hence the explicit notation. Therefore, by Newton's third law, the tensions must be equal. The force applied by the losing team on the rope must be of the same magnitude as, but opposite in direction to, the force applied by the winning team on the rope.


Figure 4.2 Question 1 free-body diagrams.

The vertical forces will not influence the horizontal interaction, so we look at the remaining force to determine how the winning team wins. The frictional forces must not be equal: The winning team exerts a larger frictional force on the ground than the losing team does in order to accelerate the losing team across the centerline.

EVALUATE We see that free-body diagrams and Newton's third law were crucial in our solution. The free-body diagram helped reduce the complexity of the problem and helped show that the frictional
force was responsible for the win. After establishing that the tensions were an action-reaction pair, we saw from Newton's third law that the tensions were equal and opposite.

Note that we assumed that the frictional force was between the team and the ground and that each team was able to grip the rope without sliding. There is also a frictional force between the teams' hands and the rope. Differences between the hand and rope frictional forces could also have led to the win.

## 2: Flying groceries

What force causes a bag of groceries to fly forward when you come to an abrupt stop in a car?

## Solution

IDENTIFY, SET UP, AND EXECUTE Suppose that, before you come to an abrupt stop, you are moving at a constant velocity. Then no net force must act on you, the car, or the bag of groceries, according to Newton's first law. As you cause an abrupt stop by hitting the brakes, you increase the frictional force between the car and the road, creating a net force on the car. When you brake, the force on the bag of groceries doesn't change, so the bag of groceries continues at its initial velocity. (We're assuming that the frictional force between the bag of groceries and the car seat is small.) Therefore, no force causes the bag of groceries to fly forward when you come to an abrupt stop in a car.

EVALUATE The solution may seem a bit illogical, for consider how the situation would appear to someone outside of the car: The bag of groceries continues moving at a constant velocity after the brakes are applied. This scenario should be more plausible and is a clearer way to imagine the situation.

This is one example of a noninertial frame of reference. The slowing car has negative acceleration and hence is an accelerated frame of reference. Newton's laws don't apply to noninertial frames of reference, so we cannot apply our new force techniques to this problem.

From inside the car, you may try to explain the situation by invoking a "force of inertia." This would be a fictitious force, however, and must be avoided. All of the forces we've encountered (and all of those we will later encounter) arise from known interactions.

## 3: Does an Apple Accelerate When Placed on a Table?

An apple is placed on a table. Can we describe the apple as having an acceleration of $9.8 \mathrm{~m} / \mathrm{s}^{2}$ toward earth and a second acceleration of $9.8 \mathrm{~m} / \mathrm{s}^{2}$ upward due to the table, thus resulting in a net acceleration of zero?

## Solution

IDENTIFY, SET UP, AND EXECUTE We have seen how forces can cause accelerations, have heard $F=m a$ often, and know that an object's weight is $m g$, so it may appear logical to replace forces with mass times acceleration in equations. However, Newton's laws apply to combining forces, not accelerations. Newton's second law states that a net force on an object will lead to an acceleration equal to the net force divided by the object's mass.

EVALUATE This question points up a common misconception about accelerations and forces. At times, replacing forces with mass times acceleration may lead to the same results as following the correct procedures, but doing so often leads to confusion. An object that is stationary is not accelerating, because there is no net force acting on the object.

## 4: Forces and Moving Objects

Does a force cause an object to move? Does a moving object "have" a force?

## Solution

IDENTIFY, SET UP, AND EXECUTE A force does not necessarily cause an object to move. Your textbook is acted upon by gravity when placed on a desk, but it does not move. A net force can cause an object to acquire velocity through acceleration.

An object moving at a constant velocity has no net force acting on it; therefore, the fact that an object is in motion does not indicate that a force is acting upon it. The fact that an object is accelerating, however, would certainly indicate that at least one force is acting upon it.

EVALUATE Acceleration and motion are not equivalent. Acceleration is motion during which the velocity changes over time. An object can also have a constant velocity, which is motion without acceleration. You must carefully distinguish between motion and acceleration in order to grasp physics.

## 5: Definition of equilibrium

Can a moving object be in equilibrium?

## Solution

IDENTIFY, SET UP, AND EXECUTE Equilibrium occurs when the net force acting on an object is zero. Newton's first law states that objects with no net force acting on them remain at rest or continue with constant velocity. An object moving at constant velocity is in equilibrium.

EVALUATE Equilibrium has a precise definition in physics, even though the word may have connotations of a stationary object. Physics relies upon precise definitions to build representations of physical processes. You must apply physics definitions carefully to build your physics understanding.

## Problems

## 1: Combining several forces to find the resultant

A mover uses a cable to drag a crate across the floor as shown in Figure 4.3. The mover provides a 300 N force and pulls the cable at an angle of $30.0^{\circ}$. The crate weighs 500 N , and the floor provides a 350 N normal force on the crate and opposes his pull with a 150 N frictional force. Find the resulting force acting on the crate. Will the crate accelerate?


Figure 4.3 Problem 1.

## Solution

IDENTIFY We will combine the forces acting on the crate to find the net force. If the net force is not zero, then there will be acceleration.

SET UP We find the resultant force by adding the forces acting on the crate. Four forces act on the crate: the tension force due to the mover's pull ( $T$ ), the crate's weight ( $w$ ), and the normal force ( $n$ ) and friction force $(f)$ due to the floor. We represent the four forces as vectors in the free-body diagram of the crate in Figure 4.4.


Figure 4.4 Problem 1 free-body diagram.

We have added an $x y$ coordinate system to the free-body diagram as the forces act in two dimensions. We've also resolved the tension force into its $x$ and $y$ components.

EVALUATE We add the four forces together by adding their components, writing separate equations for the $x$ and $y$ components. There are two $x$ components, due to the horizontal components of the tension force and the friction force:

$$
\sum F_{x}=T \cos 30^{\circ}+(-f)
$$

The $x$ component of the tension force is to the right and is thus assigned a positive value, while the friction force is to the left and is assigned a negative value. We proceed to the $y$ components. There are three $y$ components, one due to the normal force, a second due to the weight of the crate, and, finally, the vertical component of the tension force:

$$
\sum F_{y}=n+(-w)+T \sin 30^{\circ}
$$

The $y$ component of the tension force and the normal force are directed upward and are assigned positive values, while the weight of the crate is directed downward and is assigned a negative value. We now substitute the values for the variables to find the net force along both axes:

$$
\begin{aligned}
& \sum F_{x}=T \cos 30^{\circ}+(-f)=(300 \mathrm{~N}) \cos 30^{\circ}+(-150 \mathrm{~N})=+110 \mathrm{~N} \\
& \sum F_{y}=n+(-w)+T \sin 30^{\circ}=(350 \mathrm{~N})+(-500 \mathrm{~N})+(300 \mathrm{~N}) \sin 30^{\circ}=0 \mathrm{~N}
\end{aligned}
$$

The resultant force on the crate has an $x$ component of +110 N and no $y$ component, (i.e., the resulting force is horizontal and points to the right). There is also a resulting acceleration of the crate to the right, as there is a net force.

EVALUATE This is a typical force problem in which we have used our vector addition skills to find the resultant force. We see that there is no net force in the vertical direction; therefore, the crate remains on the bed of the truck.

PRACTICE PROBLEM At what rate does the crate accelerate? Answer: $2.2 \mathrm{~m} / \mathrm{s}^{2}$.

## 2: Using Newton's second law to find the mass of a cruise ship

A tugboat pulls a cruise ship out of port. (See Figure 4.5.) You estimate the acceleration by noting that the tugboat takes 60 s to move the cruise ship 100 m , starting from rest. If the tugboat exerts $3 \times 10^{6} \mathrm{~N}$ of thrust, what is the mass of the cruise ship? Ignore drag due to the water, and assume that the tugboat accelerates uniformly.


Figure 4.5 Problem 2.

## Solution

IDENTIFY We will use Newton's second law to find the mass of the cruise ship, given the net force and acceleration acting on the ship.

SET UP The problem tells us the net force provided by the tugboat, and the acceleration can be determined from the kinematics information. We ignore drag or friction with the water, so the only horizontal force acting on the cruise ship is due to the tugboat.

EXECUTE Newton's second law relates the net force to the mass and resulting acceleration:

$$
\sum F_{x}=m a .
$$

The net force acting on the cruise ship is $3 \times 10^{6} \mathrm{~N}$. The acceleration is found from the equation for position as a function of time with constant acceleration:

$$
x=x_{0}+v_{0} t+\frac{1}{2} a t^{2}
$$

Here, the initial velocity is zero and we take the initial position to be zero. Substituting these values into the equation gives

$$
x=\frac{1}{2} a t^{2} .
$$

Solving for the acceleration yields

$$
a=\frac{2 x}{t^{2}}
$$

Replacing the distance and time with the given values produces

$$
a=\frac{2 x}{t^{2}}=\frac{2(100 \mathrm{~m})}{(60 \mathrm{~s})^{2}}=0.056 \mathrm{~m} / \mathrm{s}^{2}
$$

We now use Newton's second law to find the mass. Solving for the mass gives

$$
m=\frac{F}{a}=\frac{\left(3 \times 10^{6} \mathrm{~N}\right)}{\left(0.056 \mathrm{~m} / \mathrm{s}^{2}\right)}=54,000,000 \mathrm{~kg}=54 \text { kilotonnes. }
$$

Our estimate shows that the cruise ship has a mass of 54 kilotonnes ( 1 kilotonne $=10^{6} \mathrm{~kg}$ ). More correctly, the cruise ship has a mass of 50 kilotonnes, as the values stated in the problem have only one significant figure.

EVALUATE This problem shows how we can combine Newton's law with observations to make interesting conclusions about the mass of an object.

## 3: Drawing free-body diagrams

Draw a free-body diagram for each of the following situations:
(a) A box slides down a smooth ramp. (See Figure 4.6.)


Figure 4.6 Problem 3a
(b) A box slides down a rough ramp. (See Figure 4.7.)


Figure 4.7 Problem 3b
(c) A block is placed on top of a crate, and the crate is placed on a horizontal surface. (See Figure 4.8.) Draw a free-body diagram of the crate.


Figure 4.8 Problem 3c
(d) A block is placed on top of a crate, and the crate is pulled horizontally across a rough surface. (See Figure 4.9.) The surface between the crate and the block is rough, and the block is held at rest by a string. Draw a free-body diagram of the crate.


Figure 4.9 Problem 3d

## Solution

IDENTIFY We will draw free-body diagrams that show all of the forces acting on the object, representing the forces as vectors.

SET UP The first step is to identify the object and then find the forces acting on it. We'll look at the contact tension, normal and frictional forces, and the noncontact gravitational force.

EXECUTE In part (a), there is no friction, since the ramp is smooth. The only contact force acting on the box is the normal force due to the ramp. Gravity also acts on the box. The free-body diagram includes two forces acting on the box: the normal ( $n$ ) force, directed perpendicular to the ramp; and the weight $(w)$ of the box, directed downward. The free-body diagram of the box is shown in Figure 4.10.


Figure 4.10 Problem 3a
free-body diagram
In part (b), there is friction, since the ramp is rough. The contact forces acting on the box are the normal and frictional forces due to the ramp. Gravity also acts on the box. The free-body diagram includes three forces acting on the box: the normal force ( $n$ ), directed perpendicular to the ramp; the frictional force $(f)$, directed upward along the ramp (opposing the motion of the box); and the weight $(w)$ of the box, directed downward. The free-body diagram of the box is shown in Figure 4.11.


Figure 4.11 Problem 3b
free-body diagram
In part (c), two contact forces act on the crate: the normal force due to the surface and the normal force due to the block. There are no frictional forces, as neither the crate nor the block is moving. Gravity acts on the box. The free-body diagram includes three forces acting on the crate: the normal force due to the surface $\left(n_{\text {surface }}\right)$, directed upward; the normal force due to the block ( $n_{\text {block }}$ ), directed downward; and the weight $(w)$, of the crate, directed downward. The free body diagram of the crate is shown in Figure 4.12.


Figure 4.12 Problem 3c
free-body diagram

In part (d), five contact forces act on the crate: the normal forces due to the surface and the block, the frictional forces due to the surface and the block, and the tension force provided by the pull. The sixth force acting on the crate is gravity. The free-body diagram includes six forces acting on the crate: the normal force due to the surface $\left(n_{\text {surface }}\right)$, directed upward; the normal force due to the block $\left(n_{\text {block }}\right)$, directed downward; the frictional forces due to the surface $\left(f_{\text {surface }}\right)$ and the block $\left(f_{\text {block }}\right)$, both directed to the right; the tension force $(T)$, directed to the left; and the weight $(w)$, of the crate, directed downward. The free-body diagram for the crate is shown in Figure 4.13.


Figure 4.13 Problem 3d
free-body diagram
EVALUATE Drawing free-body diagrams should become second nature to you. We will see their importance when we solve problems in Chapter 5. Free-body diagrams help catch mistakes by identifying all the forces acting on an object, as well as help identify action-reaction pairs.

## 4: Tension in a string connecting two blocks

Two blocks are connected by a massless string, as shown in Figure 4.14. A cable is attached to the upper block and is pulled upward with a 250 N force. Find the tension in the string connecting the two blocks. The upper block has a mass of 7.5 kg and the lower block has a mass of 12 kg .


Figure 4.14
Problem 4

## Solution

IDENTIFY We will use Newton's laws to solve this problem.
SET UP We cannot determine whether the system is in equilibrium or accelerating from the statement of the problem; therefore, we do not know whether to apply Newton's first law for a system in equilibrium or Newton's second law for an accelerating system. Our first step, therefore, will be to determine whether the system is in equilibrium or accelerating. Then we will apply the appropriate one of Newton's laws to find the tension in the string.

We'll use free-body diagrams to solve the problem. We can determine whether the blocks are accelerating by considering the two blocks as one system. The left panel of Figure 4.15 shows a free-body diagram of the system with the two blocks combined. We can find the tension in the string by considering
the two blocks separately. The middle and right panels of Figure 4.15 show the free-body diagrams of the two blocks separately. The top block is designated $A$, the bottom block $B$, to reduce confusion.


Figure 4.15 Problem 4 free-body diagram
The forces in the free-body diagrams are identified by their magnitudes. The combined diagram includes the tension of the cable ( $T_{\text {cable }}$ ) and the weight of the two blocks $\left(w_{A+B}\right)$. The other diagrams also include the tension of the string $\left(T_{\text {string }}\right)$, and the weights of the blocks ( $w_{A}$ and $\left.w_{B}\right)$. The upwardpointing vectors will be taken to be positive, the downward-pointing vectors negative.

EXECUTE To determine whether the blocks are accelerating, we examine the net force acting on them. From the left-hand free-body diagram, we see that there are two forces, the tension of the cable and the weight of the blocks, acting on the combination of blocks:

$$
\sum F_{y}=T_{\text {cable }}+\left(-w_{A+B}\right) .
$$

The weight is the combined mass of the blocks times the gravitational constant. The net force is found by replacing the weight and tension by the given values:

$$
\sum F_{y}=T_{\text {cable }}+\left(-g\left(m_{A}+m_{B}\right)\right)=250 \mathrm{~N}+\left(-\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(7.5 \mathrm{~kg}+12 \mathrm{~kg})\right)=58.9 \mathrm{~N} .
$$

The net force is not zero; therefore, the blocks are accelerating. We find the acceleration from Newton's second law applied to the combined blocks:

$$
\sum F_{y}=\left(m_{A}+m_{B}\right) a .
$$

Solving for the acceleration yields

$$
a=\frac{\sum F_{y}}{\left(m_{A}+m_{B}\right)}=\frac{58.9 \mathrm{~N}}{(7.5 \mathrm{~kg}+12 \mathrm{~kg})}=3.02 \mathrm{~m} / \mathrm{s}^{2} .
$$

We now apply Newton's second law to the lower block to find the tension in the string. Two forces are acting on the lower block: the tension due to the string (upward) and gravity (downward). Hence,

$$
\sum F_{y}=T_{\text {string }}+\left(-m_{B} g\right)=m_{B} a .
$$

Solving for the tension in the string and substituting the value for acceleration gives

$$
T_{\text {string }}=m_{B} g+m_{B} a=m_{B}(g+a)=(12 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}+3.02 \mathrm{~m} / \mathrm{s}^{2}\right)=150 \mathrm{~N} .
$$

The tension in the string is 150 N .
EVALUATE We see that the tension force due to the string is less than the tension force due to the cable. This is expected, as the string provides force to accelerate the lower block, whereas the cable
provides force to accelerate both blocks. It is important not to assume that tensions are equal in problems; you must consider each tension independently.

## Try It Yourself!

## 1: Rock suspended by wire

A rock of mass 4.0 kg is suspended by a wire. When a horizontal force of 29.4 kg is applied to the rock, it moves to the side such that the wire makes an angle $\theta$ with the vertical. Find the angle $\theta$ and the tension in the wire.

## Solution Checkpoints

IDENTIFY AND SET UP The net force on the rock is zero. Three forces act on the rock. By drawing a free-body diagram, we can see how to set the horizontal and vertical components of force to zero to solve the component force equations.

EXECUTE The net horizontal and vertical forces acting on the rock are

$$
\begin{aligned}
& \sum F_{x}=F_{\mathrm{H}}-T \sin \theta=0 \\
& \sum F_{y}=T \cos \theta-m g=0
\end{aligned}
$$

Dividing one equation by the other leads to an angle of $37^{\circ}$. Substituting the angle into either equation leads to a tension of 49 N .

EVALUATE We will often break the net force into horizontal and vertical components and solve each separately, much as we did in our two-dimensional kinematics problems.

## 2: Tension in an elevator cable

A 1000.0 kg elevator rises with an upward acceleration equal to $g$. What is the tension in the supporting cable?

## Solution Checkpoints

IDENTIFY AND SET UP There is a net force on the elevator, so Newton's second law will be used to find the tension.

EXECUTE The net vertical force acting on the elevator is

$$
\sum F_{y}=T-m g=m a=m g .
$$

Solving for the tension gives $19,600 \mathrm{~N}$.
EVALUATE We see that the tension in the cable is larger than the force of gravity on the elevator. Does this make physical sense?

## 3: Acceleration in an elevator

A 100.0 kg man stands on a bathroom scale in an elevator. What is the acceleration of the elevator when the scale reading is (a) 150 kg , (b) 100 kg , and (c) 50 kg ?

## Solution Checkpoints

IDENTIFY AND SET UP Two forces act on the man: the force of the scale and the force of gravity. Draw a free-body diagram to guide you.

EXECUTE The net vertical force acting on the man is

$$
\sum F_{y}=F_{\text {scale }}-m g=m a
$$

Solving for the acceleration gives (a) $a=4.9 \mathrm{~m} / \mathrm{s}^{2}$, (b) $a=0$, and (c) $a=-4.9 \mathrm{~m} / \mathrm{s}^{2}$.
EVALUATE In which case(s) does the man feel heavier than normal? In which case(s) does he feel lighter than normal? In which case(s) does he feel as if he has normal weight? What is the significance of the signs in answers (a) and (c)?

## Applying Newton's Laws

## Summary

In this chapter, we will apply Newton's laws of motion to bodies that are in equilibrium (at rest or in uniform motion) and to bodies that are not in equilibrium (in accelerated motion). We'll develop a consistent problem-solving strategy that utilizes a free-body diagram to identify the relevant forces acting on a body. We'll also expand our catalog of forces by quantifying contact forces and friction forces, as well as examine forces on bodies in uniform circular motion. By the end of the chapter, you will have built a foundation for solving equilibrium and nonequilibrium problems involving any combination of forces, including forces that we'll discover in later chapters.

## Objectives

After studying this chapter, you will understand

- How to efficiently represent forces acting on a body by using a free-body diagram.
- How to use the free-body diagram as a guide in writing force equations for Newton's laws.
- How to use Newton's first law to solve problems involving bodies in equilibrium.
- How to use Newton's second law to solve problems involving accelerating bodies.
- How to apply contact forces and various friction forces to a variety of situations.
- How to recognize action-reaction pairs and use Newton's third law to quantify their magnitudes.
- How to apply Newton's laws of motion to bodies moving in uniform circular motion.
- How to use Newton's laws to solve problems proficiently.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Using Newtons' First Law | A body in equilibrium (either at rest or moving with constant velocity) is <br> acted upon by no net force: The vector sum of the forces acting on the object <br> must be zero according to Newton's first law of motion, $\sum \vec{F}=0$. <br> When solving equilibrium problems, one starts with free-body diagrams, <br> finds the net forces along two perpendicular components, and then solves by <br> using the equations $\quad \sum F_{x}=0 \quad \sum F_{y}=0$. |
| Applying Newton's Second Law | A body that is acted upon by a nonzero net force accelerates. The acceleration <br> is given by Newton's second law of motion, |
| $\sum \vec{F}=m \overrightarrow{\mathbf{a}}$. |  |

When solving nonequilibrium problems, one identifies the forces acting on the body with free-body diagrams, finds the net forces along two perpendicular components, and determines the equations of motion given by

$$
\sum F_{x}=m a_{x} \quad \sum F_{y}=m a_{y}
$$

In problems with multiple interacting bodies, it may be necessary to apply Newton's laws to each body individually and solve the equations simultaneously.

## Friction Force

## Kinetic Friction Force

## Static Friction Force

## Forces in Circular Motion

A friction force is that component of the contact force between two objects which is parallel to the surfaces in contact. Friction forces, denoted by $\vec{f}$, are generally proportional to the normal force and include kinetic friction forces (when the objects move relative to each other), static friction forces (when there is no motion between the objects), viscosity and drag forces (for motion involving liquids and gases), and rolling frictional forces (for rolling objects).

The kinetic friction force is the friction force between two objects moving relative to each other and is generally proportional to the normal force between the objects. The proportionality constant is the coefficient of kinetic friction $\left(\mu_{k}\right)$, which depends on the objects' surface characteristics and has no units.The direction of the kinetic friction force is always opposite the direction of motion. Mathematically,

$$
f_{k}=\mu_{k} n .
$$

The static friction force is the friction force between two objects that are not moving relative to each another. The maximum static friction is generally proportional to the normal force between the objects, where the proportionality constant is the coefficient of static friction $\left(\mu_{s}\right)$. Often, $\mu_{s}$ is greater than $\mu_{k}$ for a given surface. The static friction can vary from zero to the maximum value; its magnitude depends on the component of the applied forces parallel to the surface. The direction of the static frictional force is opposite that of the parallel component of the net applied force. Mathematically,

$$
f_{s} \leq \mu_{s} n
$$

For a body in uniform circular motion, the acceleration is directed toward the center of the circle. The motion is determined by Newton's second law,

$$
\sum \vec{F}=m \vec{a}
$$

The body's acceleration is

$$
a_{\mathrm{rad}}=\frac{v^{2}}{R}=\frac{4 \pi^{2} R}{T^{2}}
$$

## Conceptual Questions

## 1: Finding errors in a free-body diagram

Two weights are suspended from the ceiling and each other by ropes as shown in Figure 5.1a. A freebody diagram is shown in Figure 5.1 b for the upper block $(A)$. Find the error in the free-body diagram and draw the correct diagram.


Figure 5.1 Question 1.

## Solution

IDENTIFY, SET UP, AND EXECUTE Three forces act on block $A$ : two tension forces due to the ropes and the gravitational force on block $A$. The gravitational force on block $B$ has been incorrectly included in the diagram. Block $B$ is not in direct contact with block $A$; only the rope is in contact with block $A$. The corrected free-body diagram is shown in Figure 5.2.


Figure 5.2 Question 1 corrected free-body diagram.

EVALUATE When drawing free-body diagrams, one must include only those forces acting on the object. Identifying the forces acting on an object is necessary to apply Newton's laws correctly.

## 2: Investigation of the normal force

For which of the following figures is the normal force not equal to the object's weight?
(a)

(b)

(c)

(d)

(e)


Figure 5.3 Question 2

## Solution

IDENTIFY, SET UP, AND EXECUTE The normal force is not equal to the object's weight in figures (a), (c), and (e). In (a), the normal force is equal to the book's weight plus the force pushing down on the book. In (c), the book's weight is directed downward and the normal force is directed upward and to the left, perpendicular to the ramp's surface. Here, the normal force is equal to the weight multiplied by the cosine of the ramp angle. In (d), the normal force is equal to the book's weight, but the applied force is along the surface; thus, it does not affect the vertical forces. In (e), a component of the applied force is parallel to the normal force, thus increasing the normal force by the amount of that component. The two forces are directed downward.

EVALUATE Often, the normal force is not equal to an object's weight. A common mistake initially encountered in force problems is assuming that the normal force is always equal to some object's weight. You must analyze all problems carefully to determine the proper normal force.

## 3: Acceleration and tension in blocks connected by a rope

Consider the situation shown in Figure 5.3. Cart A is placed on a table and is connected to block B by a rope that passes over a frictionless pulley.


Figure 5.4 Question 3.
(a) How does the acceleration of cart A compare with that of block B?

## Solution

IDENTIFY, SET UP, AND EXECUTE Both objects must accelerate at the same rate, since they are connected by the rope (as long as the rope doesn't stretch). To get a better intuitive grasp of this statement, note that cart A will move 10 cm when block B moves 10 cm . If block B moves the 10 cm in 1 second, then cart A moves 10 cm in 1 second; their velocities are the same. If block B's velocity changes by $2 \mathrm{~m} / \mathrm{s}$ in 1 second, then block A's velocity must change by $2 \mathrm{~m} / \mathrm{s}$ in 1 second; their accelerations are the same. We say that the rope constrains both objects to accelerate at the same rate.
(b) How does the tension force acting on cart A compare with the weight of block B as the system accelerates?

## Solution

IDENTIFY, SET UP, AND EXECUTE The tension force in the string is constant along the string, so the tension force is the same on block B as it is on cart A. Therefore, we compare the tension at block B with block B's weight. Newton's second law tells us that the net force on an object is equal to its mass times its acceleration. Two forces act on block B: B's weight and the tension force. The net force on block B is its weight minus the tension force. This net force must be equivalent to the acceleration multiplied by block B's mass; therefore, the tension must be less than the weight. The tension force acting on cart $A$ is less than the weight of block $B$.

EVALUATE Solving this problem gives us two important results that we'll apply repeatedly to later problems: First, objects connected by a rope are constrained to have the same magnitude of acceleration; second, the tension force in a rope connected to an object is not always equal to the object's weight if the object is accelerating.

## 4: What can a hanging ball indicate

A ball hangs on a string attached to the top of a box, as shown in Figure 5.4. The box is placed on a horizontal truck bed and the truck moves over a flat roadway. You observe the ball and find that it remains in the position shown in the figure for a long time. By looking only inside the box, what can be determined about the truck's motion?


Figure 5.5 Question 4.

## Solution

IDENTIFY, SET UP, AND EXECUTE We see that the ball has swung to the left, so we may suspect that the truck is moving. Let's look at the forces acting on the ball: to investigate the motion. Two forces act on the ball: gravity and the tension force due to the string. Figure 5.5 shows the free-body diagram.


Figure 5.6 Question 4 free-body diagram.

We see that the tension force has components in both the vertical and horizontal direction. The vertical component of the tension force must be equivalent to the force of gravity, since the box moves horizontally and there is no vertical acceleration. There is only one horizontal force-the horizontal component of the tension force-so there is horizontal acceleration to the right.

We can conclude that the truck is accelerating to the right in the horizontal direction and not at all in the vertical direction. We cannot determine the velocity of the truck; the horizontal velocity components could be zero or nonzero. For example, the truck could be moving to the left with decreasing velocity or it could be accelerating to the right from rest. Both of these motions would result in the ball's being swung to the left.

EVALUATE This problem shows that the absence of a net force results in zero acceleration of an object and the presence of a net force results in a nonzero acceleration of an object. We cannot determine the velocity of an object by knowing only its acceleration; we need additional information.

Practice Problem: How would the ball appear if the box moved with constant velocity? Answer: The ball would hang vertically, since no net force would act on it.

## 5: Motion of a box on a rough surface

A constant horizontal force is applied to a box on a rough floor. With a 15 N applied force, the box begins to slide. What is the motion of the box after it begins to slide, assuming that the applied force remains constant?

## Solution

IDENTIFY, SET UP, AND EXECUTE Before the box slides, there is static friction. Once it begins to slide, the static friction becomes kinetic friction. Kinetic friction is smaller in magnitude than static friction; therefore, the applied force must be larger than the kinetic friction force, and the box accelerates.

EVALUATE This problem helps illustrate the fact that kinetic friction is generally less than static friction. The reason is that the coefficient of kinetic friction is less than the coefficient of static friction.

## 6: Frictional forces

A box is placed on a rough floor. When you push horizontally against the box with a 20 N force, the box just begins to slide. What is the magnitude of the frictional force when you push against the box with a 10 N force? With a 15 N force?

## Solution

IDENTIFY, SET UP, AND EXECUTE Since the box just begins to slide when the 20 N force is applied, the maximum static friction force is 20 N . When you push against the box with a 10 N force, you are pushing with less than the maximum static friction force. By Newton's third law, the box must push back with the same force; therefore, the static friction force must have a magnitude of 10 N . For the same reason, when you push with a 15 N force, the static friction force has a magnitude of 15 N .

EVALUATE Static friction varies from zero to its maximum value. The static friction force equals the net force acting against the friction force. Be careful not to assume that static friction is always at its maximum.

## 7: A vertical frictional force

A block is placed against the vertical front of an accelerating cart as shown in Figure 5.6. What condition must hold in order to keep the block from falling?


Figure 5.7 Question 7.


Figure 5.8 Question 7 free-body diagram

## Solution

IDENTIFY, SET UP, AND EXECUTE The free-body diagram shown in Figure 5.7 indicates that three forces act on the box: gravity ( $m g$, downwards), the normal force due to the cart ( $n$, to the right), and the friction force at the block-cart surface $(f s)$. To keep the block from falling, the friction force must be equal and opposite to the gravitational force. The condition for equilibrium in the vertical direction gives

$$
\sum F_{y}=f_{s}-m g=0, \quad f_{s}=m g
$$

The friction force must be due to static friction, in order to prevent the block from moving. The static friction force is given by

$$
f_{s} \leq \mu_{s} n
$$

The gravitational force can then be related to the normal force and the coefficient of static friction:

$$
m g \leq \mu_{s} n
$$

The block accelerates to the right with acceleration $a$, so we can use Newton's second law to find an expression for the normal force:

$$
\sum F_{x}=n=m a .
$$

Replacing the normal force with mass and acceleration gives

$$
\begin{aligned}
& m g \leq \mu_{s} m a, \quad g \leq \mu_{s} a . \\
& \mu_{s} \geq \frac{g}{a} .
\end{aligned}
$$

The last inequality tells us that the coefficient of static friction must be equal to or greater than the gravitational constant divided by the cart's acceleration.

EVALUATE Problems that initially appear complicated often have relatively straightforward solutions. Following a consistent problem-solving procedure helps identify the key points that you will need to solve the problem.

## 8: Turning while riding in a car

As you make a right turn in your car, what pushes you against the car door?

## Solution

IDENTIFY, SET UP, AND EXECUTE As your car turns, your body tends to continue moving in a straight line; therefore, you push up against the car door. So you can say that no force actually pushes you against the door and your body tries to maintain a constant velocity due to Newton's first law. Once the car door comes in contact with you, it pushes you in the direction of the turn, accelerating you to the right.

EVALUATE From within the car, you may wonder what force pushes you to the side. However, since the car is turning to the right, it is accelerating, and the car is not an inertial reference frame. Therefore, we cannot apply Newton's laws inside the car. If we consider how the situation would appear to someone outside of the car (in an inertial reference frame), we can apply Newton's laws: You would appear to continue moving in a straight line while the car moves to the right.

## 9: Free-body diagram for a car on a hill

Draw a free-body diagram for a car going over the top of a round hill at a constant speed. Is there a nonzero net force acting on the car if it is moving at constant speed?

## Solution

SET UP AND SOLVE Two forces act on the car at the top of the hill: the normal force due to the road and gravity. The normal force is directed upward and gravity is directed downward. Figure 5.8 shows the free-body diagram.


Figure 5.9 Question 9 free-body diagram.
The force vectors in the free-body diagram are not drawn to have equal length: The gravitational force is larger than the normal force. This is because there is a nonzero net force acting on the car: The car's velocity is changing direction, so the car has a centripetal acceleration. The centripetal acceleration is downward, so the net force must be downward.

There is a net force acting on the car even though the car is moving at constant speed.
REFLECT Constant speed does not necessarily imply constant velocity. You must carefully interpret problems involving circular motion and constant speed.

## Problems

## 1: Equilibrium in two dimensions

A 322 kg block hangs from two cables as shown in Figure 5.9. Find the tension in cables A and B.


Figure 5.10 Problem 1.

## Solution

IDENTIFY The block is in equilibrium, so we can apply Newton's first law to it. The cables have tensions in two dimensions, so we will have to apply the first law to two axes. The target variables are the two tension forces (labeled $T_{A}$ for cable A and $T_{B}$ for cable B).

SET UP The free-body diagram of the block is shown in Figure 5.10. Three forces act on the block: the two tension forces $\left(T_{A}\right.$ and $\left.T_{B}\right)$ and gravity ( $m g$ ). The tensions act in two dimensions and gravity acts in the vertical direction.


Figure 5.11 Problem 1 free-body diagram.

We have added an $x y$ coordinate system to the figure to illustrate in what directions the forces act. We've also resolved the two tension forces into their $x$ and $y$ components.

EXECUTE We apply the equilibrium conditions to the block, writing separate equations for the $x$ and $y$ components:

$$
\begin{array}{ll}
\sum F_{x}=0, & T_{B} \cos 48^{\circ}+\left(-T_{A} \cos 35^{\circ}\right)=0 \\
\sum F_{y}=0, & T_{B} \sin 48^{\circ}+T_{A} \sin 35^{\circ}+(-m g)=0
\end{array}
$$

Note that components directed to the left and downward are negative, consistent with our coordinate system. We can rewrite the first equation as

$$
T_{B}=T_{A} \frac{\cos 35^{\circ}}{\cos 48^{\circ}}
$$

Substituting for $T_{B}$ in the second equation

$$
T_{A} \frac{\cos 35^{\circ}}{\cos 48^{\circ}} \sin 48^{\circ}+T_{A} \sin 35^{\circ}-m g=0
$$

Solving for $T_{A}$ yields

$$
\begin{gathered}
T_{A}\left(\cos 35^{\circ} \tan 48^{\circ}+\sin 35^{\circ}\right)=m g \\
T_{A}=\frac{m g}{\left(\cos 35^{\circ} \tan 48^{\circ}+\sin 35^{\circ}\right)}=\frac{(322 \mathrm{~kg})\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right)}{\left(\cos 35^{\circ} \tan 48^{\circ}+\sin 35^{\circ}\right)}=2130 \mathrm{~N}
\end{gathered}
$$

Substituting the value for $T_{A}$ into the first equation gives

$$
T_{B}=T_{A} \frac{\cos 35^{\circ}}{\cos 48^{\circ}}=2130 \mathrm{~N} \frac{\cos 35^{\circ}}{\cos 48^{\circ}}=2600 \mathrm{~N}
$$

The tension in cable A is 2130 N and the tension in cable B is 2600 N .
EVALUATE The sum of the magnitudes of the two tension forces ( 4730 N ) is larger than the weight of the block $(3160 \mathrm{~N})$. This is consistent because the tension forces are in two dimensions and their magnitudes are greater than their components. In addition, we check that the two horizontal components of the tension forces are equal in magnitude. Substituting into the terms in the first equation gives a magnitude of 1740 N for each component, each having opposite signs.

## 2: Accelerated motion of a block on an inclined plane

A block with mass 3.00 kg is placed on a frictionless inclined plane inclined at $35.0^{\circ}$ above the horizontal and is connected to a second hanging block with mass 7.50 kg by a cord passing over a small, frictionless pulley (See Figure 5.11). Find the acceleration (magnitude and direction) of the 3.00 kg block.


Figure 5.12 Problem 2.

## Solution

IDENTIFY The blocks are accelerating, so we apply Newton's second law to both blocks. The target variable is the acceleration of the blocks (a).

SET UP Both blocks accelerate, so we'll apply Newton's second law to each block to find two equations and solve these equations simultaneously to determine the acceleration. The free-body diagrams of both blocks are shown in Figure 5.12.


Figure 5.13 Problem 2 free-body diagrams.
The forces are identified by their magnitudes. Acting on the left-hand block (block A) are gravity $\left(m_{A} g\right)$, the normal force $(n)$, and the tension force $(T)$. The right-hand block (block B) is acted upon only by the tension force $(T)$ and gravity $\left(m_{B} g\right)$. The tension forces must be equal in magnitude and the accelerations must be equal in magnitude, since the cord connects the two blocks. (See Conceptual Question 3.) As block B accelerates downward, block A accelerates up the ramp. We've added an xy coordinate system separately to each free-body diagram, with the positive axes aligned with the direction of acceleration. Using two different coordinate systems is preferred in these situations. The coordinate system for the right-hand block is rotated to coincide with the inclined plane. A rotated axis simplifies the analysis for ramp problems. This rotated axis requires resolving the gravity force into two components, one parallel, and one perpendicular, to the incline.

EXECUTE We apply Newton's second law to each block. Block A (with mass $m_{A}$ ) accelerates in the $x$ direction (along the ramp), so

$$
\sum F_{x}=T+\left(-m_{A} g \sin 35^{\circ}\right)=m_{A} a .
$$

Block B (with mass $m_{B}$ ) accelerates at the same rate in the $y$ direction; hence,

$$
\sum F_{y}=m_{B} g+(-T)=m_{B} a .
$$

Both equations include the tension force, so we solve for the tension force in the second equation and substitute into the first. Our second equation becomes

$$
T=m_{B} g-m_{B} a .
$$

Replacing the tension force in the first equation yields

$$
\left(m_{B} g-m_{B} a\right)+\left(-m_{A} g \sin 35^{\circ}\right)=m_{A} a,
$$

Solving for the acceleration gives

$$
\begin{aligned}
& m_{B} g+\left(-m_{A} g \sin 35^{\circ}\right)=\left(m_{A}+m_{B}\right) a, \\
& a=\frac{g\left(m_{B}-m_{A} \sin 35^{\circ}\right)}{\left(m_{A}+m_{B}\right)}=\frac{\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right)\left((7.50 \mathrm{~kg})-(3.00 \mathrm{~kg}) \sin 35^{\circ}\right)}{((7.50 \mathrm{~kg})+(3.00 \mathrm{~kg}))}=5.39 \mathrm{~m} / \mathrm{s}^{2} .
\end{aligned}
$$

The 3.00 kg block accelerates up the ramp at $5.39 \mathrm{~m} / \mathrm{s}^{2}$. The positive value of acceleration confirms the block's acceleration up the ramp.

EVALUATE The value of acceleration is less than $g$, consistent with expectations. If the cord were cut, block B would accelerate at $g$. When block B is connected to block A through the cord, block B accelerates with an acceleration less than $g$. We say that block B has additional inertia when connected to block A.

What would have happened if we chose the direction of acceleration incorrectly? We would have found a negative acceleration, indicating that the acceleration was down the incline. In this problem, the forces do not depend on the direction of motion, and a negative acceleration would not indicate an error. It does, however, serve as a checkpoint for our calculation: A negative result with our choice of axes would cause suspicion because the right mass is larger and we expect it to accelerate downward.

Practice Problem: What mass must block A have for the system to remain at rest? Will that mass simply be 7.50 kg ? Answer: $m_{1}=13.1 \mathrm{~kg}$, no.

## 3: Frictional force on an accelerating block

Two blocks are connected to each other by a light cord passing over a small, frictionless pulley as shown in Figure 5.13 . Block A has mass 5.00 kg and block B has mass 4.00 kg . If block B descends at a constant acceleration of $2.00 \mathrm{~m} / \mathrm{s}^{2}$ when set in motion, find the coefficient of kinetic friction between block A and the table.


Figure 5.14 Problem 3.

## Solution

IDENTIFY The target variable is the coefficient of kinetic friction. The blocks are accelerating, so we apply Newton's second law to both blocks. We'll find the friction force and determine the coefficient from that.

SET UP Both blocks accelerate, so we'll apply Newton's second law to each block to find two equations and solve those equations simultaneously. The free-body diagrams of the two blocks are shown in Figure 5.14.


Figure 5.15 Problem 3 free-body diagram.
The forces are identified by their magnitudes. For block A , there is kinetic friction $\left(f_{k}\right)$, gravity $\left(m_{A} g\right)$, the normal force $(n)$, and the tension force $(T)$. For block B there is the tension force $(T)$ and gravity $\left(m_{B} g\right)$. The tension forces are equal in magnitude and the magnitudes of the acceleration are equal, as we have seen. As block B accelerates downward, block A accelerates to the right. To ensure that the acceleration of each block is in the positive direction, we added an $x y$ coordinate system separately to each free-body diagram, with the positive axes aligned with the direction of acceleration. All forces act along the coordinate axes, so we will not need to break the forces into components.

EXECUTE We now apply Newton's second law to each block to find the friction force. Block A (with mass $m_{A}$ ) accelerates in the $x$ direction, so

$$
\sum F_{x}=T+\left(-f_{k}\right)=m_{A} a .
$$

Block B (with mass $m_{B}$ ) accelerates at the same rate in the $y$ direction; thus,

$$
\sum F_{y}=m_{B} g+(-T)=m_{B} a .
$$

Both equations include the tension force, so we solve for the tension force in the second equation and substitute into the first. Our second equation is then

$$
T=m_{B} g-m_{B} a=m_{B}(g-a)
$$

Replacing the tension force in the first equation gives

$$
m_{B}(g-a)+\left(-f_{k}\right)=m_{A} a .
$$

Solving for the friction force yields

$$
f_{k}=m_{B}(g-a)-m_{1} a=(4.00 \mathrm{~kg})\left(9.80 \mathrm{~m} / \mathrm{s}^{2}-2.00 \mathrm{~m} / \mathrm{s}^{2}\right)-(5.00 \mathrm{~kg})\left(2.00 \mathrm{~m} / \mathrm{s}^{2}\right)=21.2 \mathrm{~N} .
$$

The friction force is related to the coefficient of kinetic friction through the normal force. We find the normal force by examining the vertical components of the forces acting on block A . There is no acceleration in the vertical direction for block A, so we can apply the equilibrium condition to block A:

$$
\sum F_{y}=n-m_{A} g=0, \quad n=m_{A} g
$$

Since there are no other vertical forces acting on block $A$, the normal force equals the weight of block $A$. The kinetic frictional force is given by

$$
f_{k}=\mu_{k} n
$$

which we can solve for $\mu_{k}$ :

$$
\mu_{\mathrm{k}}=\frac{f_{k}}{n}=\frac{f_{k}}{m_{A} g}=\frac{(21.2 \mathrm{~N})}{(5.00 \mathrm{~kg})\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right)}=0.43 .
$$

We find the coefficient of kinetic friction between the block and the table to be 0.43 .

EVALUATE A coefficient of kinetic friction equal to 0.43 compares reasonably well with values we've seen previously for smooth surfaces. Note that the tension is not equal to the weight of block B. That it is is a common misconception arising from examining the free-body diagram for block B without realizing that the block is accelerating. If we look at the rearranged second equation, the relation between the tension and the weight of block $B$ becomes clearer:

$$
T=m_{B}(g-a)
$$

The tension force is equal to the weight only when block B's acceleration is zero. Calculating the tension for this problem, we obtain a value of $31.2 \mathrm{~N}, 20 \%$ less than the weight of block $\mathrm{B}(39.2 \mathrm{~N})$.

## 4: Motion of a crate up a rough inclined plane at constant velocity

A student pushes a crate up a rough inclined plane as shown in Figure 5.15. Find the magnitude of the horizontal force the student must apply for the crate to move up the incline at constant velocity. The crate has a mass of 15.0 kg , the incline is sloped at $30.0^{\circ}$, and the coefficient of kinetic friction between the crate and the incline is 0.600 .


Figure 5.16 Problem 4.

## Solution

IDENTIFY There is no acceleration, so we apply the equilibrium condition to the crate to find the applied force.

SET UP The free-body diagram of the crate is shown in Figure 5.16. The forces are identified by their magnitudes: kinetic frictional force $\left(f_{k}\right)$, gravity $(m g)$, the normal force $(n)$, and the applied force $(F)$. Kinetic friction opposes the motion up the incline and thus is directed down the incline. The rotated $x y$ coordinate system is indicated in the diagram. This rotated axis requires resolving the gravity and applied forces into components parallel and perpendicular to the incline, as shown in the diagram.


Figure 5.17 Problem 4 free-body diagram.

EXECUTE We apply the equilibrium condition to the crate. In the $x$-direction, along the incline, we have

$$
\sum F_{x}=F \cos 30^{\circ}+\left(-f_{k}\right)+\left(-m g \sin 30^{\circ}\right)=0
$$

We must apply the equilibrium condition in the $y$-direction to find the normal force in order to quantify the friction force. Thus,

$$
\begin{aligned}
& \sum F_{y}=n+\left(-m g \cos 30^{\circ}\right)+\left(-F \sin 30^{\circ}\right)=0 \\
& n=m g \cos 30^{\circ}+F \sin 30^{\circ}
\end{aligned}
$$

The kinetic friction is then

$$
f_{k}=\mu_{k} n=\mu_{k} m g \cos 30^{\circ}+\mu_{k} F \sin 30^{\circ} .
$$

We now substitute this result into the first equation:

$$
F \cos 30^{\circ}+\left(-\mu_{k} m g \cos 30^{\circ}-\mu_{k} F \sin 30^{\circ}\right)+\left(-m g \sin 30^{\circ}\right)=0 .
$$

Solving for the applied force gives

$$
\begin{aligned}
& F \cos 30^{\circ}-\mu_{k} F \sin 30^{\circ}=\mu_{k} m g \cos 30^{\circ}+m g \sin 30^{\circ} \\
& F\left(\cos 30^{\circ}-\mu_{k} \sin 30^{\circ}\right)=m g\left(\mu_{k} \cos 30^{\circ}+\sin 30^{\circ}\right) \\
& F=\frac{m g\left(\mu_{k} \cos 30^{\circ}+\sin 30^{\circ}\right)}{\left(\cos 30^{\circ}-\mu_{k} \sin 30^{\circ}\right)}=\frac{(15.0 \mathrm{~kg})\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right)\left((0.600) \cos 30^{\circ}+\sin 30^{\circ}\right)}{\left(\cos 30^{\circ}-(0.600) \sin 30^{\circ}\right)}=265 \mathrm{~N} .
\end{aligned}
$$

The student must push with a horizontal force of 265 N to move the crate up the incline at constant velocity.

EVALUATE A force of 265 N is roughly equivalent to the weight of a 27 kg object. Would it be easier to push the crate up the incline by pushing parallel to the incline? Yes, it would be easier to push along the ramp. In this problem, the component of the applied force directed into the incline $\left(F \sin 30^{\circ}\right)$ does nothing to move the crate up the ramp. In fact, this component increases the normal force and therefore the friction force.

Practice Problem: If the student pushed along the incline, the problem would be simplified. What force would be necessary along the incline to maintain the crate at constant velocity? Answer: 150 N , $56 \%$ of the required horizontal force.

## 5: Two blocks suspended by a pulley

Two blocks are connected by a rope that passes over a small, frictionless pulley as shown in Figure 5.17. Find the tension in the rope and the acceleration of the blocks. Block 1 has a mass of 15.0 kg and block 2 has a mass of 8.0 kg .


Figure 5.18 Problem 5.

## Solution

IDENTIFY Our target variables are the tension ( $T$ ) and acceleration (a). We will apply Newton's second law to the problem to find the tension and acceleration.

SET UP The two blocks are separate objects, so we draw free-body diagrams of each block, shown in Figure 5.18. Two forces act on each block: gravity $(\mathrm{mg})$ and the tension force $(T)$. There is no friction in the pulley and the string is considered to be massless, so the tension in the rope is the same throughout.

We assume that the rope doesn't stretch, so the magnitudes of the two accelerations are the same, but the directions are opposite. Included in the diagrams are separate $x y$ coordinate axes with the positive $y$-axis in the direction of the acceleration for both blocks (upward for block 2 and downward for block 1). This choice of axes simplifies our analysis. All of the forces act along the $y$-axes, so we will not need to break the forces into components.


Figure 5.19 Problem 5 free-body diagrams.
EXECUTE Applying Newton's second law to both blocks gives

$$
\begin{aligned}
& \sum F_{y}=m_{1} g-T=m_{1} a \\
& \sum F_{y}=T-m_{2} g=m_{2} a
\end{aligned}
$$

We add both equations to eliminate $T$, leaving

$$
\begin{aligned}
& m_{1} g-m_{2} g=m_{1} a+m_{2} a \\
& a=\frac{\left(m_{1}-m_{2}\right) g}{m_{1}+m_{2}}=\frac{((15.0 \mathrm{~kg})-(8.0 \mathrm{~kg}))\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{(15.0 \mathrm{~kg})+(8.0 \mathrm{~kg})}=2.98 \mathrm{~m} / \mathrm{s}^{2}
\end{aligned}
$$

We find the tension by substituting into the second-law equation for block 2, giving

$$
T=m_{2}(g+a)=(8.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}+2.98 \mathrm{~m} / \mathrm{s}^{2}\right)=102.2 \mathrm{~N}
$$

Both blocks accelerate at $2.98 \mathrm{~m} / \mathrm{s}^{2}$, block 1 downward and block 2 accelerates upwards. The tension in the rope is 102.2 N .

EVALUATE We check that we get the same value for the tension by using the second-law equation for block 1 . We find that the equation gives a tension of 102 N , so the result is the same. We also see that the acceleration is less than $9.8 \mathrm{~m} / \mathrm{s}^{2}$, as is expected, since the net force on either block is less than its weight.

## 6: Friction force between two boxes

Two boxes, one on top of the other, are being pulled up a ramp at constant speed by an applied force, as shown in Figure 5.19. The coefficient of kinetic friction between box $A$ and the ramp is 0.35 , and the
coefficient of static friction between the two boxes is 0.80 . Box $A$ has a mass of 3.00 kg and box $B$ has a mass of 8.00 kg . What is the applied force?


Figure $\mathbf{5 . 2 0}$ Problem 6.

## Solution

IDENTIFY Our target variable is the applied force $F$. We will use Newton's first law to find the forces acting on the two boxes and then solve for $F$.

SET UP The two boxes are separate objects, so we draw free-body diagrams of each box, shown in Figure 5.20. Three forces act on box $A$ : gravity $\left(m_{A} g\right)$, the normal force due to box $B\left(n_{B \text { on } A}\right)$, and static friction $\left(f_{s}\right)$. Static friction must point up the ramp in order for the net force on box $A$ to be zero. Six forces act on box $B$ : the applied force $(F)$, gravity $\left(m_{B} g\right)$, the normal force due to box $A\left(n_{A}\right.$ on $\left.B\right)$, the normal force due to the ramp $\left(n_{\text {ramp }}\right)$, kinetic friction $\left(f_{k}\right)$, and static friction $\left(f_{s}\right)$. Static friction and the normal force due to box $A$ are action-reaction pairs; we set their directions opposite those of the forces acting on box $A$.

Included in the diagrams are a separate $x y$ coordinate axes with the positive $y$-axis in the direction of the motion of both boxes (up along the ramp). We will need to break the forces into components to solve the problem.


Figure 5.21 Problem 6 free-body diagrams.
EXECUTE Applying Newton's first law along both axes to both boxes gives
$\sum F_{x}=f_{\mathrm{s}}-m_{A} g \sin 30.0^{\circ}=0 \quad$ (box $A$ )
$\sum F_{y}=n_{B \text { on } A}-m_{A} g \cos 30.0^{\circ}=0$
$\sum F_{x}=F-f_{k}-f_{\mathrm{s}}-m_{B} g \sin 30.0^{\circ}=0$
(box B)
$\sum F_{y}=n_{\text {ramp }}-n_{A \text { on } B}-m_{B} g \cos 30.0^{\circ}=0$ (box B).
We have four equations and five unknowns. To solve these equations we need the magnitude of the kinetic friction, which is

$$
f_{k}=\mu_{k} n_{\text {ramp }}
$$

We begin by solving the first two equations for $f_{s}$ and $n_{B \text { on } A}$,

$$
\begin{aligned}
& f_{s}=m_{A} g \sin 30.0^{\circ}=(3.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \sin 30.0^{\circ}=14.7 \mathrm{~N} \\
& n_{B \text { on } A}=m_{A} g \cos 30.0^{\circ}=(3.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \cos 30.0^{\circ}=25.5 \mathrm{~N} .
\end{aligned}
$$

Next, we use the equation for the net force along the $y$-axis for box $B$ to solve for the normal force due to the ramp, giving

$$
n_{\text {ramp }}=n_{A \text { on } B}+m_{B} g \cos 30.0^{\circ}=(25.5 \mathrm{~N})+(8.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \cos 30.0^{\circ}=93.4 \mathrm{~N} .
$$

We can finally find the applied force by using the equation for the net force along the $x$-axis for box $B$, along with the kinetic friction. This gives

$$
\begin{aligned}
F & =f_{k}+f_{\mathrm{s}}+m_{B} g \sin 30.0^{\circ} \\
& =\mu_{k} n_{\mathrm{ramp}}+f_{\mathrm{s}}+m_{B} g \sin 30.0^{\circ} \\
& =(0.35)(93.4 \mathrm{~N})+(14.7 \mathrm{~N})+(8.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \sin 30.0^{\circ}=86.6 \mathrm{~N} .
\end{aligned}
$$

The applied force is 86.6 N .
EVALUATE Is the coefficient of static friction large enough to keep box $A$ from sliding off box $B$ ? We can find out by computing the maximum static friction

$$
f_{k}^{\max }=\mu_{s} n_{B \text { on } A}=(0.80)(25.5 \mathrm{~N})=20.4 \mathrm{~N} .
$$

We see that the applied static friction force $(14.7 \mathrm{~N})$ is less than the maximum static friction force ( 20.4 N ), so the box remains in place.

We saw how Newton's third-law force pairs can be useful in identifying forces and their directions. In this problem, we used those pairs to find the directions of forces on box $B$. There are also cases in which the force pairs help us identify missing force pairs. For example, the normal force acting on box $B$ due to box $A$ is often omitted in solutions. If we omitted $n_{A \text { on } B}$, but had labeled the force pairs carefully, we would have realized that there was a missing normal force.

## 7: Acceleration in a two-pulley system

A mass is attached to a rope that is connected to the ceiling and passes through two light, frictionless pulleys. One pulley is attached to the ceiling and a second mass is attached to the other pulley, as shown in Figure 5.21 . Mass 1 is 5.0 kg and mass 2 is 20.0 kg . Find the acceleration of each mass.


Figure 5.22 Problem 7.

## Solution

IDENTIFY Our target variables are the two accelerations, $a_{1}$ and $a_{2}$, of mass 1 and mass 2 , respectively. We will apply Newton's second law to find these accelerations.

SET UP The free-body diagrams for each mass are shown in Figure 5.22. Two forces act on mass 1: gravity $\left(m_{1} g\right)$ and the tension in the rope $(T)$. Three forces act on mass 2: gravity ( $\left.m_{1} g\right)$ and the tension in the rope on both sides of the pulley (two factors of $T$ ).

Included in the diagrams are $x y$ coordinate axes with the positive $y$-axis in the direction of motion for each mass (upward for mass 2 and downward for mass 1). Also shown are the accelerations $a_{1}$ and $a_{2}$ of the two masses.


Figure 5.23 Problem 7 free-body diagrams.
EXECUTE Applying Newton's second law to each mass gives
$\sum F_{y}=m_{1} g-T=m_{1} a_{1} \quad$ (mass 1)
$\sum F_{y}=T+T-m_{2} g=m_{2} a_{2} \quad$ (mass 2 ).
We have three unknowns in these two equations; thus, we will need more information to solve the problem. Let's examine the accelerations. As mass 1 moves down a distance $L$, mass 2 moves up a distance $L / 2$. The change in position of mass 1 is twice the change in position of mass 2 , so the velocity and acceleration of mass 1 are twice the velocity and acceleration of mass 2 ; therefore,

$$
a_{1}=2 a_{2}
$$

We now have enough information to solve the system of equations. We begin by replacing $a_{1}$ by $a_{2}$ :

$$
\begin{aligned}
& m_{1} g-T=m_{1} 2 a_{2} \\
& 2 T-m_{2} g=m_{2} a_{2}
\end{aligned}
$$

Doubling the first equation and adding the two equations together gives

$$
\begin{aligned}
& 2 m_{1} g-m_{2} g=4 m_{1} a_{2}+m_{2} a_{2} \\
& a_{2}=\frac{\left(2 m_{1}-m_{2}\right) g}{4 m_{1}+m_{2}}=\frac{(2(20.0 \mathrm{~kg})-(5.0 \mathrm{~kg}))\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{4(20.0 \mathrm{~kg})+(5.0 \mathrm{~kg})}=4.04 \mathrm{~m} / \mathrm{s}^{2}
\end{aligned}
$$

from which we obtain

$$
a_{1}=2 a_{2}=2\left(4.04 \mathrm{~m} / \mathrm{s}^{2}\right)=8.08 \mathrm{~m} / \mathrm{s}^{2}
$$

Mass 1 accelerates downward at $8.08 \mathrm{~m} / \mathrm{s}^{2}$ and mass 2 accelerates upward at $4.04 \mathrm{~m} / \mathrm{s}^{2}$.
EVALUATE This problem illustrates the fact that not all accelerations are equal when objects are connected by ropes. We must always evaluate the situation carefully to determine the relation between the accelerations.

## 8: Coefficient of friction in a banked curve

A circular section of road with a radius of 150 m is banked at an angle of $12^{\circ}$. What should be the minimum coefficient of friction between the tires and the road if the roadway is designed for a speed of $25 \mathrm{~m} / \mathrm{s}$ ?

## Solution

IDENTIFY Our target variable is the coefficient of static friction, $\mu_{s}$. We will use Newton's second law to find $\mu_{s}$ by finding the friction force.

SET UP Figure 5.23 is a free-body diagram of the car tire on the road, showing the three forces acting on the tire: the normal force due to the road $(n)$, static friction with the road $(f)$, and gravity $(m g)$. For the tire not to slip, the vertical forces must be in equilibrium and there must be a net horizontal force toward the center.


Figure 5.24 Problem 8 free-body diagram.
We have added an $x y$ coordinate system to the figure, since the forces act in two dimensions. Note that we aligned the axes horizontally and vertically to coincide with the directions of the net forces.

SOLVE We apply the force equations to the tire, writing separate equations for the $x$ and $y$ components. In the vertical direction, there is no net force:

$$
\sum F_{y}=0, \quad n \cos 12^{\circ}-f \sin 12^{\circ}-m g=0
$$

In the horizontal direction, we use Newton's second law with centripetal acceleration:

$$
\sum F_{x}=m a_{\mathrm{rad}}, \quad n \sin 12^{\circ}+f \cos 12^{\circ}=m \frac{v^{2}}{r}
$$

Note that the downward components are negative, consistent with our coordinate system. The static friction force can be replaced with $\mu_{s} n$ in our two equations:

$$
\begin{aligned}
& n \cos 12^{\circ}-\mu_{s} n \sin 12^{\circ}-m g=0 \\
& n \sin 12^{\circ}+\mu_{s} n \cos 12^{\circ}=m \frac{v^{2}}{r}
\end{aligned}
$$

We now rewrite the first equation in terms of $n$ and substitute into the second equation:

$$
\begin{aligned}
& n=\frac{m g}{\cos 12^{\circ}-\mu_{s} \sin 12^{\circ}} \\
& \frac{m g}{\cos 12^{\circ}-\mu_{s} \sin 12^{\circ}} \sin 12^{\circ}+\mu_{s} \frac{m g}{\cos 12^{\circ}-\mu_{s} \sin 12^{\circ}} \cos 12^{\circ}=m \frac{v^{2}}{r}
\end{aligned}
$$

The mass cancels and we can solve for $\mu_{s}$ :

$$
\mu_{s}=\frac{v^{2} \cos 12^{\circ}-r g \sin 12^{\circ}}{v^{2} \sin 12^{\circ}+r g \cos 12^{\circ}}=\frac{(25 \mathrm{~m} / \mathrm{s})^{2} \cos 12^{\circ}-(150 \mathrm{~m})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \sin 12^{\circ}}{(25 \mathrm{~m} / \mathrm{s})^{2} \sin 12^{\circ}+(150 \mathrm{~m})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \cos 12^{\circ}}=0.20
$$

The minimum coefficient of static friction between the tire and road is 0.20 .

EVALUATE The technique for solving this problem is similar to those set forth earlier in this chapter. The differences here were the inclusion of centripetal acceleration and the choice of axes that corresponded to our knowledge of the net forces.
Practice Problem: What speed would require no frictional force? Answer: $v=18 \mathrm{~m} / \mathrm{s}$.

## 9: Normal force on a roller coaster

A roller coaster has a vertical loop of radius 45 m . If the roller coaster operates at a constant speed of $35 \mathrm{~m} / \mathrm{s}$ while in the loop, what normal force does the seat provide for a 75 kg passenger at the top of the loop? The roller coaster is upside down at the top of the loop.

## Solution

IDENTIFY Our target variable is the normal force $n$. We will apply Newton's second law to find the force.

SET UP Figure 5.24 shows a free-body diagram of the passenger on the roller coaster. The figure shows the two forces acting on the passenger: the normal force due to the seat ( $n$ ) and gravity ( $m g$ ). At the top of the loop, the net force is downward and the person is accelerating toward the center. We have added an $x y$ coordinate system to the figure, with positive forces directed downward.


Figure 5.25 Problem 9 free-body diagram
SOLVE The net force on the passenger is directed downward. Newton's second law with centripetal acceleration gives

$$
\sum F_{y}=m a_{r a d}, \quad n+m g=m \frac{v^{2}}{r} .
$$

Solving for the normal force, we obtain

$$
n=m \frac{v^{2}}{r}-m g=(75 \mathrm{~kg}) \frac{(35 \mathrm{~m} / \mathrm{s})^{2}}{(45 \mathrm{~m})}-(75 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=1300 \mathrm{~N}
$$

The seat exerts a force of 1300 N on the passenger.
REFLECT We see that the seat provides a force nearly twice the passenger's weight. A seat belt would not be needed to prevent a fall from the roller coaster at the top of the loop. Amusement park rides get their reputation for excitement from their ability to rapidly change the magnitudes and directions of forces applied to passengers. The normal force of the seat on the passenger is even larger at the bottom of the loop.

Practice Problem: What normal force does the seat provide at the bottom of the loop? Answer: $N=2800 \mathrm{~N}$.

## 10: Investigating a tetherball

A tetherball is attached to a vertical pole with a 2.0 m length of rope, as shown in Figure 5.25. If the rope makes an angle of $25.0^{\circ}$ with the vertical pole, find the time required for one revolution of the tetherball.


Figure 5.26 Problem 10.

## Solution

IDENTIFY Our target variable is the period of one revolution, $T$. The period will be found from the ball's velocity and circumference. The velocity will be found by applying Newton's second law to the ball as it undergoes centripetal acceleration.

SET UP The free-body diagram of the tetherball is shown in Figure 5.26. Two forces act on the ball: gravity ( $m g$ ) and the tension in the rope ( $T$ ). An $x y$ coordinate axis is included in the diagram.


Figure 5.27 Problem 10 free-body diagram.
EXECUTE The tetherball undergoes centripetal acceleration in the horizontal direction and no net force in the vertical direction. Applying Newton's first law in the vertical direction gives

$$
\begin{aligned}
\sum F_{y} & =T \cos 25.0^{\circ}-m g=0 \\
T & =\frac{m g}{\cos 25.0^{\circ}}
\end{aligned}
$$

Applying Newton's second law in the horizontal direction gives

$$
\begin{aligned}
\sum F_{x} & =T \sin 25.0^{\circ}=m a_{\mathrm{rad}}=m \frac{v^{2}}{R} \\
v^{2} & =T \frac{R}{m} \sin 25.0^{\circ}=\frac{m g}{\cos 25.0^{\circ}} \frac{R}{m} \sin 25.0^{\circ}=R g \tan 25.0^{\circ}
\end{aligned}
$$

We now have the velocity in terms of the radius, angle, and $g$. To find the radius, we use the length of the rope and the sine:

$$
R=l \sin 25.0^{\circ}=(2.0 \mathrm{~m}) \sin 25.0^{\circ}=0.845 \mathrm{~m}
$$

Thus,

$$
v=\sqrt{R g \tan 25.0^{\circ}}=\sqrt{(0.845 \mathrm{~m})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \tan 25.0^{\circ}}=1.97 \mathrm{~m} / \mathrm{s}
$$

The period is the time required for one re volution. The ball travels the circumference of a circle of radius $R$ in one period. The period is then

$$
T=\frac{2 \pi R}{v}=\frac{2 \pi(0.845 \mathrm{~m})}{(1.97 \mathrm{~m} / \mathrm{s})}=2.69 \mathrm{~s}
$$

The ball completes one revolution in 2.69 s .
EVALUATE This problem illustrates the inclusion of centripetal acceleration into force problems. We see that Newton's laws remain unchanged. We simply set the acceleration equal to the radial acceleration in these problems.
Practice Problem: Does the period increase or decrease with larger angles? Find the period when the rope makes an angle of $50.0^{\circ}$ with respect to the vertical. Answer: 2.28 s ; the period decreased slightly.

## 11: A rotating mass

A mass attached to a vertical post by two strings rotates in a circle of constant velocity $v$. (See Figure 5.27.) At high velocities, both strings are taut. Below a critical velocity, the lower string slackens. Find the critical velocity. The mass is 1.0 kg and is 1.5 m from the post.


Figure 5.28 Problem 11.

## Solution

IDENTIFY Our target variable is the critical velocity, at which the tension in the lower string is zero. We will find the tension in the lower string and the velocity by applying Newton's second law. Then we'll set the tension in the lower string to zero and solve for the critical velocity.

SET UP The free-body diagram of the mass is shown in Figure 5.28. Three forces act on the mass: gravity ( mg ), the tension in the upper rope $\left(T_{1}\right)$, and the tension in the lower rope $\left(T_{2}\right)$. An xy coordinate axis is included in the diagram.


Figure 5.29 Problem 11 free-body diagram.

EXECUTE The mass undergoes centripetal acceleration in the horizontal direction and no net force in the vertical direction. Applying Newton's first law in the vertical direction gives

$$
\sum F_{y}=T_{1} \sin 30^{\circ}-T_{2} \sin 30^{\circ}-m g=0
$$

Applying Newton's second law in the horizontal direction gives

$$
\sum F_{x}=T_{1} \cos 30^{\circ}+T_{2} \cos 30^{\circ}=m a_{\mathrm{rad}}=m \frac{v^{2}}{R}
$$

We can solve for $T 1$ by multiplying the first equation by $\cos 30^{\circ}$, multiplying the second equation by $\sin 30^{\circ}$, and adding the two resulting equations. Doing this gives

$$
2 T_{1} \sin 30^{\circ} \cos 30^{\circ}=m g \cos 30^{\circ}+m \frac{v^{2}}{R} \sin 30^{\circ}
$$

which reduces to

$$
T_{1}=\frac{m g}{2 \sin 30^{\circ}}+\frac{m v^{2}}{2 R \cos 30^{\circ}}
$$

Solving for $T_{2}$ produces

$$
T_{2}=\frac{-m g}{2 \sin 30^{\circ}}+\frac{m v^{2}}{2 R \cos 30^{\circ}}
$$

At the critical velocity, $T_{2}$ is zero. Solving for $v$, we obtain

$$
\begin{aligned}
& \frac{m g}{2 \sin 30^{\circ}}=\frac{m v^{2}}{2 R \cos 30^{\circ}}, \\
& v=\sqrt{\frac{g R}{\tan 30^{\circ}}}=\sqrt{\frac{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(1.5 \mathrm{~m})}{\tan 30^{\circ}}}=5.04 \mathrm{~m} / \mathrm{s} .
\end{aligned}
$$

The critical velocity is $504 \mathrm{~m} / \mathrm{s}$.
EVALUATE For velocities less than the critical velocity, the bottom string is slack. Strings always have positive tensions. When they are slack, they provide no support.

Practice Problem: What is the tension in the upper string at the critical velocity? Answer: 20 N .

## Try It Yourself!

## 1: Constant velocity on an incline

Two weights are attached by a light cord that passes over a light, frictionless pulley as shown in Figure 5.29. The left weight moves up a rough ramp. Find the weight $w_{2}$ necessary to keep $w_{1}(15.0 \mathrm{~N})$ moving up the ramp at a constant rate once it is put in motion. The coefficient of kinetic friction is 0.25 and the ramp is inclined at $30.0^{\circ}$.


Figure 5.30 Try it yourself 1.

## Solution Checkpoints

IDENTIFY AND SET UP To move at constant velocity, the net force on each weight must be zero. The rough surface indicates that there is friction between the weight and the ramp. Draw a free-body diagram and apply Newton's first law to solve.
EXECUTE The net force acting on weight 1 along the ramp is

$$
\sum F_{x}=T-f-w_{1} \sin \theta=0
$$

The net force acting on weight 2 in the vertical direction is

$$
\sum F_{y}=T-w_{2}=0
$$

These two equations have three unknowns. You need to find another expression to solve for $w_{2}$. The expression will lead to

$$
w_{2}=\mu_{k} w_{1} \cos \theta+w_{1} \sin \theta=11.1 \mathrm{~N}
$$

EVALUATE Can you explain why $w_{2}$ has less weight than $w_{1}$ ?

## 2: Box sliding across rough floor

A box is kicked, giving it an initial velocity of $2.0 \mathrm{~m} / \mathrm{s}$. It slides across a rough, horizontal floor and comes to rest 1.0 m from its initial position. Find the coefficient of friction.

## Solution Checkpoints

IDENTIFY AND SET UP Begin with a sketch and free-body diagram. In the horizontal direction, the only force acting on the box is friction. Use the acceleration of the box and Newton's second law to find the friction.

EXECUTE The net horizontal force acting on the box is

$$
\sum F_{x}=f=m a .
$$

To find the acceleration, apply

$$
v^{2}=v_{0}^{2}+2 a \Delta x
$$

one of the kinematics relations for constant acceleration that we've used in previous chapters. You should find that the coefficient of kinetic friction is 0.20 .

EVALUATE We see how we can combine kinematics with our force problems in this problem. Did we find a reasonable coefficient of kinetic friction?

## 3: Three connected masses

Three blocks are attached to each other with ropes that pass over pulleys as shown in Figure 5.30. The masses of the ropes and pulleys can be ignored, and there is no friction on the surface over which the blocks slide or in the pulleys. Find the tension in the two ropes and the acceleration of the blocks.


Figure 5.31 Try it yourself 2.

## Solution Checkpoints

IDENTIFY AND SET UP Three objects are in motion, so you must draw three free-body diagrams. How does the acceleration of all three objects compare? You can apply Newton's laws to solve. Assume that the system accelerates clockwise.

EXECUTE The net forces acting in the direction of the acceleration of the blocks are

$$
\begin{aligned}
& \sum F_{A}=T-T^{\prime}-m_{A} g \sin \theta=m_{A} a \\
& \sum F_{B}=m_{B} g-T=m_{B} a \\
& \sum F_{C}=T^{\prime}-m_{C} g=m_{C} a .
\end{aligned}
$$

We can also find the net force acting on block $A$ normal to the ramp. Using these equations to solve for the tensions gives

$$
\begin{aligned}
T & =\left(\frac{2 m_{C} m_{B}+m_{B} m_{A}(1+\sin \theta)}{m_{A}+m_{B}+m_{C}}\right) g \\
T^{\prime} & =\left(\frac{2 m_{C} m_{B}+m_{C} m_{A}(1-\sin \theta)}{m_{A}+m_{B}+m_{C}}\right) g .
\end{aligned}
$$

Solving for the acceleration yields

$$
a=\left(\frac{m_{B}-m_{A} \sin \theta-m_{C}}{m_{A}+m_{B}+m_{C}}\right) g .
$$

EVALUATE How do we interpret the results if we find that the acceleration is negative? Do we need to rework the solution? Is the acceleration greater or less than $g$ ?

## 4: Tension along a heavy rope

A heavy rope of mass 10.0 kg and length 5.0 m lies on a frictionless horizontal surface. If a certain horizontal force is applied, the rope accelerates at $1.5 \mathrm{~m} / \mathrm{s}^{2}$. Find the tension in the rope at any point along its length.

## Solution Checkpoints

IDENTIFY AND SET UP To find the tension at any point in the rope, you must break the rope into many pieces and find the tension of any piece. It is easiest to pick a piece of the rope at the end where the force is applied, as shown in Figure 5.31. By applying Newton's second law to this piece, you can find the tension anywhere in the rope.

(a)

(b)

Figure 5.32 Try it yourself 3 .
EXECUTE The net force required to accelerate the chain is 15.0 N . The mass of the length $x$ of the rope is

$$
m_{x}=\frac{x}{L} m .
$$

The net force on the piece of rope is

$$
\sum F_{x}=F-T=m_{x} a=\frac{x}{L} m a .
$$

Solving for $T$ to find the tension as a function of $x$ :

$$
T=F\left(1-\frac{x}{L}\right)=(10.0 \mathrm{~N})\left(1-\frac{x}{5.0 \mathrm{~m}}\right)
$$

EVALUATE This example illustrates how to apply Newton's second law to more complicated problems. Does the tension in the rope increase or decrease as you move away from the end where the force is applied?

## 5: Riding a carousel

A 100.0 kg man stands on the outer edge of a carousel of 4.0 m radius. The coefficient of static friction between his shoes and the carousel is 0.30 . What is the minimum period of rotation required for the man to remain on the carousel?

## Solution Checkpoints

IDENTIFY AND SET UP Three forces act on the man: gravity, the normal force, and friction. The condition required for him to slip is that the centripetal force exceed the static friction. The period is determined from the velocity.

EXECUTE The net horizontal force acting on the man is

$$
\sum F_{x}=f=m a_{\mathrm{rad}}=m \frac{v^{2}}{R}
$$

This yields the critical velocity when the friction force is equal to the centripetal force. For velocities larger than the critical velocity, the man will slide off. The period is found from the velocity, given that the time required for one revolution is the distance traveled divided by the velocity. The minimum period is 7.4 s .

EVALUATE We see that the period and velocity are inversely related: Larger velocities result in shorter periods. How does the friction force vary when the period is greater than 7.4 s ?

## Problem Summary

Chapters 4 and 5 have examined a variety of problems with applied forces in diverse applications, but they share a common problem-solving foundation. For all problems, we

- Identified the general procedure for finding the solution.
- Sketched the situation when no figure was provided.
- Identified the forces acting on the objects of interest.
- Drew free-body diagrams of forces acting on the objects.
- Added appropriate coordinate systems to the free-body diagrams.
- Applied the equilibrium condition, Newton's second law, or both to the objects in order to find relations among the forces, masses, and accelerations.
- Solved the equations through algebra, trigonometry, and calculus.
- Reflected on the results, thus checking for inconsistencies.

This problem-solving foundation can be applied to all problems involving forces. Following this procedure enables one to master Newton's laws.

## 6 <br> Work and Kinetic Energy

## Summary

We introduce two new concepts in this chapter: work and energy. Our investigation begins by learning how work can be used to solve problems with variable forces. We will see how work is a form of energy transfer, leading us to learn about energy, one of the most important concepts in physics. We'll learn how work can be used to change a body's kinetic energy (the energy of motion), how to determine the work expended in many situations, and how power is the rate of change of work with respect to time. In the next chapter, we will introduce the law of conservation of energy and discover other forms of energy. The problem-solving skills we develop in these two chapters will prepare us for additional forms of energy that we'll encounter in later chapters.

## Objectives

After studying this chapter, you will understand

- The definition of work and how to calculate the work done by a force on a body.
- The definition and interpretation of kinetic energy.
- How to apply the work-energy theorem to problems.
- How to use kinetic energy and work in problems involving varying forces applied along curved paths.
- How to analyze springs and the elastic force.
- The definition of power and how to calculate power for bodies performing work or on which work is performed.

Concepts and Equations

| Term | Description |
| :---: | :---: |
| Work Done by a Force | A constant force acting on and displacing a body does work. For a constant force $\vec{F}$ acting on a particle causing a straight-line displacement $\vec{s}$ at an angle $\phi$ with respect to the force, the work done by the force on the body is $W=\vec{F} \cdot \vec{s}=F s \cos \phi .$ <br> The SI unit of work is 1 joule $=1$ newton-meter $(1 \mathrm{~J}=1 \mathrm{~N} \cdot \mathrm{~m})$. |
| Kinetic Energy | Kinetic energy $K$ is the energy of motion of a particle with mass. Kinetic energy is equal to the amount of work required to accelerate a particle from rest to a speed $v$. A particle of mass $m$ and velocity $v$ has kinetic energy $K=\frac{1}{2} m v^{2} .$ |
| Work-Energy Theorem | The work-energy theorem states that the total work done by a net external force on a particle as it undergoes a displacement is equal to the change in kinetic energy of the particle: $W_{\mathrm{tot}}=K_{2}-K_{1}=\Delta K .$ |
| Work Done by a Varying Force along a Curved Path | The work done by a varying force on a particle as it follows a curved path is determined by $W=\int_{P_{1}}^{P_{2}} \vec{F} \cdot d \vec{l}=\int_{P_{1}}^{P_{2}} F \cos \phi d l=\int_{P_{1}}^{P_{2}} F_{s} d l .$ |
| Elastic Force | An elastic force is a force that restores a body to its original equilibrium position after deformation. For a spring, the deformation is approximately proportional to the applied force, as given by Hooke's law, $F_{\mathrm{spr}}=k x,$ |

where $k$ is the force constant and $x$ is the displacement of the spring from its equilibrium position.

## Power

Power is the rate of change of work with respect to time. Average power is defined as

$$
P_{\mathrm{av}}=\frac{\Delta W}{\Delta t},
$$

where $\Delta W$ is the quantity of work performed during the time interval $\Delta t$.
Instantaneous power is defined as

$$
P=\lim _{\Delta t \rightarrow 0} \frac{\Delta W}{\Delta t}=\frac{d W}{d t} .
$$

For a force acting on a moving particle, the instantaneous power is

$$
P=\vec{F} \cdot \vec{v} .
$$

The SI unit of power is 1 watt $=1$ joule/second $(1 \mathrm{~W}=1 \mathrm{~J} / \mathrm{s})$.

## Conceptual Questions

## 1: Work done by the normal force

How much work does the normal force do on a box sliding across the floor?

## Solution

IDENTIFY, SET UP, AND EXECUTE Work is produced when a force acts on an object in the direction the object is displaced. As a box slides across the floor, the normal force is perpendicular to its motion. Therefore, the normal force does no work on the sliding box.

EVALUATE Work has a strict definition in physics. The normal force prevents the box from falling into the floor, but it does no work thereby. You also do no work as you carry your backpack, even though you arm tires.

## 2: Which force does the most work?

Rank the four situations shown in Figure 6.1 from most to least work done by the force. The displacement is the same in each case.


Figure 6.1 Conceptual Question 2.

## Solution

IDENTIFY, SET UP, AND EXECUTE Work is the dot product of force and displacement, or displacement times the component of force parallel to the displacement. All four situations have the same displacement, so we rank the components of the force parallel to the displacement. The parallel components of the forces are the magnitudes of the forces times the cosine of the angle between the force and the displacement vectors.

We find that the parallel components for the four situations are, in order, $3.47 \mathrm{~N}, 4.00 \mathrm{~N}, 5.36 \mathrm{~N}$, and 5.00 N . Therefore, the ranking from greatest to least amount of work is (c), (d), (b), and (a).

EVALUATE We see that larger forces do not necessarily produce more work. Work depends on both the magnitude of the force and the direction of the force with respect to the displacement. The 20 N force produces much less work than the 5 N force.

## 3: Ranking stopping distance

Five identical shipping crates slide down a ramp onto a rough horizontal floor. Each crate carries a different mass and has a different velocity at the bottom of the ramp, given in Table 1. Rank the distances required, from least to greatest, for each crate to stop.

TABLE 1: Conceptual Problem 3.

|  | Mass (kg) | Velocity at bottom of ramp $(\mathbf{m} / \mathbf{s})$ |
| :--- | :---: | :---: |
| Crate 1 | 20.0 | 10.0 |
| Crate 2 | 30.0 | 20.0 |
| Crate 3 | 20.0 | 15.0 |
| Crate 4 | 100.0 | 10.0 |
| Crate 5 | 200.0 | 5.0 |

## Solution

IDENTIFY, SET UP, AND EXECUTE At the bottom of the ramp, each crate has kinetic energy. As they slow, the crates lose kinetic energy because friction does work on them. The change in kinetic energy is equal to the work done by friction, which in turn is equal to the friction force ( $\mu \mathrm{mg}$ ) times the displacement $(x)$. Since the crates are identical, their coefficient of friction is the same. Algebraically, the work done is

$$
\begin{aligned}
W & =\Delta K=K_{2}-K_{1}, \\
\mu m g x & =0-\frac{1}{2} m v^{2} \\
\mu g x & =-\frac{1}{2} v^{2} .
\end{aligned}
$$

We see that the displacement is proportional to the velocity squared. To rank the stopping distances, we compare the velocities at the bottom of the ramp.

Crate 5 has the smallest velocity, crates 1 and 4 have the next-largest velocity, crate 3 has the nextlargest velocity, and crate 2 has the largest velocity. The ranking of stopping distance for the crates is $(5),(1)=(4),(3)$, and (2).

EVALUATE We see that in this case the results do not depend on the mass of the object. Crate 4 has much more mass than crate 1 , so crate 4 has more initial kinetic energy. However, the effect of friction on crate 4 is greater than the effect of friction on crate 1 . These two effects cancel, resulting in the same stopping distance.

We also see that the work done by friction is negative, because the change in kinetic energy is negative.

## 4: Work in carnival ride

A swing ride at a carnival consists of chairs attached by a cable to a vertical pole. The vertical pole rotates, causing the chairs to swing in a circle as shown in Figure 6.2. How much work is done by the tension in the cable as one chair (with a mass of 70 kg ) makes one complete revolution? The length of the rope and angle are shown in the figure.


## Solution

IDENTIFY, SET UP, AND EXECUTE As the chair swings, it is displaced along the circle. The tension in the cable is directed perpendicular to the chair's displacement. Since work is the dot product of the force and displacement, the work done by the tension is zero.

EVALUATE Does the kinetic energy change? The chair moves at a constant rate, so the kinetic energy remains constant. No work is needed to keep the chair moving at the same speed.

## 5: Work with equal forces

A force is applied to several boxes. Rank the five situations shown in Figure 6.3 from least to most work done by the force. The same magnitude of force is applied to all boxes, and the mass of each box is given.

> (a)

(b)

(d)


## Solution

IDENTIFY, SET UP, AND EXECUTE Work is the dot product of force and displacement. All five situations have the same magnitude of applied force, so we rank them on the basis of the displacement and the cosine of the angle between the force and the resulting displacement.

In cases (a) and (b), the displacement is in the direction of the force. The box in (b) is displaced further, so more work is done by the force. In case (c), the force is applied at an angle of $45^{\circ}$ with respect to the displacement and the box is displaced 10.0 m , resulting in less work than in (b), but more than in (a). The work in case (d) is negative, since the force is opposite the displacement. In case (e), no work is done, since the force is perpendicular to the displacement.

The amount of work done by the force, from least to most, is (d), (e), (a), (c), and (b).
EVALUATE We see that the mass of the box does not influence the results: The work done by the force depends only on the force, the displacement, and the angle between them.

How much work is done on the box in (e)? The box could be moving at constant speed on a frictionless surface, in which case no work is done. Or another force may be acting on the box, creating work. From the information we have, we cannot determine which is correct.

We also see that we can have negative work in problems. Negative work indicates that energy is being removed from the system, perhaps as the box slows to a stop.

## Problems

## 1: Work done in pushing a crate up an inclined plane.

A student pushes a crate 3.50 m up a rough inclined plane with a constant horizontal force of 225 N , starting from rest as shown in Figure 6.4. Find the work done by the student, the work done by friction, the work done by gravity, and the change in the crate's kinetic energy. How does the work done by the student, friction, and gravity compare with the change in kinetic energy? The crate has a mass of 15.0 kg , the incline is sloped at $30.0^{\circ}$, and the coefficient of kinetic friction between the crate and the incline is 0.400 .


Figure 6.4 Problem 1.

## Solution

IDENTIFY Each force is constant and the displacement is along a straight line, so we can find the work and kinetic energy from their definitions.

SET UP We'll begin with a free-body diagram to find the work done by the three forces. Figure 6.5 shows the free-body diagram with a rotated coordinate system that coincides with the incline to simplify the analysis. The forces are identified by their magnitudes: kinetic friction $\left(f_{k}\right)$, gravity $(m g)$, the normal force ( $n$ ), and the force applied by the student $(F)$. Kinetic friction opposes the motion up the incline and thus is directed down the incline. Gravity and the applied force are resolved into components parallel and perpendicular to the incline.


Figure 6.5 Problem 1 freebody diagram.

EXECUTE The work done by the student pushing the crate is

$$
W_{\text {student }}=F s \cos \phi=(225 \mathrm{~N})(3.50 \mathrm{~m}) \cos 30^{\circ}=682 \mathrm{~J}
$$

where the angle between the force and the displacement is $30^{\circ}$. To find the work done by friction, we need to know the friction force. We apply the equilibrium condition in the $y$ direction to find the normal force in order to quantify the friction force:

$$
\begin{aligned}
& \sum F_{y}=n+\left(-m g \cos 30^{\circ}\right)+\left(-F \sin 30^{\circ}\right)=0 \\
& n=m g \cos 30^{\circ}+F \sin 30^{\circ}=(15.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)\left(\cos 30^{\circ}\right)+(225 \mathrm{~N})\left(\sin 30^{\circ}\right)=240 \mathrm{~N}
\end{aligned}
$$

The kinetic friction force is then

$$
f_{k}=\mu_{k} n=(0.400)(240 \mathrm{~N})=95.9 \mathrm{~N}
$$

Friction is directed opposite to the displacement, so the angle between them is $180^{\circ}$. The work done by friction is

$$
W_{f_{k}}=f_{k} s \cos \phi=(95.9 \mathrm{~N})(3.50 \mathrm{~m}) \cos 180^{\circ}=-336 \mathrm{~J}
$$

The component of gravity along the displacement is also opposite to the displacement. The work done by gravity is

$$
W_{\mathrm{grav}}=F_{g} s \cos \phi=m g s \cos 120^{\circ}=(15.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(3.50 \mathrm{~m}) \cos 120^{\circ}=-257 \mathrm{~J}
$$

To find the change in kinetic energy, we need the initial and final velocities. The initial velocity is zero. The final velocity is found by applying Newton's second law and kinematics. In the $x$-direction, along the incline, Newton's second law gives

$$
\sum F_{x}=F \cos 30^{\circ}+\left(-f_{k}\right)+\left(-m g \sin 30^{\circ}\right)=m a .
$$

The acceleration of the box is then

$$
\begin{aligned}
a & =\frac{F \cos 30^{\circ}+\left(-f_{k}\right)+\left(-m g \sin 30^{\circ}\right)}{m} \\
& =\frac{\left.(225 \mathrm{~N}) \cos 30^{\circ}-(95.9 \mathrm{~N})-(15.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \sin 30^{\circ}\right)}{(15.0 \mathrm{~kg})}=1.70 \mathrm{~m} / \mathrm{s}^{2}
\end{aligned}
$$

Constant-acceleration kinematics gives the final velocity:

$$
\begin{aligned}
v^{2} & =v_{x 0}^{2}+2 a_{x}\left(x-x_{0}\right) \\
v & =\sqrt{0+2\left(1.70 \mathrm{~m} / \mathrm{s}^{2}\right)(3.50 \mathrm{~m})}=3.45 \mathrm{~m} / \mathrm{s}
\end{aligned}
$$

The change in kinetic energy is then

$$
\Delta K=K_{2}-K_{1}=\frac{1}{2} m v^{2}-0=\frac{1}{2}(15.0 \mathrm{~kg})(3.45 \mathrm{~m} / \mathrm{s})^{2}=89.3 \mathrm{~J}
$$

In sum, we found that the student did 682 J of work on the crate, friction did -336 J of work on the crate, gravity did -257 J of work, and the kinetic energy increased by 89 J . When we add the work due to the three forces together, the total work is 89 J . The total work is equal to the change in kinetic energy.

EVALUATE This problem affords a thorough investigation of work and kinetic energy. It illustrates how to combine the work due to several forces into the total work and how the total work on a system is used to increase the kinetic energy of the system.

CAUTION Watch Signs! You must evaluate the signs carefully when determining work and energy. Negative work indicates that the force is directed opposite to the displacement and often slows the object, as it does in this problem.

## 2: Spring force between two blocks

Two blocks are placed on a horizontal, frictionless surface and attached to each other by a spring with force constant $4500 \mathrm{~N} / \mathrm{m}$. If the right-hand block is pulled with a force of 150.0 N , find the displacement of the spring as the blocks accelerate. The left-hand block has a mass of 5.00 kg , and the righthand block has a mass of 3.00 kg .

## Solution

IDENTIFY Our target variable is the displacement of the spring. We can solve this problem with Newton's second law.

SET UP The displacement of the spring is proportional to the spring force. We find the spring force by applying Newton's second law to the blocks and then use Hooke's law to find the displacement. The first task is to sketch the situation, as shown in Figure 6.6.


Figure 6.6 Problem 2.
Examining the sketch, we see the two blocks interact through the spring. We draw free-body diagrams for the two blocks, shown in Figure 6.7.


Figure 6.7 Problem 2 free-body diagram.

The forces are identified by their magnitudes: force due to the spring $\left(F_{\mathrm{spr}}\right)$, gravity ( $m g$ ), the normal force ( $n$ ), and the applied force $\left(F^{\prime}\right.$ ). Knowledge of the vertical forces will not be necessary to solve this problem. The blocks are connected to each other; therefore, both accelerate at the same rate and are acted upon by the same magnitude of spring force. The diagrams include a common $x y$ coordinate axis, with the positive $x$-axis in the direction of the acceleration. All of the forces act along the coordinate axes, so we will not need to break the forces into components.

EXECUTE We apply Newton's second law to the horizontal forces acting on each block to determine the force due to the spring. For the right-hand block (with mass $m_{R}$ ),

$$
\sum F_{x}=F+\left(-F_{\mathrm{spr}}\right)=m_{R} a .
$$

For the left-hand block (with mass $m_{L}$ ),

$$
\sum F_{x}=F_{\mathrm{spr}}=m_{L} a .
$$

Examining these two equations, we find two unknowns: $F_{\text {spr }}$ and $a$. We wish to find the spring force, so we rewrite the second equation in terms of acceleration:

$$
a=\frac{F_{\mathrm{spr}}}{m_{L}}
$$

Substituting for the acceleration in the first equation yields

$$
\begin{aligned}
& F+\left(-F_{\mathrm{spr}}\right)=m_{R} \frac{F_{\mathrm{spr}}}{m_{L}} \\
& F=F_{\mathrm{spr}}+F_{\mathrm{spr}} \frac{m_{R}}{m_{L}}=F_{\mathrm{spr}}\left(1+\frac{m_{R}}{m_{L}}\right) \\
& F_{\mathrm{spr}}=\frac{F}{\left(1+\frac{m_{R}}{m_{L}}\right)}=F \frac{m_{L}}{m_{R}+m_{L}}=(150) \frac{(5.00 \mathrm{~kg})}{(3.00 \mathrm{~kg})+(5.00 \mathrm{~kg})}=93.8 \mathrm{~N}
\end{aligned}
$$

This gives us the magnitude of the force due to the spring. The direction is opposite to the displacement. We can now use Hooke's law,

$$
F_{\mathrm{spr}}=k x,
$$

to solve for the displacement:

$$
x=\frac{F_{\mathrm{spr}}}{k}=\frac{(93.8 \mathrm{~N})}{4500 \mathrm{~N} / \mathrm{m}}=0.0208 \mathrm{~m}=2.08 \mathrm{~cm}
$$

The spring is displaced 2.08 cm when the blocks are pulled.
EVALUATE We see that the force due to the spring is less than the applied force in this problem. That is reasonable, as the spring must provide force to accelerate only the right-hand mass, while the applied force must accelerate both masses.

## 3: Block stopped by spring and friction

A 5.00 kg block is moving along a rough horizontal surface toward a spring with force constant $500 \mathrm{~N} / \mathrm{m}$. The velocity of the block just before it contacts the spring is $12.0 \mathrm{~m} / \mathrm{s}$. If the coefficient of kinetic friction between the block and the surface is 0.400 , what is the maximum compression of the spring?

## Solution

IDENTIFY We can apply the work-energy theorem to the problem. Two forces (friction and the spring force) do work to slow the block, taking away the kinetic energy from the block. The target variable is the spring's maximum compression.

SET UP Figure 6.8 shows a sketch of the situation. Just before contacting the spring, the block has kinetic energy and the spring is uncompressed. At maximum compression, the spring has been compressed a distance $X$, the block is not moving $(K=0)$, and friction has done work on the block as it
moves the same distance $X$. The work-energy theorem tells us that the change in kinetic energy is equal to the total work done.


Figure 6.8 Problem 3 sketch.
Four forces act on the block as it slows: the force due to the spring ( $F_{\text {spr }}$ ), gravity ( $m g$ ), the normal force $(n)$, and the force of kinetic friction $\left(f_{k}\right)$. These forces are shown in the free-body diagram in Figure 6.9. The normal force and gravity do no work on the block as it slows.


Figure 6.9 Problem 3 free-body diagram.

EXECUTE The change in kinetic energy of the block is

$$
\Delta K=K_{2}-K_{1}=0-\frac{1}{2} m v^{2}=-\frac{1}{2} m v^{2} .
$$

The work done on the spring by the block is

$$
W_{\text {Block on Spring }}=\frac{1}{2} k x_{2}^{2}-\frac{1}{2} k x_{1}^{2}=\frac{1}{2} k X^{2} .
$$

The work done by the spring on the block is the negative of this value:

$$
W_{\text {Spring }}=-\frac{1}{2} k X^{2} .
$$

The work done by friction on the block is

$$
W_{f}=F s \cos \phi=(\mu m g) s \cos 180^{\circ}=-\mu m g X
$$

The total work is the work due to the spring and the work due to friction. Setting the sum of these two quantities equal to the change in kinetic energy gives

$$
-\frac{1}{2} k X^{2}-\mu m g X=-\frac{1}{2} m v^{2}
$$

This is a quadratic equation. Solving for $X$ yields

$$
\begin{aligned}
X & =\frac{-(\mu m g) \pm \sqrt{(\mu m g)^{2}-4\left(\frac{1}{2} k\right)\left(-\frac{1}{2} m v^{2}\right)}}{2\left(\frac{1}{2} k\right)} \\
& =\frac{-(19.6) \pm \sqrt{(19.6)^{2}+4(15.0)(22.5)}}{2(15.0)} m=0.735,-2.04 \mathrm{~m} .
\end{aligned}
$$

In this case, we require the positive root, 0.735 m . The spring compresses 0.735 m when the block comes to a momentary stop.

EVALUATE This problem illustrates how to use energy to work with varying forces. The analysis would have been more difficult had we used Newton's second law, as we would then have had to integrate the force equation.

CAUTION Watch signs in work done by a spring! As you pull on a spring, you do positive work on it and the spring does negative work on you. Understanding the signs for work done on or by a spring will help you understand the proper signs for work and energy.

## 4: A novel spring

Suppose you have invented a novel spring that will slow a 2.5 kg toy car. The spring exerts a force that depends on position such that $F_{x}=\left[10.0 \mathrm{~N}+\left(16.0 \mathrm{~N} / \mathrm{m}^{2}\right) x^{2}\right]$. What maximum speed of the toy car will the spring stop with a compression of 0.50 m ?

## Solution

IDENTIFY The maximum work due to the spring occurs when the displacement is directed along the force, so we will take that direction as the direction of displacement. Our target variable is the speed of the toy car, and we will find it by using the work-energy theorem.

SET UP Only the force due to the novel spring acts on the toy car. Initially, there is just the kinetic energy of the car. At the maximum compression, there is no kinetic energy.

EXECUTE The force of the spring is in the direction of the motion. The work that the spring does on the toy car as the spring is compressed a distance $X$ is

$$
W=\int_{0}^{x} F_{x} d x
$$

Substituting for the force and finding the work when $X=0.50 \mathrm{~m}$ yields

$$
\begin{aligned}
W & =\int_{0}^{x}\left[10.0 \mathrm{~N}+\left(16.0 \frac{\mathrm{~N}}{\mathrm{~m}^{2}}\right) x^{2}\right] d x \\
& =\left.\left[(10.0 \mathrm{~N}) x+\left(16.0 \frac{\mathrm{~N}}{\mathrm{~m}^{2}}\right) \frac{x^{3}}{3}\right]\right|_{0} ^{X} \\
& =\left[(10.0 \mathrm{~N}) X+\left(16.0 \frac{\mathrm{~N}}{\mathrm{~m}^{2}}\right) \frac{X^{3}}{3}\right] \\
& =\left[(10.0 \mathrm{~N})(0.50 \mathrm{~m})+\left(16.0 \frac{\mathrm{~N}}{\mathrm{~m}^{2}}\right) \frac{(0.50 \mathrm{~m})^{3}}{3}\right] \\
& =6.33 \mathrm{~J} .
\end{aligned}
$$

The spring can stop a toy car with up to 6.33 J of energy. The velocity of that car is determined from the kinetic energy:

$$
\begin{aligned}
& K=\frac{1}{2} m v^{2}, \\
& v=\sqrt{\frac{2 K}{m}}=\sqrt{\frac{2(6.33 \mathrm{~J})}{(2.5 \mathrm{~kg})}}=2.25 \mathrm{~m} / \mathrm{s} .
\end{aligned}
$$

The maximum velocity that the spring can stop is $2.25 \mathrm{~m} / \mathrm{s}$.
EVALUATE In this problem, the work is negative because the force due to the spring acts opposite to the direction of the displacement in order to reduce the kinetic energy. The change in kinetic energy is also negative, so the work-energy theorem is satisfied.

## 5: Average power to run an escalator

What average power does an escalator require to lift twenty 100.0 kg people 3.0 m high in 1 minute?

## Solution

IDENTIFY The target variable is the power, or the amount of work done per unit time.
SET UP The escalator provides work equivalent to the amount of work due to gravity as the people are lifted. We will find the work done by gravity and divide by the time.

EXECUTE The work done by gravity is the weight of the people, multiplied by their change in height:

$$
W=m g h=(20 \times 100 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(3.0 \mathrm{~m})=58,800 \mathrm{~J} .
$$

The power is the work of gravity divided by time:

$$
P=\frac{\Delta W}{\Delta t}=\frac{(58,800 \mathrm{~J})}{(60 \mathrm{~s})}=980 \mathrm{~W} .
$$

The power needed to run the escalator is 980 W .
EVALUATE We see that the power is independent of the angle of the escalator. This is due to the fact that the gravitational work depends only on the change in vertical elevation.

## Try It Yourself!

## 1: Box on a smooth incline

A 10.0 kg box is pushed 2.0 m up a smooth inclined plane of angle $30.0^{\circ}$ by a 100 N horizontal force, starting from rest. Find the work done on the box and the change in kinetic energy.

## Solution Checkpoints

IDENTIFY AND SET UP Start with a free-body diagram to identify the forces acting on the box. The work due to each force is equal to the dot product of the force and the displacement. The change in kinetic energy is equal to the total work done.

EXECUTE Three forces act on the box. The work done by each is

$$
\begin{aligned}
W_{\text {applied }} & =F s \cos \phi=173 \mathrm{~J}, \\
\mathrm{~W}_{\text {normal }} & =0, \\
W_{g} & =-98 \mathrm{~J} .
\end{aligned}
$$

The change in kinetic energy is 75 J .
EVALUATE You can check your results by using Newton's second law to find the acceleration, which leads to the final velocity of the box and final kinetic energy of the box. Do they agree?

## 2: Body sliding on a rough surface

A body slides on a rough surface. If the body is given an initial velocity of $3.0 \mathrm{~m} / \mathrm{s}$, it comes to a stop in 1.0 m . Find the coefficient of kinetic friction between the body and the surface.

## Solution Checkpoints

IDENTIFY AND SET UP Use the work-energy theorem to relate the change in kinetic energy to the work done by friction. Only one force acts in the direction of motion.

EXECUTE Set the change in kinetic energy equal to the work done by friction. This results in a coefficient of kinetic friction equal to 0.50 .

EVALUATE Is the change in kinetic energy and in the work negative or positive? Why?

## 3: Drag on an automobile

An automobile has a 150 hp engine and a top speed of 100 mph . If you assume that half of the power of the engine is delivered to the tires on the road, find the net drag (air resistance and other dissipative forces) on the automobile.

## Solution Checkpoints

IDENTIFY AND SET UP At constant velocity, the net force on the car must be zero and the force acting on the tires must be equal to the drag forces. Force is related to power.

EXECUTE The power is equivalent to the force multiplied by the velocity. The units need to be converted to solve the problem:

$$
100 \mathrm{mph}=147 \mathrm{ft} / \mathrm{s}
$$

The force is

$$
F=\frac{\frac{1}{2} P}{v}=281 \mathrm{lb}=1250 \mathrm{~N}
$$

EVALUATE Why was the power divided by 2 to get the force?

## Potential Energy and Energy Conservation

## Summary

In this chapter, we'll continue our investigation of energy by defining potential energy and learning about conservation of energy. Potential energy is a form of energy storage that applies to gravitational and elastic forces. Conservation of energy is one of the most fundamental concepts in physics, and we will learn how it can be applied to problems. We will learn the difference between conservative and nonconservative forces. We'll conclude by learning how to find forces, given a potential-energy function. By the end of the chapter, we'll be able to apply energy concepts to the analysis of problems and be prepared to extend our methods to additional forms of energy that we'll encounter in later chapters.

## Objectives

After studying this chapter, you will understand

- The definition of potential energy.
- How to use gravitational potential energy and elastic potential energy in a variety of problems.
- The definitions of conservative and nonconservative forces.
- How to apply conservation of energy to problems.
- How to find the force, given a potential-energy function.

Concepts and Equations

| Term | Description |
| :--- | :--- |
| Gravitational Potential Energy | Gravitational potential energy is the potential energy associated with the posi- <br> tion of a particle relative to earth. For a particle of mass $m$ at a vertical dis- <br> tance $y$ above the origin in a uniform gravitational field $g$, the gravitational <br> potential energy of the system is <br> $\quad U_{\text {grav }}=m g y$. <br> Gravitational potential energy does not depend upon the location of the ori- <br> gin; only differences in gravitational potential energy are significant. |
| Elastic Potential Energy | Elastic potential energy is the potential energy associated with an ideal <br> spring. For a spring of force constant $k$ stretched or compressed a distance $x$ <br> from equilibrium, the elastic potential energy is <br> $U_{\text {el }}=\frac{1}{2} k x^{2}$. |
| Conservation of Mechanical Energy | When only conservative forces act on a particle, the total mechanical energy <br> is constant; that is, $\quad K_{1}+U_{1}=K_{2}+U_{2}$, |

Nonconservation of Mechanical Energy

|  | c |
| :--- | :---: |
| Conservative Forces and | F |
| Conservation of Energy |  |

where $U$ is the sum of the gravitational and elastic potential energies.
When forces other than gravitation or elastic forces do work on a particle, the work $W_{\text {other }}$ done by these other forces equals the change in the total mechanical energy:

$$
K_{1}+U_{1}+W_{\text {other }}=K_{2}+U_{2} .
$$

Forces are either conservative or nonconservative. Conservative forces are forces for which the work-kinetic-energy theorem is completely reversible and the work can be represented by potential-energy functions. Work done by nonconservative forces manifests itself as changes in the internal energy of the object. The sum of the kinetic, potential, and internal energy is always conserved:

$$
\Delta K+\Delta U+\Delta U_{\mathrm{int}}=0
$$

## Determining Force from

Potential Energy

A conservative force is the negative derivative of its potential-energy function in one, two, or three dimensions:

$$
\begin{aligned}
F_{x}(x) & =-\frac{d U(x)}{d x}, \\
F_{x} & =-\frac{\partial U}{\partial x}, \\
F_{y} & =-\frac{\partial U}{\partial y}, \\
F_{z} & =-\frac{\partial U}{\partial z}, \\
\vec{F} & =-\left(\frac{\partial U}{\partial x} \hat{i}+\frac{\partial U}{\partial y} \hat{j}+\frac{\partial U}{\partial z} \hat{k}\right) .
\end{aligned}
$$

## Conceptual Questions

## 1: Launching a ball

A compressed spring is used to shoot a ball straight up into the air. Compressing the spring a distance of 10 cm results in a maximum height of 3.2 m . How high does the ball go if the spring is compressed 5.0 cm ?

## Solution

IDENTIFY, SET UP, AND EXECUTE The spring stores elastic potential energy that is converted to gravitational potential energy at the top of the ball's flight. Elastic potential energy is proportional to the displacement squared, and gravitational potential energy is proportional to the height. One-half the compression reduces the potential energy of the spring by a factor of four, so the ball reaches onefourth the height, or 0.8 m .

EVALUATE How does the velocity just above the spring compare for the two cases? Just above the spring, the elastic potential energy has been transformed to kinetic energy. Kinetic energy depends on the velocity squared; therefore, the velocity is proportional to the compression. One-half of the compression results in one-half the velocity just above the spring.

## 2: An accelerating car

A car accelerates from zero to 30 mph in 2.0 s . How long does it take to accelerate from zero to 60 mph ?

## Solution

IDENTIFY, SET UP, AND EXECUTE We assume that the power provided to the wheels is constant and that there is no friction. Kinetic energy is proportional to velocity squared, so a doubling of the final speed requires four times the energy. Power is energy per unit time; therefore, the time required to reach the final speed will increase by a factor of four, assuming constant power. The car will take 8.0 s to accelerate from zero to 60 mph .

EVALUATE This problem illustrates how energy principles provide alternative solutions to our earlier kinematics problems. How would you solve the problem by using forces?

## 3: Multiple routes to bottom of a hill

Figure 7.1 shows four different routes to the bottom of a hill, all starting from the same initial height. If you and three of your friends slide down the four routes, how do the four speeds at the bottom of the hill compare? Each of the paths is frictionless, and everyone starts from rest.


Figure 7.1 Question 3.

## Solution

IDENTIFY, SET UP, AND EXECUTE Gravitational potential energy depends only on changes in height. Therefore, the change in gravitational potential energy is the same for all four routes, as all four have the same change in height. The kinetic energy at the bottom of the hill will be the same for all of the friends; therefore, the speeds of all four friends will be the same at the bottom.

EVALUATE If the four speeds are the same at the bottom, what quantity differs? You can see that the four routes have different lengths and are shaped differently. If you compare path 2 with path 3 , you see that path 3 has a steep initial drop-off while path 2 has a shallower initial drop-off. The friend on path 3 will accelerate faster initially than the friend on path 2 , will have a larger speed throughout, and will arrive at the bottom first. Therefore, the time to reach the bottom differs for the different paths.

## 4: Swinging by vines

Tarzan crosses a river gorge by starting from rest and swinging across the gorge on a vine. Can he ever reach a height above his starting point with this method?

## Solution

IDENTIFY, SET UP, AND EXECUTE From an energy standpoint, Tarzan converts gravitational potential energy into kinetic energy as he swings across the gorge. After he passes the low point of his path, his speed slows as his gravitational potential energy increases. At his starting height, he will come to a momentary stop after all of his kinetic energy has converted to potential energy. To get to a greater final height, he needs additional energy. He could increase his initial energy by starting with an initial velocity, using a running start, for example.

EVALUATE Without additional energy, Tarzan's final height cannot be greater than his initial height.

## Problems

## 1: Velocity of a mass on a string

A mass $m$ is attached by a string of length $l$ to the ceiling and is released from rest at an angle of $60^{\circ}$ from the vertical. Find the velocity as a function of the angle.

## Solution

IDENTIFY Only gravity and tension act on the mass. Tension does no work, so we can use energy conservation to solve the problem. The target variable is the velocity.

SET UP A sketch of the problem is shown in Figure 7.2. At the initial angle, the mass has only gravitational potential energy. As the mass falls, the potential energy transforms to kinetic energy. At any point, the sum of the potential and kinetic energy is equal to the initial gravitational potential energy. We will need to relate the height of the mass to the angle. The origin is placed at the bottom of the mass' path, below the anchor point.


Figure 7.2 Problem 1.
EXECUTE Energy is conserved, so the initial energy is equal to the energy at any other point:

$$
U_{1}+K_{1}=U_{2}+K_{2}
$$

The initial kinetic energy is zero (the mass starts from rest), and the initial potential energy is

$$
U_{1}=m g y_{1}=m g\left(l-l \cos \theta_{0}\right),
$$

where the initial angle $\theta_{0}=60.0^{\circ}$. At any other point, the kinetic and potential energies are

$$
\begin{aligned}
& K_{2}=\frac{1}{2} m v^{2}, \\
& U_{2}=m g y_{2}=m g(l-l \cos \theta) .
\end{aligned}
$$

Equating the energies gives

$$
m g\left(l-l \cos \theta_{0}\right)=\frac{1}{2} m v^{2}+m g(l-l \cos \theta) .
$$

Rearranging to solve for $v$ as a function of the angle gives

$$
\begin{aligned}
v & =\sqrt{2 g l\left(\cos \theta-\cos \theta_{0}\right)} \\
& =\sqrt{2 g l\left(\cos \theta-\frac{1}{2}\right)}
\end{aligned}
$$

EVALUATE What is the maximum velocity? The maximum velocity occurs at the bottom of the swing, where $\theta=0^{\circ}$, and is equal to $v=\sqrt{g l}$.

## 2: Professor landing on spring platform

Your professor, with a mass of 60.0 kg , falls from a height of 2.50 m onto a platform mounted on a spring. As the springs compresses, she compresses the spring a maximum distance of 0.240 m . What is the force constant of the spring? Assume that the spring and platform have negligible mass.

## Solution

IDENTIFY Energy is conserved, as the only forces acting on the professor are gravity and the spring force. The target variable is the force constant of the spring.

SET UP Figure 7.3 shows a sketch of the situation. Initially, the professor has only $U_{\text {grav }}$, since her velocity is zero $(K=0)$ and the spring is uncompressed $\left(U_{\mathrm{el}}=0\right)$. As she falls to the top of the platform, her


Figure 7.3 Problem 2 sketch.
kinetic energy increases and gravitational potential energy decreases. As she starts to compress the spring, she slows down as energy is transformed to the spring's elastic potential energy. At maximum compression, she comes to a momentary stop $(K=0)$. At this point, she is below the origin, so she has negative gravitational potential energy. We'll use energy conservation to solve for the spring's force constant.

EXECUTE Energy conservation relates the initial and final energies:

$$
K_{1}+U_{1}=K_{2}+U_{2}
$$

Initially, there is only $U_{\text {grav }}$. At the maximum compression, there are two potential-energy terms: $U_{\text {grav }}$ and $U_{\text {el }}$. Also,

$$
U_{\mathrm{grav}, 1}=U_{\mathrm{grav}, 2}+U_{\mathrm{el}, 2}
$$

Substituting the expressions for the energies yields

$$
m g y_{1}=m g y_{2}+\frac{1}{2} k x^{2}
$$

The initial height is 2.50 m , and the final height and compression is -0.240 m . Solving for $k$ gives

$$
k=\frac{2 m g\left(y_{1}-y_{2}\right)}{x^{2}}=\frac{2(60.0 \mathrm{~kg})\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right)(2.50 \mathrm{~m}-(-0.240 \mathrm{~m}))}{(0.240 \mathrm{~m})^{2}}=55,900 \mathrm{~N} / \mathrm{m}
$$

The force constant of the spring is $55,900 \mathrm{~N} / \mathrm{m}$.
EVALUATE Our choice of origin gave a negative $y_{2}$, but only differences in gravitational potential energies influence the result. This problem would have been much more challenging to solve with our force techniques, as the force of the spring varies with position.

CAUTION You set zero for gravitational potential energy! Only differences in gravitational potential energy are useful in energy problems. You may set the zero at any point. It is best to choose one that simplifies the solution.

## 3: Designing a bungee jump

You are entering the bungee-jumping business and must design the bungee cord. The jump will be from a bridge that is 100.0 m above a river. The design calls for 2.00 seconds of free fall before the cord begins to slow the fall, and the person just touches the water after jumping. Find the force constant and length of the bungee cord for a 100.0 kg person.

## Solution

IDENTIFY The forces acting on the jumper are gravity and the spring force. There is no mechanical work done on the system, so we will use energy conservation to solve the problem. Our target variables are the force constant and the length of the bungee cord.

SET UP Figure 7.4 shows a sketch of the situation with the coordinate origin at the river. On the bridge, the jumper has gravitational potential energy. After he jumps, the energy transforms to kinetic and elastic potential energies. At the river, the jumper momentarily stops and all the energy has transformed into elastic potential energy.


Figure 7.4 Problem 3.

We'll also need to recall our kinematics for freely falling objects to find the length of the bungee cord. The length of the bungee cord is found by determining its length when it becomes taut. We know that the cord becomes taut after 2.00 s , so we can use free-fall kinematics to solve for the distance the person falls in 2.00 s , which is equal to the length of the bungee cord. We'll ignore air resistance and any friction in the bungee cord.

EXECUTE Energy conservation relates the initial and final energies:

$$
K_{1}+U_{1}=K_{2}+U_{2}
$$

Initially, there is only $U_{\text {grav }}$ at the bridge. At the river, there is only $U_{\text {el }}$. Hence,

$$
U_{\mathrm{grav}}=U_{\mathrm{el}} .
$$

Replacing the energies gives

$$
m g y_{1}=\frac{1}{2} k x^{2},
$$

where $m$ is the mass of the jumper, $y_{1}$ is height of the bridge, $x$ is the stretch length of the bungee cord, and $k$ is the force constant of the bungee cord. We need the amount of stretch in the bungee cord, so we first use constant $=$ acceleration kinematics for freely falling objects to find the position where the bungee cord becomes taut:

$$
y_{\mathrm{taut}}=y_{0}+v_{0 y} t+\frac{1}{2} a_{y} t^{2} .
$$

Here, $v_{0 y}$ is zero as the person starts from rest, $y_{0}=100 \mathrm{~m}, a_{y}=-g$, and $t$ is 2.00 s . Solving for the position where the bungee cord becomes taut, we have

$$
y_{\text {taut }}=100 \mathrm{~m}+\frac{1}{2}\left(-9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(2.00 \mathrm{~s})^{2}=80.4 \mathrm{~m} .
$$

Thus, $y_{\text {taut }}=80.4 \mathrm{~m}$ is the vertical position above the river where the bungee cord becomes taut. The starting point was at $y_{0}=100 \mathrm{~m}$, so the length of the bungee cord is the difference between $y_{0}$ and $y_{\text {taut }}: 100 \mathrm{~m}-80.4 \mathrm{~m}=19.6 \mathrm{~m}$. The stretch length is how much the cord is stretched from its original length, or 80.4 m in this case (i.e., the distance from where the cord becomes taut to the river). Substituting into our energy expression to solve for the force constant yields

$$
k=\frac{2 m g\left(y_{0}\right)}{x^{2}}=\frac{2(100.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(100.0 \mathrm{~m})}{(80.4 \mathrm{~m})^{2}}=30.3 \mathrm{~N} / \mathrm{m}
$$

You will need a bungee cord that is 19.6 m long with a $30.3 \mathrm{~N} / \mathrm{m}$ force constant.
EVALUATE The spring constant was found to be relatively small, indicating that the person will be slowed gently. What will happen to a person with a mass of less than 100 kg ? What about a person with a mass greater than 100 kg ? The lighter person has less initial energy and so stops above the river. The heavier person has more initial energy and so stops under the surface of the river. (So the cord should be changed!)

## 4: Toy car loop-the-loop

A toy car is released from a spring launcher onto a horizontal track that leads to a vertical loop-theloop, as shown in Figure 7.5. What is the minimum compression needed for the launcher so that, when released, the car remains on the track throughout the loop? The mass of the car is 10.0 g , the force constant of the launcher is $20.0 \mathrm{~N} / \mathrm{m}$, the loop has a radius of 20.0 cm , and you may assume that the car moves along the track without friction.


Figure 7.5 Problem 4.

## Solution

IDENTIFY We will use both energy conservation and Newton's second law to solve the problem. The target variable is minimum compression of the spring.

SET UP The forces acting on the car are gravity, the spring force of the launcher while the car is in contact with the launcher, and the normal force of the ground or loop. There is no mechanical work done on the system, so energy is conserved.

For the car to remain in contact with the track at the top of the loop, it must have sufficient velocity to maintain centripetal force. A free-body diagram for the car at the top of the loop is shown in Figure 7.6.


Figure 7.6 Problem 4 free-body diagram.

We place the origin at ground level. Our initial point (1) will be when the car is at rest and the launcher is compressed, storing all the energy in the spring. After the car is released, the elastic potential energy is transformed into kinetic energy and then into a combination of kinetic and gravitational potential energy when the car enters the loop. The final point (2) will be at the top of the loop, where there are both kinetic and gravitational potential energies. We'll use energy conservation to solve for the spring's compression.

EXECUTE Energy conservation relates the initial and final energies:

$$
K_{1}+U_{1}=K_{2}+U_{2} .
$$

Initially, there is only $U_{\text {el }}$ stored in the spring. At the top of the loop, both $K$ and $U_{\text {grav }}$ are stored, and

$$
U_{\mathrm{el}}=U_{\mathrm{grav}}+K
$$

Substituting with the expressions for the energies gives

$$
\frac{1}{2} k x^{2}=m g y_{2}+\frac{1}{2} m v^{2},
$$

where $x$ is the spring compression, $k$ is the force constant of the spring, $m$ is the mass of the car, $y_{2}$ is twice the radius of the loop (the height at the top of the loop), and $v$ is the speed of the car at the top of the loop. We find the velocity at the top of the loop by applying Newton's second law. To find the minimum compression of the spring, we need the minimum velocity at the top of the loop. The minimum velocity corresponds to the minimum force on the car at the top; therefore, the only force acting on the car at the top is gravity, so

$$
\sum F_{y}=m g=m a_{\mathrm{rad}}=\frac{m v^{2}}{r} .
$$

Solving for $v$ yields the velocity at the top of the loop:

$$
v=\sqrt{g r}=\sqrt{\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right)(0.200 \mathrm{~m})}=1.40 \mathrm{~m} / \mathrm{s} .
$$

Combining the results and solving for the displacement of the spring gives

$$
x=\sqrt{\frac{m\left(4 g r+v^{2}\right)}{k}}=\sqrt{\frac{0.0100 \mathrm{~kg}\left(4\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right)(0.200 \mathrm{~m})+(1.40 \mathrm{~m} / \mathrm{s})^{2}\right)}{20.0 \mathrm{~N} / \mathrm{m}}}=0.0700 \mathrm{~m} .
$$

The minimum spring compression necessary to keep the car on the track throughout the loop is 7.00 cm .
EVALUATE We found the minimum compression of the spring. Additional compression would have resulted in greater total energy after the car is launched, which would also keep the car on the track.

This problem illustrates how we'll sometimes combine our knowledge of previous materials (e.g., forces) with our current topics. As we progress through the text, we will add to our knowledge base and not merely exchange one concept for another.

## 5: Losing contact with the hill

A frictionless puck slides down a large, round dome of radius 2.0 m . If the puck starts at the top of the dome with a very small initial velocity, how far below the starting point does the puck lose contact with the dome?

## Solution

IDENTIFY Gravitation and the normal force are the only forces acting on the puck. The normal force does no work, so we will use energy conservation. At the point where the puck loses contact with the dome, the normal force is zero. The target variable is the height at which the puck loses contact.

SET UP A sketch of the problem is shown in Figure 7.7. The origin is at the center of the dome and the target variable is the change in height, $y$. No mechanical work is done on the system, so energy is conserved.


Figure 7.7 Problem 5.
The forces acting on the puck are gravity and the normal force, and the net force is a centripetal force directed toward the center of the dome. When the puck loses contact with the dome, the normal force is zero. A free-body diagram of the puck is shown in Figure 7.8.


Figure 7.8 Problem 5 free-body diagram.

Our initial point (1) will be when the puck is at the top of the dome, where there is only gravitational potential energy. (We ignore the small quantity of kinetic energy at the top of the dome, to simplify the solution.) Our second point (2) will be the moment the puck leaves the dome, when there are both kinetic and gravitational potential energies.

EXECUTE We write the change in height in terms of the radius and angle:

$$
y=R-R \cos \theta
$$

We need to find the angle at which the puck leaves the dome. We start with the forces acting on the puck. The net force along the radius at any point is given by Newton's second law:

$$
\sum F_{\mathrm{rad}}=m g \cos \theta-n=m a_{\mathrm{rad}}=m \frac{v^{2}}{R}
$$

When the puck loses contact, the normal force goes to zero, producing the following relation between the angle and the velocity:

$$
v^{2}=g R \cos \theta
$$

We now have $\theta$ in terms of velocity. The $v^{2}$ reminds us of energy, so we write the expression for energy conservation:

$$
U_{1}+K_{1}=U_{2}+K_{2}
$$

The potential energy at the top of the dome is $m g R$ and the kinetic energy is zero. At a later point, the puck has both kinetic energy and potential energy. Substituting in the expressions for the energies gives

$$
m g R=m g(R \cos \theta)+\frac{1}{2} m v^{2} .
$$

Substituting the expression from the forces results in

$$
m g R=m g(R \cos \theta)+\frac{1}{2} m g R \cos \theta .
$$

Simplifying this equation yields

$$
\frac{2}{3} R=R \cos \theta \text {. }
$$

Solving for $y$, we obtain

$$
\begin{aligned}
y & =R-R \cos \theta \\
& =R-\frac{2}{3} R \\
& =\frac{1}{3} R=\frac{1}{3}(2.0 \mathrm{~m})=0.67 \mathrm{~m} .
\end{aligned}
$$

The puck leaves the dome 0.67 m below the top.
EVALUATE We see that the result depends on neither the puck's mass nor gravity. If the puck and dome were transported to the moon, the puck would lose contact at the same position as on earth.

## 6: Force from potential-energy function

The potential-energy function of a particle is

$$
U(x, y)=a x y-3 b x^{2}+2 c y^{2}
$$

where $a, b$, and $c$ are constants. What is the force on the particle?

## Solution

IDENTIFY AND SET UP Given the potential-energy function, we find the force by taking partial derivatives. The potential-energy function depends on $x$ and $y$, so we will find the negative partial derivative of the potential-energy function with respect to $x$ and $y$.

EXECUTE The $x$ component of the force is

$$
F_{x}=-\frac{\partial U}{\partial x} .
$$

Substituting the expression $U$ and solving gives

$$
F_{x}=-\frac{\partial}{\partial x}\left(a x y-3 b x^{2}+2 c y^{2}\right)=a y-6 b x .
$$

The $y$ component of the force is

$$
F_{y}=-\frac{\partial U}{\partial y} .
$$

Substituting again and solving yields

$$
F_{y}=-\frac{\partial}{\partial y}\left(a x y-3 b x^{2}+2 c y^{2}\right)=a x+4 c y
$$

The force is

$$
\vec{F}=(a y-6 b x) \hat{i}+(a x+4 c y) \hat{j}
$$

EVALUATE This problem illustrates how we can find the force, given a potential-energy function. We see that it is easier to find the force from the potential energy than the potential energy from the force.

CAUTION Use both forces and energy! This problem shows how you can combine your knowledge of both forces and energy to solve complex problems. Without using the combined technique, the solution would have been more difficult.

## Try It Yourself!

## 1: Mass on a spring

A 0.5 kg box hangs from a spring whose unstretched length is 1.0 m . The box stretches the spring 0.5 m . The box is pulled 0.5 m from its equilibrium position and is released from rest. Find its maximum velocity and height.

## Solution Checkpoints

IDENTIFY AND SET UP Start by finding the spring constant by using information from the first part of the problem. What forces act on the box? Can you use energy conservation to find the position and velocity of the box at any point? The initial kinetic energy of the box is zero when it is released. What is the initial elastic potential energy?

EXECUTE Energy conservation gives

$$
\frac{1}{2} k(1.0 \mathrm{~m})^{2}=m g y+\frac{1}{2} m v^{2}+\frac{1}{2} k(1.0 \mathrm{~m}-y)^{2}
$$

when the gravitational potential energy is initially zero. This gives a maximum velocity of $2.2 \mathrm{~m} / \mathrm{s}$ when $y=0.5 \mathrm{~m}$ and a maximum height of 1.0 m when $v=0$.

EVALUATE Do you get the same result when you use a different origin?

## 2: Two masses connected by a pulley

For the frictionless system shown in Figure 7.9, find the velocity of the mass $m$ when it hits the floor if it is released from a height $L$ above the floor.


Figure 7.9 Problem 1.

## Solution Checkpoints

IDENTIFY AND SET UP What forces act on the masses? Can you use energy conservation in this case? The initial kinetic energy is zero when mass $m$ is released.

EXECUTE Energy conservation gives

$$
m g L=\frac{1}{2}\left(m+m^{\prime}\right) v_{2}^{2}+m^{\prime} g L \sin \theta
$$

when the origin is at the bottom of the incline. This equation results in a velocity of

$$
v_{2}=\sqrt{2 \frac{\left(m-m^{\prime}\right)}{\left(m+m^{\prime}\right)} g L .}
$$

EVALUATE What is the velocity of $m$ as $\theta$ approaches $90^{\circ}$ (an Atwood's machine)? What is the velocity as $\theta$ approaches $0^{\circ}$ (a flat table)? Do these results agree with those of earlier problems in Chapter 5?

## 3: Force from the potential-energy function

The potential-energy function of a particle is

$$
U(x)=a x+\frac{1}{2} k x^{2} .
$$

What is the force on the particle? What is the equilibrium position of the particle?

## Solution Checkpoints

IDENTIFY AND SET UP How do you find the force, given a potential-energy function? The equilibrium position is where the force is zero and is found by setting the force equal to zero and solving for position.

EXECUTE The force on the particle is the negative derivative with respect to position:

$$
F_{x}=-(a+k x) .
$$

The force is zero when $x=-a / k$.
EVALUATE We see that this force is an elastic force.

## Problem Summary

The previous two chapters have augmented our knowledge of forces and kinematics with our newly formed knowledge of energy analysis. Energy analysis shares many of the problem-solving principles we have encountered. In these problems, we

- Identified the general procedure to find the solution.
- Sketched the situation when no figure was provided.
- Identified the energies involved and the forces acting in the system.
- Applied energy principles, including conservation of energy (when possible).
- Drew free-body diagrams of the objects when appropriate.
- Applied Newton's laws when appropriate.
- Solved the equations through algebra and substitutions.
- Reflected on the results, checking for inconsistencies.

This problem-solving foundation can be applied to all problems, including those that could be solved by force analysis. Following the procedure set forth here leads to mastery of many physics problems.

## Momentum, Impulse, and Collisions

## Summary

In this chapter, we'll introduce two new concepts-momentum and impulse-which together will serve as our third major analysis technique in mechanics. Like energy analysis, momentum and impulse analysis will expand our problem-solving repertoire and allow us to tackle collision problems that would be challenging with Newton's laws. Also, like energy, momentum is a conserved quantity that has important consequences throughout physics. We'll be able to apply momentum and impulse analyses to a wide variety of problems by the end of the chapter, and, when those analyses are combined with force and energy analyses, we will be able apply a powerful set of tools to investigate many natural phenomena.

## Objectives

After studying this chapter, you will understand

- The definition of momentum and the distinction between momentum and velocity.
- The definition of impulse and the distinction between impulse and force.
- How to restate Newton's law in terms of momentum and impulse.
- How to use conservation of momentum to solve a variety of collision problems.
- How to identify elastic, inelastic, and totally inelastic collisions and how to apply conservation of energy appropriately to each of these situations.
- The definition of the center of mass of a system and how to use the center of mass to solve problems.
- How to apply momentum conservation to rocket propulsion problems in which the rocket's mass changes.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Momentum | The momentum $\vec{p}$ of a particle of mass $m$ moving with velocity $\vec{v}$ is defined as |
| $\qquad \vec{p}=m \vec{v}$. |  |
|  | Newton's second law states that the net force on a particle is equal to the rate | of change of momentum of the particle:

$$
\sum \vec{F}=\frac{d \vec{p}}{d t}
$$

The total momentum $\vec{P}$ of a system of particles is the vector sum of the individual momenta:

$$
\vec{P}=\vec{p}_{A}+\vec{p}_{B}+\cdots=m_{A} \vec{v}_{A}+m_{B} \vec{v}_{B}+\cdots
$$

## Impulse

## Conservation of Momentum

| $\qquad \vec{P}=$ constant if $\sum \vec{F}=0$. |  |
| :--- | :--- |
| Elastic Collision | Each component of momentum is separately conserved. | | In an elastic collision between two bodies, the initial and final kinetic energies |
| :--- | :--- |
| are equal and the initial and final relative velocities have equal magnitudes. |,

$$
\vec{r}_{\mathrm{cm}}=\frac{m_{1} \vec{r}_{1}+m_{2} \vec{r}_{2}+\cdots}{m_{1}+m_{2}+\cdots}=\frac{\sum_{i} m_{i} \vec{r}_{i}}{\sum_{i} m_{i}}
$$

## Motion of Center of Mass

The total momentum of a system equals its total mass multiplied by the velocity of the center of mass:

$$
\vec{P}=M \vec{v}_{\mathrm{cm}}=m_{1} \vec{v}_{1}+m_{2} \vec{v}_{2}+\cdots
$$

| $\qquad$The center of mass of a system moves as if all the mass were located at the <br> center of mass: |  |
| :--- | :--- |
| Rocket Motion | If the net external force on a system is zero, the velocity $\vec{v}_{\mathrm{cm}}=M \vec{a}_{\mathrm{cm}}$. <br> mass of the system is constant. |
|  | Any analysis of rocket motion must include the momentum carried away by <br> the spent fuel and the momentum of the rocket itself. |

## Conceptual Questions

## 1 : Jumping off a wall

If you fall off a 2-m-high wall, would you prefer to land on concrete or grass? Why?

## Solution

Your speed at the ground will be the same in both cases. The change in your momentum, or impulse, as you come to rest will also be the same. Given the same impulse, the average force will be less if the time interval is longer. Since you will be in contact with the grass longer as you land, grass is the preferred landing material.

This conceptual question illustrates why cushioning is used to reduce the average force exerted in a collision by increasing the duration of the collision. Consider this principle when you buy your next pair of running shoes or feel your car's padded dashboard.

## 2: Beanbag versus tennis ball

You wish to close your bedroom door from across the room. You can toss either a beanbag or a tennis ball at the door. (Both have the same mass). Which should you choose?

## Solution

Consider how much momentum each object can impart to the door. Both the beanbag and the tennis ball have the same mass, and you give each the same velocity, so their initial momenta are the same. After colliding with the door, the beanbag falls to the floor while the tennis ball bounces back toward you. The beanbag ends with zero momentum, whereas the tennis ball has momentum in the direction opposite that of its original momentum. The change in momentum of the tennis ball is larger than the change in momentum for the beanbag, so you should use the tennis ball to close the door.

Even though both objects have the same initial momentum, we see that the change in momentum determines the best choice.

## 3: Getting off the ice

You're standing on a frictionless ice rink. If you toss your physics book vertically upwards, will you move?

## Solution

You will not move, since the book carries away no component of momentum parallel to the ice rink. You should toss your physics book horizontally to move along the ice.

Is momentum conserved in this case? Initially, there is no momentum. When you toss the book, you and the earth must move in the opposite direction. Since the mass of the earth is extremely large, the velocity imparted to the earth is tiny.

## 4: Car-train collision

If a train engine and a compact car collide, which exhibits the greatest impact force? Which exhibits the greatest change in momentum? Which exhibits the greatest acceleration?

## Solution

Both the train engine and the compact car exhibit the same impact force, due to Newton's third law. Both also exhibit the same change in momentum, since the momentum is conserved. (The friction force with the ground is very small compared with the impact force and may be ignored.) However, the car exhibits the greatest acceleration, since it has the least mass.

The preferred vehicle to be in during a crash is a larger vehicle, because you will experience less acceleration (and less force will be acting on you). Of course, it is far better to avoid the crash altogether.

## 5: Collisions on a pool table

A billiard ball can stop when it collides head-on with another ball of the same mass that is at rest. Can the ball stop if the collision is at a slight angle?

## Solution

If the balls collide head-on, then all of the momentum is in one direction (the direction of the initial velocity) and momentum is conserved if the second ball moves away with the first ball's initial velocity. If the balls collide at an angle, the impact gives the second ball a component of momentum perpendicular to the first ball's initial direction of motion. For the component of momentum perpendicular to the initial velocity to be conserved, the first ball must also have a perpendicular component of momentum and cannot stop.

## 6: Walking in a canoe

You are standing in a canoe. If you walk to the other end of the canoe, what happens to the canoe? You may neglect the resistance of the water.

## Solution

No external forces act on you or the canoe (ignoring the resistance of the water). The center of mass of you and the canoe remains at rest, and its location is constant. When you walk to the other end of the canoe, the canoe must move in the opposite direction to preserve the location of the center of mass.

This conceptual question also illustrates how the frictional force is necessary for walking. As you walk, the frictional force between your shoes and the canoe pushes you in one direction while pushing the canoe in the opposite direction.

## Problems

## 1: Tossing bubble gum

You throw your bubble gum at a stationary puck on a frictionless air-hockey table. The gum sticks to the puck, and both move away with a velocity of $1.2 \mathrm{~m} / \mathrm{s}$. If the puck has a mass of 0.30 kg and the bubble gum has a mass of 0.020 kg , find the initial speed of the bubble gum.

## Solution

IDENTIFY There are no horizontal external forces, so the $x$ component of total momentum is the same before and after the collision. The target variable is the bubble gum's initial speed.

SET UP Figure 8.1 shows the before and after sketches of the situation, including axes. Our system consists of the bubble gum and puck. Before the collision, the gum has a velocity $v_{1 x}$ and the puck is stationary. After the collision, the gum and puck move away at velocity $v_{2 x}$. All of the velocities and momenta have only $x$ components.



Figure 8.1 Problem 1.
EXECUTE We solve the problem by using conservation of momentum. The initial momentum is that of the gum:

$$
p_{1 x}=m_{g} v_{1 x} .
$$

The final momentum is that of the gum and puck together:

$$
p_{2 x}=\left(m_{g}+m_{p}\right) v_{2 x} .
$$

Momentum is conserved, so we set the momenta equal to each other:

$$
m_{g} v_{1 x}=\left(m_{g}+m_{p}\right) v_{2 x} .
$$

Solving for $v_{1 x}$ yields

$$
v_{1 x}=\frac{\left(m_{g}+m_{p}\right) v_{2 x}}{m_{g}}=\frac{((0.020 \mathrm{~kg})+(0.30 \mathrm{~kg}))(1.2 \mathrm{~m} / \mathrm{s})}{(0.020 \mathrm{~kg})}=19 \mathrm{~m} / \mathrm{s}
$$

The initial speed of the bubble gum is $19 \mathrm{~m} / \mathrm{s}$.
EVALUATE This problem illustrates how we can determine a projectile's velocity by examining its collision with a larger object and applying conservation of momentum.

Was the collision elastic? No, the collision was not elastic: The initial kinetic energy was 3.6 J , and the final kinetic energy was 0.23 J . This is an example of a totally inelastic collision, since the masses stuck together after the collision.

## 2: Ball hits the floor

A golf ball of mass 0.045 kg bounces off a tile floor. The velocity of the ball just before it hits the floor is $6.2 \mathrm{~m} / \mathrm{s}$. If the ball is in contact with the floor for 0.012 s and the floor exerts an average force of 40.0 N , find the maximum height of the ball after its impact with the floor.

## Solution

IDENTIFY We can find the maximum height $h$ by using conservation of energy, but we need to know the velocity of the ball just after impact with the floor. We are given the initial velocity and can determine the impulse imparted by the floor. Since the impulse is the change in momentum, we can find the momentum (and velocity) after the impact.

SET UP The motion is purely vertical, so we will use a single vertical axis with the positive direction taken to be upward as shown in Figure 8.2. We will first use the impulse-momentum theorem to find the velocity $v_{2 y}$ just after the ball's impact with the floor. We will then use energy conservation to find our target variable, the height $h$ the ball reaches. The origin is located at the ground, so the golf ball only has kinetic energy immediately after the bounce. At the maximum height, the ball has pure gravitational potential energy.

$$
-y=h
$$



Figure 8.2 Problem 2.
EXECUTE Impulse is change in momentum, which is equal to the average force times the contact interval:

$$
J_{y}=p_{2 y}-p_{\mathrm{ly}}=\left(F_{\mathrm{av}}\right)_{y} \Delta t
$$

We solve this equation for the final momentum:

$$
\begin{aligned}
p_{2 y} & =\left(F_{\mathrm{av}}\right)_{y} \Delta t+p_{1 y} \\
& =\left(F_{\mathrm{av}}\right)_{y} \Delta t+m v_{1 y} \\
& =(40.0 \mathrm{~N})(0.012 \mathrm{~s})+(0.045 \mathrm{~kg})(-6.2 \mathrm{~m} / \mathrm{s}) \\
& =+0.201 \mathrm{~kg} \cdot \mathrm{~m} / \mathrm{s} .
\end{aligned}
$$

Note that the initial velocity is downward, or negative, with our choice of axis. The velocity just after impact is

$$
v_{2 y}=\frac{p_{2 y}}{m}=\frac{(+0.201 \mathrm{~kg} \cdot \mathrm{~m} / \mathrm{s})}{(0.045 \mathrm{~kg})}=4.47 \mathrm{~m} / \mathrm{s}
$$

We can now apply conservation of energy to find the maximum height of the ball. We equate the kinetic energy of the ball just after impact with the gravitational potential energy gained by the ball when it reaches maximum height:

$$
\begin{aligned}
& \frac{1}{2} m v_{2 y}^{2}=m g h, \\
& h=\frac{\frac{1}{2} m v_{2 y}^{2}}{m g}=\frac{v_{2 y}^{2}}{2 g}=\frac{(4.47 \mathrm{~m} / \mathrm{s})^{2}}{2\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}=1.02 \mathrm{~m} .
\end{aligned}
$$

The maximum height reached by the golf ball after bouncing off of the tile floor is 1.02 m .

EVALUATE This problem illustrates the definitions of momentum and impulse in a relatively straightforward way. In addition, it reminds us that we may need to recall problem-solving skills from earlier chapters-conservation of energy in this case.

We also see how we must evaluate signs carefully, since momentum is a vector. The initial momentum is directed downward, which is negative, in our example. Had we omitted the negative sign, we would have calculated a final velocity of $16.9 \mathrm{~m} / \mathrm{s}$. Clearly, this velocity is nonsensical, since it is greater than the velocity just before the bounce!

## 3: Stepping off a sled

A 60.0 kg sled is traveling across an ice rink at $4.5 \mathrm{~m} / \mathrm{s}$. Irene, riding on the sled, jumps off and lands on the ice with a velocity of $1.5 \mathrm{~m} / \mathrm{s}$ in the opposite direction. What is the velocity of the sled after Irene jumps? Irene has a mass of 45.0 kg .

## Solution

IDENTIFY There are no external horizontal forces acting on the sled-plus-Irene system, so momentum is conserved.

SET UP We take the $x$-axis to lie along the direction of motion of the sled, with the positive direction taken to be in the initial direction of the sled as shown in Figure 8.3. We are given the masses of the sled and Irene, the initial velocity of the sled (with Irene on it), and Irene's final velocity. Our target variable is $v_{S 2 x}$, the final velocity of the sled.


Figure 8.3 Problem 3.
EXECUTE The $x$-component of total momentum before Irene jumps off is

$$
\begin{aligned}
p_{1 x} & =m_{S} v_{S I x}+m_{I} v_{\mathrm{IIx}} \\
& =(60.0 \mathrm{~kg})(4.5 \mathrm{~m} / \mathrm{s})+(45.0 \mathrm{~kg})(4.5 \mathrm{~m} / \mathrm{s}) \\
& =472.5 \mathrm{~kg} \cdot \mathrm{~m} / \mathrm{s}
\end{aligned}
$$

(Note again that both the sled and Irene are moving with the same initial velocity.) After Irene jumps off, the $x$ component of total momentum has the same value, so

$$
p_{2 x}=m_{\mathrm{S}} v_{\mathrm{S} 2 x}+m_{\mathrm{I}} v_{\mathrm{I} 2 x} .
$$

Solving for $v_{S 2 r}$, gives

$$
\begin{aligned}
v_{S 2 x} & =\frac{p_{1 x}-m_{\mathrm{I}} v_{\mathrm{I} 2 x}}{m_{S}} \\
& =\frac{(472.5 \mathrm{~kg} \cdot \mathrm{~m} / \mathrm{s})-(45.0 \mathrm{~kg})(-1.5 \mathrm{~m} / \mathrm{s})}{(60.0 \mathrm{~kg})} \\
& =9.0 \mathrm{~m} / \mathrm{s},
\end{aligned}
$$

where Irene's final velocity is negative, since she ends up moving away from the sled. The final velocity of the sled is $9.0 \mathrm{~m} / \mathrm{s}$ in the positive direction.

EVALUATE We check the result by noting that the final velocity of the sled is greater than the initial velocity. For Irene to move away from the sled, she had to give the cart an impulse that resulted in a larger final velocity.

This problem illustrates how we can apply conservation of momentum to problems involving no net external force. Here, there is no explicit collision, but simply someone jumping off a sled.

CAUTION Watch out for signs! In the last two problems, we paid careful attention to the direction, and therefore the sign, of the momentum. Note how an incorrect sign could have led to nonsensical results. Momentum is a vector, and you must always check and recheck the direction to ensure accurate results.

## 4: Colliding pucks in two dimensions

Two pucks collide on a frictionless air-hockey table. Initially, puck $A$ is traveling at $3.50 \mathrm{~m} / \mathrm{s}$ and puck $B$ is at rest. After the collision, puck $A$ moves away at a speed of $2.50 \mathrm{~m} / \mathrm{s}$ and an angle of $30.0^{\circ}$ from the initial direction. Find the final velocity of puck $B$. Puck $A$ has a mass of 3.0 kg and puck $B$ has a mass of 5.0 kg .

## Solution

IDENTIFY There are no horizontal external forces, so both the $x$ component and the $y$ component of the total momentum are conserved in the collision.

SET UP Figure 8.4 shows the before and after sketches of the situation, including axes. Our system consists of the two pucks. Before the collision, puck $A$ moves with velocity $v_{A 1}$ to the right. After the collision, puck $A$ moves away at velocity $v_{A 2} 30.0^{\circ}$ above the $x$-axis and puck $B$ moves away at velocity $v_{B 2}$ at an angle $\theta$ below the $x$-axis (to conserve momentum). The velocities are not along a single axis, so we will have to solve for both the $x$ and $y$ components of momentum. Our target variable is the final velocity $v_{B 2}$ of puck $B$.


Figure 8.4 Problem 4.
EXECUTE Starting with the $x$ components of momentum before and after the collision, we have

$$
\begin{aligned}
& p_{1 x}=m_{A} v_{A 1 x}, \\
& p_{2 x}=m_{A} v_{A 2 x}+m_{B} v_{B 2 x}=m_{A} v_{A 2} \cos 30.0^{\circ}+m_{B} v_{B 2 x} .
\end{aligned}
$$

The $x$ component of momentum is conserved, so we set the expressions for $p_{1 x}$ and $p_{2 x}$ equal to each other:

$$
m_{A} v_{A 1 x}=m_{A} v_{A 2} \cos 30.0^{\circ}+m_{B} v_{B 2 x}
$$

Solving for $v_{B 2 x}$ yields

$$
\begin{aligned}
v_{B 2 x} & =\frac{m_{A} v_{A 1 x}-m_{A} v_{A 2} \cos 30.0^{\circ}}{m_{B}} \\
& =\frac{(3.0 \mathrm{~kg})(3.50 \mathrm{~m} / \mathrm{s})-(3.0 \mathrm{~kg})(2.50 \mathrm{~m} / \mathrm{s}) \cos 30.0^{\circ}}{(5.0 \mathrm{~kg})} \\
& =0.801 \mathrm{~m} / \mathrm{s}
\end{aligned}
$$

We follow the same procedure for the $y$ components of momentum. The initial $y$ component of momentum is zero. The $y$ components of the final momentum must sum to zero; that is,

$$
0=p_{2 y}=m_{A} v_{A 2 y}+m_{B} v_{B 2 y}=m_{A} v_{A 2} \sin 30.0^{\circ}+m_{B} v_{B 2 y}
$$

Solving for $v_{B 2 y}$ gives

$$
\begin{aligned}
v_{B 2 y} & =\frac{-m_{A} v_{A 2} \sin 30.0^{\circ}}{m_{B}} \\
& =\frac{-(2.50 \mathrm{~m} / \mathrm{s})(3.0 \mathrm{~kg}) \sin 30.0^{\circ}}{(5.0 \mathrm{~kg})} \\
& =-0.75 \mathrm{~m} / \mathrm{s} .
\end{aligned}
$$

The $x$ and $y$ components of puck $B$ 's velocity are $0.80 \mathrm{~m} / \mathrm{s}$ and $-0.75 \mathrm{~m} / \mathrm{s}$, respectively. Thus, puck $B$ must travel into the fourth quadrant, as expected. We find the magnitude and direction of puck $B$ 's velocity from

$$
\begin{aligned}
& v_{B 2}=\sqrt{v_{B 2 x}^{2}+v_{B 2 y}^{2}}=\sqrt{(0.80 \mathrm{~m} / \mathrm{s})^{2}+(-0.75 \mathrm{~m} / \mathrm{s})^{2}}=1.10 \mathrm{~m} / \mathrm{s}, \\
& \phi=\tan ^{-1} \frac{v_{B 2 y}}{v_{B 2 x}}=\tan ^{-1} \frac{(-0.75 \mathrm{~m} / \mathrm{s})^{2}}{(0.80 \mathrm{~m} / \mathrm{s})}=-43.2^{\circ},
\end{aligned}
$$

Puck $B$ 's final velocity is $1.10 \mathrm{~m} / \mathrm{s}$, directed at an angle of $43.2^{\circ}$ below the positive $x$-axis.
EVALUATE We can check our answer by examining the momentum components before and after the collision. Initially, puck $A$ has an $x$ component of momentum equal to $10.5 \mathrm{~kg} \mathrm{~m} / \mathrm{s}$. After the collision, the $x$ component of momentum of puck $A$ is $6.5 \mathrm{~kg} \mathrm{~m} / \mathrm{s}$ and that of puck $B$ is $4.0 \mathrm{~kg} \mathrm{~m} / \mathrm{s}$, or a total of $10.5 \mathrm{~kg} \mathrm{~m} / \mathrm{s}$, as expected. There is no initial momentum along the $y$-axis. After the collision, the $y$ component of momentum of puck $A$ is $3.75 \mathrm{~kg} \mathrm{~m} / \mathrm{s}$ and that of puck $B$ is $-3.75 \mathrm{~kg} \mathrm{~m} / \mathrm{s}$, also as expected.

Solving momentum problems in two dimensions follows from the one-dimensional cases. You just need to remember that momentum is a vector that can have multiple components, each of which can be conserved individually.

## 5: Elastic collision in one dimension

Two gliders collide elastically on a frictionless, linear air track. Glider $A$ has a mass of 0.60 kg and initially moves to the right at $3.0 \mathrm{~m} / \mathrm{s}$. Glider $B$ has mass of 0.40 kg and initially moves to the left at $4.0 \mathrm{~m} / \mathrm{s}$. What are the final velocities of the two gliders after the collision?

## Solution

IDENTIFY There are no net external forces acting on the system, so the momentum of the system is conserved. The collision is elastic; therefore, energy is conserved as well.

SET UP Figure 8.5 shows the before and after sketches of the situation including axes. Our system consists of the two gliders. Before the collision, glider $A$ moves with velocity $v_{A 1 x}=3.0 \mathrm{~m} / \mathrm{s}$ to the right and glider $B$ moves with velocity $v_{B 1 x}=-4.0 \mathrm{~m} / \mathrm{s}$ to the left. After the collision, glider $A$ moves away with velocity $v_{A 2 x}$ and glider $B$ moves away with velocity $v_{B 2 x}$. Our target variables are the final velocities of the two gliders. We'll need to use the relative velocity relation for elastic collisions in our solution.

## Before



Af ter


Figure 8.5 Problem 5.

EXECUTE From conservation of momentum,

$$
\begin{aligned}
& m_{A} v_{A 1 x}+m_{B} v_{B 1 x}=m_{A} v_{A 2 x}+m_{B} v_{B 2 x}, \\
& (0.60 \mathrm{~kg})(3.0 \mathrm{~m} / \mathrm{s})+(0.40 \mathrm{~kg})(-4.0 \mathrm{~m} / \mathrm{s})=(0.60 \mathrm{~kg}) v_{A 2 x}+(0.40 \mathrm{~kg}) v_{B 2 x}, \\
& 0.20 \mathrm{~m} / \mathrm{s}=0.60 v_{A 2 x}+0.40 v_{B 2 x} .
\end{aligned}
$$

The last equation has two unknowns, and we need more information to solve for the velocities. Since this is an elastic collision, we can apply the relative velocity relation to solve for the velocities:

$$
\begin{aligned}
& v_{B 2 x}-v_{A 2 x}=-\left(v_{B 1 x}-v_{A 1 x}\right) \\
& v_{B 2 x}-v_{A 2 x}=-((-4.0 \mathrm{~m} / \mathrm{s})-(3.0 \mathrm{~m} / \mathrm{s}))=7.0 \mathrm{~m} / \mathrm{s} \\
& v_{B 2 x}=v_{A 2 x}+7.0 \mathrm{~m} / \mathrm{s} .
\end{aligned}
$$

Substituting the last expression into the earlier momentum conservation relation gives

$$
\begin{aligned}
& 0.60 v_{A 2 x}+0.40\left(v_{A 2 x}+7.0 \mathrm{~m} / \mathrm{s}\right)=0.20 \mathrm{~m} / \mathrm{s} \\
& v_{A 2 x}=(0.20 \mathrm{~m} / \mathrm{s})-(0.40(7.0 \mathrm{~m} / \mathrm{s}))=-2.6 \mathrm{~m} / \mathrm{s} \\
& v_{B 2 x}=v_{A 2 x}+7.0 \mathrm{~m} / \mathrm{s}=(-2.6 \mathrm{~m} / \mathrm{s})+7.0 \mathrm{~m} / \mathrm{s}=4.4 \mathrm{~m} / \mathrm{s}
\end{aligned}
$$

After the collision, puck $A$ moves to the left at $2.6 \mathrm{~m} / \mathrm{s}$ ( $v_{A 2 x}$ is negative) and puck $B$ moves to the right at $4.4 \mathrm{~m} / \mathrm{s}$.

EVALUATE Both gliders reversed their directions in this elastic collision. Are the kinetic energies equivalent before and after the collision? The initial kinetic energy is

$$
K_{\mathrm{i}}=\frac{1}{2} m_{A} v_{A \mid x}^{2}+\frac{1}{2} m_{B} v_{B 1 x}^{2}=\frac{1}{2}(0.60 \mathrm{~kg})(3.0 \mathrm{~m} / \mathrm{s})^{2}+\frac{1}{2}(0.40 \mathrm{~kg})(-4.0 \mathrm{~m} / \mathrm{s})^{2}=5.9 \mathrm{~J} .
$$

The final kinetic energy is

$$
K_{\mathrm{f}}=\frac{1}{2} m_{A} v_{A 2 x}^{2}+\frac{1}{2} m_{B} v_{B 2 x}^{2}=\frac{1}{2}(0.60 \mathrm{~kg})(-2.6 \mathrm{~m} / \mathrm{s})^{2}+\frac{1}{2}(0.40 \mathrm{~kg})(4.4 \mathrm{~m} / \mathrm{s})^{2}=5.9 \mathrm{~J} .
$$

The initial and final kinetic energies are equivalent, as expected.

## 6: Ballistic pendulum

A 6.00 g bullet is shot through a 2.00 kg block suspended on a $1.00-\mathrm{m}$-long string. If the initial speed of the bullet is $650 \mathrm{~m} / \mathrm{s}$ and it emerges with a velocity of $175 \mathrm{~m} / \mathrm{s}$, find the maximum angle through which the block swings after it is hit.

## Solution

IDENTIFY We will analyze this problem in two stages. First, we'll look at the interaction of the bullet with the block. Second, we'll examine the swinging of the block on the string after the bullet passes through the block.

In the first stage, the bullet passes through the block, giving the block an impulse. The bullet passes through the block very quickly, so the block has no time to swing any significant distance from its initial position. The only force acting on the bullet-and-block system is between the bullet and block; there are no appreciable external forces acting on the system. We conclude that, during the first stage, the horizontal component of momentum is conserved.

In the second stage, the block moves away from its initial position. The only forces acting on the block are gravity (a conserved force) and tensions due to the strings (which do no work as the block swings). Mechanical energy is conserved as the block swings.

SET UP Figure 8.6 shows sketches of the two stages of the problem. We take the positive $x$-axis to be to the right and the positive $y$-axis upward. Our target variable for the first stage is the velocity $V_{2 x}$ of the block after the bullet emerges from the block. We will use momentum conservation to find $V_{2 x}$. Our target variable for the second stage is $\theta$, the angle the string makes with the vertical when the block stops momentarily after swinging to the right. We will use energy conservation to find the height $h$ the block rises to after the collision and relate $h$ to $\theta$.


Figure 8.6 Problem 6.
EXECUTE In the first stage, conservation of momentum gives

$$
m v_{1 x}=m v_{2 x}+M V_{2 x} .
$$

The velocity of the block is

$$
\begin{aligned}
V_{2 \mathrm{x}} & =\frac{m\left(v_{1 x}-v_{2 x}\right)}{M} \\
& =\frac{(0.006 \mathrm{~kg})(650 \mathrm{~m} / \mathrm{s}-175 \mathrm{~m} / \mathrm{s})}{(2.0 \mathrm{~kg})} \\
& =1.43 \mathrm{~m} / \mathrm{s} .
\end{aligned}
$$

At the beginning of the second stage, the block has kinetic energy. Afterwards, the block swings up and comes to rest momentarily. At this point, the block's kinetic energy is zero and its gravitational potential energy has increased by $m g h$. Energy conservation gives

$$
\frac{1}{2} M V_{2 x}^{2}=M g h .
$$

The block reaches a height

$$
h=\frac{V_{2 x}^{2}}{2 g}=0.104 \mathrm{~m}
$$

The angle between the vertical and the string when the block stops momentarily is

$$
\theta=\cos ^{-1} \frac{l-h}{l}=26.4^{\circ}
$$

The maximum angle the block swings to after the collision is $26.4^{\circ}$.
EVALUATE Where did most of the energy go? A quick check of the kinetic energies gives an initial energy of the bullet (before striking the block) of 1270 J , a final energy of the bullet of 92 J , and a final energy of the block of 2.0 J . Most of the initial energy was dissipated in the deformation and heating of the bullet and block.

## 7: Boat on a lake

Luka is standing in a boat on a calm lake. His position is 5.00 m from the end of the pier. As he walks towards the pier, the boat moves away from it. When Luka stops walking, he finds that the boat has moved 2.00 m away from the pier. If you ignore the boat's resistance to motion in the water, how far from the end of the pier is Luka when he stops walking? The mass of the boat is 80.0 kg , and Luka's mass is 60.0 kg .

## Solution

IDENTIFY Since we can ignore friction between the boat and the water, the net external force on the boat-and-Luka system is zero. Momentum is therefore conserved. There is no initial motion, so the total momentum is zero and the velocity of the center of mass of the system is zero. The center of mass will remain at rest as Luka walks in the boat.

SET UP We take the origin to be the end of the pier, since the measurements are given with respect to that point. Figure 8.7 shows a sketch of the situation. Our target variable is the final position of Luka, $x_{L 2}$.


Figure 8.7 Problem 7.
EXECUTE The initial coordinate of the center of mass of the boat-and-Luka system is given by the formula

$$
x_{\mathrm{cm}}=\frac{x_{L 1} m_{L}+x_{B 1} m_{B}}{m_{L}+m_{B}}
$$

The final coordinate of the center of mass of the system is given by the equation

$$
x_{\mathrm{cm}}=\frac{x_{L 2} m_{L}+x_{B 2} m_{B}}{m_{L}+m_{B}} .
$$

Since the center of mass doesn't move, we set the two right-hand expressions equal to each other:

$$
\frac{x_{L 1} m_{L}+x_{B 1} m_{B}}{m_{L}+m_{B}}=\frac{x_{L 2} m_{L}+x_{B 2} m_{B}}{m_{L}+m_{B}} .
$$

The boat moves 2.0 m away from the end of the pier, so $x_{B 2}=x_{B 1}+2.0 \mathrm{~m}$. Substituting and solving gives

$$
\begin{aligned}
& x_{L 1} m_{L}+x_{B 1} m_{B}=x_{L 2} m_{L}+\left(x_{B 1}+2.0 \mathrm{~m}\right) m_{B} \\
& (5.0 \mathrm{~m})(60.0 \mathrm{~kg})+x_{B 1}(80.0 \mathrm{~kg})=x_{L 2}(60.0 \mathrm{~kg})+\left(x_{B 1}+2.0 \mathrm{~m}\right)(80.0 \mathrm{~kg}), \\
& x_{L 2}=2.3 \mathrm{~m} .
\end{aligned}
$$

EVALUATE Since the boat and Luka's mass are similar, Luka moves a distance similar to the distance the boat moves. Luka's mass is less, so he moves a little farther than the boat moves.

## 8: Falling sand

Sand is dropped from a height of 1.0 m onto a kitchen scale at the uniform rate of $100.0 \mathrm{~g} / \mathrm{s}$. Find the force on the scale and its reading if the scale is calibrated in kg .

## Solution

IDENTIFY We need to find the force on the scale due to the falling sand. We can find the change in momentum of the falling sand and relate it to the force through impulse. The target variable is the force on the scale.

SET UP We will first find the velocity of the sand as it hits the scale. We will then use that velocity to find the change in momentum as the sand strikes the scale. We will incorporate the change in momentum into the impulse formula and solve for the force on the scale. We will complete the problem by converting the force into a calibrated scale reading. We will take our vertical axis to be positive downward.

EXECUTE As the sand falls 1.0 m , it acquires a velocity

$$
v_{1 y}=\sqrt{2 g h}
$$

according to energy conservation. The sand stops after striking the scale. The change in momentum as the sand strikes the scale is

$$
\Delta p=p_{2 y}-p_{1 y}=-\Delta m v=-\Delta m \sqrt{2 g h}
$$

The change in momentum is negative (directed upward) in our coordinate system. The force on the sand due to the scale is directed upward to stop the sand. The impulse imparted to the sand by the scale is

$$
J=-F \Delta t=\Delta p=-\Delta m \sqrt{2 g h}
$$

Rearranging terms and dividing both sides by $\Delta t$ gives

$$
\begin{aligned}
F & =\frac{\Delta m}{\Delta t} \sqrt{2 g h} \\
& =(0.100 \mathrm{~kg} / \mathrm{s}) \sqrt{2\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(1.0 \mathrm{~m})} \\
& =0.44 \mathrm{~N} .
\end{aligned}
$$

Here, the rate of change of mass per unit time is the given rate at which the sand falls. The scale is calibrated in terms of mass, so we divide the force by $g$ to get the reading on the scale:

$$
m=\frac{F}{g}=\frac{0.44 \mathrm{~N}}{9.8 \mathrm{~m} / \mathrm{s}^{2}}=0.045 \mathrm{~kg}
$$

The scale reads 0.045 kg as the sand falls.
EVALUATE This problem illustrates how we can use force, momentum, and impulse together to solve problems.

## Try It Yourself!

## 1: Rocket motion

A two-stage rocket traveling at $350 \mathrm{~m} / \mathrm{s}$ through space separates, with one stage having twice the mass of the other. If the final velocity of the larger stage is $120 \mathrm{~m} / \mathrm{s}$ in a direction opposite its initial direction, find the final velocity of the smaller stage.

## Solution Checkpoints

IDENTIFY AND SET UP Confirm that there are no net external forces acting on the rocket. Apply conservation of momentum along one axis. The target variable is the velocity of the smaller rocket stage.

EXECUTE Momentum conservation gives

$$
m_{A} v_{1 \mathrm{x}}+m_{B} v_{1 \mathrm{x}}=m_{A} v_{A 2 x}+m_{B} v_{B 2 x} .
$$

Solving for $v_{A 2 x}$ yields

$$
v_{A 2 x}=\frac{\left(m_{A}+m_{B}\right) v_{1 \mathrm{x}}-m_{B} v_{B 2 x}}{m_{A}}=1300 \mathrm{~m} / \mathrm{s}
$$

EVALUATE Do you expect the smaller stage to have a final velocity that is larger or smaller than the other velocities in the problem? Do you expect energy to be conserved in this case?

## 2: Car collision

A 2000 kg car moving at $30 \mathrm{~km} / \mathrm{hr}$ collides with a stopped 1000 kg car, and the two lock bumpers. What is their common velocity after the collision? What fraction of the initial energy is dissipated in the collision?

## Solution Checkpoints

IDENTIFY AND SET UP Confirm that there are no net external forces acting on the cars just before and just after the collision. Apply conservation of momentum along one axis. The target variables are the velocity of the combined cars and the fraction of energy lost in the collision.

EXECUTE Momentum conservation gives

$$
m_{A} v_{A 1 x}=\left(m_{A}+m_{B}\right) v_{2 x}
$$

Solving for the final velocity gives $20 \mathrm{~km} / \mathrm{hr}$. The initial and final energies are

$$
\begin{aligned}
& E_{1}=\frac{1}{2} m_{A} v_{A l x}^{2} \\
& E_{2}=\frac{1}{2}\left(m_{A}+m_{B}\right) v_{2 x}^{2} .
\end{aligned}
$$

Dividing the two demonstrates that two-thirds of the initial energy remains kinetic after the collision.
EVALUATE We will contrast this problem with the next one, in which the collision is elastic.

## 3: Car collision revisited

A 2000 kg car moving at $30 \mathrm{~km} / \mathrm{hr}$ collides with a stopped 1000 kg car, and the two have spring bumpers so that the collision is perfectly elastic. What is the velocity of each car after the collision?

## Solution Checkpoints

IDENTIFY AND SET UP Confirm that there are no net external forces acting on the cars just before and just after the collision. Apply conservation of momentum along one axis. Apply conservation of energy, since the collision is elastic. The target variables are the velocities of the two cars.

EXECUTE Momentum conservation gives

$$
m_{A} v_{A 1 x}=m_{A} v_{A 2 x}+m_{B} v_{B 2 x} .
$$

Energy conservation gives

$$
\frac{1}{2} m_{A} v_{A 1 x}^{2}=\frac{1}{2} m_{A} v_{A 2 x}^{2}+\frac{1}{2} m_{B} v_{B 2 x}^{2} .
$$

These two equations can be solved to find $v_{A 2 x}=10 \mathrm{~km} / \mathrm{hr}$ and $v_{2 B x}=40 \mathrm{~km} / \mathrm{hr}$. You may wish to consult the text for a derivation of velocities in elastic collisions.

EVALUATE In this elastic collision, we see that the lighter car ends up with a velocity greater than the heavier car after the collision. You can also check that the initial and final momenta and energies are equal.

## 4: Putty hitting mass on spring

A 0.75 kg mass of putty is hurled with a velocity of $2.0 \mathrm{~m} / \mathrm{s}$ against a 0.50 kg mass attached to a spring, as shown in Figure 8.8. The mass attached to the spring slides across the frictionless horizontal surface, depressing the spring a maximum distance of 12.0 cm . Find the spring constant if the putty sticks to the mass on the spring.


Figure 8.8 Try It Yourself! 4.

## Solution Checkpoints

IDENTIFY AND SET UP Is the collision elastic or inelastic if the putty sticks to the mass? You can apply conservation of momentum to the initial collision (before the spring begins to compress). Then apply conservation of energy to the interval after the spring begins to compress.

EXECUTE Momentum conservation gives

$$
m_{A} v_{A \mid x}=\left(m_{A}+m_{B}\right) v_{2 x} .
$$

The velocity of the putty plus mass is $1.2 \mathrm{~m} / \mathrm{s}$ just after the collision. Energy conservation applied to the interval after the collision gives

$$
\frac{1}{2}\left(m_{A}+m_{B}\right) v_{2 x}^{2}=\frac{1}{2} k x^{2} .
$$

Substituting yields a spring constant of $125 \mathrm{~N} / \mathrm{m}$.
EVALUATE Energy and momentum conservation do not apply to the whole process; rather, the process must be broken into intervals during which energy and momentum can be applied. Can you repeat the problem, replacing the putty for a rubber ball that bounces elastically off the mass attached to the spring?

## 5: Collision on a football field

A 135 kg football player traveling at $5.0 \mathrm{~m} / \mathrm{s}$ collides with an 85 kg player at rest. The two slide 2.0 m on wet grass. Their collision is completely inelastic. What is the coefficient of friction between the grass and the players? How much energy is dissipated in the collision?

## Solution Checkpoints

IDENTIFY AND SET UP Confirm that there are minimal external forces acting on the players just before and just after the collision. Apply conservation of momentum to the two players as they collide. Then use the work-energy theorem to find the friction force. The energy dissipated is found by comparing the energy just before and just after the collision.

EXECUTE Momentum conservation applied to the collision gives

$$
m_{A} v_{A \mid x}=\left(m_{A}+m_{B}\right) v_{2 x}
$$

The work-energy theorem applied to the two players as they slide on the grass after they collide results in

$$
-\mu\left(m_{A}+m_{B}\right) g x=0-\frac{1}{2}\left(m_{A}+m_{B}\right) v_{2 x}^{2} .
$$

Solving these equations, we obtain $\mu=0.24$. The change in energy during the collision is

$$
\Delta E=\frac{1}{2}\left(m_{A}+m_{B}\right) v_{2 x}^{2}-\frac{1}{2} m_{A} v_{A \mid x}^{2}
$$

or 650 J dissipated during the collision.
EVALUATE In this collision, we had to split the process up into two intervals, to which we separately applied conservation of momentum and the work-energy theorem. We also found that almost one-third of the energy is dissipated in the totally inelastic collision of the players.

## Problem Summary

This chapter has augmented our knowledge of energy, forces, and kinematics with our newly formed knowledge of momentum analysis. Momentum analysis shares many of the problem-solving principles we have encountered. In the problems presented, we

- Identified the general procedure to find the solution.
- Sketched the situation, including before and after views.
- Identified the momenta, energies, and forces in the system.
- Applied conservation of momentum to the system.
- Applied energy principles, including conservation of energy (when possible).
- Drew free-body diagrams of the objects when appropriate.
- Applied Newton's laws when appropriate.
- Solved for the target variable(s) algebraically from the equations we derived.
- Reflected on the results, checking for inconsistencies.
- 


## 9 <br> Rotation of Rigid Bodies

## Summary

In this chapter, we will investigate the rotational motion of rigid bodies-objects that don't change size or shape as they move. We'll describe first the kinematics of rotation for a rigid body and then its rotational kinetic energy. We'll see how these quantities are analogous to linear kinematics and translational kinetic energy. Finally, we'll learn about the moment of inertia, how to calculate it, and how to use it to measure rotational inertia. We'll use our new knowledge of rotational motion in the next chapter as we learn how to cause such motion.

## Objectives

After studying this chapter, you will understand

- How to use radians in angular measurements.
- The definition and application of angular displacement, velocity, and acceleration.
- How to solve problems involving constant angular acceleration.
- How to define moment of inertia and apply it to systems of varying shapes.
- How to calculate the moment of inertia.
- How to solve conservation-of-energy problems that include rotational kinetic energy.
- How to draw analogies between translational and rotational motion and energy.


## Concepts and Equations

| Term | Description |
| :---: | :---: |
| Rigid Body | A rigid body is an object that maintains an unchanging size and shape. We neglect squeezing, stretching, and twisting in our analysis of a rigid body. |
| Radian | Angular displacements are usually measured in radians. A displacement $\theta$ measured in radians is the ratio of the arc length $s$ to the radius $r$ : $\theta=\frac{s}{r}$ <br> There are $2 \pi$ radians in one revolution $\left(360^{\circ}\right)$. |
| Angular Velocity | The instantaneous angular velocity about the $z$-axis is the rate of change of angular displacement with respect to time: $\omega_{z}=\lim _{\Delta t \rightarrow 0} \frac{\Delta \theta}{\Delta t}=\frac{d \theta}{d t} .$ <br> The term angular velocity refers to the instantaneous angular velocity. All pieces of a rigid object have the same angular velocity at any given instant. The direction of the angular velocity is given by the right-hand rule. |
| Angular Acceleration | The instantaneous angular acceleration about the $z$-axis is the rate of change of angular velocity with respect to time: $\alpha_{z}=\lim _{\Delta t \rightarrow 0} \frac{\Delta \omega_{z}}{\Delta t}=\frac{d \omega_{z}}{d t}=\frac{d^{2} \theta}{d t^{2}}$ <br> The term angular acceleration refers to the instantaneous angular acceleration. All pieces of a rigid object have the same angular acceleration at any given instant. The direction of the angular acceleration is the same as that of the angular speed when the object is speeding up and opposite that of the angular speed when the object is slowing down. |
| Rotation with Constant Angular Acceleration | When an object moves with constant angular acceleration, the angular displacement, velocity, acceleration, and time are related by the formulas $\begin{aligned} & \theta=\theta_{0}+\omega_{0 z} t+\frac{1}{2} \alpha_{z} t^{2}, \\ & \omega_{z}=\omega_{0 z}+\alpha_{z} t \\ & \omega_{z}^{2}=\omega_{0 z}^{2}+2 \alpha_{z}\left(\theta-\theta_{0}\right), \\ & \theta-\theta_{0}=\frac{1}{2}\left(\omega_{z}+\omega_{0 z}\right) t, \end{aligned}$ |

where $\theta_{0}$ and $\omega_{0 x}$ are the initial values of the angular position and velocity, respectively.

## Connecting Linear

 and Angular QuantitiesThe tangential speed $v$ of a particle rotating in a rigid body at a distance $r$ from the axis of rotation is

$$
v=r \omega .
$$

The particle's acceleration $\vec{a}$ has a tangential component

$$
a_{\tan }=\frac{d v}{d t}=\frac{r d \omega}{d t}=r \alpha
$$

and a radial component

$$
a_{\mathrm{rad}}=\frac{v^{2}}{r}=\omega^{2} r .
$$

| Moment of Inertia | The moment of inertia, $I$, of a body is a measure of its rotational inertia and depends on how the body's mass is distributed relative to the axis of rotation. The moment of inertia is given by $\begin{aligned} I & =m_{1} r_{1}^{2}+m_{2} r_{2}^{2}+m_{3} r_{3}^{2}+\cdots \\ & =\sum_{i} m_{i} r_{i}^{2} \end{aligned}$ <br> For arbitrarily distributed masses, the moment of inertia is given by $I=\int r^{2} d m$ <br> Moments of inertia for common shapes are given in Table 9.2 in the textbook The parallel-axis theorem can be used to find the moment of inertia, $I_{P}$, about another parallel axis; that is, $I_{P}=I_{\mathrm{cm}}+M d^{2}$ <br> where $I_{\mathrm{cm}}$ is the moment of inertia about the center of mass, $M$ is the mass of the body, and $d$ is the distance between the two axes. |
| :---: | :---: |
| Energy of Rotating Body | The kinetic energy of a rigid body rotating about a fixed axis is $K=\frac{1}{2} I \omega^{2} .$ <br> This quantity is the sum of the kinetic energies of all of the particles that make up the rigid body. |

## Conceptual Questions

## 1: Rolling versus sliding down a hill

A ball travels down a hill. Will the ball reach the bottom of the hill faster if it rolls or if it slides without friction down the hill?

## Solution

IDENTIFY, SET UP, AND EXECUTE No energy is lost as the ball travels down the hill, so we can use energy conservation to answer this question. At the top of the hill, the ball has gravitational potential energy. As the ball descends, the gravitational potential energy is transformed into kinetic energy. When the ball rolls down the hill, the kinetic energy is shared between translational kinetic energy and rotational kinetic energy. When the ball slides down the hill without friction, the gravitational potential energy transforms into translational kinetic energy alone. There is no rotational kinetic energy in this case. Therefore, more energy is transformed into translational kinetic energy if the ball slides without friction than if it rolls down the hill. Since it acquires more translational kinetic energy, its velocity is higher and it reaches the bottom faster when it slides without friction.

EVALUATE This question shows how we must include both translational and rotational kinetic energy in our energy analyses. We'll practice using rotational kinetic energy in the problem section.

## 2: Comparing moments of inertia

A light rod of length $L$ has two lead weights of mass $M$ attached at both ends of the rod. How does the system's moment of inertia compare when the rod is spun about an axis at its center as opposed to when it is spun around a point one-quarter along the length of the rod?

IDENTIFY, SET UP, AND EXECUTE The moment of inertia of an object is

$$
I=m_{1} r_{1}^{2}+m_{2} r_{2}^{2}+\cdots
$$

The moment of inertia when the axis is at the center of the rod is then

$$
I_{\text {center axis }}=M\left(\frac{L}{2}\right)^{2}+M\left(\frac{L}{2}\right)^{2}=M\left(\frac{L^{2}}{2}\right)=\frac{1}{2} M L^{2},
$$

since each mass is positioned half of the length of the rod away from the axis. When the axis is onequarter along the length of the rod, the moment of inertia is

$$
I_{1 / 4 \text { along rod }}=M\left(\frac{L}{4}\right)^{2}+M\left(\frac{3 L}{4}\right)^{2}=M\left(\frac{L^{2}}{16}\right)+M\left(\frac{9 L^{2}}{16}\right)=\frac{10}{16} M L^{2},
$$

since one mass is $\frac{1}{4} L$ from the axis and the other is $\frac{3}{4} L$ from the axis. The moment of inertia when the rod is spun one-quarter along its length is $25 \%$ larger than the moment of inertia when the rod is spun at the center.

EVALUATE The moment of inertia depends on both mass and the location of the mass. Since the location from the axis enters as the square of the distance, moving the axis changes the moment of inertia. Do you get the same result if you apply the parallel-axis theorem?

Practice Problem: What axis of rotation provides the largest moment of inertia? Answer: The axis located at one end of the rod gives the largest moment of inertia, $I=M L^{2}$.

## 3: Rolling up a ramp

A solid sphere and a thin-walled sphere roll without slipping along a horizontal surface. The two spheres roll with the same translational speed. The surface leads to a ramp. Which sphere rises to the greatest height on the ramp before stopping momentarily?

IDENTIFY AND SET UP We'll use conservation of energy and ignore air drag. Both spheres have initial translational and rotational kinetic energies that are transformed completely into gravitational potential energy when they stop momentarily on the ramp. The sphere with the greatest initial total kinetic energy will rise to the greatest height.

EXECUTE The initial kinetic energy of each sphere is

$$
K_{i}=\frac{1}{2} m_{\text {sphere }} v_{\mathrm{cm}}^{2}+\frac{1}{2} I_{\text {sphere }} \omega^{2} .
$$

The angular velocity is related to the velocity of the center of mass, since both spheres roll without slipping. Thus,

$$
\omega=\frac{v_{\mathrm{cm}}}{\mathrm{R}}
$$

where $R$ is the radius of the sphere. The moment of inertia of the solid sphere is $I_{\text {solid }}=2 / 5 m_{\text {solid }} R_{\text {solid }}^{2}$. The kinetic energy of the solid sphere is then

$$
K_{\text {solid }}=\frac{1}{2} m_{\text {solid }} v_{\mathrm{cm}}^{2}+\frac{1}{2} I_{\text {solid }} \omega^{2}=\frac{1}{2} m_{\text {solid }} v_{\mathrm{cm}}^{2}+\frac{1}{2} \frac{2}{5} m_{\text {solid }} R_{\text {solid }}^{2}\left(\frac{v_{\mathrm{cm}}}{R_{\text {solid }}}\right)^{2}=\frac{7}{10} m_{\text {solid }} v_{\mathrm{cm}}^{2} .
$$

The moment of inertia of the thin-walled sphere (shell) is $I_{\text {shell }}=2 / 3 m_{\text {shell }} R_{\text {shell }}^{2}$. The kinetic energy of the shell is then

$$
K_{\text {shell }}=\frac{1}{2} m_{\text {shell }} v_{\mathrm{cm}}^{2}+\frac{1}{2} I_{\text {shell }} \omega^{2}=\frac{1}{2} m_{\text {shell }} v_{\mathrm{cm}}^{2}+\frac{1}{2} \frac{2}{3} m_{\text {shell }} R_{\text {shell }}^{2}\left(\frac{v_{\mathrm{cm}}}{R_{\text {shell }}}\right)^{2}=\frac{5}{6} m_{\text {shell }} v_{\mathrm{cm}}^{2} .
$$

Since the final gravitational potential energy depends on mass ( $U=m g h$ ), the masses will cancel and we compare the leading fractions in the kinetic-energy terms to determine which sphere has the greatest initial kinetic energy. We see that the shell has more initial kinetic energy; therefore, the thin-walled sphere rises to the greatest height.

EVALUATE It is interesting to find that the results don't depend on either the mass or the radius of the two spheres. The results depend only on how the mass is distributed in the object.

## 4: Racing down a ramp

A thin-walled hollow cylinder, a solid cylinder, a solid sphere, and a thin-walled sphere start from rest at the same height and roll without slipping down a wide ramp. Rank the velocities of the four objects, from first to last.

IDENTIFY AND SET UP We'll use conservation of energy. All of the objects have gravitational potential energy that is transformed into translational and rotational kinetic energies. We will find the velocity at the bottom of the ramp in terms of the height and other factors.

EXECUTE For any of the four objects, energy conservation applied to the starting point and the bottom of the ramp gives

$$
m g h=\frac{1}{2} m v^{2}+\frac{1}{2} I \omega^{2}
$$

The angular velocity is related to the velocity of the center of mass, since all of the objects roll without slipping; thus,

$$
\omega=\frac{v}{\mathrm{R}},
$$

where $R$ is the radius of the object. Replacing the angular velocity yields

$$
m g h=\frac{1}{2} m v^{2}+\frac{1}{2} I\left(\frac{v}{R}\right)^{2}
$$

We now solve for the velocity of each of the four objects. The thin-walled hollow cylinder (TWHC) gives

$$
\begin{aligned}
& m_{\mathrm{TWHC}} g h=\frac{1}{2} m_{\mathrm{TWHC}} v_{\mathrm{TWHC}}^{2}+\frac{1}{2}\left(m_{\mathrm{TWHC}} R^{2}\right)\left(\frac{v_{\mathrm{TWHC}}}{R}\right)^{2}=m_{\mathrm{TWHC}} v_{\mathrm{TWHC}}^{2} \\
& v_{\mathrm{TWHC}}=\sqrt{g h} .
\end{aligned}
$$

The solid cylinder (SC) results in

$$
\begin{aligned}
& m_{\mathrm{SC}} g h=\frac{1}{2} m_{\mathrm{SC}} v_{\mathrm{SC}}^{2}+\frac{1}{2}\left(\frac{1}{2} m_{\mathrm{SC}} R^{2}\right)\left(\frac{v_{\mathrm{SC}}}{R}\right)^{2}=\frac{3}{4} m_{\mathrm{SC}} v_{\mathrm{SC}}^{2}, \\
& v_{\mathrm{SC}}=\sqrt{\frac{4}{3}} \sqrt{g h} .
\end{aligned}
$$

The solid sphere (SS) produces

$$
\begin{aligned}
& m_{\mathrm{SS}} g h=\frac{1}{2} m_{\mathrm{SS}} v_{\mathrm{SS}}^{2}+\frac{1}{2}\left(\frac{2}{5} m_{\mathrm{SS}} R^{2}\right)\left(\frac{v_{\mathrm{SS}}}{R}\right)^{2}=\frac{7}{10} m_{\mathrm{SS}} v_{\mathrm{SS}}^{2} \\
& v_{\mathrm{SS}}=\sqrt{\frac{10}{7}} \sqrt{g h}
\end{aligned}
$$

The thin-walled hollow sphere (TWHS) gives

$$
\begin{aligned}
& m_{\mathrm{TWHS}} g h=\frac{1}{2} m_{\mathrm{TWHS}} v_{\mathrm{TWHS}}^{2}+\frac{1}{2}\left(\frac{2}{3} m_{\mathrm{TWHS}} R^{2}\right)\left(\frac{v_{\mathrm{TWHS}}}{R}\right)^{2}=\frac{5}{6} m_{\mathrm{TWHS}} v_{\mathrm{TWHS}}^{2}, \\
& v_{\mathrm{TWHS}}=\sqrt{\frac{6}{5}} \sqrt{g h} .
\end{aligned}
$$

Comparing the leading factors, we see that the solid sphere has the largest velocity, followed by the solid cylinder, the thin-walled hollow sphere, and the thin-walled hollow cylinder.

EVALUATE The largest moment of inertias resulted in the smallest velocities at the bottom of the ramp, since energy went into rotational kinetic energy. We see that the radii and masses cancel in this problem: The velocity at the bottom is dependent only upon the distribution of mass in the object.

In what order do the four objects reach the bottom? Since all the velocities depend on the square root of the height, those with the fastest velocity reach the bottom first.

## Problems

## 1: Constant angular acceleration in a pottery wheel

A pottery wheel is rotating with an initial angular velocity $\omega_{0}$ when the wheel's drive motor is turned on. The wheel increases to a final angular velocity of 125 rpm while making 30.0 revolutions in 25.0 seconds. Find the initial angular velocity and angular acceleration, assuming that the latter is constant. A pottery wheel is essentially a cylinder rotating about a vertical axis driven by a motor.


Figure 9.1 Problem 1 sketch.

## Solution

IDENTIFY The wheel exhibits constant angular acceleration, so we use the equations for constant angular acceleration to solve the problem. The target variables are the initial angular velocity and the angular acceleration.

SET UP Figure 9.1 shows a sketch of the pottery wheel. For consistency, we'll use radians and seconds as our units and convert the given quantities.

We are given the final angular velocity, the angular displacement, and the time. We need to find the initial angular velocity and the angular acceleration. None of the equations for constant angular acceleration allow us to solve for both unknowns at once, so we'll solve for the initial angular velocity first and then use the results to solve for the angular acceleration.

EXECUTE The angular displacement can be written in terms of the average angular velocity and the time interval as

$$
\theta-\theta_{0}=\frac{1}{2}\left(\omega_{z}+\omega_{0 z}\right) t
$$

Solving for the initial angular velocity gives

$$
\omega_{0 z}=\frac{2\left(\theta-\theta_{0}\right)}{t}-\omega_{z} .
$$

The quantity $\left(\theta-\theta_{0}\right)$ is our angular displacement, 30.0 revolutions. The quantity $\omega_{z}$ is the final angular velocity, 125 rpm . We convert the revolutions to radians by multiplying by ( $2 \pi \mathrm{rad} / \mathrm{rev}$ ) and the rpm to radians $/$ second by multiplying by $(2 \pi \mathrm{rad} / \mathrm{rev})(1 \mathrm{~min} / 60 \mathrm{~s})$. Solving for $\omega_{0 z}$ gives

$$
\omega_{0 z}=\frac{2(30.0 \mathrm{rev})(2 \pi \mathrm{rad} / \mathrm{rev})}{25.0 \mathrm{~s}}-(125 \mathrm{rpm})(2 \pi \mathrm{rad} / \mathrm{rev})(1 \mathrm{~min} / 60 \mathrm{~s})=1.99 \mathrm{rad} / \mathrm{s}
$$

The initial angular velocity is $1.99 \mathrm{rad} / \mathrm{s}$, or 19.0 rpm . To find the angular acceleration, we use the relationship between angular velocity, angular acceleration, and time:

$$
\omega_{z}=\omega_{0 z}+\alpha_{z} t
$$

Solving for the angular acceleration, we obtain

$$
\alpha_{z}=\frac{\omega_{z}-\omega_{0 z}}{t}=\frac{(125 \mathrm{rpm})(2 \pi \mathrm{rad} / \mathrm{rev})(1 \mathrm{~min} / 60 \mathrm{~s})-(1.99 \mathrm{rad} / \mathrm{s})}{25.0 \mathrm{~s}}=0.444 \mathrm{rad} / \mathrm{s}^{2}
$$

The angular acceleration is $0.444 \mathrm{rad} / \mathrm{s}^{2}$.
EVALUATE This problem reminds us of the problems involving constant linear acceleration we first encountered in Chapter 2. The same problem-solving strategy applies: Draw a diagram, check for constant acceleration, find one or more equations that can be used to solve for the unknowns, and reflect upon the results. As we' ve seen here, we may need to use more than one equation for the solution, and we must watch our units carefully.

## 2: A slowing pottery wheel

A pottery wheel rotating at 30 revolutions per minute is shut off and slows uniformly, coming to a stop in 2 complete revolutions. Find the angular acceleration and the time it takes to come to a stop.

## Solution

IDENTIFY The wheel exhibits constant angular acceleration as it slows uniformly, so we use the equations for constant angular acceleration to solve the problem. The target variables are the angular acceleration and the time required to stop.

SET UP We are given the initial angular velocity, the final angular velocity (zero), and the angular displacement for the wheel to come to a stop. We will use two of the equations for constant angular acceleration to solve for the two unknowns.

EXECUTE The angular displacement, velocity, and acceleration are related by the formula

$$
\omega_{z}^{2}=\omega_{0 z}^{2}+2 \alpha_{z}\left(\theta-\theta_{0}\right) .
$$

Solving for the angular acceleration gives

$$
\alpha_{z}=\frac{\omega_{z}^{2}-\omega_{0 z}^{2}}{2\left(\theta-\theta_{0}\right)}=\frac{(0)^{2}-(30 \mathrm{rev} / \mathrm{min})^{2}}{2(2 \mathrm{rev})}=-225 \mathrm{rev} / \mathrm{min}^{2}\left(\frac{2 \pi \mathrm{rad}}{\mathrm{rev}}\right)\left(\frac{1 \mathrm{~min}}{60 \mathrm{~s}}\right)^{2}=-0.393 \mathrm{rad} / \mathrm{s}^{2}
$$

To solve for time, we use the acceleration-velocity relation:

$$
\omega_{z}=\omega_{0 z}+\alpha_{z} t
$$

Solving for the time yields

$$
t=\frac{\omega_{z}-\omega_{0 z}}{\alpha_{z}}=\frac{(0)-(30 \mathrm{rev} / \mathrm{min})}{-225 \mathrm{rev} / \mathrm{min}^{2}}=0.133 \mathrm{~min}=8.0 \mathrm{~s}
$$

The wheel slows to a stop at a rate of $-0.393 \mathrm{rad} / \mathrm{s}^{2}$, stopping in 8.0 s .
EVALUATE This problem reminds us of our earlier problems involving constant linear acceleration. You should be able to become proficient at these problems rather quickly. Just make sure that you watch your units!

## 3: Energy in a wheel-stone system

A thin light string is wrapped around the rim of a spoked wheel that can rotate without friction around its center axle. An 8.00 kg stone is attached to the end of the string as shown in Figure 9.2. If the stone is released from rest, how far does it travel before attaining a speed of $4.80 \mathrm{~m} / \mathrm{s}$ ? The spoked wheel is made of a central hub (a solid uniform cylinder of radius 7.50 cm and mass 22.0 kg ) attached to a rim (a thin-walled hollow cylinder of radius 30.0 cm mass 12.0 kg ) by spokes of negligible mass.


Figure 9.2 Problem 3.

## Solution

IDENTIFY No work is done by external forces, so mechanical energy is conserved. The target variable is the height the stone falls.

SET UP Initially, there is only gravitational potential energy. When the stone is released, the gravitational potential energy is transformed into kinetic energy of the stone and wheel. We'll set the origin at the point where the stone reaches a speed of $4.80 \mathrm{~m} / \mathrm{s}$; the starting position is a distance $H$ above the origin, as we see in Figure 9.2. We'll find the moment of inertia of the wheel by combining the moment of inertia of a solid cylinder with the moment of inertia of a thin-walled hollow cylinder.

EXECUTE Energy conservation relates the initial and final energies:

$$
K_{1}+U_{1}=K_{2}+U_{2}
$$

Initially, there is only $U_{\text {grav }}$. At the origin, the gravitational potential energy has been transformed into the kinetic energies of the stone and wheel:

$$
U_{\text {grav }}=K_{\text {stone }}+K_{\text {wheel }} .
$$

Replacing the energies yields

$$
m_{\text {stone }} g H=\frac{1}{2} m_{\text {stone }} v^{2}+\frac{1}{2} I_{\text {wheel }} \omega^{2}
$$

We need to find the moment of inertia of the wheel and its angular velocity when the stone reaches its final velocity. The moment of inertia of the wheel is the algebraic sum of the moment of inertia of the central hub plus the moment of inertia of the outer rim (the thin-walled cylinder.) Using Table 9.2 of the text, we find the total moment of inertia of the wheel:

$$
I_{\text {wheel }}=I_{\text {solid cylinder }}+I_{\text {thin-walled cylinder }}=\frac{1}{2} M_{\text {hub }} R_{\text {hub }}^{2}+M_{\text {rim }} R_{\text {rim }}^{2} \text {. }
$$

The speed of the stone is the tangential speed of the wheel, so we can find the angular speed of the wheel from the equation

$$
\omega=\frac{v}{R_{\mathrm{rim}}} .
$$

Substituting the two expressions we found into the energy relation gives

$$
m_{\text {stone }} g H=\frac{1}{2} m_{\text {stone }} v^{2}+\frac{1}{2}\left(\frac{1}{2} M_{\text {hub }} R_{\text {hub }}^{2}+M_{\text {rim }} R_{\text {rim }}^{2}\right)\left(\frac{v}{R_{\text {rim }}}\right)^{2} .
$$

Solving for $H$ results in

$$
\begin{gathered}
H=\frac{1}{m_{\text {stone }} g}\left[\frac{1}{2} m_{\text {stone }} v^{2}+\frac{1}{2}\left(\frac{1}{2} M_{\text {hub }} R_{\text {hub }}^{2}+M_{\text {rim }} R_{\text {rim }}^{2}\right)\left(\frac{v}{R_{\text {rim }}}\right)^{2}\right], \\
H=\frac{1}{(8.00 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}\left[\frac{1}{2}(8.00 \mathrm{~kg})(4.80 \mathrm{~m} / \mathrm{s})^{2}+\frac{1}{2}\left(\frac{1}{2}(22.0 \mathrm{~kg})(0.0750 \mathrm{~m})^{2}+(12.0 \mathrm{~kg})(0.300 \mathrm{~m})^{2}\right)\left(\frac{4.80 \mathrm{~m} / \mathrm{s}}{0.300 \mathrm{~m}}\right)^{2}\right]=3.04 \mathrm{~m} .
\end{gathered}
$$

After the stone falls 3.04 m , its speed will be $4.80 \mathrm{~m} / \mathrm{s}$.
EVALUATE If the stone had fallen freely, it would have attained the final speed when $h=v / \sqrt{2 g}$ $(1.08 \mathrm{~m})$. Why does it take almost three times this distance to reach the final speed? Looking at the energy conservation equation, we see that the energy is shared between the kinetic energies of the stone and wheel. Roughly two-thirds of the energy goes into the kinetic energy of the wheel.

Practice Problem: Repeat the problem with only the inner hub of the wheel. Answer: 1.28 m .

## 4: Energy in a falling cylinder

A thin, light string is wrapped around a solid uniform cylinder of mass $M$ and radius $R$ as shown in Figure 9.3. The string is held stationary and the cylinder is released from rest. What is the cylinder's radius if it reaches an angular speed of 350.0 rpm after it falls 3.00 m ?


## Solution

IDENTIFY Gravity and tension are the only forces acting in this problem, so energy is conserved. The target variable is the cylinder's radius.

SET UP Initially, there is only gravitational potential energy. When the cylinder is released, the gravitational potential energy is transformed into rotational and translational kinetic energy. We'll set the origin 3.00 m below the initial position.

EXECUTE Energy conservation relates the initial and final energies:

$$
K_{1}+U_{1}=K_{2}+U_{2} .
$$

Initially, there is only $U_{\text {grav }}$. At the origin, the gravitational potential energy has been transformed into the total kinetic energy of the cylinder:

$$
U_{\text {grav }}=K_{\text {translational }}+K_{\text {rotational }} .
$$

Replacing the energies, we obtain

$$
M g h=\frac{1}{2} M v^{2}+\frac{1}{2} I \omega^{2} .
$$

Now, recall that the moment of inertia of a solid uniform cylinder is $\frac{1}{2} M R^{2}$. (See Table 9.2 in the text.) The speed of the cylinder is the tangential speed of the wheel, so we can find the angular speed of the wheel from the relation

$$
v=\omega R .
$$

Substituting the preceding expressions into the energy relation gives

$$
M g h=\frac{1}{2} M v^{2}+\frac{1}{2} I \omega^{2}=\frac{1}{2} M(\omega R)^{2}+\frac{1}{2}\left(\frac{1}{2} M R^{2}\right) \omega^{2}=\frac{3}{4} M(\omega R)^{2} .
$$

Solving for $R$, we get

$$
R=\frac{\sqrt{\frac{4}{3} g h}}{\omega}=\frac{\sqrt{\frac{4}{3}\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right)(3.00 \mathrm{~m})}}{(350 \mathrm{rpm})(2 \pi \mathrm{rad} / \mathrm{rev})(1 \mathrm{~min} / 60 \mathrm{~s})}=0.171 \mathrm{~m} .
$$

The cylinder's radius is 17.1 cm .

EVALUATE We see that the mass of the cylinder cancels in the conservation-of-energy equation; the results apply to a cylinder of any mass.

Practice Problem: Repeat the problem with a thin hoop replacing the cylinder. Answer: 14.8 cm .

## 5: Cylinder rolling down a ramp

A uniform hollow cylinder rolls along a horizontal floor and then up a flat ramp without slipping. The ramp is inclined at $20.0^{\circ}$. How far along the ramp does the cylinder roll before stopping if its initial forward speed is $12.0 \mathrm{~m} / \mathrm{s}$ ? The hollow cylinder has a mass of 6.5 kg , an inner radius of 0.13 m , and an outer radius of 0.25 m .


Figure 9.4 Problem 5.

## Solution

IDENTIFY There are no nonconservative forces (ignoring air drag), so we use energy conservation. The target variable is the distance up the ramp the cylinder travels before stopping momentarily.

SET UP Figure 9.4 shows a sketch of the problem. Initially, the cylinder has translational and rotational kinetic energy. When the cylinder stops momentarily at the top of the ramp, the kinetic energy has been totally transformed to gravitational potential energy. We'll set the origin at the base of the ramp.

EXECUTE Energy conservation relates the initial and final energies:

$$
K_{1}+U_{1}=K_{2}+U_{2}
$$

Initially, the cylinder has only kinetic energy. At its highest point, the cylinder has only gravitational potential energy. Thus,

$$
K_{\text {translational }}+K_{\text {rotational }}+0=0+U_{\text {grav }} .
$$

Replacing the energies with their equivalent expressions yields

$$
\left.\frac{1}{2} M v^{2}+\frac{1}{2} I \omega\right)^{2}=M g y .
$$

Recall that the moment of inertia of a hollow uniform cylinder is $\frac{1}{2} M\left(R_{1}^{2}+R_{2}^{2}\right)$. (See Table 9.2 in the text.) The speed of the cylinder is the tangential speed of the wheel at the outer radius, so we can replace the speed with

$$
\omega=\frac{v}{R_{2}} .
$$

Substituting these expressions into the energy relation gives

$$
\frac{1}{2} M v^{2}+\frac{1}{2} I \omega^{2}=\frac{1}{2} M v^{2}+\frac{1}{2}\left(\frac{1}{2} M\left(R_{1}^{2}+R_{2}^{2}\right)\right)\left(\frac{v}{R_{2}}\right)^{2}=M g y .
$$

Solving for $y$, the maximum vertical height, we obtain

$$
y=\frac{v^{2}}{g}\left[\frac{1}{2}+\frac{1}{4}\left(\frac{R_{1}^{2}+R_{2}^{2}}{R_{2}^{2}}\right)\right]=\frac{(12.0 \mathrm{~m} / \mathrm{s})^{2}}{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}\left[\frac{1}{2}+\frac{1}{4}\left(\frac{(0.13 \mathrm{~m})^{2}+(0.25 \mathrm{~m})^{2}}{(0.25 \mathrm{~m})^{2}}\right)\right]=12.0 \mathrm{~m}
$$

The cylinder's maximum vertical height is 12.0 m . To find the distance along the ramp, we use the sine relation:

$$
L=\frac{y}{\sin 20^{\circ}}=\frac{(12.0 \mathrm{~m})}{\sin 20^{\circ}}=35.1 \mathrm{~m}
$$

The hollow cylinder rolls 35.1 m up along the ramp.
EVALUATE How far up the ramp would the cylinder travel without friction? It would move 21.5 m along the ramp without friction. This result shows how the added initial rotational kinetic energy results in a greater distance along the ramp.

## 6: Moment of inertia of a rectangular sheet

Find the moment of inertia of a uniform thin rectangular sheet of metal with mass $M$, length $L$, and width $W$ about the $x$-axis in Figure 9.5.


Figure 9.5 Problem 6.

## Solution

IDENTIFY The sheet is a continuous distribution of mass, so we must integrate to find the moment of inertia. We break the sheet up into thin strips along the $x$-axis as shown. The target variable is the moment of inertia.

SET UP The mass density of the sheet is the total mass divided by the volume, or

$$
\rho=\frac{M}{W L t},
$$

where $t$ is the thickness of the sheet. The volume of the thin strip along the $x$-axis is

$$
d V=W t d y
$$

The mass of the strip is the density multiplied by the volume. We will integrate the product of the mass and $r^{2}$ to find the moment of inertia.

EXECUTE The moment of inertia is given by

$$
I=\int r^{2} d m
$$

We will integrate from $-L / 2$ to $+L / 2$, since the sheet is centered on the origin. Replacing $d m$ and adding the limits of integration gives

$$
\begin{aligned}
I & =\int_{-L / 2}^{L / 2} r^{2} \rho d V \\
& =\int_{-L / 2}^{L / 2} y^{2} \frac{M}{W L t} W t d y \\
& =\frac{M}{L} \int_{-L / 2}^{L / 2} y^{2} d y \\
& =\left.\frac{M}{L} \frac{y^{3}}{3}\right|_{-L / 2} ^{L / 2}=\frac{M}{L}\left[\frac{(L / 2)^{3}}{3}-\frac{(-L / 2)^{3}}{3}\right] \\
& =\frac{M L^{2}}{12} .
\end{aligned}
$$

The moment of inertia of the sheet rotating about the $x$-axis is $M L^{2} / 12$.
EVALUATE We see that the moment of inertia does not depend on the width or thickness of the sheet when it is rotated about the $x$-axis. If we look at other shapes in Table 9.2, we see that neither does the moment of inertia depend upon dimensions along the axis of rotation.

Practice Problem: Use the parallel-axis theorem to check that this result agrees with the moment of inertia of a thin rectangular plate rotated about the edge.

## Try It Yourself!

## 1: Motion of flywheel

A flywheel is a disk-shaped mass that rotates about its central perpendicular axis. A flywheel 1.0 m in diameter rotates with an initial velocity of 500 rpm . It increases its speed to 1000 rpm in 20.0 s . Assuming constant acceleration, find the angular acceleration and the angular displacement of the flywheel as it increases its speed from 500 to 1000 rpm.

## Solution Checkpoints

IDENTIFY AND SET UP Determine the target variables and identify the appropriate constant-angularacceleration equations needed to find the target variables.

EXECUTE The angular acceleration is found from the relation

$$
\omega_{z}=\omega_{0 z}+\alpha_{z} t .
$$

This gives an angular acceleration of $1510 \mathrm{rev} / \mathrm{min}^{2}$. The angular displacement is found from the formula

$$
\omega_{z}^{2}=\omega_{0 z}^{2}+2 \alpha_{z}\left(\theta-\theta_{0}\right) .
$$

This gives an angular displacement of 248 rev.

EVALUATE We see that this problem closely parallels the linear kinematics problems we have seen throughout the previous chapters. We will continue to use angular kinematics as we investigate rotations.

## 2: Energy in a grinding wheel

How much energy is dissipated when a 2.0 kg grinding wheel of radius 0.1 m is brought to rest from an initial velocity of 3000 rpm ? What is the average power dissipated if the wheel stops in 10 rev ? Assume constant angular acceleration.

## Solution Checkpoints

IDENTIFY AND SET UP The grinding wheel loses all of its energy as it stops, so you must find the initial kinetic energy. The moment of inertia of the wheel is that of a disk. To find the power, you must find the time it takes the wheel to stop and divide the energy by the time. The target variables are the energy and the power.

EXECUTE The initial kinetic energy of the grinding wheel is given by

$$
K=\frac{1}{4} M R^{2} \omega^{2}
$$

The final kinetic energy is zero, so the wheel loses 493 J .
We find the time it takes the wheel to stop by directly combining two kinematics equations:

$$
\begin{aligned}
\omega_{z}^{2} & =\omega_{0 z}^{2}+2 \alpha_{z}\left(\theta-\theta_{0}\right) \text { and } \\
\omega_{z} & =\omega_{0 z}+\alpha_{z} t .
\end{aligned}
$$

These equations result in a time of 0.40 s . The power is then the energy lost divided by the time, giving 1.23 kW .

EVALUATE This problem combines energy, kinematics, and power. How can we best check our results?

## Dynamics of Rotational Motion

## Summary

In this chapter, we will investigate the dynamics of rotational motion to learn what gives an object angular acceleration. We will define torque-the turning or twisting effort of a force-and learn how to apply it to both equilibrium and nonequilibrium situations. Work and power for rotating systems will also be investigated. Angular momentum will be introduced and become the basis of an important new conservation law that will lead to an analysis of a spinning gyroscope and the motion called precession. The linear dynamics foundation we have developed throughout the text will help build our intuition about rotational dynamics.

## Objectives

After studying this chapter, you will understand

- The definition and meaning of torque.
- How to identify torques acting on a body.
- The equation of motion for rotational systems and how to apply it to problems.
- How to apply work and power to rotational dynamics problems.
- The definition of angular momentum and how it changes with time.
- How to apply conservation of angular momentum to problems.
- The concept of precession as it applies to gyroscopes.

Concepts and Equations

| Term | Description |
| :--- | :--- |
| Torque | Torque is the tendency of a force to cause or change rotational motion about a <br> chosen axis. The magnitude of torque is the magnitude of the force $(F)$ times <br> the moment arm $(l)$, which is the perpendicular distance between the axis and <br> the line of force: |
| $\quad \tau=F l$. |  |

For a force $\vec{F}$ applied at point $O$ and a vector $\vec{r}$ from the chosen axis to point $O$, the torque is given by

$$
\vec{\tau}=\vec{r} \times \vec{F} .
$$

The SI unit of torque is the newton-meter ( Nm ).

## Combined Translation and Rotation

The motion of a rigid body, moving through space and rotating, can be regarded as translational motion of the center of mass plus rotational motion about the center of mass. The kinetic energy, net force, and net torque are given respectively, by

$$
\begin{aligned}
& K=\frac{1}{2} M v_{\mathrm{cm}}^{2}+\frac{1}{2} I_{\mathrm{cm}} \omega^{2}, \\
& \sum \vec{F}_{\mathrm{ext}}=M \vec{a}_{\mathrm{cm}}, \\
& \sum \vec{\tau}_{z}=I_{\mathrm{cm}} \alpha_{z} .
\end{aligned}
$$

In the case of rolling without slipping, the motion of the center of mass is related to the angular velocity by

$$
v_{\mathrm{cm}}=R \omega
$$

The work done by a torque is given by

$$
W=\int_{\theta_{1}}^{\theta_{2}} \tau_{z} d \theta
$$

For a constant torque,

$$
W=\tau_{z} \Delta \theta .
$$

The power provided by a torque is the product of the torque and the angular velocity of the body:

$$
P=\tau_{z} \omega_{z} .
$$

The angular momentum $L$ of a particle with respect to point $O$ is the vector product of the position vector and momentum of the particle:

$$
\vec{L}=\vec{r} \times \vec{p}
$$

The angular momentum of a symmetrical body rotating about a stationary axis of symmetry is

$$
\vec{L}=I \vec{\omega} .
$$

Our convention is that counterclockwise rotations have positive $L$ and clockwise rotations have negative $L$.

## Rotational Dynamics and Angular Momentum

The net external torque on a system is equal to the rate of change of the angular momentum of the system:

$$
\sum \vec{\tau}=\frac{d \vec{L}}{d t}
$$

If the net external torque acting on a system is zero, the total angular momentum is conserved and remains constant.

## Conceptual Questions

## 1: Ranking torques

In the following diagram, each rod pivots about the indicated axis with the indicated force:


Figure 10.1
Rank the diagrams in order of increasing torque.

## Solution

IDENTIFY, SET UP, AND EXECUTE Torque is given by $\tau=r F \sin \phi$, so we must examine the magnitude of the force, the location at which the force is applied, and the direction in which the force is applied. Comparing (a) and (b), we see that both forces act in the same direction, the force is larger in (a), and the moment arm is larger in (b). We estimate the moment arm in (a) as half the moment arm of (b). Since the force in (a) is less than double the force in (b), the torque in (b) is greater than that in (a). No torque is generated in (c), because the force acts along the moment arm. The force and moment arm are the same in (d) and (e), but the directions are different. However, since both forces act $30^{\circ}$ from the horizontal, the components of the forces perpendicular to the moment arm are the same for both and the torque in (d) and (e) are the same. The vertical component of the force in (d) and (e) is 2.5 N , which is less than the force in (b) and less than half of the force in (a). Diagrams (d) and (e) both show less torque than diagrams (a) and (b).

Ranked in order of increasing torque, the diagrams are therefore (c), (d) = (e), (a), (b).
EVALUATE Torque is the most complicated quantity that we have discussed thus far. It depends on the magnitude and direction of the force responsible for it, as well as where that force acts relative to the rotation axis. Gaining intuition about torque will help guide you through problems.

## 2: Massless versus massive pulleys

Why were massless pulleys used in problems from the previous chapters?

## Solution

IDENTIFY, SET UP, AND EXECUTE Consider the net torque acting on a pulley with a rope resting on the pulley. The left segment of rope has a tension $T_{L}$, and the right segment of rope has a tension $T_{R}$. The net torque is then

$$
\sum \tau=T_{R} R-T_{L} R=I \alpha
$$

Each tension creates a torque $T R$, since the rope will be perpendicular to the radius and the two torques are opposite in direction. Massless pulleys have no moment of inertia (since their mass is zero); therefore, the net torque acting on the pulley is zero. If the net torque is zero, the torques in each segment must be equal and the tensions are equal in each rope.

When we include the pulley's mass, the pulley has a moment of inertia. If the pulley accelerates, there must be a difference in the left and right torques and the tensions must also be different. Only when the angular acceleration is zero are the two tensions equal.

Massless pulleys were used in earlier chapters to avoid having to include torque in our analysis. Using massless pulleys simplified our analysis and let us focus on learning about forces.

EVALUATE From now on, we must assume that the tensions in segments of rope may vary and we must identify each segment's tension separately.

How does the acceleration of each segment of rope compare when we include a pulley's mass? There is no change: Objects connected by the rope are constrained to have the same magnitude of acceleration.

## 3: Spinning on a roundabout

You are standing at the center of a rotating playground roundabout (a round, horizontal plate that spins about its center axis). As you move to the edge, will your angular speed increase, decrease, or stay the same?

## Solution

IDENTIFY, SET UP, AND EXECUTE To simplify our analysis, we ignore friction in the roundabout. Without friction, there is no torque to slow the roundabout, so angular momentum is conserved. As you move to the edge, the moment of inertia of the system increases. (Your mass moves to a greater radius.) For angular momentum to be conserved, the angular speed must be reduced.

EVALUATE Like linear momentum and energy, angular momentum is a conserved quantity. We could solve this problem numerically by picking initial and final angular momenta before and after you moved and setting them equal to each other.

## 4: Standing on a turntable

You are standing on a small frictionless turntable with your arms outstretched and spinning about the axis of the turntable. As you pull your arms in, you spin faster. Does your rotational kinetic energy increase, decrease, or stay the same?

## Solution

IDENTIFY AND SET UP There are no external torques acting on you or the turntable, so angular momentum is conserved. You spin faster, since you decrease your moment of inertia as you bring your arms in, leading to a greater angular velocity. Rotational kinetic energy is given by

$$
K=\frac{1}{2} I \omega^{2} .
$$

Since kinetic energy depends on both moment of inertia and angular velocity, and since both change, we'll have to consider more information. Angular momentum is conserved, so $I \omega$ is constant.

EXECUTE We include angular momentum in the kinetic energy:

$$
K=\frac{1}{2} L \omega .
$$

Since $L$ is constant and $\omega$ increases as you bring your arms in, the kinetic energy increases.
EVALUATE Where does the increased energy come from? You must do work to pull your arms in; doing work on the system increases its kinetic energy.

This example also applies to a figure skater spinning on ice.

## Problems

## 1: Balancing a food tray

A waiter balances a tray of food on his hand. On the tray is a 0.40 kg drink and a 2.0 kg lobster dinner. The drink is placed 6.5 cm from one edge of the tray, and the lobster dinner is placed 8.0 cm from the opposite edge. The tray has a mass of 1.2 kg and a diameter of 42 cm . Where should the waiter hold the tray so that it doesn't tip over?

## Solution

IDENTIFY Figure 10.2 shows a sketch, and a free-body diagram, of the food tray. The tray should be held so that it is in equilibrium and the net torque on it is zero. The target variable is the location at which the waiter's hand holds the tray.

SET UP The forces on the food tray include the force of the waiter's hand holding the tray, the weight of the tray, and the normal forces due to the lobster dinner and drink. To find the location of the hand, we need to consider the torque acting on the tray. When the tray is in equilibrium, the net torque must be zero about any axis. We'll take the axis to be the left edge, marked by an $X$.


Figure 10.2 Problem 1 sketch and free-body diagram.
When we choose the left edge as the axis, we note that four torques act on the tray, corresponding to the four forces on the tray. Each of the four forces is applied perpendicular to the moment arm (the plane of the tray); each torque is the magnitude of the force times the distance from the axis. We'll take counterclockwise torques to be positive.

EXECUTE Since the tray is in equilibrium, the net torque is zero:

$$
\sum \tau=\tau_{\text {tray }}+\tau_{\text {lobster }}+\tau_{\text {drink }}+\tau_{\text {hand }}=0
$$

Writing the four torques explicitly, we have

$$
\sum \tau=\tau_{\text {tray }}+\tau_{\text {lobster }}+\tau_{\text {drink }}+\tau_{\text {hand }}=-m_{\text {tray }} g x_{\text {tray }}-n_{\text {lobster }} x_{\text {lobster }}-n_{\text {drink }} x_{\text {drink }}+n_{\text {hand }} x_{\text {hand }}=0
$$

The first term is the torque due to the weight of the tray; the moment arm is half the tray diameter (at the center of mass). The second and third terms are the torques due to the normal forces of the lobster dinner and drink, respectively. The normal force is equal to the weight of the objects, and the moment arm is the distance from the left edge of the tray. The first three terms are negative, since they are all in the clockwise direction. The last term is the torque due to the normal force of the waiter's hand and is the only positive term, because it is counterclockwise. We need to find this normal force of his hand to solve the problem. We find the force by using Newton's first law:

$$
\sum F=0
$$

In the vertical direction, four forces act on the tray. We have

$$
\sum F=-m_{\text {tray }} g-n_{\text {lobster }}-n_{\text {drink }}+n_{\text {hand }}=0
$$

Solving yields

$$
n_{\text {hand }}=\left(m_{\text {tray }}+m_{\text {lobster }}+m_{\text {drink }}\right) g=((1.2 \mathrm{~kg})+(2.0 \mathrm{~kg})+(0.40 \mathrm{~kg}))\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=35.3 \mathrm{~N} .
$$

We can now solve for the location of the waiter's hand:

$$
\begin{aligned}
x_{\text {hand }} & =\frac{m_{\text {tray }} g x_{\text {tray }}+m_{\text {lobsterg }} g x_{\text {lobster }}+m_{\text {drink }} g x_{\text {drink }}}{n_{\text {hand }}} \\
& =\frac{((1.2 \mathrm{~kg})(0.21 \mathrm{~m})+(2.0 \mathrm{~kg})(0.080 \mathrm{~m})+(0.40 \mathrm{~kg})(0.42 \mathrm{~m}-0.065 \mathrm{~m}))\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{(35.3 \mathrm{~N})} \\
& =0.15 \mathrm{~m} .
\end{aligned}
$$

The waiter should hold the tray 15 cm from the left edge (closest to the lobster dinner).
EVALUATE This equilibrium problem required us to apply both the equilibrium torque and the equilibrium force conditions to solve the problem. You should be familiar with the equilibrium force condition and should need to gain expertise only in torque to solve similar problems.

Experience shows that we may be able to simplify similar problems by picking an axis that coincides with the location at which a force is applied. This choice will reduce the number of torques in the problem. In the current problem, for example, we could have chosen the axis to be at the location of the lobster dinner, thus removing the torque due to the normal force of the lobster dinner.

## 2: Tension in string attached to a falling cylinder

A thin, light string is wrapped around the outer rim of a uniform hollow cylinder of mass 12.0 kg , inner radius 15.0 cm , and outer radius 30.0 cm . The cylinder is released from rest. What is the tension in the string as the cylinder falls?

## Solution

IDENTIFY We will apply Newton's second law and its rotational analog to solve for the tension, the target variable.

SET UP Figure 10.3 shows a sketch, and a free-body diagram, of the situation. As the cylinder falls, it will accelerate downward and rotate about its central axis. In falling, the cylinder will rotate faster and undergo angular acceleration. The cylinder has both a net force and a net torque acting on it.


Figure 10.3 Problem 2 sketch and free-body diagram.
EXECUTE We first apply Newton's second law to the translational motion of the center of mass in the vertical direction. The only forces acting in the vertical direction are gravity and tension. We have

$$
\sum F_{y}=M g-T=M a_{\mathrm{cm}, y}
$$

The moment of inertia of the hollow cylinder is $I=\frac{1}{2} M\left(R_{\text {inner }}^{2}+R_{\text {outer }}^{2}\right)$. The one torque acting on the cylinder as it rotates about its central axis is due to the tension force. Gravity acts on the center of mass, but creates no torque about the central axis. The torque acts perpendicular to the outer radius. Thus,

$$
\sum \tau_{z}=T R_{\text {outer }}=I \alpha_{z}=\frac{1}{2} M\left(R_{\text {inner }}^{2}+R_{\text {outer }}^{2}\right) \alpha_{z} .
$$

We can relate the two accelerations, since the cylinder falls without slipping:

$$
a_{y}=R_{\text {outer }} \alpha_{z} .
$$

Solving for the acceleration of the center of mass in the first equation yields

$$
\boldsymbol{a}_{y}=\frac{M g-T}{M}
$$

Substituting the last two equations into the torque result gives

$$
\begin{aligned}
& T R_{\text {outer }}=\frac{1}{2} M\left(R_{\text {inner }}^{2}+R_{\text {outer }}^{2}\right) \frac{a_{y}}{R_{\text {outer }}}=\frac{1}{2} M\left(R_{\text {inner }}^{2}+R_{\text {outer }}^{2}\right) \frac{M g-T}{M R_{\text {outer }}} \\
& T=\frac{\left(R_{\text {inner }}^{2}+R_{\text {outer }}^{2}\right) M g}{R_{\text {inner }}^{2}+3 R_{\text {outer }}^{2}}=\frac{\left((15.0 \mathrm{~cm})^{2}+(30.0 \mathrm{~cm})^{2}\right)(12.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{(15.0 \mathrm{~cm})^{2}+3(30.0 \mathrm{~cm})^{2}}=45.2 \mathrm{~N} .
\end{aligned}
$$

The tension in the string is 45.2 N .
EVALUATE This problem resembles our earlier Newton's-law problems, but with the addition of torque. Once torque is included, the problem becomes a relatively straightforward algebraic one. We also see that the tension in the string is less than that for a stationary cylinder. If the cylinder were stationary, the tension would have been 118 N .
Practice Problem: At what rate does the cylinder accelerate? Answer: $6.03 \mathrm{~m} / \mathrm{s}^{2}$.

## 3: Acceleration and tension of two blocks connected by a pulley

Two blocks are connected to each other by a light cord passing over a pulley as shown in Figure 10.4. Block $A$ has a mass of 5.00 kg and block $B$ has a mass of 4.00 kg . The pulley has a mass of 8.00 kg and a radius of 4.00 cm . Find the acceleration of the blocks and the tensions in the horizontal and vertical segments of the cord. Assume that the pulley is a solid, uniform disk and there is no friction between block $A$ and the table.


Figure $\mathbf{1 0 . 4}$ Problem 3.

## Solution

IDENTIFY We'll apply the net-force and net-torque equations to solve the problem. The accelerations of both blocks are the same, as we saw in Chapter 5. Our target variables are the two tensions and the acceleration of the blocks.

SET UP Figure 10.5 shows the free-body diagram of the two blocks and the pulley. The forces on the blocks include tension, gravity, and the normal force (block $A$ ). We assume that the tensions of the two segments are not equal and label them $T_{A}$ and $T_{B}$. The two tension forces lead to two torques acting on the pulley. (The axis of rotation is the center of the pulley.)


Figure 10.5 Problem 3 free-body diagrams.
EXECUTE We first apply Newton's second law to each block. Block $A$ (with mass $m_{A}$ ) accelerates in the $x$ direction due to tension $T_{A}$, so

$$
\sum F_{x}=T_{A}=m_{A} a .
$$

Gravity and tension $T_{B}$ act on block $B$ (with mass $m_{B}$ ) and that block accelerates at the same rate as block $A$ in the $y$ direction, so

$$
\sum F_{y}=m_{B} g+\left(-T_{B}\right)=m_{B} a .
$$

As block $B$ falls, the pulley's rotational speed increases. The net torque on the pulley is

$$
\sum \tau_{z}=\tau_{A}-\tau_{B}=I \alpha_{z},
$$

where we have taken the counterclockwise torque as positive and clockwise torque as negative. The moment of inertia of a uniform cylinder is $I=\frac{1}{2} M R^{2}$. We assume that the cord doesn't slip on the pulley, so we relate the angular acceleration of the pulley to the tangential acceleration of the cord (a):

$$
\alpha_{z}=-\frac{a}{R} .
$$

We included a minus sign in the equation because the pulley rotates clockwise (negative, according to our convention). The tension forces act perpendicular to the moment arm, so the torques are simply $T R$. Rewriting the net torque, we have

$$
\sum \tau_{z}=\tau_{A}-\tau_{B}=T_{A} R-T_{B} R=I \alpha_{z}=\frac{1}{2} M R^{2}\left(-\frac{a}{R}\right)
$$

Simpifying yields

$$
T_{B}-T_{A}=\frac{1}{2} M a .
$$

Our second-law equations are used to replace the tensions:

$$
m_{B} g-m_{B} a-m_{A} a=\frac{1}{2} M a .
$$

Solving for the acceleration gives

$$
a=\frac{m_{B} g}{m_{B}+m_{A}+\frac{1}{2} M}=\frac{(4.00 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{(4.00 \mathrm{~kg})+(5.00 \mathrm{~kg})+\frac{1}{2}(8.00 \mathrm{~kg})}=3.02 \mathrm{~m} / \mathrm{s}^{2} .
$$

Using this result to find the two tensions, we get

$$
\begin{aligned}
& T_{A}=m_{A} a=(5.00 \mathrm{~kg})\left(3.02 \mathrm{~m} / \mathrm{s}^{2}\right)=15.1 \mathrm{~N} \\
& T_{B}=m_{B} g-m_{B} a=(4.00 \mathrm{~kg})\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right)-(4.00 \mathrm{~kg})\left(3.02 \mathrm{~m} / \mathrm{s}^{2}\right)=27.1 \mathrm{~N} .
\end{aligned}
$$

The blocks accelerate at $3.02 \mathrm{~m} / \mathrm{s}^{2}$, the tension in the horizontal segment of the cord is 15.1 N , and the tension in the vertical segment of the cord is 27.1 N .

EVALUATE The tensions in the two segments of the cord differ by almost a factor of two. Recall from our problems in Chapter 5 that the tension was constant in both segments of the cord. What causes this difference? The current problem includes the pulley's mass, resulting in some energy spent on increasing the pulley's angular velocity, leaving less energy available for the blocks. This problem also illustrates why we let the first pulleys we encountered be massless.

## 4: Solid cylinder rolling down ramp

A solid cylinder rolls without slipping down an incline of $40^{\circ}$. Find the acceleration and minimum coefficient of friction needed to prevent slipping.

## Solution

IDENTIFY We'll apply translational and rotational dynamics to the cylinder. Since the cylinder doesn't slip, we will use the relationship between the linear and angular acceleration of the cylinder. The target variables are the acceleration and coefficient of friction.

SET UP Figure 10.6 shows a sketch, and a free-body diagram, of the cylinder on the incline. Gravity, the normal force, and the frictional force act on the cylinder. If we set the axis of rotation at the center of the cylinder, a torque due to friction acts on the cylinder. We'll apply the net-force and net-torque equations to solve the problem.

A rotated coordinate system is included in the free-body diagram to simplify our analysis.


Figure 10.6 Problem 4 sketch and free-body diagram.
EXECUTE We first apply Newton's second law to the translational motion along the $x$-axis:

$$
\sum F_{x}=M g \sin \theta-f_{\mathrm{s}}=M a_{x} .
$$

The equation of motion for the rotation about the axis is

$$
\sum \tau_{z}=f_{s} R=I_{\mathrm{cm}} \alpha_{z}=\frac{1}{2} M R^{2} \alpha_{z}
$$

where we included the moment of inertia ( $I=\frac{1}{2} M R^{2}$.) The translational and rotational accelerations are related by $a_{\mathrm{cm}}=\alpha_{z} R$, since the cylinder rolls without slipping. Combining and writing the second equation in terms of $f_{s}$ yields

$$
f_{s}=\frac{1}{2} M a_{x} .
$$

We use this result in the first equation and solve for the acceleration:

$$
\begin{aligned}
& M g \sin \theta-\frac{1}{2} M a_{x}=M a_{x}, \\
& a_{x}=\frac{2}{3} g \sin \theta=\frac{2}{3}\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right) \sin 40^{\circ}=4.20 \mathrm{~m} / \mathrm{s}^{2} .
\end{aligned}
$$

To find the minimum coefficient of static friction, we use the equilibrium equation along the $y$-axis:

$$
\sum F_{y}=n-m g \cos \theta=0 .
$$

The friction force is then

$$
f_{\mathrm{s}}=\mu_{\mathrm{s}} n=\mu_{\mathrm{s}} m g \cos \theta=\frac{1}{2} M a_{x} .
$$

Solving for $\mu_{\mathrm{s}}$ gives

$$
\mu_{\mathrm{s}}=\frac{\frac{1}{2} a_{x}}{g \cos \theta}=\frac{\frac{1}{2}\left(4.20 \mathrm{~m} / \mathrm{s}^{2}\right)}{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \cos 40^{\circ}}=0.280
$$

The cylinder accelerates at $4.20 \mathrm{~m} / \mathrm{s}^{2}$, and the minimum coefficient of static friction that will prevent slipping is 0.280 .

Evaluate This solution is a straightforward application of net-force and net-torque problem-solving techniques. We see that neither the mass nor the radius of the cylinder affect the results, which are valid for any cylinder rolling down a $40^{\circ}$ incline.

## 5: Pivoting rod

A rod of mass 1.0 kg and length 0.50 m is connected to the ceiling by a frictionless hinge. It is released from rest when it is in the horizontal position. What is its angular velocity when it is vertical?

## Solution

IDENTIFY There are no dissipative forces, so we will use conservation of energy. The rod begins with only gravitational potential energy, which is subsequently transformed to kinetic energy. The target variable is the rod's angular velocity when the rod is vertical.

SET UP Figure 10.7 shows a diagram of the rod as it falls. The center of mass of the rod is at the center of the rod; the gravitational potential energy depends on the position of the center of mass. As the rod falls, its rotational kinetic energy increases.

We set the origin of the coordinate axis at the center of the rod when it is in the vertical position. The initial height of the rod in this coordinate system is $h=L / 2$.


Figure 10.7 Problem 5 sketch and free-body diagram.
EXECUTE Conservation of energy gives

$$
K_{1}+U_{1}=K_{2}+U_{2}
$$

Initially, there is only gravitational potential energy $(M g L / 2)$. When the rod is in the vertical position, there is only kinetic energy. The kinetic energy of the rod is

$$
K=\frac{1}{2} I \omega_{z}^{2}=\frac{1}{2}\left(\frac{1}{3} M L^{2}\right) \omega_{z}^{2} .
$$

Combining these energies, we find that

$$
\begin{aligned}
& U_{1}=K_{2} \\
& M g \frac{L}{2}=\frac{1}{6} M L^{2} \omega_{z}^{2}
\end{aligned}
$$

Solving for angular velocity gives

$$
\omega_{z}=\sqrt{\frac{3 g}{L}}=\sqrt{\frac{3\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{(0.5 \mathrm{~m})}}=7.67 \mathrm{rad} / \mathrm{s} .
$$

The angular velocity of the rod when it is vertical is $7.67 \mathrm{rad} / \mathrm{s}$.

EVALUATE Could we have obtained this result by using rotational dynamics? We could have, although it would have required additional calculations. The forces and torques change with position; hence, the solution would have been more challenging.

## 6: Baggage carousel

A baggage carousel has a mass of 500.0 kg and can be approximated as a disk of radius 2.0 m . It is rotating freely at an angular velocity of $1.0 \mathrm{rad} / \mathrm{s}$ when 10 pieces of baggage, each with a mass of 20.0 kg , are dropped on the edge of the carousel. Assuming that no external torques act on the system, what is the final angular velocity of the system?

## Solution

IDENTIFY Since there are no external torques, the total angular momentum of the system remains constant. We will find the initial angular momentum and set it equal to the final angular momentum. The target variable is the carousel's angular velocity after the baggage is added.

SET UP The initial angular momentum is that of the carousel. As the bags are added, they share the angular momentum, resulting in a slower final angular velocity.

EXECUTE The initial angular momentum is

$$
L_{1}=I_{\text {carousel }} \omega_{1}=\left(\frac{1}{2} M_{\text {carousel }} R_{\text {carousel }}^{2}\right) \omega_{1} .
$$

The final angular momentum will be the angular momentum of the carousel plus the angular momentum of the 10 bags:

$$
L_{2}=I_{\text {carousel }} \omega_{2}+10 m_{\text {bag }} r_{\text {bag }}^{2} \omega_{2}=\left(\frac{1}{2} M_{\text {carousel }} R_{\text {carousel }}^{2}\right) \omega_{2}+10 m_{\text {bag }} r_{\text {bag }}^{2} \omega_{2}
$$

Equating the initial and final angular momenta gives

$$
\left(\frac{1}{2} M_{\text {carousel }} R_{\text {carousel }}^{2}\right) \omega_{1}=\left(\frac{1}{2} M_{\text {carousel }} R_{\text {carousel }}^{2}\right) \omega_{2}+10 m_{\text {bag }} r_{\text {bag }}^{2} \omega_{2}
$$

All of the radii are 2.0 m , so they cancel. Solving for the final angular velocity gives

$$
\omega_{2}=\frac{\left(\frac{1}{2} M_{\text {carousel }}\right) \omega_{1}}{\left(\frac{1}{2} M_{\text {carousel }}\right)+10 m_{\text {bag }}}=\frac{\left(\frac{1}{2}(500.0 \mathrm{~kg})\right)(1.0 \mathrm{rad} / \mathrm{s})}{\left(\frac{1}{2}(500.0 \mathrm{~kg})\right)+10(20.0 \mathrm{~kg})}=0.556 \mathrm{rad} / \mathrm{s}
$$

The final angular velocity of the system is $0.556 \mathrm{rad} / \mathrm{s}$.
EVALUATE We see how to apply conservation of angular momentum in this problem. Was energy conserved? Even though there were no external torques, there must have been external forces, since energy was not conserved.
Practice Problem: How much energy was lost as the baggage was added? Answer: 218 J.

## 7: Rotating mass on a string

A 0.10 kg block of mass is attached to a cord that passes through a hole in a horizontal frictionless surface. The block initially is rotating in a circle of radius 0.20 m at an angular velocity of $7.0 \mathrm{rad} / \mathrm{s}$. A force is applied to the cord, shortening it to 0.10 m . What is the new angular velocity of the block?

## Solution

IDENTIFY AND SET UP There is no external torque, as the force exerts no torque at the hole. Therefore, the total angular momentum of the system remains constant. The target variable is the block's final angular velocity.

EXECUTE The initial angular momentum is

$$
L_{1}=I_{1} \omega_{1}=m r_{1}^{2} \omega_{1}
$$

The final angular momentum is

$$
L_{2}=I_{2} \omega_{2}=m r_{2}^{2} \omega_{2}
$$

Setting these equal to each other gives

$$
m r_{1}^{2} \omega_{1}=m r_{2}^{2} \omega_{2} .
$$

Solving results in

$$
\omega_{2}=\frac{r_{1}^{2}}{r_{2}^{2}} \omega_{1}=4(7.0 \mathrm{rad} / \mathrm{s})=28 \mathrm{rad} / \mathrm{s}
$$

The final angular velocity of the block is $28 \mathrm{rad} / \mathrm{s}$.
EVALUATE We see that the final angular velocity is greater than the initial angular velocity, since the radius decreased. The final result does not depend on the mass of the block. We can go on to find the amount of work done by the force by comparing the initial and final kinetic energies.

Practice Problem: How much work did the force do when shortening the cord? Answer: 0.29 J .

## Try It Yourself!

## 1: Torque in a grinding wheel

How much torque is required to bring a 2.0 kg grinding wheel of radius 0.1 m to rest from an initial velocity of 3000 rpm ? The grinding wheel stops in 10 rev . How much work is done by the torque to bring the grinding wheel to a halt? Assume constant angular acceleration.

## Solution

IDENTIFY AND SET UP We apply our results from Try It Yourself! Problem 9.2 as a starting point. We begin by finding the angular acceleration and use that in combination with the moment of inertia to find the torque. The work is the torque times the angular displacement.

EXECUTE The angular acceleration is $-4.5 \times 10^{5} \mathrm{rev} / \mathrm{min}^{2}$, which, when combined with the moment of inertia, gives a torque of -7.85 Nm .

The work done by the torque is -493 J .
EVALUATE Why are the angular acceleration, torque, and work done all negative values? The negative sign indicates that the grinding wheel is slowing.

## 2: Mass on a flywheel

A cord is wrapped around the rim of a uniform flywheel of radius 02.0 m and mass 10.0 kg . A 10.0 kg mass is suspended from the cord 10.0 m above the floor. How much time does it take the mass to hit the floor? What is the tension in the rope as it falls?

## Solution

IDENTIFY AND SET UP Apply the net-force and net-torque equations to solve the problem. The length of cord pulled from the flywheel is equal to the arc length $r \theta$ at the wheel.

EXECUTE Combining the torque and force equations yields the relation

$$
a=\frac{g}{1+\frac{m_{\text {flywheel }}}{2 m_{\text {mass }}}}
$$

for the acceleration. This result can be combined with linear kinematics to find the time to fall, 1.74 s . The tension in the rope is found from the net-force equation:

$$
m_{\mathrm{mass}} g-T=m_{\mathrm{mass}} a .
$$

The tension is 32.7 N .
EVALUATE Can energy conservation be used to check the results?

## 3: Man on a turntable

A turntable with moment of inertia of $2000 \mathrm{kgm}^{2}$ makes one revolution every 5.0 s . A man of mass 100 kg standing at the center of the turntable runs out along a radius fixed on the turntable. What is the angular velocity of the turntable when the man is 3.0 m from the center?

## Solution

IDENTIFY AND SET UP There are no external torques acting on the system, so angular momentum is conserved. As the man runs out, he changes the angular momentum of the system.

EXECUTE Conservation of angular momentum gives

$$
I_{\text {turntable }} \omega_{1}=I_{\text {turntable }} \omega_{2}+M R^{2} \omega_{2}
$$

This results in a final angular velocity of $0.14 \mathrm{rev} / \mathrm{s}$, or one revolution every 7.25 s .
EVALUATE Does the man do positive or negative work on the system? Work is done on the man as he runs out.

## Problem Summary

These first 10 chapters of the book form the basis of kinematic and dynamic problem-solving techniques. Future chapters will expand to include additional forces and forms of energy, but still utilize the same problem-solving techniques. Our problem-solving methodology continues to encompass the following techniques:

- Identifying the general procedure to find the solution.
- Sketching the situation when no figure is provided.
- Identifying the forces and torques acting on the bodies.
- Identif ying the forms of energy included in the problem.
- Drawing free-body diagrams of the bodies.
- Applying appropriate coordinate systems to the diagrams.
- Applying the equations of motion to find relations among the forces, masses, and accelerations.
- Applying conservation of energy and conservation of momentum when appropriate.
- Solving the equations through algebra and substitutions.
- Reflecting on the results and checking for inconsistencies.

Expert problem solvers use this foundation at all levels of physics investigations, from introductory courses through cutting-edge research projects.

## 11 <br> Equilibrium and Elasticity

## Summary

We will explore equilibrium and elasticity in this chapter. We will focus on extended bodies in equilibrium: bodies having no net force or torque acting on them. Bodies deform when forces act on them, so we will examine deformations that describe the stretching, twisting, and compressing of a body. New concepts and principles will be introduced to quantify deformations-ideas based on the concepts and principles we encountered in previous chapters. We will learn about stress, strain, and elastic modulus, and we will further clarify Hooke's law.

## Objectives

After studying this chapter, you will understand:

- The conditions required for a body to be in equilibrium.
- The definition of center of gravity and how to apply it to a problem.
- How to solve problems when bodies are in equilibrium.
- How to analyze problems involving the deformation of bodies.
- Stress and strain with respect to tension, compression, and shear forces.
- How to use Young's, bulk, and shear moduli to predict the changes due to stress.
- The limits of stress and strain.

Concepts and Equations
Term

## Center of Mass

The torque due to the weight of a body is found by assuming that the entire weight of the body is located at the center of gravity, given by

$$
\vec{r}_{\mathrm{cm}}=\frac{m_{1} \vec{r}_{1}+m_{2} \vec{r}_{2}+m_{3} \vec{r}_{3}+\cdots}{m_{1}+m_{2}+m_{3}+\cdots}
$$

The center of gravity is equivalent to the center of mass when gravity is constant.

## Stress and Strain

Stress characterizes the strength of a force that stretches, squeezes, or twists an object. Strain is the resulting deformation. Stress and strain are often directly proportional, with the proportionality-the elastic modulus-given by Hooke's law:

$$
\text { elastic modulus }=\frac{\text { stress }}{\text { stain }}
$$

## Tensile and Compressive Stress

Tensile stress is the ratio of the perpendicular component of a force to the cross-sectional area where the force is applied:

$$
\text { Tensile stress }=\frac{F_{\perp}}{A} .
$$

The SI unit of stress is the pascal ( Pa ), equal to 1 newton per meter squared. Tensile strain is the ratio of the change in an object's length under stress to its original length:

$$
\text { Tensile Strain }=\frac{\Delta l}{l_{0}}
$$

Young's modulus $Y$ is the elastic modulus:

$$
Y=\frac{\text { Tensile stress }}{\text { Tensile strain }}=\frac{F_{\perp} / A}{\Delta l / l_{0}} .
$$

Compressive stress and strain are defined in the same manner.
Bulk Stress
The pressure in a fluid is the force per unit area of the fluid:

$$
p=\frac{F_{\perp}}{A} .
$$

Bulk stress is the change in pressure, and bulk strain is the fractional change in volume, of the fluid. The bulk modulus is the elastic modulus:

$$
B=\frac{\text { Bulk stress }}{\text { Bulk Strain }}=-\frac{\Delta p}{\Delta V / V_{0}} .
$$

| Shear Stress | Shear stress is the force tangent to an object's surface, divided by the area on <br> which the force acts. The shear modulus $(S)$ is the ratio of shear stress to <br> strain: |
| :--- | :--- |
| $\qquad S=\frac{\text { Shear stress }}{\text { Shear strain }}=\frac{F_{\\|} / A}{x / h}$. |  |
| Limits of Hooke's law | There is a maximum stress for which stress and strain are proportional, <br> beyond which Hooke's law is not valid. The elastic limit is the stress beyond <br> which irreversible deformation occurs. |

## Conceptual Questions

## 1: Body in equilibrium

A body is acted upon by no net force and no net torque. Is it at rest?

## Solution

IDENTIFY, SET UP, AND EXECUTE A body that is moving with constant velocity is in equilibrium. A body could exhibit translational motion with a constant velocity or rotate with a constant angular velocity (or both) and remain in equilibrium.

EVALUATE Just as we saw with forces, constant velocity is a state of equilibrium.

## 2: Ladder on a frictionless surface

A ladder is placed against a wall. The wall is rough, but the floor is frictionless. Can the ladder be in equilibrium?

## Solution

IDENTIFY, SET UP, AND EXECUTE Four forces act on the ladder: the normal force due to the wall, the normal force due to the floor, gravity, and friction due to the wall. Three of the forces act in the vertical direction: gravity (acting downward), friction due to the wall (acting upward), and the normal force due to the floor (acting upward). These forces can sum to zero, since they act in different directions. One force acts in the horizontal direction: the normal force due to the wall. Since there is only one force, the net force on the ladder cannot be zero and therefore the ladder cannot be in equilibrium.

EVALUATE Equilibrium is needed in order for us to use the ladder. How can we achieve equilibrium? Friction with the floor is needed to establish equilibrium. Can the ladder be in equilibrium if placed against a frictionless wall on a rough floor? Yes, the forces and torques can be in equilibrium in this case.

## Problems

## 1: Forces on a diving board

A 4.0-m-long diving board with a uniform mass of 150.0 kg is mounted as shown in Figure 11.1. Find the forces holding the board in place when a 100.0 kg man is standing on the end of the board.


Figure 11.1 Problem 1.

## Solution

IDENTIFY We'll use the conditions of equilibrium to solve the problem. The board has two forces holding it in place: a downward force at the left end and an upward force at the pivot. The target variables are the forces acting on the board.

SET UP Figure 11.2 shows the free-body diagram of the board. Forces $A$ and $B$ hold the board in place, the weight of the board acts at the board's center, and the weight of the man acts at the end. We will take counterclockwise torques as positive.


Figure 11.2 Problem 1 free-body diagram.

EXECUTE Newton's first law applied to the board gives

$$
\sum F_{y}=0=-F_{A}+F_{B}-m_{\text {board }} g-m_{\operatorname{man}} g .
$$

We have two unknowns in this equation, so we need to use the net torque equation. The net torque about the left end is zero:

$$
\sum \tau=0=F_{B}(1.0 \mathrm{~m})-m_{\text {board }} g(2.0 \mathrm{~m})-m_{\operatorname{man}} g(4.0 \mathrm{~m})
$$

Solving for the force at $B$ gives

$$
F_{B}=m_{\text {board }} g(2.0)+m_{\operatorname{man} g} g(4.0)=6860 \mathrm{~N} .
$$

Substituting and solving for the force at $A$ gives

$$
F_{A}=F_{B}-m_{\text {board }} g-m_{\operatorname{man}} g=4410 \mathrm{~N} .
$$

The force at the left end is 4410 N downward, and the force at the pivot point is 6860 N upward.
EVALUATE To simplify our analysis, we chose the axis for the net torque such that the torque due to force $A$ was zero. We can double check the result by calculating the torque about the pivot point. If we do, we find the same result.

## 2: Force on a support strut

Find the force exerted by the wall on the uniform strut shown in Figure 11.3 if the strut weighs 100.0 N .


Figure 11.3 Problem 2.

## Solution

IDENTIFY We'll use equilibrium conditions to solve the problem. The strut has four forces acting on it; both the net force and the net torque are zero. The target variable is the force acting on the strut due to the wall.

SET UP Figure 11.4 shows the free-body diagram of the strut. The weights of the strut and hanging mass, tension, and the force of the wall act on the strut. We assume that the force of the wall on the strut acts to the right and upward. The forces act in two directions, so we'll need to include net forces in those directions. We will take counterclockwise torques as positive.


Figure 11.4 Problem 2 free-body diagram.

EXECUTE Newton's first law applied in the $x$ direction gives

$$
\sum F_{x}=0=F_{W x}-T \cos \theta
$$

Newton's first law applied in the $y$ direction gives

$$
\sum F_{y}=0=F_{W y}+T \sin \theta-w_{S}-w_{M} .
$$

We have three unknowns in these equations, so we need to use the net torque equation. The net torque about the left end of the strut is zero:

$$
\sum \tau=0=-w_{S}(2.0 \mathrm{~m})-w_{M}(4.0 \mathrm{~m})+T \sin \theta(4.0 \mathrm{~m})
$$

Inspecting the figure, we see that the sine and cosine of the angle are $3 / 5$ and $4 / 5$, respectively. We solve the torque equation for tension, giving

$$
T=\frac{w_{S}(2.0 \mathrm{~m})+w_{M}(4.0 \mathrm{~m})}{\sin \theta(4.0 \mathrm{~m})}=583 \mathrm{~N} .
$$

With the tension, we can solve for the components of the force due to the wall:

$$
\begin{aligned}
& F_{W x}=T \cos \theta=467 \mathrm{~N} \\
& F_{W y}=w_{S}+w_{M}-T \sin \theta=50 \mathrm{~N} .
\end{aligned}
$$

The magnitude of the force is

$$
F_{W}=\sqrt{F_{W x}^{2}+F_{W y}^{2}}=470 \mathrm{~N},
$$

and it acts at an angle

$$
\phi=\tan ^{-1}\left(\frac{F_{W y}}{F_{W x}}\right)=83.9^{\circ}
$$

above the positive $x$-axis.
EVALUATE We see that we chose the correct directions for the force of the wall on the strut. If we hadn't guessed correctly, we would have found negative results for one or both of the force components.

You can double check the result by calculating the torque about any other point. Do you find the same result when you do?

CAUTION Pick pivots carefully! Carefully choosing your pivot point simplifies the net torque equation, as we have seen in the previous two examples. Note how the pivot point is chosen in the next two problems.

## 3: Boom in equilibrium

A horizontal wire supports a boom of length $L$. The boom supports a 200.0 N weight as shown in Figure 11.5. The boom weighs 200.0 N. Find the tension in the wire and the force exerted by the ground on the boom.


Figure 11.5 Problem 3.

## Solution

IDENTIFY We'll use equilibrium conditions to solve the problem. The boom has four forces acting on it; both the net force and the net torque are zero. The target variables are the tension and the force acting on the boom due to the ground.

SET UP Figure 11.6 shows the free-body diagram of the boom. The weights of the boom and the hanging mass, tension, and the force of the floor act on the boom. We assume that the force of the ground on the boom acts to the right and upward. The forces act in two directions, so we'll need to include net forces in those directions. We will take counterclockwise torques as positive.


Figure 11.6 Problem 3 free-body diagram.

EXECUTE Newton's first law applied in the $x$ direction gives

$$
\sum F_{x}=0=F_{x}-T .
$$

Newton's first law applied in the $y$ direction gives

$$
\sum F_{y}=0=F_{y}-w-w .
$$

We have three unknowns in these equations, so we need to use the net torque equation. The net torque about the bottom end of the boom is zero:

$$
\sum \tau=0=L T \sin \theta-L w \cos \theta-\frac{1}{2} L w \cos \theta
$$

We solve the torque equation for tension, giving

$$
T=\frac{3}{2} \frac{w \cos \theta}{\sin \theta}=\frac{3}{2} \frac{(200.0 \mathrm{~N}) \cos 60^{\circ}}{\sin 60^{\circ}}=173 \mathrm{~N} .
$$

With the tension, we can solve for the components of the force due to the ground:

$$
\begin{aligned}
& F_{x}=T=173 \mathrm{~N} \\
& F_{y}=-2 w=400.0 \mathrm{~N}
\end{aligned}
$$

The magnitude of the force is

$$
F=\sqrt{F_{x}^{2}+F_{y}^{2}}=436 \mathrm{~N}
$$

and it acts at an angle

$$
\phi=\tan ^{-1}\left(\frac{F_{y}}{F_{x}}\right)=66.6^{\circ}
$$

above the positive $x$-axis.
EVALUATE We see that we chose the correct directions for the force of the ground on the boom. If we hadn't guessed correctly, we would have found negative results for one or both of the force components.

You can double check the result by calculating the torque about any other point. Do you find the same result when you do?

## 4: Coefficient of friction for a strut

Find the minimum coefficient of friction between the weightless horizontal strut and the wall in the system shown in Figure 11.7.


Figure 11.7 Problem 4.

## Solution

IDENTIFY We'll use equilibrium conditions to solve the problem. The strut has four forces acting on it; both the net force and the net torque are zero. The target variable is the coefficient of friction at the wall.

SET UP Figure 11.8 shows the free-body diagram of the strut. The weights of the hanging mass, tension, friction, and the normal force of the wall act on the strut. The forces act in two directions, so we'll need to include net forces in those directions. We will take counterclockwise torques as positive.


Figure 11.8 Problem 4 free-body diagram.
EXECUTE Newton's first law applied in the $x$ direction gives

$$
\sum F_{x}=0=n-T \cos 30^{\circ}
$$

Newton's first law applied in the $y$ direction gives

$$
\sum F_{y}=0=f+T \sin 30^{\circ}-w
$$

We have three unknowns in these equations, so we need to use the net torque equation. The net torque about the left end of the strut is zero:

$$
\sum \tau=0=-w \frac{1}{4} L+T \sin 30^{\circ} L
$$

We solve the torque equation for tension, giving

$$
T=\frac{w}{4 \sin 30^{\circ}} .
$$

With the tension, we can solve for the normal and friction forces:

$$
\begin{aligned}
& n=T \cos 30^{\circ}=\frac{w \cos 30^{\circ}}{4 \sin 30^{\circ}} \\
& f=w-T \sin 30^{\circ}=w-\frac{1}{4} w=\frac{3}{4} w .
\end{aligned}
$$

For there to be no slipping,

$$
f<\mu n
$$

Solving for the coefficient of friction gives

$$
\mu=\frac{f}{n}=\frac{\frac{3}{4} w}{\frac{w}{4 \tan 30^{\circ}}}=3 \tan 30^{\circ}=1.73
$$

The minimum coefficient of friction is 1.73 .
EVALUATE How can you double-check the result? Do you find the same result when you do?

## 5: Strain in a steel cable

A 10.0 kg weight is hung from a steel wire having an unstretched length of 1.0 m and a diameter of 2.0 mm . How much does the wire stretch?

## Solution

IDENTIFY The force acting on the cable is the weight of the mass. Young's modulus will be used to find the change in length.

SET UP We look up Young's modulus for steel and find that it is $2 \times 10^{11} \mathrm{~Pa}$.
EXECUTE Young's modulus is

$$
Y=\frac{F / A}{\Delta L / L_{0}}
$$

Rearranging terms to find the change in length gives

$$
\Delta L=\frac{F L_{0}}{Y A}=\frac{m g L_{0}}{Y \pi(d / 2)^{2}}=\frac{(10.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(1.0 \mathrm{~m})}{\left(2 \times 10^{11} \mathrm{~N} / \mathrm{m}^{2}\right) \pi(0.002 \mathrm{~mm} / 2)^{2}}=1.56 \times 10^{-4} \mathrm{~m}=0.156 \mathrm{~mm}
$$

EVALUATE We see that the stretch is very small for the wire. This agrees with experience: Steel is a difficult material to stretch.

## 6: Strain on an elevator cable

A steel elevator cable can support a maximum stress of $9.0 \times 10^{7} \mathrm{~Pa}$. If the maximum mass of the fully loaded elevator is 2100 kg and the maximum upward acceleration is $3.0 \mathrm{~m} / \mathrm{s}^{2}$, what should the diameter of the cable be? By how much does the cable stretch when the elevator is accelerating upward at $3.0 \mathrm{~m} / \mathrm{s}^{2}$ and 120 m of cable has been released? (Young's modulus for steel is $2 \times 10^{11} \mathrm{~Pa}$.)

## Solution

IDENTIFY We'll use Newton's second law to find the tension in the elevator cable and then use the maximum stress to find the cable diameter. To solve the second part, we'll use Young's modulus. The target variables are the cable diameter and the elongation of the cable.

SET UP Figure 11.9 shows the free-body diagram of the elevator. Gravity and tension act on the elevator.


Figure 11.9 Problem 6 free-body diagram.

EXECUTE We'll apply Newton's second law to find the maximum tension in the cable:

$$
\sum F_{y}=T-m g=m a_{y}
$$

The tension is then

$$
T=m\left(g+a_{x}\right)=(2100 \mathrm{~kg})\left(\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)+\left(3.0 \mathrm{~m} / \mathrm{s}^{2}\right)\right)=26,900 \mathrm{~N} .
$$

Stress is the force per unit area, or

$$
S=\frac{F}{A}=\frac{T}{A}
$$

The area can be written in terms of the diameter as

$$
A=\frac{\pi d^{2}}{4}
$$

The diameter is then

$$
d=2 \sqrt{\frac{T}{S \pi}}=2 \sqrt{\frac{(26,900 \mathrm{~N})}{\left(9.0 \times 10^{7}\right) \pi}}=0.020 \mathrm{~m}=2.0 \mathrm{~cm}
$$

where we have replaced the stress with the maximum stress. Young's modulus then leads to the amount of cable stretch:

$$
Y=\frac{l_{0} F}{A \Delta l}=\frac{4 l_{0} T}{\pi d^{2} \Delta l} .
$$

Rearranging terms to find the cable stretch gives

$$
\Delta l=\frac{4 l_{0} T}{\pi d^{2} Y}=\frac{4(120 \mathrm{~m})(26,900 \mathrm{~N})}{\pi(0.020 \mathrm{~m})^{2}\left(2.0 \times 10^{11} \mathrm{~Pa}\right)}=0.051 \mathrm{~m}=5.1 \mathrm{~cm}
$$

The cable must have a 2.0 cm diameter and stretch 5.1 cm when 120 m of cable has been released.
EVALUATE We see that the $2-\mathrm{cm}$-thick cable stretches over 5 cm . This may appear to be a significant elongation, but it represents only $0.04 \%$ of the cable length.

## 7: Compressibility of oil

Find the compressibility of a $0.1 \mathrm{~m}^{3}$ sample of oil whose volume decreases $2.04 \times 10^{-4} \mathrm{~m}^{3}$ when subjected to an increase in pressure of $1.02 \times 10^{7} \mathrm{~Pa}$.

## Solution

IDENTIFY AND SET UP We will use the bulk modulus to find the compressibility. With the given information, the compressibility follows directly from the definition of the bulk modulus.

EXECUTE The bulk modulus is given by

$$
B=-\frac{\Delta p}{\Delta V / V_{0}}
$$

Solving, we have

$$
B=-\frac{\Delta p V_{0}}{\Delta V}=-\frac{\left(1.02 \times 10^{7} \mathrm{~Pa}\right)\left(0.1 \mathrm{~m}^{3}\right)}{\left(-2.04 \times 10^{-4} \mathrm{~m}^{3}\right)}=5.0 \times 10^{9} \mathrm{~Pa}=4.9 \times 10^{4} \mathrm{~atm}
$$

The compressibility is the inverse of the bulk modulus, or

$$
k=\frac{1}{B}=\frac{1}{4.9 \times 10^{4} \mathrm{~atm}}=2.0 \times 10^{-5} / \mathrm{atm}
$$

EVALUATE We see that the oil does not compress when subjected to pressure. It requires an increase of 500 atmospheres of pressure to change the volume by $1 \%$.

## Try It Yourself!

## 1: Tension in support cable

Find the tension in the supporting cable and the force acting on the strut due to the wall for the weightless horizontal strut shown in Figure 11.10.


Figure 11.10 Try it yourself 1.

## Solution Checkpoints

IDENTIFY AND SET UP The strut is in equilibrium, so apply equilibrium conditions to solve the problem. Start with a free-body diagram, and include the force due to the weight hanging off the end of the strut, tension, and the force of the wall. Set the net forces and net torques equal to zero.

EXECUTE Newton's first law applied in the $x$ and $y$ directions gives

$$
\begin{aligned}
& \sum F_{x}=0=F_{\text {wall } x}-T \cos \theta, \\
& \sum F_{y}=0=F_{\text {wall } y}+T \sin \theta-w .
\end{aligned}
$$

The net torque about the right end of the strut is zero:

$$
\sum \tau=0=-F_{\text {wall }}(4.0 \mathrm{~m})
$$

Determine the sine and cosine of the angle and solve.
The tension is 500.0 N . The force due to the wall is 400 N , directed perpendicular to the wall.
EVALUATE Choosing a good pivot point simplifies calculations. By calculating the torques about the right end of the strut, we immediately learned that there is no $y$ component of force due to the wall. Do you get the same result if you set the pivot on the left end of the strut?

## 2: Force acting on a ladder

A ladder of mass 25.0 kg rests against a frictionless wall as shown in Figure 11.11. Find all the forces acting on the ladder.


Figure 11.11 Try it yourself 2.

## Solution Checkpoints

IDENTIFY AND SET UP The ladder is in equilibrium, so apply equilibrium conditions to solve the problem. Start with a free-body diagram, and include the forces due to the wall, the ground, and gravity. Set the net forces and net torques equal to zero.

EXECUTE Newton's first law applied in the $x$ and $y$ directions gives

$$
\begin{aligned}
& \sum F_{x}=0=F_{\text {ground } x}-n, \\
& \sum F_{y}=0=F_{\text {ground } y}-w .
\end{aligned}
$$

The net torque about the bottom end of the ladder is zero:

$$
\sum \tau=0=-\frac{L}{2} w \cos \theta+L n \sin \theta
$$

The normal force due to the wall is 71 N . The force due to the ground is 71 N in the positive $x$ direction and 245 N upward.

EVALUATE Does the wall exert a force in the vertical direction? How can you check your results?

## 3: Friction force acting on a ladder

A ladder of mass 25.0 kg rests against a frictionless wall as shown in Figure 11.11. What is the minimum coefficient of friction between the ladder and the ground that allows the ladder to stand without slipping?

## Solution Checkpoints

IDENTIFY AND SET UP The $x$ component of the force due to the ground is the friction force in the previous problem. Use the definition of friction to solve for the coefficient.

EXECUTE The static friction force is

$$
F_{\text {ground } x}=f \leq \mu_{s} n
$$

The coefficient of static friction is 0.29 .

EVALUATE: How did you find the normal force?

## 4: Stress in wire

A copper wire of cross-sectional area $0.050 \mathrm{~cm}^{2}$ and length 5.0 m is attached end to end to a steel wire of length 3.0 m and cross-sectional area $0.020 \mathrm{~cm}^{2}$. The wires are stretched under a tension of 200.0 N . Find the stress in each wire and the total change in length for the combination.

## Solution Checkpoints

IDENTIFY Tensile stress is the force per area. Young's modulus will be used to find the change in length. You can find Young's modulus for steel and copper from Table 11.1 of the text. Then add the changes in lengths to find the total change in length.

EXECUTE The stresses are

$$
\begin{aligned}
& \text { Stress }_{\text {Copper }}=\frac{F}{A}=4.0 \times 10^{7} \mathrm{~Pa} \\
& \text { Stress }_{\text {Steel }}=\frac{F}{A}=1.0 \times 10^{8} \mathrm{~Pa}
\end{aligned}
$$

The change in length of the copper is

$$
\Delta L=\frac{F L_{0}}{Y A}=1.5 \times 10^{-3} \mathrm{~m}
$$

For steel, the change in length is $1.8 \times 10^{-3} \mathrm{~m}$. The total length is $3.3 \times 10^{-3} \mathrm{~m}$, or 3.3 mm .
EVALUATE We see that the stretch is very small for combined length of wire.

## 12 Gravitation

## Summary

In this chapter, we will delve into the gravitational interaction by learning that the gravity we experience on earth also applies to planets and celestial objects and is responsible for their motion. We will see how to apply Newton's law of gravitation and gain a better understanding of the concept of weight. We'll use this knowledge to explain the orbits of satellites and planets. We will also examine an extreme case of gravity: black holes.

## Objectives

After studying this chapter, you will understand

- How to apply Newton's law of gravitation to pairs of masses.
- The general definition of weight.
- How to use the generalized expression for gravitational potential energy.
- How satellites orbit astronomical bodies.
- How to predict the motion of satellites.
- Kepler's three laws of planetary motion.
- The definition of a black hole and the properties of black holes.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Newton's Law of Gravitation | Newton's law of gravitation states that the magnitude of the force between <br> two bodies with masses $m_{1}$ and $m_{2}$, separated by a distance $r$, is given by |
| $\qquad F_{g}=G \frac{m_{1} m_{2},}{r^{2}}$, |  |
| where $G$ denotes the gravitational constant and is equal to $6.67 \times$ |  |
| $10^{-11} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{kg}^{2}$. The gravitational force is always attractive and is directed |  |
| along the line that separates the objects. |  |, | The weight of an object is the total gravitational force exerted on the object |
| :--- |
| by all other objects in the universe. Near the surface of the earth, an object's |
| weight is very nearly equal to the gravitational force of the earth on the |
| object alone. |, rated by a distance $r$, is given by

$$
U=-\frac{G m_{E} m}{r} .
$$

The potential energy is never positive and is zero only when the two objects are infinitely far apart.

## Orbits of Satellites

For a satellite moving in a circular orbit, the gravitational attraction between the satellite and the astronomical body provides the centripetal acceleration. The velocity $v$ and period $T$ of a satellite orbiting at a radius $r$ are given, respectively, by

$$
\begin{aligned}
& v=\sqrt{\frac{G M}{r}}, \\
& T=\frac{2 \pi r^{3 / 2}}{\sqrt{G M}},
\end{aligned}
$$

where $M$ is the mass of the astronomical body.
Kepler's Laws

Kepler's three laws describe the motion of a planet or satellite around the sun or another planet. They describe the elliptical motion and the area swept out per unit time in the orbit, as well as relate the period of the planet or satellite to the major axis of its orbit.

## Black Holes

A black hole is a nonrotating spherical mass distribution with total mass $M$ contained within a radius $R_{S}$, the Schwarzschild radius, given by

$$
R_{S}=\frac{2 G M}{c^{2}} .
$$

Gravity prevents matter and light from escaping from within a sphere with radius $R_{S}$.

## Conceptual Questions

## 1 : Is the earth falling?

There is a net gravitational force between the earth and the sun, so why doesn't the earth fall into the sun?

## Solution

IDENTIFY, SET UP, AND EXECUTE The earth is constantly falling toward the sun, but the earth doesn't get closer to the sun, since the sun's surface curves away beneath the earth. Recall projectile motion from Chapter 3. If we launch an object parallel to the ground, it follows a parabolic path. If we give the object a larger initial velocity, then the object moves farther away from the launch site as it falls. The earth is round, so as the object moves farther away, the object will have a larger distance to fall. At a sufficiently high launch velocity, the object will make a complete revolution and not land on the ground. This is the same physical situation as the earth revolving around the sun. If the earth had a lower velocity, it would fall into the sun.

EVALUATE This result may seem a bit odd, but is indeed accurate. The result also shows how our understanding of one physical phenomenon helps us understand other phenomena: Our experience with projectile motion helped us interpret the motion of the earth around the sun.

## 2: Does the earth maintain a constant speed?

In the previous question, we saw that the earth is constantly falling. Does it maintain a constant speed?

## Solution

IDENTIFY, SET UP, AND EXECUTE There is a net gravitational force acting on the earth due to the sun. The direction of the net force is toward the sun; however, the earth's velocity is perpendicular to the direction of force. The gravitational force can change only the direction of the earth's velocity around the sun and not the magnitude of the velocity. Thus, the earth maintains a constant speed as it orbits the sun. The earth does not maintain a constant velocity, since its direction is always changing.

EVALUATE We've come to know that a net force causes acceleration-a change in velocity. In the previous chapters, of ten the magnitude of an object's velocity changed when the object was acted upon by a net force. This chapter examines additional consequences of the influences of forces.

## 3: Two satellites and the earth

The moon and the international space station are located on opposite sides of the earth. How does the presence of the earth influence the gravitational force between the moon and the space station?

## Solution

IDENTIFY, SET UP, AND EXECUTE The gravitational force between two bodies depends only on the mass of the bodies and their separation, according to Newton's law of gravitation. The force between the moon and the space station is not affected by the presence of the earth. There are forces between the earth and the two orbiting satellites, but those forces do not affect the force between the satellites.

EVALUATE Forces are between two bodies. The net force on a single body may include forces due to many bodies, but each force acts between two bodies.

## Problems

## 1: Gravitational force due to three masses

Three masses are arranged as shown in Figure 12.1. Find the net force acting on the top mass ( $A$ ). Each mass is 5.00 kg .


Figure 12.1 Problem 1.

## Solution

IDENTIFY The net force on $A$ is found by adding the force on $A$ due to $B$ and the force on $A$ due to $C$. The target variable is the force on $A$.

SETUP A free-body diagram illustrating the two forces on $A$ is shown in Figure 12.2. We'll need to add the two forces by using components. Newton's law of gravitation gives the magnitude of the forces. We'll use the coordinate axes provided in the figure.


Figure 12.2 Problem 1 free-body diagram.

EXECUTE To apply Newton's law of gravity, we need the distances between the masses. The distance between $A$ and $B$ is 10.0 cm . By summing the squares of the sides of the triangle and taking the square root, we find that the distance between $A$ and $C$ is 14.1 cm . The force of $B$ on $A$ is

$$
F_{B \text { on } A}=\frac{G m_{B} m_{A}}{r_{B A}^{2}}=\frac{\left(6.67 \times 10^{-11} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{kg}^{2}\right)(5.00 \mathrm{~kg})(5.00 \mathrm{~kg})}{(0.100 \mathrm{~cm})^{2}}=1.67 \times 10^{-7} \mathrm{~N}
$$

The force of $C$ on $A$ is

$$
F_{C \text { on } A}=\frac{G m_{C} m_{A}}{r_{C A}^{2}}=\frac{\left(6.67 \times 10^{-11} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{kg}^{2}\right)(5.00 \mathrm{~kg})(5.00 \mathrm{~kg})}{(0.141 \mathrm{~cm})^{2}}=8.39 \times 10^{-8} \mathrm{~N}
$$

With the magnitudes of the forces determined, we simply add the two vectors together, using components. The force of $C$ on $A$ has the only $x$ component:

$$
\sum F_{x}=F_{C \text { on } A} \sin 45^{\circ}=5.93 \times 10^{-8} \mathrm{~N}
$$

The $45^{\circ}$ angle results from the masses arranged as an isosceles triangle. Both forces have $y$ components:

$$
\sum F_{y}=-F_{B \text { on } A}-F_{C \text { on } A} \sin 45^{\circ}=-2.26 \times 10^{-7} \mathrm{~N} .
$$

The negative result indicates that the $y$ component points downward. We find the magnitude of the net force by combining the components:

$$
F=\sqrt{F_{x}^{2}+F_{y}^{2}}=2.34 \times 10^{-7} \mathrm{~N} .
$$

The direction of the net force is found by using the tangent. We want to specify the angle $\phi$ with respect to the $x$-axis:

$$
\phi=\tan ^{-1} \frac{F_{y}}{F_{x}}=\tan ^{-1} \frac{\left(-2.26 \times 10^{-7} \mathrm{~N}\right)}{\left(5.93 \times 10^{-8} \mathrm{~N}\right)}=-75.3^{\circ} .
$$

The net force on $A$ has magnitude $2.34 \times 10^{-7} \mathrm{~N}$ and points $75.3^{\circ}$ below the positive $x$-axis.
EVALUATE We see that the gravitational force between the masses is very small. To have an appreciable gravitational force, we need at least one large mass, such as the earth. Also, we can see that Newton's third law is valid: Reversing indices in the first two equations would result in a force of the same magnitude, but opposite in direction.

## 2: Orbit of a weather satellite

Imagine you are designing a new weather satellite. The goal is to have the satellite orbit the earth in a circular orbit every 6 hours. At what distance above the earth's surface should the satellite be placed to obtain the correct period?

## Solution

IDENTIFY The force acting on the satellite is the force of gravitation between the satellite and the earth. The satellite follows a circular orbit and so has a radial acceleration toward the center of the earth. The target variable is the height of the satellite's orbit.

SET UP Newton's law of gravitation gives the force on the satellite due to the earth. The acceleration of the satellite is centripetal. Combining the two equations will lead to the velocity of the satellite and the period of rotation. We solve for the distance by setting the period to 6 hours.

EXECUTE Newton's law of gravitation gives the force on the satellite, namely,

$$
F_{g}=\frac{G m m_{E}}{r^{2}}
$$

where $m$ is the mass of the satellite, $m_{E}$ is the mass of the earth, and $r$ is the distance from the center of the earth to the satellite. Newton's second law gives the net force on the satellite (the acceleration is $\left.v^{2} / r\right)$ :

$$
\sum F=m a_{\mathrm{rad}}, \sum F=F_{g}=\frac{G m m_{E}}{r^{2}}=m \frac{v^{2}}{r}
$$

Solving for $v$, we find that

$$
v=\sqrt{\frac{G m_{E}}{r}}
$$

We can also write the velocity in terms of the distance the satellite travels ( $2 \pi r$ ) in one period $(T)$ :

$$
v=\frac{2 \pi r}{T}
$$

To find the radius, we equate the last two equations and solve for $r$, obtaining

$$
r=\sqrt[3]{\frac{G m_{E} T^{2}}{4 \pi^{2}}}=\sqrt[3]{\frac{\left(6.67 \times 10^{-11} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{kg}^{2}\right)\left(5.98 \times 10^{24} \mathrm{~kg}\right)\left(2.16 \times 10^{4} \mathrm{~s}\right)^{2}}{4 \pi^{2}}}=1.68 \times 10^{7} \mathrm{~m}
$$

where we replaced the 6 hour period with the equivalent $21,600 \mathrm{~s}$. The satellite should be placed in an orbit of radius $16,800 \mathrm{~km}$. Subtracting the radius of the earth ( 6380 km ) from the radius of the satellite's orbit, we find that the satellite should be placed $10,400 \mathrm{~km}$ above the earth's surface to achieve a 6 hour orbital period.

EVALUATE We could have avoided our derivation and used the textbook's equation 12.12 to arrive at the solution directly. However, this review helps remind us how to find the solution without searching the book.

Practice Problem: What does the free-body diagram of the satellite look like? Answer: A single vector.

## 3: Velocity of rocket

What velocity must a rocket have at the surface of the earth if it is to rise to a height equal to the earth's radius before it begins to descend? Ignore air resistance.

## Solution

IDENTIFY We'll use energy conservation to solve the problem. The target variable is the initial velocity of the rocket.

SET UP The rocket is given an initial kinetic energy for lift-off and has initial gravitational potential energy. At the top of the flight, its kinetic energy drops to zero, leaving only gravitational potential energy. We'll set these equal to each other to solve for the initial velocity.

EXECUTE Energy conservation gives

$$
K_{1}+U_{1}=K_{2}+U_{2} .
$$

Replacing both sides with the expressions for the two forms of energy gives

$$
\frac{1}{2} m v^{2}-\frac{G m m_{E}}{r_{E}}=-\frac{G m m_{E}}{2 r_{E}}
$$

where $r_{E}$ is the radius of the earth and $K_{2}=0$. Solving for the initial velocity, we obtain

$$
v=\sqrt{\frac{G m_{E}}{r_{E}}}=\sqrt{\frac{r_{\text {earth }}^{2} g}{r_{E}}}=\sqrt{g r_{E}}=\sqrt{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)\left(6.36 \times 10^{6} \mathrm{~m}\right)}=7.9 \times 10^{3} \mathrm{~m} / \mathrm{s}
$$

The initial velocity of the rocket must be $7.9 \times 10^{3} \mathrm{~m} / \mathrm{s}$.
EVALUATE By examining the solution to the problem, we see why rockets must have large initial velocities to be propelled into space. The actual value is higher due to air resistance.

## 4: What if the sun were a black hole?

What would the sun's radius need to be in order for the surface escape velocity to be $c$ ?

## Solution

IDENTIFY AND SET UP A radius corresponding to an escape velocity of $c$ is the Schwarzschild radius. We can use the expression for the Schwarzschild radius to solve the problem.

EXECUTE The Schwarzschild radius is given by

$$
R_{\mathrm{S}}=\frac{2 G M}{c^{2}}
$$

Substituting gives

$$
R_{\mathrm{S}}=\frac{2 G M}{c^{2}}=\frac{2\left(6.67 \times 10^{-11} \mathrm{Nm}^{2} / \mathrm{kg}^{2}\right)\left(1.99 \times 10^{30} \mathrm{~kg}\right)}{\left(3.0 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)^{2}}=2960 \mathrm{~m}
$$

The radius would be $2,960 \mathrm{~m}$.
EVALUATE The sun would have to be compressed by a factor of over 200,000 to become a black hole.
Practice Problem: What is the Schwarzschild radius for the earth? Answer: 8.8 mm .

## Try It Yourself!

## 1: Sun's gravity on earth

The sun is a distance of $1.48 \times 10^{11} \mathrm{~m}$ from earth and has a mass of $1.99 \times 10^{30} \mathrm{~kg}$. Find the ratio of the sun's gravitational force to the earth's gravitational force on an object on the earth's surface.

## Solution Checkpoints

IDENTIFY AND SET UP Newton's law of gravitation is used to find the force of gravity due to the earth and the force of gravity due to the sun. Taking their ratio solves the problem.

Substituting the values given results in a ratio of $6.03 \times 10^{-4} \mathrm{~m} / \mathrm{s}$.
EVALUATE Do you need to include the force of gravity due to the sun in physics problems on earth?

## 2: Three masses positioned in a triangle

Three 1000.0 kg masses are at the vertices of an equilateral triangle with sides of length 1.0 m . Find the force due to any two masses on the third.

## Solution Checkpoints

IDENTIFY AND SET UP Newton's law of gravitation is used to find the force of gravitation due to the other masses. The two forces must be added as vectors. A free-body diagram should be used.

EXECUTE The force between two masses is

$$
F_{g}=\frac{G m_{1} m_{2}}{r^{2}}=6.67 \times 10^{-5} \mathrm{~N}
$$

The net force is $1.16 \times 10^{-4} \mathrm{~N}$, directed toward the line separating the other two masses.

EVALUATE Did you use symmetry arguments to determine that the force has no component parallel to the line separating the two masses?

## 3: Orbit of a communications satellite

Communications satellites revolve in orbits over the earth's equator, adjusted so that their period of rotation is the same as the period of rotation of the earth about its axis. This speed and a period of rotation together cause the satellite to remain in a fixed position in the sky. Find the height of these satellites' orbit above the earth.

## Solution Checkpoints

IDENTIFY AND SET UP You can find the height by setting the period equal to 24 hours. The solution can be found by solving Newton's second law or by using the equation in the book.

EXECUTE The period is given by

$$
T=\frac{2 \pi r^{3 / 2}}{\sqrt{g R^{2}}}
$$

The terms in this equation can be rearranged to solve for $r$. The height above earth is $r$ minus the radius of earth. The satellite must be placed $36,000 \mathrm{~km}$ above the surface of the earth to remain in geosynchronous orbit.

EVALUATE Such heights require additional power for radio signals to reach the satellites and for the signal to be redirected back to earth.

## 4: Escape from the sun

What is the escape velocity of a particle on the surface of the sun?

## Solution Checkpoints

IDENTIFY AND SET UP The escape velocity can be found from energy conservation or the equation given in the text. Parameters for the sun can be found in the chapter or appendix.

EXECUTE The escape velocity is given by

$$
v=\sqrt{\frac{2 G M}{r}}
$$

After substituting, we find the escape velocity to be $6.48 \times 10^{5} \mathrm{~m} / \mathrm{s}$.
EVALUATE Is the resulting velocity greater or less than the escape velocity of a particle on the surface of the earth? Why?

## Periodic Motion

## Summary

We will examine periodic motion, or oscillation, in this chapter. Many systems exhibit periodic motion, such as a swinging pendulum, a ball on a spring, or the membrane of a drum. We will describe the motion of oscillating bodies, characterized by amplitude, period, frequency, and angular frequency. We'll use force equations and energy concepts to analyze their motions. We will look at several models of periodic motion that can be used to represent the motion of many oscillators. Periodic motion plays a vital role in many areas of physics, and this chapter will lay the foundation for further studies.

## Objectives

After studying this chapter, you will understand

- Periodic motion and the terminology used to describe oscillations.
- How to identify and analyze simple harmonic motion.
- Energy and motion as a function of time for a particle in simple harmonic motion.
- The simple pendulum, the physical pendulum, damped and forced oscillations, and resonance.

Concepts and Equations

| Term | Description |
| :--- | :--- |
| Periodic Motion | Periodic motion is motion that repeats in a definite cycle. Periodic motion <br> occurs when an object is displaced from its equilibrium position and a restor- <br> ing force exists that tends to return the object to equilibrium. The amplitude is <br> the maximum magnitude of displacement from equilibrium. A cycle is one <br> complete round-trip. The period is the time taken to complete one cycle. Fre- <br> quency $(f)$ is the number of cycles per unit time. Angular frequency $(\omega)$ is <br> $2 \pi$ times the frequency. Period, frequency, and angular frequency are related: <br> $\qquad T=\frac{1}{f}, \quad f=\frac{1}{T}, \quad \omega=2 \pi f=\frac{2 \pi}{T}$. |
| Simple Harmonic Motion | Simple harmonic motion $(S H M)$ is periodic motion in which the restoring <br> force is directly proportional to the object's displacement. Often, SHM occurs <br> when the displacement is small. The equation of motion is <br> $x=A \cos (\omega t+\phi)$. |

A system attached to a spring with spring constant $k$ and having mass $m$ will oscillate at a frequency of

$$
\omega=2 \pi f \frac{\omega}{2 \pi}=\sqrt{\frac{k}{m}}
$$

and a period of

$$
T=\frac{1}{f}=2 \pi \sqrt{\frac{m}{k}} .
$$

The total mechanical energy remains constant in SHM and can be expressed in terms of its amplitude:

$$
E=\frac{1}{2} m v^{2}+\frac{1}{2} k x^{2}=\frac{1}{2} k A^{2} .
$$

For angular simple harmonic motion, the frequency is related to the moment of inertia and the torsion constant by

$$
\omega=\sqrt{\frac{\kappa}{I}}
$$

## Simple Pendulum

Physical Pendulum

A simple pendulum is a model of a point mass suspended by a massless string in a gravitational field. For small displacements, a pendulum of length $L$ has frequency

$$
\omega=2 \pi f=\sqrt{\frac{g}{L}} .
$$

A physical pendulum is any body suspended from an axis of rotation. The angular frequency for small-amplitude oscillations is given by

$$
\omega=\sqrt{\frac{m g d}{I}} .
$$

## Damped Oscillations

A simple harmonic oscillator impelled by a force that is proportional to velocity exhibits damped oscillations. The angular frequency becomes

$$
\omega^{\prime}=\sqrt{\frac{k}{m}-\frac{b^{2}}{4 m^{2}}} .
$$

|  | Systems are called critically damped, overdamped, and underdamped accord- <br> ing to how they return to equilibrium. |
| :--- | :--- |
| Forced Oscillations | Periodic motion in a system with a sinusoidally varying driving force is <br> called forced oscillation or driven oscillation. Resonance occurs when the <br> driving angular frequency is near the natural-oscillation angular frequency, <br> increasing the amplitude of the motion. The amplitude is given by |
| $\qquad A=\frac{F_{\max }}{\sqrt{\left(k-m \omega_{d}^{2}\right)^{2}+b \omega_{d}^{2}}}$. |  |

## Conceptual Questions

## 1: Glider in simple harmonic motion

A glider attached to a spring and set on a horizontal air track is allowed to oscillate with a 5.0 cm amplitude. How far does the glider travel in one period?

## Solution

IDENTIFY, SET UP, AND EXECUTE We answer the question by considering how the glider moves during one period. The period is the time an object takes to move from any position through one complete periodic cycle and return to the starting position. Imagine that the glider starts from an equilibrium position, moves to the right, and momentarily stops at a displacement equal to the amplitude. It has traveled a distance of one amplitude, or 5.0 cm . The glider then returns to its equilibrium position, traveling a second distance equal to the amplitude, or a total of 10.0 cm . The glider continues moving to the left until it reaches its maximum displacement on the left side, thus traveling a third distance equal to the amplitude ( 15.0 cm total). The glider then moves to the right and returns to the starting position, traveling a fourth distance equal to the amplitude ( 20.0 cm total).

In one period, the glider travels a distance equal to four amplitudes, or 20.0 cm .
EVALUATE You must distinguish between amplitude and total distance traveled. Comprehending this difference helps build an understanding of simple harmonic motion.

## 2: Gravity on the moon

You are asked to estimate the moon's gravitational acceleration by watching a video of the early lunar explorations. How could you estimate the acceleration due to gravity on the moon?

## Solution

IDENTIFY, SET UP, AND EXECUTE We've seen that, for small oscillations, the period of a simple pendulum is related to the gravitational constant and the length of the pendulum. If you can find an object that can be approximated by a simple pendulum, then you can determine the gravitational acceleration from the object's motion. One approach would be to look for a dangling object during a moonwalk. You can estimate the length of the pendulum by comparing it with the size of the astronaut on the walk and measure the time with a stopwatch or by counting video frames.

EVALUATE The moon's gravitational acceleration was estimated by physics students around the world watching the early moonwalks. The technique can also be used to estimate the sizes of objects in videos by taking the known gravitational acceleration value and combining it with the period to find the length of the pendulum.

## Problems

## 1: Mass on a spring

A spring stretches 4.7 cm from its equilibrium position when a 1.2 kg mass is hung from it. If the mass is now stretched 6.5 cm from the equilibrium position and released, find (a) the period of the motion, (b) the maximum velocity of the mass, and (c) the maximum acceleration of the mass.

## Solution

IDENTIFY Since the net force acting on the block is proportional to the displacement of the block, the motion is simple harmonic motion. The target variables are the period, maximum velocity, and maximum acceleration.

SET UP We'll use the equations of simple harmonic motion to find the solutions to the problem. We first find the spring constant, using the preliminary information.

EXECUTE We find the spring constant from Hooke's law. When the mass is initially attached to the spring, it hangs in equilibrium, so the spring force is equal to the product of the mass and the acceleration due to gravity:

$$
F_{\mathrm{s}}=k x=m g
$$

The spring constant is

$$
k=\frac{m g}{x}=\frac{(1.2 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{0.047 \mathrm{~m}}=250 \mathrm{~N} / \mathrm{m}
$$

With the spring constant, we can directly find the period:

$$
T=2 \pi \sqrt{\frac{m}{k}}=2 \pi \sqrt{\frac{(1.2 \mathrm{~kg})}{(250 \mathrm{~N} / \mathrm{m})}}=0.44 \mathrm{~s}
$$

The maximum velocity is

$$
v_{\max }=\sqrt{\frac{k}{m}} A=\sqrt{\frac{(250 \mathrm{~N} / \mathrm{m})}{(1.2 \mathrm{~kg})}}(0.065 \mathrm{~m})=0.94 \mathrm{~m} / \mathrm{s} .
$$

The maximum (positive) acceleration occurs when the mass is at its most negative position, so

$$
a_{\max }=-\frac{k}{m} x=-\frac{(250 \mathrm{~N} / \mathrm{m})}{(1.2 \mathrm{~kg})}(-0.065 \mathrm{~m})=13.5 \mathrm{~m} / \mathrm{s}^{2}
$$

The mass oscillates with a period of 0.44 s and has a maximum velocity of $0.94 \mathrm{~m} / \mathrm{s}$ and a maximum acceleration of $13.5 \mathrm{~m} / \mathrm{s}^{2}$, upward.

EVALUATE Simple harmonic motion is the most complicated motion we have studied to date. However, our previous experiences led to straightforward relationships from which we can easily extract useful information.

We could also have found the motion as a function of time, taken derivatives to find the velocity and acceleration as a function of time, and then determined the maxima-the amplitudes of the velocity and acceleration functions.

## 2: Object in SHM

A 200.0 g mass vibrates in SHM with a total energy of 25.0 J and a frequency of 5.0 Hz . Find the time it takes to move from 25.0 cm below to 25.0 cm above the equilibrium position.

## Solution

IDENTIFY We will use the simple harmonic motion relations to find the solution. The target variable is the time needed to move the specified distance.

SET UP We'll use the equation of simple harmonic motion to find the solutions to the problem. We will find the times the mass is -0.25 cm and +25.0 cm from the equilibrium position. We will need to find the spring constant and amplitude from the information given.

EXECUTE The spring constant can be found from the frequency equation:

$$
f=\frac{1}{2 \pi} \sqrt{\frac{k}{m}}
$$

Solving for $k$ gives

$$
k=m(2 \pi f)^{2}=(0.200 \mathrm{~kg})(2 \pi(5.0 \mathrm{~Hz}))^{2}=197.4 \mathrm{~N} / \mathrm{m}
$$

The total energy is given by

$$
E=\frac{1}{2} m v^{2}+\frac{1}{2} k x^{2}
$$

The amplitude is the maximum displacement. At the maximum displacement, the velocity is zero. Solving for the amplitude results in

$$
\begin{gathered}
E=\frac{1}{2} k A^{2} \\
A=\sqrt{2 E / k}=\sqrt{2(24.7 \mathrm{~J}) /(197 \mathrm{~N} / \mathrm{m})}=0.500 \mathrm{~m}
\end{gathered}
$$

The position as a function of time is given by

$$
x=A \sin \omega t=A \sin 2 \pi f t
$$

We need to solve for the time when $x=-0.25 \mathrm{~m}$ and $x=0.25 \mathrm{~m}$. Solving gives

$$
\begin{aligned}
& t_{-0.25 \mathrm{~m}}=\frac{1}{2 \pi f} \sin ^{-1}\left(\frac{x}{A}\right)=\frac{1}{2 \pi(5.0 \mathrm{~Hz})} \sin ^{-1}\left(\frac{-0.25 \mathrm{~m}}{0.50 \mathrm{~m}}\right)=-0.0167 \mathrm{~s} \\
& t_{+0.25 \mathrm{~m}}=\frac{1}{2 \pi f} \sin ^{-1}\left(\frac{x}{A}\right)=\frac{1}{2 \pi(5.0 \mathrm{~Hz})} \sin ^{-1}\left(\frac{0.25 \mathrm{~m}}{0.50 \mathrm{~m}}\right)=0.0167 \mathrm{~s}
\end{aligned}
$$

The time the mass takes to move from 25.0 cm below to 25.0 cm above the equilibrium position is 0.0333 s.
eVALUATE We see that the time the mass takes to move from half the amplitude below to half the amplitude above the equilibrium position is about $15 \%$ of the period. Does this make sense? Yes, it makes sense, since the velocity near the equilibrium point is maximal.
Practice Problem: How long would it take the mass to move from the equilibrium point to the maximum amplitude? Answer: 0.05 s , or one-fourth the period.

CAUTION Use radians! When you work with trigonometric functions, the arguments are in radians. You must either set your calculator to radian mode or convert degrees to radians after taking the inverse of the trigonometric function.

## 3: Period of a simple pendulum

A simple pendulum reaches a maximum angle of $7.2^{\circ}$ after swinging through the bottom of its path with a maximum speed of $0.35 \mathrm{~m} / \mathrm{s}$. What is the period of the pendulum's oscillation?

## Solution

IDENTIFY The period of a simple pendulum depends on its length and the gravitational constant. We can find the target variable-the length-from the velocity and maximum angle.

SET UP The maximum angle, arc length, and length are related to each other. The amplitude and the maximum speed will be used to find the length of the pendulum, and the period will be derived from the length.

EXECUTE The amplitude of a simple pendulum is the maximum arc length, which is related to the maximum angle by the length:

$$
S_{\max }=L \theta_{\max }
$$

The maximum velocity is

$$
v_{\max }=2 \pi f A=2 \pi f L \theta_{\max }
$$

For a simple pendulum, the frequency is found from the length:

$$
f=\frac{1}{2 \pi} \sqrt{\frac{g}{L}}
$$

Combining these equations, we obtain the length:

$$
\begin{aligned}
& v_{\max }=2 \pi\left(\frac{1}{2 \pi} \sqrt{\frac{g}{L}}\right) L \theta_{\max }=\sqrt{g L} \theta_{\max } \\
& L=\frac{v_{\max }^{2}}{\mathrm{~g} \theta_{\max }^{2}}=\frac{(0.35 \mathrm{~m} / \mathrm{s})^{2}}{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(0.126)^{2}}=0.79 \mathrm{~m}
\end{aligned}
$$

Note that we replaced the maximum angle of $7.2^{\circ}$ with the equivalent 0.126 radian. The period is then

$$
T=2 \pi \sqrt{\frac{L}{g}}=2 \pi \sqrt{\frac{(0.79 \mathrm{~m})}{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}}=1.8 \mathrm{~s}
$$

The pendulum has a length of 79 cm and a period of 1.8 s .
EVALUATE The period of a simple pendulum depends only on the length of the pendulum and the gravitational constant. The maximum angle and velocity provided enough information to solve the problem.

## 4: Period of a physical pendulum

A thin, uniform rod is pivoted at a point one-quarter of its length from one end and is then pivoted at a point at its end. Find the ratio of the two periods.

## Solution

IDENTIFY The period of a physical pendulum depends on its moment of inertia, its mass, the distance to its center of mass, and the gravitational constant. The target variable is the ratio of the periods for the two pivot points.

SET UP We will calculate the moment of inertia for each of the two pivot points and then combine the two moments to form the ratio.

EXECUTE The moment of inertia of a rod about its end point is

$$
I_{\mathrm{end}}=\frac{1}{3} M L^{2} .
$$

When the rod is pivoted at a point one-quarter along its length, the moment of inertia is found by the parallel-axis theorem:

$$
I_{1 / 4}=I_{\mathrm{end}}+m x^{2}=\frac{1}{3} M L^{2}+M\left(\frac{L}{4}\right)^{2}=\frac{7}{48} M L^{2}
$$

The period of a physical pendulum is given by

$$
T=2 \pi \sqrt{\frac{I}{m g d}}
$$

The ratio of the two periods is then

$$
\frac{T_{1 / 4}}{T_{\mathrm{end}}}=\frac{2 \pi \sqrt{\frac{I_{1 / 4}}{m g d_{1 / 4}}}}{2 \pi \sqrt{\frac{I_{\mathrm{end}}}{m g d_{\mathrm{end}}}}}=\sqrt{\frac{I_{1 / 4} d_{\mathrm{end}}}{I_{\mathrm{end}} d_{1 / 4}}}=\sqrt{\frac{\frac{7}{48} M L^{2} \frac{L}{2}}{\frac{1}{3} M L^{2} \frac{L}{4}}}=\sqrt{\frac{(7)(3)(4)}{(2)(48)}}=0.94
$$

The ratio of the period when the rod is pivoted at a point one-quarter of its length from one end to the period when the rod is pivoted at a point at its end is 0.94 .

EVALUATE We see that the period does not change substantially when the pivot point moves between the two positions.

## 5: Oscillating blocks

Two blocks shown in Figure 13.1 oscillate on a frictionless surface with a frequency of 0.30 Hz . The top block has a mass of 2.0 kg and the bottom block has a mass of 4.5 kg . If the amplitude is increased to 25 cm , the top block begins to slide. What is the coefficient of static friction?


Figure 13.1 Problem 5.

## Solution

IDENTIFY When the top block just begins to slide, the force applied must be equal to the maximum static frictional force. The maximum applied force occurs at the maximum displacement (equal to the amplitude). We will determine the force at the turning point and use that to solve for the coefficient of friction, the target variable.

SET UP We'll need the spring constant, which we extract from the initial frequency. The motion is simple harmonic motion, as the only horizontal force acting on the blocks is the spring force, a restoring force that is directly proportional to the displacement.

EXECUTE The frequency in simple harmonic motion depends on the spring constant according to the formula

$$
f=\frac{1}{2 \pi} \sqrt{\frac{k}{m+M}},
$$

where we include the combined mass of the oscillating blocks. Solving for $k$ gives

$$
k=(2 \pi f)^{2}(m+M)=(2 \pi(0.30 \mathrm{~Hz}))^{2}((2.0 \mathrm{~kg})+(4.5 \mathrm{~kg}))=23.1 \mathrm{~N} / \mathrm{m}
$$

Recall that the maximum static frictional force is $f_{\mathrm{s}}=\mu_{\mathrm{s}} n$. For the top block, the normal force is $m g$. The maximum force applied by the spring is $k A$. Equating these forces gives

$$
f_{\mathrm{s}}=\mu_{\mathrm{s}} n=\mu_{\mathrm{s}} m g=k A
$$

Rearranging terms to find the coefficient of static friction yields

$$
\mu_{\mathrm{s}}=\frac{k A}{m g}=\frac{(23.1 \mathrm{~N} / \mathrm{m})(0.25 \mathrm{~m})}{(2.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}=0.29
$$

The coefficient of static friction between the blocks is 0.29 .
EVALUATE This problem brings together topics from several areas we have studied throughout the text, including the normal force, frictional forces, the spring force, and simple harmonic motion. Combining our knowledge helps us understand complex phenomena.

## 6: Damped oscillation

A body with mass 0.30 kg hangs by a spring with force constant $50.0 \mathrm{~N} / \mathrm{m}$. By what factor is the frequency of oscillation reduced if the oscillation is damped and reaches $1 / e$ of its original amplitude in 100 oscillations?

## Solution

IDENTIFY The amplitude in damped SHM diminishes by a factor of $e^{-b t / 2 m}$ as a function of time. We'll set this quantity to $1 / e$ to solve. The target variable is the fractional frequency shift, which is the change in frequency divided by the undamped frequency.

SET UP We will write the fractional frequency shift in terms of the damped and undamped frequencies. We will need to expand the frequency equation to simplify the shift and substitute the exponential decay information to solve.

EXECUTE The undamped frequency is given by

$$
\omega=\sqrt{\frac{k}{m}}
$$

The damped frequency is given by

$$
\omega^{\prime}=\sqrt{\frac{k}{m}-\frac{b^{2}}{4 m^{2}}}
$$

Combining these two equations to find the fractional frequency shift produces

$$
\begin{aligned}
\frac{\Delta \omega}{\omega} & =\frac{\omega^{\prime}-\omega}{\omega}=\frac{\sqrt{\frac{k}{m}-\frac{b^{2}}{4 m^{2}}}-\sqrt{\frac{k}{m}}}{\sqrt{\frac{k}{m}}} \\
& =\sqrt{\left[1-\left(\frac{m}{k}\right) \frac{b^{2}}{4 m^{2}}\right]-1 .}
\end{aligned}
$$

For small damping, the second term in the square root is small, so we can simplify by using the approximation

$$
\sqrt{1-x} \approx 1-\frac{1}{2} x
$$

This gives

$$
\frac{\Delta \omega}{\omega} \approx 1-\frac{1}{2}\left(\frac{m}{k}\right) \frac{b^{2}}{4 m^{2}}-1=\frac{1}{2}\left(\frac{m}{k}\right) \frac{b^{2}}{4 m^{2}} .
$$

We need to eliminate the damping term. We use the exponential decay information. The amplitude in damped oscillations changes as

$$
A(t)=A_{0} e^{-(b / 2 m) t}
$$

We know that after 100 oscillations the amplitude drops to $1 / e$. This gives

$$
e^{-(b / 2 m) t}=e^{-1}
$$

Taking the logarithm of each side yields an expression for $b / 2 m$ :

$$
\begin{gathered}
-(b / 2 m)[100 T]=-1 \\
\frac{b}{2 m}=\frac{1}{100} \frac{1}{T}=\frac{1}{100} \frac{1}{2 \pi} \sqrt{\frac{k}{m}}
\end{gathered}
$$

We now solve for the shift:

$$
\frac{\Delta \omega}{\omega}=\frac{1}{2}\left(\frac{m}{k}\right) \frac{b^{2}}{4 m^{2}}=\frac{1}{2}\left(\frac{m}{k}\right)\left(\frac{1}{100} \frac{1}{2 \pi} \sqrt{\frac{k}{m}}\right)^{2}=\frac{1}{2}\left(\frac{1}{100} \frac{1}{2 \pi}\right)^{2}=1.27 \times 10^{-6}
$$

The frequency shift is $1.27 \times 10^{-6}$.
EVALUATE We see that the frequency shift for this problem is very small.

## Try It Yourself!

## 1: SHM practice

A body of mass 0.5 kg is attached to a spring with spring constant $100.0 \mathrm{~N} / \mathrm{m}$ and is allowed to oscillate on a horizontal frictionless surface. It is given an initial velocity, at $x=0$, of $5.0 \mathrm{~m} / \mathrm{s}$. Find (a) the total energy of the body, (b) the amplitude of oscillation, (c) the velocity when the displacement is half of the amplitude, (d) the displacement when the velocity is half of its initial value, (e) the displacement when the kinetic and potential energies are equal, and (f) the frequency and period of the motion.

## Solution Checkpoints

IDENTIFY AND SET UP The body is in simple harmonic motion. Use the energy equations for SHM to solve for the many target variables. Remember that energy is conserved in SHM.

EXECUTE Energy conservation is used to solve (a) through (e) by substituting the appropriate knowns to find the unknowns. Energy conservation for the system states that

$$
E=\frac{1}{2} m v^{2}+\frac{1}{2} k x^{2}=\frac{1}{2} k A^{2}
$$

This equation can be used to find (a) the total energy of the body ( 6.25 J ), (b) the amplitude of oscillation $(0.35 \mathrm{~m})$, (c) the velocity when the displacement is half of the amplitude ( $\pm 4.34 \mathrm{~m} / \mathrm{s}$, ) (d) the displacement when the velocity is half of its initial value $( \pm 0.31 \mathrm{~m})$, and (e) the displacement when the kinetic and potential energies are equal ( 0.25 m ).

The period and frequency are found from the relationships

$$
f=\frac{1}{2 \pi} \sqrt{\frac{k}{m}}, \quad T=\frac{1}{f}
$$

The period is 0.44 s and the frequency is 2.25 Hz .
EVALUATE This problem illustrates how to apply simple harmonic energy relations to find displacements, velocities, and the amplitude, period, and frequency of the motion.

## 2: SHM practice

A body in SHM with angular frequency 0.5 s is initially 10.0 cm from its equilibrium position and is moving back toward equilibrium with a velocity of $5.0 \mathrm{~cm} / \mathrm{s}$, as shown in Figure 13.2. How long does it take for the body to return to its equilibrium position?
(b)


Figure 13.2 Try It Yourself 2.

## Solution Checkpoints

IDENTIFY AND SET UP The body is in simple harmonic motion. Find the motion equation in terms of amplitude, angular frequency, and phase angle. Then solve for the time needed to move back to equilibrium, the target variable.

EXECUTE The general forms of the position and velocity equations are

$$
\begin{aligned}
& x=A \cos \left(\omega t+\theta_{0}\right) \\
& v=-\omega A \sin \left(\omega t+\theta_{0}\right)
\end{aligned}
$$

To find the two constants, we use the initial conditions. At time zero, the position and velocity are given. This yields

$$
\begin{aligned}
A & =14.1 \mathrm{~cm} \\
\theta_{0} & =0.79 \mathrm{rad}
\end{aligned}
$$

The position is solved for time when $x=0$, giving a time of 1.56 s .
EVALUATE This problem shows how to use initial conditions to solve for the equations of motion.
CAUTION Watch $\boldsymbol{f}$ and $\boldsymbol{\omega}$ ! Be careful to distinguish the frequency $f$ from the angular frequency $\omega=2 \pi f$. The last two problems involved both quantities.

## 3: Simple pendulum

A clock pendulum with mass 5.0 kg is set to swing with a 2.0 s period. How long should the pendulum be made if you approximate it as a simple pendulum?

## Solution Checkpoints

IDENTIFY AND SET UP The period of a simple pendulum is given in terms of the length of the pendulum and the gravitational constant.

EXECUTE The period of a simple pendulum is given by

$$
T=2 \pi \sqrt{\frac{L}{g}}
$$

Rearranging terms and solving yields a length of 0.99 m .
EVALUATE Mass does not affect the results for a simple pendulum.

## 4: Physical pendulum

A body of mass 2.0 kg is suspended at a point 3.0 cm from its center of mass and observed to oscillate with a 2.0 s period. Find its moment of inertia.

## Solution Checkpoints

IDENTIFY AND SET UP Is this a physical pendulum?
EXECUTE The period of a physical pendulum is given by

$$
T=2 \pi \sqrt{\frac{I}{m g d}}
$$

Rearranging terms and solving yields a moment of inertia of $5.96 \times 10^{-2} \mathrm{~kg} \mathrm{~m}^{2}$.
EVALUATE This problem shows another method of determining the moment of inertia: Set the body in oscillation and measure the period.

## 14

## Fluid Mechanics

## Summary

We interact with fluids on a continual basis, from walking through air to swimming in the ocean. This chapter examines fluids, or substances that can flow, including liquids and gases. We will begin with fluid statics and use Newton's laws to describe the behavior of fluids at rest. Density, pressure, buoyancy, and surface tension, concepts needed for our investigation, will be defined. We will also delve into fluid dynamics and see how to analyze fluids in motion. Conservation of energy and Newton's laws will guide us in this examination. Although fluid dynamics can be quite complex, several examples will give us insight into the subject.

## Objectives

After studying this chapter, you will understand

- The definition of a material's density.
- The definition of pressure in a fluid and its measurement.
- How to analyze fluids in equilibrium and find the pressure at varying depths.
- Buoyancy and how to calculate the buoyancy acting on a body.
- How to compare and contrast laminar and turbulent fluid flow.
- How to apply Bernoulli's equation to fluid dynamics problems.

| Term | Description |
| :---: | :---: |
| Density | Density is the mass per unit volume of a material. For a homogeneous material with mass $m$ and volume $V$, the density is $\rho=\frac{m}{V} .$ <br> The SI unit of density is the kilogram per cubic meter $\left(1 \mathrm{~kg} / \mathrm{m}^{3}\right)$. The cgs unit is used to express density in grams per cubic centimeter $\left(1 \mathrm{gm} / \mathrm{cm}^{3}=1000 \mathrm{~kg} / \mathrm{m}^{3}\right)$. |
| Pressure | The pressure $p$ in a fluid is the normal force per unit area: $p=\frac{d F_{\perp}}{d A} .$ <br> The SI unit of pressure is the pascal $(\mathrm{Pa}) ; 1 \mathrm{~Pa}=1 \mathrm{~N} / \mathrm{m}^{2}$. Also common are the bar ( $10^{5} \mathrm{~Pa}$ ) and millibar ( $10^{2} \mathrm{bar}$ ). |
| Pressure in a Fluid | The pressure difference between two points in a fluid with uniform density $\rho$ is proportional to the difference in elevations between the two points: $p_{2}-p_{1}=-\rho g\left(y_{2}-y_{1}\right) .$ <br> Pascal's law states that the pressure applied to a fluid is transmitted through the fluid and depends only on depth. |
| Buoyant Force | Archimedes' principle states that when an object is immersed in a fluid, the fluid exerts an upward buoyant force on the object equal in magnitude to the weight of the fluid displaced by the object. |
| Fluid Flow | An ideal fluid is incompressible and has no viscosity. Conservation of mass requires that the amount of fluid flowing through a cross section of a tube per unit time be the same for all cross sections: $\frac{\Delta V}{\Delta t}=A_{1} v_{1}=A_{2} v_{2} .$ |
| Bernoulli's Equation | Bernoulli's equation relates the pressure $p$, flow speed $v$, and elevation $y$ of an ideal fluid at any two points: $p_{1}+\rho g y_{1}+\frac{1}{2} \rho v_{1}^{2}=p_{2}+\rho g y_{2}+\frac{1}{2} \rho v_{2}^{2}=\text { constant. }$ |

## Conceptual Questions

## 1: Ice in a glass

Two glasses are filled with water to the same level. In one glass, ice cubes float on the top. If the two glasses are made of the same material and have the same shape, how do their total weights compare?

## Solution

IDENTIFY, SET UP, AND EXECUTE The glasses must have the same weight, since they are made of the same material and have the same shape. We solve the problem by comparing the mass of the water alone to the mass of the water-plus-ice mix.

Archimedes' principle states that an object will displace its own weight in a fluid. The volume of water displaced by the ice has the same weight as the ice; therefore, the water in one glass weighs the same as the water plus ice in the second glass.

EVALUATE The volume of the glass with the ice is greater than the volume of the glass without ice, but weight depends on both density and volume. As with any new physical principle, we need to develop our skills carefully and not jump to conclusions.

## 2: Energy in a hydraulic lift

A hydraulic lift is used to lift a car. The piston supporting the car has a cross-sectional area 100 times larger than the cross-sectional area of the piston driving the lift. The drive piston will therefore require a force 100 times smaller than the weight of the car to lift the car. Does this mean that energy conservation is violated?

## Solution

IDENTIFY, SET UP, AND EXECUTE Pascal's law states that the pressure is the same at both pistons. The drive piston requires a small force to create the pressure that will lift the car. A large displacement in the drive piston creates a small displacement in the lift piston, due to the differences in areas. The amount of work done in moving the drive piston a long distance is equal to the work done by the lift piston moving a small distance (ignoring friction). The amounts of work are equivalent; energy conservation is not violated.

EVALUATE If you have ever operated a hydraulic jack to lift your car or a house, you should recall that you had to pump the jack several times to move a small distance. The work produced by your small force applied over a long distance was equivalent to the work done in lifting the object.

## 3: Race car spoilers

Why do race cars have spoilers, or wings, on their bodies?

## Solution

IDENTIFY, SET UP, AND EXECUTE Spoilers are essentially inverted airplane wings. We've seen that airplane wings produce lift by reducing the pressure above the plane's wing. The inverted wing produces a downward force to help hold the race car on the pavement and maintain contact between the wheels and the road. The spoiler also helps stabilize the car as it moves around the track.

EVALUATE Spoiler design for race cars is critical: There is a careful balance between enough downward force to keep the car on the track and too much force that causes lost fuel economy and premature tire wear. Some race cars have downward forces of up to three times the force of gravity (i.e., they could operate on an upside-down track and not fall off). Many cars have spoilers; most serve only an aesthetic purpose and have no effect on the car's performance.

## Problems

## 1: How much seawater in a tank

Seawater is stored under pressure in a tank of horizontal cross-sectional area $4.5 \mathrm{~m}^{2}$. The pressure above the seawater in the tank is $7.2 \times 10^{5} \mathrm{~Pa}$, and the pressure at the bottom of the tank is $1.2 \times 10^{6} \mathrm{~Pa}$. What is the mass of the seawater in the tank?

## Solution

IDENTIFY We will use the relations among pressure, density, and height to solve the problem. The target variable is the mass of the seawater.

SET UP The tank is sketched in Figure 14.1. We can find the mass by first finding the volume of the seawater and then multiplying by the density. To find the volume, we multiply the height by the crosssectional area of the tank. To find the height, we use the variation in the pressures due to the height of seawater. We assume that the seawater is incompressible.


Figure 14.1 Problem 1 sketch.
EXECUTE We start by finding the height of the seawater in the container. The difference in pressure is related to the height by

$$
p_{2}-p_{1}=\rho g\left(y_{2}-y_{1}\right)
$$

where we set $p_{1}$ as the pressure at the bottom of the tank and $p_{2}$ as the pressure at the top of the tank. Solving for the height, we get

$$
h=y_{2}-y_{1}=\frac{p-p_{0}}{\rho g}=\frac{\left(1.2 \times 10^{6} \mathrm{~Pa}\right)-\left(7.2 \times 10^{5} \mathrm{~Pa}\right)}{\left(1.03 \times 10^{3} \mathrm{~kg} / \mathrm{m}^{3}\right)\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}=47.6 \mathrm{~m}
$$

where we used the density of seawater given in Table 14.1 in the text $\left(1.03 \times 10^{3} \mathrm{~kg} / \mathrm{m}^{3}\right)$. We now find the volume of seawater in the tank. The volume is the height times the cross-sectional area:

$$
V=h A=(47.6 \mathrm{~m})\left(4.5 \mathrm{~m}^{2}\right)=214 \mathrm{~m}^{3} .
$$

The mass is the volume times the density we found from the table:

$$
m=\rho V=\left(1.03 \times 10^{3} \mathrm{~kg} / \mathrm{m}^{3}\right)\left(214 \mathrm{~m}^{3}\right)=221,000 \mathrm{~kg}
$$

The tank holds $221,000 \mathrm{~kg}$ of seawater.
EVALUATE This problem shows how to determine height from change in pressure. Altimeters find changes in altitude by monitoring the change in pressure of air.

## 2: Water pressure in a town

Water pressure in a town is maintained by a water tower 35.0 m high, open to the atmosphere at the top. (a) What is the gauge pressure at ground level? (b) A $1.75-\mathrm{cm}$-diameter garden hose at the bottom of the tower is open and spilling water. How much force is needed at the end of the hose to seal the end?

## Solution

IDENTIFY We will use the relations among pressure, density, and height to solve the problem. The target variables are the gauge pressure at the ground and the force needed to seal the hose.

SET UP The gauge pressure is found by multiplying the height by the density by $g$, the acceleration due to gravity. The force is found by finding the pressure differential at the hose and multiplying by the area of the opening. We assume that the water is incompressible.

EXECUTE We start by finding the pressure at the surface. The gauge pressure is given by

$$
\begin{aligned}
p_{\mathrm{g}} & =p-p_{\mathrm{a}}=\rho g\left(y_{2}-y_{1}\right) \\
& =\left(10^{3} \mathrm{~kg} / \mathrm{m}^{3}\right)\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(35.0 \mathrm{~m}) \\
& =3.43 \times 10^{5} \mathrm{~Pa} \\
& =3.40 \mathrm{~atm} .
\end{aligned}
$$

The pressure difference between the inside and the outside of the hose is the gauge pressure, equivalent to the force per unit area needed to seal the hose. The force needed is then

$$
F=\left(p-p_{\mathrm{a}}\right) A=\left(3.43 \times 10^{5} \mathrm{~Pa}\right)\left(\pi(0.0175 / 2 \mathrm{~m})^{2}\right)=83 \mathrm{~N}
$$

The gauge pressure at the ground level is 3.40 atm , and the force required to seal the hose is 83 N .
EVALUATE This problem illustrates how force, pressure, and height are related in a fluid.

## 3: Velocity of water exiting a fire hose

Water enters a round fire hose of diameter 3.5 cm and exits from a round, $0.60-\mathrm{cm}$-diameter nozzle. If the water enters the hose at $2.0 \mathrm{~m} / \mathrm{s}$, what is the velocity of the exiting water? What is the maximum horizontal range of the water leaving the hose?

## Solution

IDENTIFY The continuity equation relates the velocities and cross-sectional areas of incompressible fluids in a tube. The target variables are the velocity of the water at the outlet of the nozzle and the maximum range of the water.

SET UP We use the continuity equation to find the velocity of the water exiting the nozzle. To find the range, we employ projectile motion. We treat the water as incompressible.

EXECUTE The amount of fluid flowing through a tube per unit time is constant. The flow through the hose is equal to the flow through the nozzle:

$$
A_{\text {hose }} v_{\text {hose }}=A_{\text {nozzle }} v_{\text {nozzle }}
$$

The area of the hose or nozzle is $\pi$ times the square of half the diameter. Solving for the velocity of the nozzle gives

$$
v_{\text {nozzle }}=\frac{A_{\text {hose }} v_{\text {hose }}}{A_{\text {nozzle }}}=\frac{\pi\left(D_{\text {hose }} / 2\right)^{2} v_{\text {hose }}}{\pi\left(D_{\text {nozzle }} / 2\right)^{2}}=\frac{(3.5 \mathrm{~cm})^{2}(2.0 \mathrm{~m} / \mathrm{s})}{(0.6 \mathrm{~cm})^{2}}=68 \mathrm{~m} / \mathrm{s},
$$

where the factors $\pi$ and 2 cancel. The water molecules leaving the hose have a velocity of $68 \mathrm{~m} / \mathrm{s}$ and undergo acceleration due to gravity. We use the kinematic relations for two-dimensional motion to find the range. Recall that the horizontal range of a projectile in terms of the launch angle $\theta_{0}$ and the initial velocity $v_{0}$ is

$$
R=\frac{v_{0}^{2} \sin 2 \theta_{0}}{g}
$$

The maximum range occurs when the launch angle is $45^{\circ}$. Substituting our values, we obtain

$$
R=\frac{v_{0}^{2} \sin 2\left(45^{\circ}\right)}{g}=\frac{(68 \mathrm{~m} / \mathrm{s})^{2}(1)}{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}=470 \mathrm{~m}
$$

The maximum horizontal range of the water leaving the nozzle at $68 \mathrm{~m} / \mathrm{s}$ is 470 m .

EVALUATE This problem illustrates why nozzles are placed at the ends of hoses. The reduced diameter of the nozzle increases the exit velocity and thereby increases the range of the water. Next time you wash your car, compare the velocity and range of the water leaving the hose with and without the nozzle attached.

## 4: Ice cube in glycerine

What fraction of an ice cube is submerged when floating in glycerine?

## Solution

IDENTIFY We will use Newton's law and the definition of buoyancy to solve the problem. The target variable is the fraction of the ice cube submerged.

SET UP The buoyant force acts upward and gravity acts downward. The ice cube is in equilibrium, so the forces sum to zero.

EXECUTE The weight of the ice cube is

$$
w_{w}=\rho_{w} g V_{w} .
$$

The buoyant force is equal to the amount of glycerine displaced by the ice cube, which is equal to the weight of the ice cube. Combining terms yields

$$
F_{B}=\rho_{g} g V_{g}=\rho_{w} g V_{w}
$$

The fraction of the ice cube that is submerged is the volume of glycerine displaced divided by the volume of the ice cube. Rearranging terms in the previous equation yields the fraction

$$
\frac{V_{\mathrm{g}}}{V_{\mathrm{w}}}=\frac{\rho_{\mathrm{w}}}{\rho_{\mathrm{g}}}=\left(\frac{0.92 \mathrm{~g} / \mathrm{cm}^{2}}{1.26 \mathrm{~g} / \mathrm{cm}^{2}}\right)=0.73
$$

Thus, $73 \%$ of the ice cube is submerged when floating in glycerine. Note that the densities of ice and glycerine were taken from Table 14.1 in the text.

EVALUATE We see how we can determine the fraction of ice located below the surface when the ice is placed in a liquid. You can use the same procedure to find out how much of an iceberg is below the surface in the ocean.

Practice Problem: Draw the free-body diagram of the problem.

## 5: Examining the buoyant force

A $4.5-\mathrm{cm}$-radius sphere of wood is held in fresh water below the surface by a spring. If the spring's force constant is $55 \mathrm{~N} / \mathrm{m}$, by how much is the spring stretched from its equilibrium position? Take the density of wood to be $700 \mathrm{~kg} / \mathrm{m}^{3}$.

## Solution

IDENTIFY We will use Newton's law and the definition of buoyancy to solve for the stretch of the spring, the target variable. The wood is in equilibrium.

SET UP The free-body diagram of the block of wood is shown in Figure 14.2. The buoyant force is directed upward, and both gravity and tension due to the spring are directed downward. The block of wood is in equilibrium, so the forces sum to zero. The size and density of the wood determine its vol-
ume and mass, needed for the buoyant force and gravity. Hooke's law will be used to determine the amount of stretch.


Figure 14.2 Problem 5 sketch and free-body diagram.
EXECUTE We sum the three forces acting on the block of wood. The forces act only in the vertical direction and add to zero:

$$
\sum F_{y}=F_{B}-m g-F_{s}=0
$$

The buoyant force is equal to the amount of water displaced by the wooden sphere. The volume of a sphere is $4 / 3 \pi r^{2}$. Combining terms yields

$$
F_{B}=\rho_{\text {water }} V_{\text {sphere }} g=\rho_{\text {water }}\left(\frac{4}{3} \pi r_{\text {sphere }}^{3}\right) g .
$$

The spring force is equal to the spring constant times its displacement, $k x$. The mass of the sphere is its volume times its density. Inserting these expressions into the equilibrium equation gives

$$
\sum F_{y}=F_{B}-m g-F_{s}=\rho_{\text {water }}\left(\frac{4}{3} \pi r_{\text {sphere }}^{3}\right) g-\rho_{\text {wood }}\left(\frac{4}{3} \pi r_{\text {sphere }}^{3}\right) g-k x=0
$$

Rearranging terms, we solve for $x$ :

$$
\begin{aligned}
& x=\frac{\rho_{\text {water }}\left(\frac{4}{3} \pi r_{\text {sphere }}^{3}\right) g-\rho_{\text {wood }}\left(\frac{4}{3} \pi r_{\text {sphere }}^{3}\right) g}{k}=\frac{\left(\rho_{\text {water }}-\rho_{\text {wood }}\right)\left(\frac{4}{3} \pi r_{\text {sphere }}^{3}\right) g}{k}, \\
& x=\frac{\left(\left(1 \times 10^{3} \mathrm{~kg} / \mathrm{m}^{3}\right)-\left(700 \mathrm{~kg} / \mathrm{m}^{3}\right)\right)\left(\frac{4}{3} \pi(0.045 \mathrm{~m})^{3}\right)\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{(55 \mathrm{~N} / \mathrm{m})}=0.020 \mathrm{~m} .
\end{aligned}
$$

The spring is stretched 0.020 m , or 2.0 cm , from its equilibrium position.
EVALUATE This problem illustrates how to incorporate the buoyant force with previously encountered forces to solve equilibrium problems. We've used the same procedure as in the past: starting with a free-body diagram and setting the net force equal to zero. Working with fluids requires conversions among volume, mass, and density.

## 6: Water from a tank

A 10.0-m-high cylindrical tank of cross-sectional area $0.75 \mathrm{~m}^{2}$ is filled with water. (a) Find the velocity of discharge as a function of the height of water remaining in the tank when a hole of area $0.40 \mathrm{~m}^{2}$ is opened at the bottom of the tank. (b) Find the initial discharge velocity. (c) Find the initial volume rate of discharge.

## Solution

IDENTIFY Bernoulli's equation and the continuity equation will be used to relate the pressure difference, height, and velocity of the flowing water. The target variables are the discharge velocity and volume rate of discharge.

SET UP Bernoulli's equation will be used to relate the velocities at the top and bottom of the tank to the change in height of the water. The continuity equation also relates the two velocities. Combining both equations will yield the velocity at the bottom of the tank, from which we can determine the solutions to parts (b) and (c).

EXECUTE Bernoulli's equation applied to the top and bottom of the cylinder gives

$$
p_{\text {top }}+\rho g y_{\text {top }}+\frac{1}{2} \rho v_{\text {top }}^{2}=p_{\text {bottom }}+\rho g y_{\text {bottom }}+\frac{1}{2} \rho v_{\text {bottom }}^{2} .
$$

Both sides are open to atmospheric pressure, so the pressures are the same. We derive an expression for the velocities:

$$
v_{\text {bottom }}^{2}-v_{\text {top }}^{2}=2 g\left(y_{\text {top }}-y_{\text {bottom }}\right)=2 g h .
$$

The continuity equation yields another relation between the two velocities:

$$
v_{\text {bottom }} A_{\text {bottom }}=v_{\text {top }} A_{\text {top }}
$$

or

$$
v_{\text {top }}=\frac{A_{\text {botom }}}{A_{\text {top }}} v_{\text {bottom }} .
$$

Placing the right-hand side of the latter equation into the previous equation and solving for the velocity at the bottom gives

$$
v_{\text {bottom }}^{2}-\left(\frac{A_{\text {bottom }}}{A_{\text {top }}} v_{\text {botom }}\right)^{2}=2 g h,
$$

or

$$
v_{\text {botom }}=\sqrt{2 g h}\left(1-\left(\frac{A_{\text {bottom }}}{A_{\text {top }}}\right)^{2}\right)^{-1 / 2} .
$$

The discharge velocity as a function of the height of the water remaining in the tank is

$$
v_{\text {bottom }}=1.20 \sqrt{g h} .
$$

The initial discharge velocity is when the tank begins to drain $(h=10.0 \mathrm{~m})$ and is equal to $11.8 \mathrm{~m} / \mathrm{s}$. The discharge rate when the tank begins to drain is

$$
\text { volume discharge rate }=A_{\text {botom }} v_{\text {botom }}=4.73 \mathrm{~m}^{3} / \mathrm{s} \text {. }
$$

EVALUATE We must carefully check the units in this problem. Do they cancel correctly?

## 7: Lift on a car on a highway

As a car travels down the highway, the speed of the air flowing over the top of the car is higher than the speed of the air flowing under the car, thus creating lift. Estimate the lift on a car as it travels at 100 kph .

Take the density of air to be $1.20 \mathrm{~kg} / \mathrm{m}^{3}$, the car's area to be $6 \mathrm{~m}^{2}$, and the height of the car to be 1.0 m . Assume that air travels under the car at 100 kph and over the top of the car at 140 kph .

## Solution

IDENTIFY Bernoulli's equation gives the pressure difference between the top and bottom of the car. We will use that equation to find the target variable, the lift on the car.

SET UP We find the lift force acting on the car from the definition of pressure as force per unit surface area.

EXECUTE The car is moving through a fluid (air), so Bernoulli's equation can be applied. We'll compare the pressures below and above the car to find the pressure difference. Bernoulli's equation is

$$
p_{\text {above }}+\rho g y_{\text {above }}+\frac{1}{2} \rho v_{\text {above }}^{2}=p_{\text {below }}+\rho g y_{\text {below }}+\frac{1}{2} \rho v_{\text {below. }}^{2} .
$$

The pressure difference is then

$$
\Delta p=\rho g y_{\text {above }}+\frac{1}{2} \rho v_{\text {above }}^{2}-\frac{1}{2} \rho v_{\text {below }}^{2},
$$

where we have set the origin below the car $\left(y_{\text {below }}=0\right)$. The pressure difference is

$$
\begin{aligned}
\Delta p & =\rho\left(g y_{\text {above }}+\frac{1}{2}\left(v_{\text {above }}^{2}-v_{\text {below }}^{2}\right)\right) \\
& =\left(1.20 \mathrm{~kg} / \mathrm{m}^{3}\right)\left(\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(1.0 \mathrm{~m})+\frac{1}{2}\left((38.9 \mathrm{~m} / \mathrm{s})^{2}-(27.7 \mathrm{~m} / \mathrm{s})^{2}\right)\right) \\
& =459 \mathrm{~Pa},
\end{aligned}
$$

where we replaced 100 kpm with $27.7 \mathrm{~m} / \mathrm{s}$ and 140 kph with $38.9 \mathrm{~m} / \mathrm{s}$. We find the force by multiplying the pressure by the area of the car:

$$
F=P A=(459 \mathrm{~Pa})\left(6.0 \mathrm{~m}^{2}\right)=2750 \mathrm{~N} .
$$

The lift on the car is 2750 N .
EVALUATE We see that the lift is significant in this case-roughly equivalent to a weight of 280 kg . It is not enough to lift the car off the highway, since most cars weigh over 1000 kg . We assumed that the flow of air around the car was smooth and that air is incompressible. Neither are valid assumptions and should be modified in a careful examination. Our results show the maximum lift of the car.

## 8: Pressure in a water system

Water is discharged from a closed system, reaching a maximum height of 10.0 m . What is the gauge pressure of the water system at the hose nozzle?

## Solution

IDENTIFY Bernoulli's equation gives the pressure difference between the tank and the top of water in flight. The target variable is the gauge pressure in the tank.

SET UP Kinematics is used to find the initial velocity of the water, taking into account the maximum height. Bernoulli's equation is then used to find the pressure in the tank.

EXECUTE From kinematics, for the water to reach a height $h$, it must have an initial velocity of

$$
v_{\mathrm{noz}}^{2}=2 g h .
$$

Bernoulli's equation is applied to two points, one inside the tank and one at the nozzle, both at the same height, to find the pressure:

$$
p_{\mathrm{in}}+\rho g y+\frac{1}{2} \rho v_{\mathrm{in}}^{2}=p_{\mathrm{noz}}+\rho g y+\frac{1}{2} \rho v_{\mathrm{noz}}^{2} .
$$

The outside pressure is atmospheric pressure, so the difference between the two pressures is the gauge pressure, our target variable. The velocity inside the tank is zero, giving

$$
p_{\text {gauge }}=p_{\text {in }}-p_{\mathrm{a}}=\frac{1}{2} \rho v_{\mathrm{noz}}^{2} .
$$

Combining the results produces

$$
p_{\text {gauge }}=\frac{1}{2} \rho(2 g h)=\left(10^{3} \mathrm{~kg} / \mathrm{m}^{3}\right)\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(10.0 \mathrm{~m})=9.8 \times 10^{4} \mathrm{~Pa} .
$$

The tank is at a gauge pressure of $9.8 \times 10^{4} \mathrm{~Pa}$.
EVALUATE We've combined kinematics and fluid dynamics in this problem. We'll continue building our physics models and call upon older material to help us as we move forward.

## Try It Yourself!

## 1: Leak in a submarine

A submarine descends 35.0 m into the ocean and springs a leak. The hole out of which water is leaking has a diameter of 2.5 cm . What force must be used to plug the hole?

## Solution Checkpoints

IDENTIFY AND SET UP We will use the relations among pressure, density, and depth to solve the problem. The target variable is the force needed to seal the hole. We must find the pressure at depth first. The density of seawater is $1.03 \times 10^{3} \mathrm{~kg} / \mathrm{m}^{3}$.

EXECUTE The pressure difference between the water and the submarine (assuming that the latter is at a pressure of 1 atmosphere inside) is given by

$$
p-p_{\mathrm{a}}=\rho g\left(y_{2}-y_{1}\right) .
$$

This equation yields a pressure difference of 3.50 atm . The force needed to plug the hole is

$$
F=\left(p-p_{a}\right) A .
$$

The force required to seal the hole is 173 N .
EVALUATE Do you think it is possible for someone to push with a force of 173 N and seal the hole?

## 2: Floating dumpling

A dumpling floats two-thirds submerged in water. What is its density?

## Solution Checkpoints

IDENTIFY AND SET UP The dumpling is in equilibrium. What forces are acting on it? How do you quantify the buoyant force?

EXECUTE The fraction of the dumpling that is submerged is given by

$$
\frac{V_{D}}{V_{w}}=\frac{\rho_{w}}{\rho_{D}}
$$

This equation can be rearranged to solve for the density of the dumpling. Doing so yields a density of $0.66 \mathrm{~g} / \mathrm{cm}^{2}$.

EVALUATE This problem illustrates how we can find the density of an object by examining how it floats.

## 3: Weighing a sphere under water

A sphere of volume $10 \mathrm{~cm}^{3}$ displaces a spring scale and is found to weigh 50.0 g when submerged in water. What are the sphere's mass and density?

## Solution Checkpoints

IDENTIFY AND SET UP The sphere is in equilibrium. What three forces are acting on it?
EXECUTE The net force acting on the sphere is

$$
\sum F_{y}=F_{s}+F_{B}-m g=0 .
$$

The weight and buoyant force are written in terms of density, volume, and $g$, the acceleration due to gravity. These quantities can be rearranged to yield

$$
\rho=\frac{F_{s}}{V g}+\rho_{w}
$$

From this equation, the density of the sphere is $6.0 \mathrm{~g} / \mathrm{cm}^{3}$ and the mass of the sphere is 60.0 g .
EVALUATE How did you decide the direction that the spring force acted in?

## 4: Water from a tank

Water inside an enclosed tank is subjected to a pressure of two atmospheres at the top of the tank. What is the velocity of discharge from a small hole 3.0 m below the surface of the water?

## Solution Checkpoints

IDENTIFY AND SET UP Use Bernoulli's equation and the continuity equation to solve this problem. The target variable is the discharge velocity. The small hole indicates that the ratio of the hole to the surface area is small.

EXECUTE Bernoulli's equation applied to the top of the water and hole gives

$$
p_{\text {top }}+\rho g y_{\text {top }}+\frac{1}{2} \rho v_{\text {top }}^{2}=p_{\text {hole }}+\rho g y_{\text {hole }}+\frac{1}{2} \rho v_{\text {hole }}^{2}
$$

What is the pressure outside and at the top? Can you assume that the velocity at the top surface is small? The continuity equation indicates that

$$
v_{\mathrm{top}}=\frac{A_{\mathrm{hole}}}{A_{\mathrm{top}}} v_{\mathrm{hole}} \approx 0
$$

Solving for the velocity at the hole gives

$$
v_{\text {hole }}^{2}=\frac{2}{\rho}\left(p_{\mathrm{a}}+\rho g h\right),
$$

or a velocity of $16 \mathrm{~m} / \mathrm{s}$.
EVALUATE How can you check these results?

## 5: Water from a rocket

A toy rocket of diameter 2.0 in consists of water under the pressure of compressed air pumped into the nose chamber. When the gauge air pressure is $60 \mathrm{lb} / \mathrm{in}^{2}$, the water is ejected through a hole of diameter 0.2 in . Find the propelling force, or thrust, of the rocket.

## Solution Checkpoints

IDENTIFY AND SET UP Use Bernoulli's equation and the continuity equation to solve this problem. The target variable is the thrust.

EXECUTE Bernoulli's equation applied to points inside and outside of the rocket, at the same height, gives

$$
p_{\mathrm{in}}+\rho g y+\frac{1}{2} \rho v_{\mathrm{in}}^{2}=p_{\mathrm{out}}+\rho g y+\frac{1}{2} \rho v_{\mathrm{out}}^{2} .
$$

What is the pressure $p_{a}$ outside the rocket? What is the velocity of water inside the rocket? The thrust is given by

$$
F=v_{\text {out }} \frac{d m_{\text {out }}}{\mathrm{dt}}=v_{\text {out }} \rho \frac{d V_{\text {out }}}{\mathrm{dt}}=v_{\text {out }} \rho A_{\text {out }} v_{\text {out }} \text {. }
$$

Solving for the force, we find that it is 3.8 lb .
EVALUATE Did you check units?

## 15 <br> Mechanical Waves

## Summary

In this chapter, we expand the concept of the periodic motion of an object to the periodic motion of many particles connected together as a medium. The periodic motion of a medium is a mechanical wave. Waves occur in many forms, including ocean waves, sound, light, earthquakes, and television transmission. This chapter will form the foundation for studying a variety of waves. We'll begin with the description of transverse and longitudinal waves and their amplitudes, periods, frequencies, and wavelengths. We'll see how waves move, interact, transmit energy, reflect, and combine in a variety of ways and how to describe their frequencies.

## Objectives

After studying this chapter, you will understand

- How to identify longitudinal and transverse waves and their media.
- The relations among the period, velocity, frequency, and wavelength of a wave.
- How a wave function that satisfies the wave equation describes a wave.
- The concepts of superposition, standing waves, nodes, and antinodes.
- The allowed frequencies for standing waves.
- How waves interact and interfere.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Mechanical Wave | A mechanical wave is a disturbance from equilibrium that propagates from <br> one region of space to another through a medium. In a transverse wave, the <br> particles in the medium are displaced perpendicular to the direction of travel. <br> In a longitudinal wave, the particles in the medium are displaced parallel to <br> the direction of travel. |
| Periodic Mechanical Waves | In a periodic wave, particles in the medium exhibit periodic motion. The <br> speed, wavelength, period, and frequency of a periodic wave are related by |

$$
v=\lambda f=\frac{\lambda}{T} .
$$

The speed of a transverse wave in a string under tension is given by

$$
v=\sqrt{\frac{F_{T}}{\mu}}
$$

where $F_{T}$ is the tension in the rope and $\mu$ is the mass per unit length.

## Wave functions

The wave function $y(x, t)$ describes the displacements of individual particles in the medium. For a sinusoidal wave traveling in the $+x$ direction, the equation

$$
\begin{aligned}
y(x, t) & =A \cos \left[\omega\left(\frac{x}{v}-t\right)\right] \\
& =A \cos (k x-\omega t)
\end{aligned}
$$

describes the wave. The wave functions are solutions of the wave equation

$$
\frac{\partial^{2} y(x, t)}{\partial x^{2}}=\frac{1}{v^{2}} \frac{\partial^{2} y(x, t)}{\partial t^{2}} .
$$

## Wave Power

## Principle of Superposition

The principle of superposition states that when two waves overlap, the net displacement at any point at any time is found by taking the sum of the displacements of the individual waves:

$$
y(x, t)=y_{1}(x, t)+y_{2}(x, t) .
$$

## Standing Waves

A standing wave is that combination of sinusoidal waves which produces a stationary sinusoidal pattern. Nodes are points where the standing wave pattern does not change with time. Antinodes are positions half way between nodes; the amplitude is maximum at an antinode. The distance between successive nodes or antinodes is one-half of the wavelength. A string of length $L$ held stationary at both ends can have standing waves only with frequencies such that

$$
f_{n}=n \frac{v}{2 L}=n f_{1} \quad(n=1,2,3, \ldots) .
$$

The fundamental frequency is given by

$$
f_{1}=\frac{1}{2 L} \sqrt{\frac{F}{\mu}} .
$$

The multiples of $f_{1}$ are the harmonics. Each frequency and its associated vibration pattern is called a normal mode.

## Conceptual Questions

## 1: Waves in a jump rope

Your younger sister is playing with a jump rope. She ties one end to a fence post and moves the other end up and down, observing waves in the rope. She sees waves moving down the rope and becomes confused. She asks you why the rope doesn't move toward the fence, since that is how the waves move. How do you answer her question?

## Solution

IDENTIFY, SET UP, AND EXECUTE Having just learned about mechanical waves, you explain that the rope is made up of lots of little pieces (particles) and the pieces at the end she holds move up and down as she moves her hand up and down. The pieces of rope next to her end are connected to the pieces she is moving and so are pulled up and down at the same time. These pieces, which take a little bit longer to move than the ones she holds, are connected to other pieces, which also move up and down. Each successive piece takes a bit longer to start moving than the piece before it, which is why there is a wave pattern. This wave pattern is what your sister observes moving down the rope. All of the individual pieces of rope move only up and down. Since they are connected to each other, their moving creates a disturbance in the rope that appears to move down the rope. The rope doesn't move toward the fence because none of the pieces of the rope move toward the fence.

EVALUATE Remember that when you pluck a string, you pull the string to the side, so you give the string a velocity to the side of, or perpendicular to, the string. During the pluck, you impart a force perpendicular to the string, not along the string.

## 2: Velocity of a wave

A taut string is plucked and a wave travels down the string at speed $v$. How can you double the speed of the wave?

## Solution

IDENTIFY, SET UP, AND EXECUTE The speed of a transverse wave on a string is proportional to the square root of the tension on the string and inversely proportional to the square root of the mass per unit length. To double the speed, you can quadruple the tension in the string or decrease the mass per unit length by a factor of 4 .

EVALUATE This problem illustrates the dependence of the speed of a wave in a string on the mass per length and the tension in the string.

## 3: Heavy rope

A heavy rope is suspended vertically and is stretched taut by a $10.0-\mathrm{kg}$ mass attached to the bottom of the rope. The top end of the rope is plucked, creating a wave. Does the speed of the wave change as it propagates down the rope? If so, how?

## Solution

IDENTIFY, SET UP, AND EXECUTE The speed of a transverse wave on a string depends on the tension and the mass per unit length in the rope. We can assume that the mass per unit length is constant. Because the rope is heavy, the top end must have greater tension than the bottom end, since the top end supports both the mass at the end and the rope itself. The tension, therefore, decreases toward the bottom of the rope. The decreasing tension will slow the speed of the wave as it travels down the rope.

The wave speed is not constant and slows as it approaches the bottom of the rope.
EVALUATE We will normally not encounter the varying tension and speed found in this problem. We'll focus on light strings to understand wave propagation better.

## Problems

## 1: An unusual scale

Your strange physics professor builds an unusual scale by hanging an object from a $3.0-\mathrm{m}$-long wire attached to the ceiling. She plucks the string just above the object and finds that the pulse takes 0.50 s to propagate up and down the wire. What is the mass of the object? The mass of the wire is 0.50 kg .

## Solution

IDENTIFY AND SET UP The speed of a wave in a wire under tension is related to the tension in the wire. By finding the speed of the wave in the wire, we'll determine the mass of the object.

EXECUTE The speed of the wave in the wire is given by

$$
v=\sqrt{\frac{F_{T}}{\mu}}
$$

The tension at the bottom of the wire is equal to the gravitational force on the object, since the object is in equilibrium. To calculate the tension, we first need the velocity and the mass per unit length. The velocity is found by noting that the wave takes 0.50 s to travel 6.0 m (up and down the wire):

$$
v=\frac{\Delta d}{\Delta t}=\frac{6.0 \mathrm{~m}}{0.5 \mathrm{~s}}=12 \mathrm{~m} / \mathrm{s}
$$

The mass per unit length is

$$
\mu=\frac{m}{L}=\frac{0.50 \mathrm{~kg}}{3.0 \mathrm{~m}}=0.167 \mathrm{~kg} / \mathrm{m} .
$$

Substituting to find the tension, we obtain

$$
F_{T}=v^{2} \mu=(12 \mathrm{~m} / \mathrm{s})^{2}(0.167 \mathrm{~kg} / \mathrm{m})=24 \mathrm{~N}
$$

The tension force is equal to the weight, so the mass of the object is

$$
m=\frac{F_{T}}{g}=\frac{24 \mathrm{~N}}{9.8 \mathrm{~m} / \mathrm{s}^{2}}=2.4 \mathrm{~kg} .
$$

The object's mass is 2.4 kg .
EVALUATE This unusual scale illustrates how we can use mechanical waves to measure mass, but it is impractical for several reasons. First, the scale requires a high ceiling and a method of accurately measuring the speed of waves in the wire. Also, we have omitted the mass of the wire, which is roughly $15 \%$ higher at the ceiling, in our calculation of the tension.

## 2: Write a wave equation

Write the wave equation for a traveling transverse wave that propagates in the $+x$ direction, has a maximum disturbance from equilibrium of 1.0 cm , and has a wavelength of 2.0 m and a period 0.02 s . At $x=0.5 \mathrm{~m}$ and $t=0$, the instantaneous particle velocity is $\pi / 2 \mathrm{~m} / \mathrm{s}$ downward.

## Solution

IDENTIFY AND SET UP We'll begin with the general form of the wave function and use the given conditions to determine the constants.

EXECUTE The general form of a wave equation is

$$
y(x, t)=A \sin (\omega t-k x+\phi)
$$

Here, we used a minus sign in front of the wave number to ensure that the wave propagates toward positive $x$, and we included a phase angle. The frequency is found from the period:

$$
\omega=2 \pi \frac{1}{T}=314 \mathrm{rad} / \mathrm{s} .
$$

The wave number is related to the wavelength:

$$
k=\frac{2 \pi}{\lambda}=3.14 / \mathrm{m}
$$

The amplitude is the maximum displacement from zero, so $A=0.01 \mathrm{~m}$. To find the phase angle, we'll have to use the given velocity at the specified point. The velocity is the first derivative:

$$
v_{y}=\frac{\partial y}{\partial t}=A \omega \cos (\omega t-k x+\phi)
$$

At $x=0.5 \mathrm{~m}$ and $t=0$, the instantaneous particle velocity is $\pi / 2 \mathrm{~m} / \mathrm{s}$ downward, or negative. This gives

$$
\begin{gathered}
v_{y}(0.5 \mathrm{~m}, 0)=-\pi / 2 \mathrm{~m} / \mathrm{s} \\
(0.01 \mathrm{~m})(314 \mathrm{rad} / \mathrm{s}) \cos ((314 \mathrm{rad} / \mathrm{s})(0)-(3.14 / \mathrm{m})(0.5 \mathrm{~m})+\phi)=-\pi / 2 \mathrm{~m} / \mathrm{s} \\
(\pi) \cos (-\pi / 2+\phi)=-\pi / 2 \\
\sin (\phi)=-1 / 2
\end{gathered}
$$

The phase angle must be $-30^{\circ}$, or $-\pi / 6$. The complete wave function is

$$
y(x, t)=(0.01 \mathrm{~m}) \sin ((314 \mathrm{rad} / \mathrm{s}) t-(3.14 / \mathrm{m}) x-\pi / 6)
$$

EVALUATE This problem illustrates how to construct a wave function, given the properties of the wave.

## 3: Wave function check

Does $y(x, t)=A e^{-k x} \sin \omega t$ satisfy the wave function equation?

## Solution

IDENTIFY AND SET UP We will take the second derivatives of the function with respect to both time and position. We will then substitute the results into the wave equation and see if it is satisfied.

EXECUTE The wave equation is given by

$$
\frac{\partial^{2} y(x, t)}{\partial x^{2}}=\frac{1}{v^{2}} \frac{\partial^{2} y(x, t)}{\partial t^{2}}
$$

We start by taking the derivative of the given function with respect to position. We have

$$
\frac{\partial y(x, t)}{\partial x}=\frac{\partial A e^{-k x} \sin \omega t}{\partial x}=A(-k) e^{-k x} \sin \omega t .
$$

Next, we take the second derivative with respect to position:

$$
\frac{\partial^{2} y(x, t)}{\partial x^{2}}=\frac{\partial}{\partial x}\left(A(-k) e^{-k x} \sin \omega t\right)=A k^{2} e^{-k x} \sin \omega t
$$

We now switch to the time derivatives. The first derivative with respect to time is

$$
\frac{\partial y(x, t)}{\partial t}=\frac{\partial A e^{-k x} \sin \omega t}{\partial t}=A(\omega) e^{-k x} \cos \omega t
$$

The second derivative is

$$
\frac{\partial^{2} y(x, t)}{\partial t^{2}}=\frac{\partial^{2} A e^{-k x} \sin \omega t}{\partial t^{2}}=A\left(-\omega^{2}\right) e^{-k x} \sin \omega t
$$

Combining the results gives

$$
\begin{gathered}
A\left(-\omega^{2}\right) e^{-k x} \sin \omega t=A\left(k^{2}\right) e^{-k x} \sin \omega t \\
A e^{-k x} \sin \omega t=-\frac{1}{v^{2}} A e^{-k x} \sin \omega t
\end{gathered}
$$

The latter formula would satisfy the wave equation were it not for the minus sign. The function does not satisfy the wave equation and is not a valid wave function.

EVALUATE Constructing valid wave functions takes practice and experience. We'll see that there are several common types of equations that describe most waves and satisfy the wave equation.

## 4: Combining strings

Two strings of mass per unit length $\mu_{1}$ and $\mu_{2}$ are joined together at their ends. The tension in the two strings is the same. If the wavelength in the first string with $\mu_{1}=5.0 \mathrm{~g} / \mathrm{m}$ is 3.0 cm , what is the mass per unit length in the second string if the wavelength in that string is 5.0 cm ? Assume that the wave frequencies are the same in the two strings.

## Solution

IDENTIFY We'll use the fact that the tension is the same in both strings to determine the target variable, the mass per unit length in the second string.

SET UP The speed of propagation depends on the mass per unit length and the tension. The tension and frequency are the same in both strings. We'll set these equal to each other to solve.

EXECUTE The speed is given by

$$
v=\sqrt{\frac{T}{\mu}}
$$

Squaring both sides of this equation and solving for the tension gives

$$
v^{2} \mu=T
$$

Since both tensions are the same, we have

$$
v_{1}^{2} \mu_{1}=v_{2}^{2} \mu_{2}
$$

The speeds may not be the same, but we know the frequencies and wavelengths in both strings. The frequencies are the same in both strings. Rewriting the speed in terms of frequency and wavelength gives

$$
\lambda_{1}^{2} f^{2} \mu_{1}=\lambda_{2}^{2} f^{2} \mu_{2}
$$

Solving for the mass per unit length yields

$$
\mu_{2}=\frac{\lambda_{1}^{2} \mu_{1}}{\lambda_{2}^{2}}=\frac{(3.0 \mathrm{~cm})^{2}(5.0 \mathrm{~g} / \mathrm{m})}{(5.0 \mathrm{~cm})^{2}}=1.8 \mathrm{~g} / \mathrm{m}
$$

The second string has a mass per unit length of $1.8 \mathrm{~g} / \mathrm{m}$.
EVALUATE Do we expect a smaller mass per unit length in the second string? Yes, since the wavelength is larger in the second string, the mass per length must be less in order for the tensions to be the same.

Practice Problem: How do the speeds of the waves compare in the two strings? Answer: The speed of the wave in string 2 must be $\frac{5}{3}$ the speed in string 1 .

## 5: Modes in a string

A uniform string of length 0.50 m is fixed at both ends. Find the wavelength of the fundamental mode of vibration. If the wave speed is $300 \mathrm{~m} / \mathrm{s}$, find the frequency of the fundamental and next possible modes.

## Solution

IDENTIFY We'll use the properties and definitions of modes for a standing wave on a string to solve the problem.

SET UP The fundamental mode has a wavelength twice the length of the string. The frequency of the fundamental mode is the velocity divided by the wavelength. Higher frequencies are integer multiples of the fundamental frequency.

EXECUTE The wavelength of the fundamental mode is twice the length of the string. The wavelength is 1.0 m . For a wave speed of $300 \mathrm{~m} / \mathrm{s}$, the frequency of the fundamental mode is

$$
f_{1}=\frac{v}{2 L}=300 \mathrm{~Hz}
$$

The next possible mode will have half of the fundamental mode's wavelength, or a wavelength of 0.5 m . Its frequency will be

$$
f_{2}=2 f_{1}=\frac{v}{\lambda_{2}}=600 \mathrm{~Hz}
$$

EVALUATE This problem gives us practice understanding the properties of standing waves.

## Try It Yourself!

## 1: Wave function check

Does $y(x, t)=A(x-v t)^{n}$ satisfy the wave function equation? $A$ is a constant and $n>1$.

## Solution Checkpoints

IDENTIFY AND SET UP Take the second derivatives of the function with respect to both time and position. Then substitute the results into the wave equation and see if it is satisfied.

EXECUTE The second derivative with respect to position is

$$
\frac{\partial^{2} y(x, t)}{\partial x^{2}}=n(n-1) A(x-v t)^{n-2}
$$

The second derivative with respect to time is

$$
\frac{\partial^{2} y(x, t)}{\partial t^{2}}=n(n-1) A v^{2}(x-v t)^{n-2}
$$

Combining the results, we see that the function does satisfy the wave equation and is a valid wave function.

EVALUATE What type of wave does the function represent? Is it periodic?

## 2: Changing the diameter

Waves propagate through a rope under tension, provided by a hanging mass of 20.0 kg , with a speed of $30.0 \mathrm{~m} / \mathrm{s}$. The rope is replaced by ropes made of the same material but different diameters. For the velocity to remain the same, what mass should be hung from the end of the rope if the replacement rope has (a) half the diameter and (b) twice the diameter of the original rope?

## Solution Checkpoints

IDENTIFY AND SET UP The speed of propagation depends on the tension and the mass per unit length.
EXECUTE How must the mass per unit length and the tension relate if the speed is to remain constant? How does the mass change with (a) half the diameter rope and (b) twice the diameter rope?

With half the diameter, the volume decreases by a factor of four, so the mass should be 5.0 kg . With twice the diameter, the volume increases by a factor of four, so the mass should be 80.0 kg .

EVALUATE How do you confirm these results?

## 3: Modes in a free string

A uniform string of length 0.50 m is fixed at one end and free at the other end. Find the wavelength of the fundamental mode of vibration. If the wave speed is $300 \mathrm{~m} / \mathrm{s}$, find the frequency of the fundamental and next possible modes.

## Solution Checkpoints

IDENTIFY Use the properties and definitions of modes for a standing wave on a string to solve the problem.

SET UP How do the fundamental mode's frequency and wavelength for a standing wave on a string with only one end fixed compare with the fundamental mode's frequency and wavelength for a string with both ends fixed?

EXECUTE The wavelength of the fundamental mode is four times the length, or 2.0 m . The frequency of the fundamental mode is then 150 Hz .

The next mode has a wavelength of $\frac{4}{3} L$, or 0.67 m , and a frequency of 450 Hz .
EVALUATE Sketch the standing waves on the string for the fundamental and next possible modes. Does your sketch agree with the results you just calculated?

## Sound and Hearing

## Summary

In this chapter, we expand the concept of mechanical waves in order to understand sound and hearing. We'll begin with a description of longitudinal sound waves and their amplitudes, periods, frequencies, and wavelengths. We'll see how waves propagate through gases, liquids, and solids; how to determine the speed and intensity of sound waves; and how sound is produced by musical instruments. We'll also examine how the frequency of sound waves changes relative to the motion of the source and listener, summarized in the Doppler effect.

## Objectives

After studying this chapter, you will understand

- How sound waves are formed and propagate through media.
- How to apply the concepts of superposition, standing waves, nodes, and antinodes to sound waves
- How to calculate the intensity of a sound wave.
- The allowed frequencies for longitudinal standing sound waves.
- The definition and how to calculate sound beats.
- How to apply the Doppler effect to moving sources and listeners.
- How to apply acoustics to a variety of systems.

| Term | Description |
| :---: | :---: |
| Sound Waves | Sound consists of longitudinal waves propagating through a mediu pressure amplitude is given by $p_{\max }=B k A,$ <br> where $B$ is the bulk modulus of the medium, $k$ is the wave number, displacement amplitude. The speed of the sound wave depends on th $\begin{array}{ll} v=\sqrt{\frac{B}{\rho}} & \text { (longitudinal wave in a fluid) } \\ v=\sqrt{\frac{\gamma R T}{M}} & \text { (longitudinal wave in an ideal gas) } \\ v=\sqrt{\frac{Y}{\rho}} & \text { (longitudinal wave in a solid rod) } \end{array}$ |

## Intensity

The intensity of a sound wave is the rate at which energy is transported per unit area per unit time. The intensity of a sinusoidal wave is given by

$$
I=\frac{1}{2} \sqrt{\rho B} \omega^{2} A^{2}=\frac{p_{\max }^{2}}{2 \rho \mathrm{v}}=\frac{p_{\max }^{2}}{2 \sqrt{\rho \mathrm{~B}}}
$$

The intensity $\beta$ of a sound wave is a logarithmic measure given by

$$
\beta=(10 \mathrm{~dB}) \log \frac{I}{I_{0}}
$$

where $I_{0}$ is the reference intensity $\left(10^{-12} \mathrm{~W} / \mathrm{m}^{2}\right)$. The units of $\beta$ are decibels (dB).

## Standing Sound Waves

## Interference

|  |  |
| :--- | :--- |
| Beats | Ba |

## Beats

Standing sound waves that propagate in a fluid in a pipe can reflect and form longitudinal standing waves. The closed end of a pipe is a displacement node and a pressure antinode; the open end of a pipe is a displacement antinode and a pressure node. For a pipe of length $L$ with an open end, the fundamental frequency and harmonics are

$$
f_{n}=n \frac{v}{2 L}=n f_{1} \quad(n=1,2,3, \ldots)
$$

For a pipe of length $L$ with a closed end, the fundamental frequency and harmonics are

$$
f_{n}=n \frac{v}{4 L}=n f_{1} \quad(n=1,3,5, \ldots)
$$

When waves overlap in the same region of space, the waves are said to interfere. When the waves combine to form a wave with a larger amplitude, the waves interfere constructively, or reinforce one another. When the waves differ by a half cycle, their sum results in a wave with a smaller amplitude, and the waves interfere destructively, or cancel.

Beats are heard when two tones of slightly different frequencies are sounded together, creating a beat frequency that is the difference of the original two frequencies. The beat frequency given by

$$
f_{\text {beat }}=f_{a}-f_{b} .
$$

The Doppler effect is the frequency shift that occurs when the listener is in motion relative to the source of sound. The listener's frequency $f_{\mathrm{L}}$ is related to the source frequency $f_{\mathrm{s}}$ by

$$
f_{\mathrm{L}}=\frac{v+v_{\mathrm{L}}}{v+v_{\mathrm{S}}} f_{\mathrm{S}} .
$$

where $v$ is the speed of sound and $v_{\mathrm{L}}$ and $v_{\mathrm{S}}$ are the $x$ components of the speed of the listener and source, respectively.

## Conceptual Questions

## 1: Threshold of pain

By what factor must you amplify the intensity of a normal conversation to make it reach the threshold of pain?

## Solution

IDENTIFY, SET UP, AND EXECUTE The intensity of a normal conversation is 65 dB , and the intensity at the threshold of pain is 120 dB , according to Table 16.2 in the text. These two intensities differ by 55 dB , or $10^{5.5}=3.1 \times 10^{5}$. Therefore, you must amplify the normal conversation by $3.1 \times 10^{5}$ to reach the threshold of pain.

EVALUATE A normal conversation is five orders of magnitude less than the pain threshold. Remember this fact the next time your physics professor lectures: His lectures are not painful, because his voice hasn't reached the sound pain threshold!

## 2: Explaining Doppler shift

Your younger brother asks you to explain why the sound from train whistles changes from a high pitch to a low pitch when a train passes. How do you explain the change?

## Solution

IDENTIFY, SET UP, AND EXECUTE You first explain that sound comes from vibrations, or changing pressure. A high pitch comes from faster vibrations and a low pitch comes from slower vibrations. The train whistle produces a steady number of vibrations. When the train approaches, the vibrations become compressed, effectively increasing the number of vibrations that reach your ear per unit time. When the train leaves, the vibrations become expanded, effectively slowing the vibrations.

EVALUATE This description provides an alternative explanation of the Doppler shift. As the listener and source move toward each other, the wave fronts become closer together and the frequency increases.

## 3: An orchestra warming up

While you wait for an orchestral performance, you hear the musicians tuning their instruments. You observe that several musicians play the same note for a few seconds to check the tune. What are they doing?

## Solution

IDENTIFY, SET UP, AND EXECUTE The musicians are trying to play the same frequency when they tune their instruments. If one or more instruments vibrate at a slightly different frequency, the different frequencies interfere and produce beats. When tuning, the musicians listen for beats and readjust their instruments until the beats are removed.

EVALUATE You can listen for beats when you hear music. Beats may be intentional; for example, some pipe organs have a slow beat to create an undulating effect.

## 4: Frequencies in a water bottle

Estimate the two lowest frequencies that you can achieve by blowing across the top of a plastic water bottle.

## Solution

IDENTIFY, SET UP, AND EXECUTE We can treat the bottle as a pipe that is open at one end and closed at the other. The normal-mode frequencies of a closed-end pipe are given by

$$
f_{n}=\frac{n v}{4 L}
$$

where we want frequencies corresponding to $n=1$ and 3 . The water bottle is approximately $8^{\prime \prime}$, or 20.3 cm , long. Taking the speed of sound to be $344 \mathrm{~m} / \mathrm{s}$, we find that the two frequencies are 424 Hz and 1271 Hz .

EVALUATE Note that only the odd $n$ 's are valid for the closed-end pipe.
You can now build a pipe organ from recycled soda cans and soda bottles.

## Problems

## 1: Finding a plug in a tube

In an attempt to find where a plug is in a tube containing air, a plumber blows air across the opening of the tube and hears a resonance at a frequency of 80 Hz . If this is the fundamental mode, how far away from the end of the pipe is the plug? Take the velocity of sound in air to be $345 \mathrm{~m} / \mathrm{s}$.

## Solution

IDENTIFY We'll use the relationship between pipe length and normal-mode frequencies to find the distance the plug is from the end of the tube-the target variable.
SET UP We will use the normal-mode relationship to find the fundamental frequency of a closed pipe, since the plug effectively closes the pipe.

EXECUTE The fundamental frequency of the closed pipe is

$$
f_{1}=\frac{v}{4 L} .
$$

We know the frequency and speed of sound, so we solve for the length:

$$
L=\frac{v}{4 f_{1}}=\frac{(345 \mathrm{~m} / \mathrm{s})}{4(80 \mathrm{~Hz})}=1.08 \mathrm{~m}
$$

The plug is 1.08 from the end of the pipe.

EVALUATE This problem illustrates how to count overtones carefully and how to interpret integer results.

## 2: Overtones in a pipe

An open pipe of length 1.5 m is played on a day when the speed of sound in air is $345 \mathrm{~m} / \mathrm{s}$. How many overtones can be heard by a person with good hearing?

## Solution

IDENTIFY The number of overtones is the number of frequencies above the fundamental frequency. We'll use the relationship between pipe length and normal-mode frequencies to find the number of overtones-the target variable.

SET UP A person with good hearing can hear in the range from 20 Hz to $20,000 \mathrm{~Hz}$. We'll find the number of frequencies in that range for the pipe, and then we'll subtract the fundamental frequency to solve the problem.

EXECUTE The frequencies of standing waves in an open pipe are

$$
f_{n}=n f_{1},
$$

where $n$ is an integer. The fundamental frequency of the pipe is

$$
f_{1}=\frac{v}{2 L}
$$

For this pipe,

$$
f_{1}=\frac{v}{2 L}=\frac{(345 \mathrm{~m} / \mathrm{s})}{2(1.5 \mathrm{~m})}=115 \mathrm{~Hz}
$$

The fundamental frequency is above 20 Hz , so it can be heard. The highest frequency that can be heard by the human ear is $20,000 \mathrm{~Hz}$. This frequency corresponds to

$$
n=\frac{20,000 \mathrm{~Hz}}{f_{1}}=\frac{20,000 \mathrm{~Hz}}{115 \mathrm{~Hz}}=173.9 .
$$

Since we cannot hear nine-tenths of a frequency, we truncate $n$ to 173 . Thus, 173 frequencies can be heard: the fundamental frequency and 172 overtones.

EVALUATE This problem illustrates how to count overtones carefully and how to interpret integer results.

## 3: Making notes

A 1-meter-long tube open at one end and closed at the other contains water to a depth $d$. Assuming that the sound waves have a displacement node at the water surface and an antinode at the open end, find the depth of liquid that makes the tube resonate at middle $\mathrm{C}(264 \mathrm{~Hz})$ and one octave below middle C $(132 \mathrm{~Hz})$. Take the speed of sound in air to be $345 \mathrm{~m} / \mathrm{s}$.

## Solution

IDENTIFY We'll use the relationship between pipe length and normal-mode frequencies to find the distance the water is from the top of the tube. We'll subtract that distance from the length of the pipe to find the height of the water-the target variable.

SET UP We will use the normal-mode relationship to find the fundamental frequency of a closed pipe, since there is a displacement node at the water. We will find the lowest mode, that corresponding to $n=1$.

EXECUTE The fundamental frequency of a closed pipe is

$$
f_{1}=\frac{v}{4 L}
$$

where $L$ is the distance from the top of the tube to the top of the water. We solve for $L$ for the two frequencies:

$$
\begin{aligned}
L_{\text {middle C }} & =\frac{v}{4 f_{1}}=\frac{(345 \mathrm{~m} / \mathrm{s})}{4(264 \mathrm{~Hz})}=0.327 \mathrm{~m} \\
L_{\text {below C }} & =\frac{v}{4 f_{1}}=\frac{(345 \mathrm{~m} / \mathrm{s})}{4(132 \mathrm{~Hz})}=0.653 \mathrm{~m}
\end{aligned}
$$

The water must be at a height of $1.0 \mathrm{~m}-0.327 \mathrm{~m}=0.673 \mathrm{~m}$ for middle C and at a height of $1.0 \mathrm{~m}-$ $0.653 \mathrm{~m}=0.347 \mathrm{~m}$ for one octave below middle C .

EVALUATE This problem illustrates how to design a pipe organ that is made by filling several pipes of the same length with water. The pipe, however, would be a tough organ to tune, because the water will evaporate.

## 4: Power at a concert

You are given the task of determining how much power is needed in the sound system at the new stadium. There is a single set of speakers on the stage. If the design calls for an intensity of 100 dB at the farthest seats ( 120 m from the speakers), how much power is required?

## Solution

IDENTIFY AND SET UP In order to determine the intensity from the power, we assume that the sound is distributed over a sphere of radius 120 m . We'll use the definition of intensity to relate the design intensity to the power.

EXECUTE The intensity is given by

$$
\beta=(10 \mathrm{~dB}) \log \frac{I}{I_{0}}
$$

Taking the logarithm of both sides gives

$$
I=I_{0} 10^{(\beta / 10 \mathrm{~dB})}
$$

The intensity is the power per unit area, where the area in this case is that of a sphere $\left(4 \pi r^{2}\right)$. Combining the various equations gives

$$
P=A I=\left(4 \pi r^{2}\right)\left(I_{0} 10^{(\beta / 10 \mathrm{~dB})}\right)=\left(4 \pi(120 \mathrm{~m})^{2}\right)\left(\left(10^{-12} \mathrm{~W} / \mathrm{m}^{2}\right) 10^{((120 \mathrm{~dB}) / 10 \mathrm{~dB})}\right)=180 \mathrm{~kW}
$$

The required power is 180 kW .

EVALUATE Our stadium requires a substantial sound system in order for all visitors to hear the concert. The amount of power required would harm the hearing of those near the speaker. Stadiums are designed with multiple speakers placed around the stadium and closer to the visitors, to reduce the maximum volume.

Practice Problem: What is the intensity of the sound for persons seated 20 m from the speakers? Answer: 135 dB , above the threshold for permanent hearing damage.

## 5: Speed of approaching train

You are driving along a country road at $20.0 \mathrm{~m} / \mathrm{s}$. A train approaches on a rail that parallels the road. The train whistle blasts at a frequency of 800 Hz , but you hear a $950-\mathrm{Hz}$ whistle. What is the speed of the approaching train?

## Solution

IDENTIFY Our target variable is the speed of the approaching train.
SET UP The Doppler effect describes the frequency shift for moving sources, so we'll use the Doppler formula to determine the speed of the approaching train. Our coordinate system is shown in Figure 16.1; positive velocities are taken to be from the listener toward the source. Both the source and listener are moving. The listener's velocity is positive and the source's velocity is negative in our coordinate system. The speed of sound is taken to be $340 \mathrm{~m} / \mathrm{s}$.


Figure 16.1 Problem 5 sketch.
EXECUTE The listener's frequency is related to the source frequency by the Doppler shift,

$$
f_{\mathrm{L}}=\frac{v+v_{\mathrm{L}}}{v+v_{\mathrm{S}}} f_{\mathrm{S}}
$$

where $v$ is the speed of sound and $v_{L}$ and $v_{S}$ are the $x$ components of the speed of the listener and source, respectively. We can rearrange to solve for $v_{S}$ :

$$
v_{S}=\frac{v+v_{L}}{f_{L}} f_{S}-v
$$

In this case, $v_{L}$ is $+20.0 \mathrm{~m} / \mathrm{s}, f_{L}$ is $950 \mathrm{~Hz}, f_{S}$ is 800 Hz , and $v$ is $340 \mathrm{~m} / \mathrm{s}$. Substituting yields

$$
v_{\mathrm{S}}=\frac{v+v_{\mathrm{L}}}{f_{\mathrm{L}}} f_{\mathrm{S}}-v=\frac{(340 \mathrm{~m} / \mathrm{s})+(20.0 \mathrm{~m} / \mathrm{s})}{(950 \mathrm{~Hz})}(800 \mathrm{~Hz})-(340 \mathrm{~m} / \mathrm{s})=-36.8 \mathrm{~m} / \mathrm{s}
$$

The train is approaching at $36.8 \mathrm{~m} / \mathrm{s}$, or 132 kilometers per hour.
EVALUATE This problem illustrates how to use the Doppler shift to find the speed of an object. We expected and found a negative speed, indicating that the source was moving toward the listener. We see that proper Doppler-shift solutions require a coordinate system and careful interpretation of the directions of the velocities.

Practice Problem: What frequency would you hear if you were moving away from the train at $20.0 \mathrm{~m} / \mathrm{s}$ ? Answer: 844 Hz .

## Try It Yourself!

## 1: Designing an organ pipe

You are asked to design an organ pipe that will produce a middle C on the "even-tempered scale." Middle C is equivalent to a frequency of 261.6 Hz . (a) If the tube is open at both ends, how long should it be? (b) If the tube is open at one end and closed at the other, how long should it be? Take the speed of sound in air to be $345 \mathrm{~m} / \mathrm{s}$.

## Solution Checkpoints

IDENTIFY AND SET UP Use the normal-mode frequency relationships for open and closed tubes to find the lengths.

EXECUTE The fundamental frequency of an open pipe is

$$
f_{1}=\frac{v}{2 L}
$$

The fundamental frequency of a closed pipe is

$$
f_{1}=\frac{v}{4 L} .
$$

The two lengths are 0.659 m and 0.330 m .
EVALUATE Which pipe is more desirable for a compact pipe organ?

## 2: Designing a organ pipe, part 2

Find the allowed normal-mode frequencies of the two organ pipes in the previous problem. Take the speed of sound in air to be $345 \mathrm{~m} / \mathrm{s}$.

## Solution Checkpoints

IDENTIFY AND SET UP Use the normal-mode frequency relationships for open and closed tubes to find the frequencies.

EXECUTE The normal-mode frequencies of an open pipe are

$$
f_{n}=\frac{n v}{2 L}
$$

Substituting $n=1,2,3 \ldots$, we find that $f_{1}=261.6 \mathrm{~Hz}, f_{2}=523.2 \mathrm{~Hz}, f_{3}=784.8 \mathrm{~Hz}, \ldots$
The fundamental frequency of a closed pipe is

$$
f_{1}=\frac{v}{4 L}
$$

The normal-mode frequencies are integer multiples of the fundamental frequency; that is, $f_{n}=n f_{1}$. So, substituting $n=1,3,5 \ldots$ we find that $f_{1}=261.6 \mathrm{~Hz}, f_{2}=784.8 \mathrm{~Hz}, f_{3}=1308 \mathrm{~Hz}, \ldots$

EVALUATE Which pipe is more desirable for the number of frequencies it can produce?

## 3: Doppler-shift practice

Consider a source that produces a sound with a frequency of 500 Hz . If the speed of sound in air is $345 \mathrm{~m} / \mathrm{s}$, and the source and listener both move along the line joining them at speeds of $25 \mathrm{~m} / \mathrm{s}$, what frequencies can be heard by the listener for all possible directions of velocities?

## Solution Checkpoints

IDENTIFY AND SET UP Use the Doppler-shift equation to find the possible frequencies while varying the direction, or sign, of the velocities.

EXECUTE The listener's frequency is related to the source frequency by the Doppler shift,

$$
f_{\mathrm{L}}=\frac{v+v_{\mathrm{L}}}{v+v_{\mathrm{S}}} f_{\mathrm{S}}
$$

where $v$ is the speed of sound and $v_{L}$ and $v_{S}$ are the speeds of the listener and source, respectively. When the two velocities are in the same direction, either positive or negative, there is no Doppler shift and the listener hears a sound with a frequency of 500 Hz .

If the source and listener are moving away from each other, (say, $v_{L}$ is negative and $v_{S}$ is positive,) then the frequency is reduced to 432 Hz . If the source and listener are moving toward each other (say, $v_{L}$ is positive and $v_{S}$ is negative), then the frequency is increased to 579 Hz .

EVALUATE This problem illustrates the amount of Doppler shift for four possible combinations of velocity between two objects.

## 17 <br> Temperature and Heat

## Summary

In this chapter, we will begin a four-chapter investigation of thermodynamics. We lay the groundwork for the upcoming chapters with an initial definition of temperature, and then we see how materials change size with temperature. Heat will be introduced as a method of energy transfer due to temperature differences, and the rate of heat transfer will be calculated. We will also learn about the amount of heat required to change the phase of matter and about the three types of heat transfer: conduction, convection, and radiation.

## Objectives

After studying this chapter, you will understand

- The definition of temperature and thermal equilibrium.
- The three temperature scales and how to measure temperature.
- How thermal expansion describes the change in length and volume of materials due to temperature changes.
- About heat, phase changes, and calorimetry and how to apply these concepts to problems.
- How heat is transferred by conduction, convection, and radiation.


## Concepts and Equations

| Term |
| :--- |
| Thermal Equilibrium |
| Temperature Scales |
|  |
| Thermal Expansion and Thermal |
| Stress |

## Heat

Heat is energy transferred from one object to another due to changes in temperature. The quantity of heat $Q$ needed to raise the temperature of a mass $m$ of material by an amount $\Delta T$ is

$$
Q=m c \Delta T,
$$

where $c$ is the specific heat capacity of the material. The SI unit of heat capacity is the joule per kilogram per kelvin ( $\mathrm{J} /(\mathrm{kg} \mathrm{K})$ ).

## Phase Change

A phase transition is the change from one phase of matter to another. Phases include solid, liquid, and gas. The heat of fusion, $L_{f}$, is the heat per unit mass required to change a solid material to liquid. The heat of vaporization, $L_{v}$, is the heat per unit mass required to change a liquid material to gas. The heat of sublimation, $L_{s}$, is the heat per unit mass required to change a solid material to gas.

## Calorimetry

Calorimetry is the measurement of heat in a system. For an isolated system, the algebraic sum of the quantities of heat must add to zero:

$$
\sum Q=0 .
$$

## Heat Transfer

duction is the transfer of energy within a material without bulk motion of the material. Convection is the transfer of energy due to the motion of mass from one region to another. Radiation is the transfer of energy through electromagnetic waves. The heat current $H$ for an area $A$ and length $L$ through which the heat flows is given by

$$
H=\frac{d Q}{d t}=k A \frac{T_{\mathrm{H}}-T_{\mathrm{C}}}{L},
$$

where $T_{H}$ and $T_{C}$ are, respectively, the temperatures of the hot and cold sides of the material and $k$ is the thermal conductivity. The heat current $H$ due to radiation is

$$
H=A e \sigma T^{4}
$$

where $A$ is the surface area, $e$ is the emissivity of the surface (a pure number between 0 and 1 ), $T$ is the absolute temperature, and $\sigma$ is the Stefan-Boltzmann constant $\left(5.6705 \times 10^{-8} \mathrm{~W} / \mathrm{m}^{2} / \mathrm{K}^{4}\right)$.

## Conceptual Questions

## 1: Do holes expand or contract?

Your younger brother knows that solids expand as they heat up. He thinks that the metal surrounding the hole in a cookie sheet will expand into the hole as the cookie sheet heats up. Is he right or wrong?

## Solution

IDENTIFY, SET UP, AND EXECUTE Figure 17.1 shows a sketch of a cookie sheet with a hole in it. After considering the problem, you realize that the hole will enlarge as the cookie sheet heats up, since all dimensions of an object enlarge with temperature. The challenge is how to best explain this phenomenon to him.


Figure 17.1 Question 1 sketch.
If you give the problem a bit more thought, you come up with a convincing argument. If the cookie sheet had no hole, the whole sheet would increase with temperature. If you punch out a hole in the same cookie sheet and consider the piece of metal that was removed, this piece expands as its temperature rises. Therefore, the hole in the cookie sheet must also expand, just as it did when the cookie sheet was holeless.

EVALUATE Thermal expansion must be considered carefully. Here, we see that a confusing point can be clarified by imagining what happens to the piece that was once the hole.

CAUTION Holes expand when heated! Keep the results of this problem in mind when you encounter similar problems. Holes don't shrink when heated.

## 2: Cooler after a shower

When you step out of the shower, you often feel cold. After drying off, you feel warmer, even though the room's temperature is the same as when you stepped out of the shower. Why?

## Solution

IDENTIFY, SET UP, AND EXECUTE When you step out of the shower, water on your body evaporates. Evaporation requires heat energy (the heat of vaporization), much of which energy comes from heat leaving your body. You feel cold because your body is transferring its heat to evaporate the water. When you are dry, there is little heat lost due to evaporation.

EVALUATE Evaporation also explains why one feels cooler in a dry climate than in a humid climate: Your sweat evaporates more rapidly in a dry climate, taking away more heat, than in a humid climate.

## 3: Cold water versus cold air

Would you prefer to spend 10 minutes in a $40^{\circ} \mathrm{F}\left(4^{\circ} \mathrm{C}\right)$ room or in a $40^{\circ} \mathrm{F}$ pool? Why?

## Solution

IDENTIFY, SET UP, AND EXECUTE Both the room and the pool are at the same temperature, but the $40^{\circ} \mathrm{F}$ room would be much more comfortable. The reason is that the specific heat of air is much less than the specific heat of water (i.e., air will carry away less heat from your body than the water would in any time interval). Since the air carries away less heat, you are more comfortable in the room.

EVALUATE Specific heat is the amount of heat needed to change the temperature of a material per unit mass and per unit temperature. Larger specific heats mean that more heat is carried away from an object.

## Problems

## 1: Volume of a copper cup

A copper cup is filled to the brim with ethanol at $0^{\circ} \mathrm{C}$. When the cup and ethanol are heated to $35^{\circ} \mathrm{C}, 4.7 \mathrm{~cm}^{3}$ of ethanol spills from the cup. What is the initial volume of the cup?

## Solution

IDENTIFY We will use temperature expansion to find the change in volume of the cup and the ethanol. Their difference will lead to the initial volume of the cup-the target variable.

SET UP Both the cup and the ethanol expand as the temperature rises; the difference in their expansion is equal to the volume of the spilled ethanol. We'll apply the volume expansion equation to both the cup and the ethanol, setting their difference equal to the volume of the spill. The coefficient of volume expansion is $5.1 \times 10^{-5} / \mathrm{K}$ for copper and $75 \times 10^{-5} / \mathrm{K}$ for ethanol.

EXECUTE For any material, the change in volume due to temperature is

$$
\Delta V=\beta V_{0} \Delta T
$$

We are given the volume of the spill, which is the change in volume of the ethanol minus the change in volume of the cup:

$$
V_{\text {spill }}=\Delta V_{\text {ethanol }}-\Delta V_{\text {cup }} .
$$

The initial volumes of the cup and the ethanol are the same. We'll call their common volume $V_{0}$. The temperature of both materials is $35^{\circ} \mathrm{C}$. Replacing the changes in volumes yields

$$
V_{\text {spill }}=\beta_{\text {ethanol }} V_{0} \Delta T-\beta_{\text {copper }} V_{0} \Delta T .
$$

Solving for $V_{0}$, we obtain

$$
V_{0}=\frac{V_{\text {spill }}}{\left(\beta_{\text {ethanol }}-\beta_{\text {copper }}\right) \Delta T}=\frac{\left(4.7 \mathrm{~cm}^{3}\right)}{\left(\left(75 \times 10^{-5} / \mathrm{K}\right)-\left(5.1 \times 10^{-5} / \mathrm{K}\right)\right)\left(35^{\circ} \mathrm{C}\right)}=190 \mathrm{~cm}^{3}
$$

The original volume of the cup is $190 \mathrm{~cm}^{3}$.
EVALUATE This is a straightforward application of volume thermal expansion. We did need to note carefully that both the copper and the ethanol expanded, so the spillage was the difference in the changes in volumes. We could almost ignore the change in volume of the copper, since the coefficient of thermal expansion is much smaller for copper than for ethanol.

Practice Problem: What is the final volume of the cup? Answer: $190.3 \mathrm{~cm}^{3}$.

## 2: Stress in a wire

An aluminum wire is stretched across a large steel frame. Initially, the wire is at $20^{\circ} \mathrm{C}$ and is unstressed. The system (wire plus frame) is cooled by $50^{\circ} \mathrm{C}$. If the area of contact for the wire is $9.0 \times 10^{-6} \mathrm{~m}^{2}$, what force is exerted on the wire?

## Solution

IDENTIFY The differences in the expansion of the wire and frame will lead to tensile stress acting on the wire. The target variable is the force on the wire.

SET UP We will first use temperature expansion to find the change in lengths of the wire and frame. We will then find the tensile stress on the wire. The coefficient of linear expansion is $2.4 \times 10^{-5} / \mathrm{K}$ for aluminum and $1.2 \times 10^{-5} / \mathrm{K}$ for steel. Young's modulus for the aluminum is $0.7 \times 10^{11} \mathrm{~N} / \mathrm{m}^{2}$.

EXECUTE The percent change in length for a material due to temperature is

$$
\Delta L=\alpha L_{0} \Delta T
$$

When the aluminum is cooled by $50^{\circ} \mathrm{C}$, the change in its length is

$$
\left(\frac{\Delta L}{L_{0}}\right)_{\mathrm{a} 1}=\alpha_{\mathrm{al}} \Delta T=\left(2.4 \times 10^{-5}\right)(50)=1.2 \times 10^{-3}
$$

For the steel, the change in length is

$$
\left(\frac{\Delta L}{L_{0}}\right)_{\mathrm{st}}=\alpha_{\mathrm{st}} \Delta T=\left(1.2 \times 10^{-5}\right)(50)=0.60 \times 10^{-3}
$$

Both of these changes are decreases, since the temperature has decreased. We see the aluminum changes more than the steel. Because the steel decreases less, stress is induced in the aluminum. The stress is given by

$$
\frac{F}{A}=Y_{\mathrm{al}}\left(\frac{\Delta L}{L_{0}}\right)
$$

The stress is proportional to the net change in length of the wire. The frame shrinks, so the net change in length is the difference in the changes of the wire and the frame. The force is then

$$
F=A Y_{\mathrm{al}}\left(\frac{\Delta L}{L_{0}}\right)=\left(9.0 \times 10^{-6} \mathrm{~m}^{2}\right)\left(0.7 \times 10^{11} \mathrm{~N} / \mathrm{m}^{2}\right)\left(1.2 \times 10^{-3}-0.60 \times 10^{-3}\right)=380 \mathrm{~N} .
$$

The stress on the wire is 380 N .
EVALUATE This is an application of linear thermal expansion. To find the stress, we did need to note carefully the difference in how the two materials contracted. We could also have used the relation in the book that directly provides the stress as a function of temperature.

## 3: Ice to steam

A copper calorimeter of mass 2.0 kg initially contains 1.5 kg of ice at $-10^{\circ} \mathrm{C}$. How much heat energy must be added to convert all of the ice to water and then half of the water into steam?

## Solution

IDENTIFY We will use heat capacity and heat of fusion to determine how the ice melts and turns to steam. The target variable is the amount of heat needed to convert the ice to water and the water to steam.

SET UP We will solve the problem in several steps. We'll first find the heat required to raise the temperature of the ice to $0^{\circ} \mathrm{C}$, then find the heat required to melt the ice, then find the heat required to warm the water to $100^{\circ} \mathrm{C}$, and finally find the heat required to vaporize half of the water. We know that the final temperature of the remaining water must be $100^{\circ} \mathrm{C}$, since the water remains in equilibrium throughout the process. We must also add the heat required to heat the copper pot to $100^{\circ} \mathrm{C}$.

EXECUTE The heat required to heat the ice to $0^{\circ} \mathrm{C}$ is

$$
Q_{1}=m_{\mathrm{ice}} c_{\mathrm{ice}} \Delta T=(1.5 \mathrm{~kg})(2100 \mathrm{~J} / \mathrm{kg} / \mathrm{K})\left(10.0^{\circ} \mathrm{C}\right)=31,500 \mathrm{~J}
$$

We used the specific heat of ice $(2010 \mathrm{~J} / \mathrm{kg} / \mathrm{K})$ to find $Q_{1}$. The heat required to melt the ice is the heat of fusion for ice:

$$
Q_{2}=m_{\mathrm{ice}} L_{\mathrm{f}}=(1.5 \mathrm{~kg})\left(3.34 \times 10^{5} \mathrm{~J} / \mathrm{kg}\right)=501,000 \mathrm{~J}
$$

The melted ice must warm to the final temperature $\left(100.0^{\circ} \mathrm{C}\right)$ :

$$
Q_{3}=m_{\text {ice }} c_{\text {water }} \Delta T=m_{\text {ice }}(4190 \mathrm{~J} / \mathrm{kg} / \mathrm{K})\left(100.0^{\circ} \mathrm{C}\right)=628,500 \mathrm{~J}
$$

Here, we used the heat capacity of water $(4190 \mathrm{~J} / \mathrm{kg} / \mathrm{K})$. Half of the mass turns to steam. Using the heat of vaporization $\left(2.256 \times 10^{6} \mathrm{~J} / \mathrm{kg}\right)$, we find that the heat required is

$$
Q_{4}=\frac{1}{2} m_{\mathrm{ice}} L_{\mathrm{f}}=\frac{1}{2}(1.5 \mathrm{~kg})\left(2.256 \times 10^{6} \mathrm{~J} / \mathrm{kg}\right)=1,692,000 \mathrm{~J} .
$$

The copper pot also increases in temperature, from $-10^{\circ} \mathrm{C}$ to $100^{\circ} \mathrm{C}$. The heat required to bring about this increase is

$$
Q_{\text {copper }}=m_{\text {copper }} c_{\text {copper }} \Delta T=(2.0 \mathrm{~kg})(390 \mathrm{~J} / \mathrm{kg} / \mathrm{K})\left(110.0^{\circ} \mathrm{C}\right)=85,800 \mathrm{~J} .
$$

The total heat is the sum of the five quantities of heat:

$$
Q_{\mathrm{t}}=Q_{1}+Q_{2}+Q_{3}+Q_{4}+Q_{\text {copper }}=2.94 \times 10^{6} \mathrm{~J}
$$

So $2.94 \times 10^{6} \mathrm{~J}$ are needed to heat the pot and ice from $-10^{\circ} \mathrm{C}$ to $100^{\circ} \mathrm{C}$ and then to vaporize half of the water that is produced.

EVALUATE We see how we must solve calorimetry problems in multiple steps. We need to include the latent heat when materials change phase and the heat capacity when materials heat up.

CAUTION Temperature is not heat! Temperature characterizes the state of an object. Heat is the flow of energy. The two terms may be synonymous in everyday language, but they are not synonymous in physics.

## 4: Cooling hot tea

You wish to chill your freshly brewed tea with the minimum amount of ice that will avoid watering it down too much. What is the minimum amount of ice you should add to 2.0 kg of freshly brewed tea at $95^{\circ} \mathrm{C}$ to cool it to $5.0^{\circ} \mathrm{C}$ ? The ice is initially at a temperature of $-5.0^{\circ} \mathrm{C}$.

## Solution

IDENTIFY We'll set the heat lost from the tea equal to the heat gained by the ice. The target variable is the amount of ice needed to cool the tea.

SET UP The amount of heat lost by the tea is given by the specific heat capacity equation, since the tea doesn't go through a phase change. The ice melts, so, in calculating the heat gain of the ice, we need to include the latent heat of fusion, plus the changes due to the ice warming to $0^{\circ} \mathrm{C}$, and the changes due to the melted ice warming to $5.0^{\circ} \mathrm{C}$.

EXECUTE The heat transfer from the hot tea as it cools to $5.0^{\circ} \mathrm{C}$ is negative:

$$
Q_{\text {tea }}=m_{\text {tea }} c_{\text {water }} \Delta T_{\text {tea }}=(2.0 \mathrm{~kg})(4190 \mathrm{~J} / \mathrm{kg} / \mathrm{K})\left(5.0^{\circ} \mathrm{C}-95^{\circ} \mathrm{C}\right)=-754,000 \mathrm{~J}
$$

Here, we used the heat capacity of water $(4190 \mathrm{~J} / \mathrm{kg} / \mathrm{K})$ for the tea. The ice must warm to $0^{\circ} \mathrm{C}$, then melt, and then heat to $5.0^{\circ} \mathrm{C}$. We find the heat required for each segment of the ice warming. For the ice to heat to $0^{\circ} \mathrm{C}$, we use the specific heat of ice $(2010 \mathrm{~J} / \mathrm{kg} / \mathrm{K})$ :

$$
Q_{\text {ice }}=m_{\text {ice }} \mathrm{i}_{\text {ice }} \Delta T_{\text {ice }}=m_{\text {ice }}(2010 \mathrm{~J} / \mathrm{kg} / \mathrm{K})\left[0.0^{\circ} \mathrm{C}-\left(-5.0^{\circ} \mathrm{C}\right)\right]=m_{\text {ice }}(10,000 \mathrm{~J} / \mathrm{kg}) .
$$

The heat needed to melt the ice is the heat of fusion for ice:

$$
Q_{\mathrm{melt}}=m_{\mathrm{ice}} L_{\mathrm{f}}=m_{\text {ice }}\left(3.34 \times 10^{5} \mathrm{~J} / \mathrm{kg}\right)
$$

The melted ice must warm to the final temperature $\left(5.0^{\circ} \mathrm{C}\right)$ :

$$
Q_{\text {melted ice }}=m_{\text {ice }} e_{\text {water }} \Delta T_{\text {melted ice }}=m_{\text {ice }}(4190 \mathrm{~J} / \mathrm{kg} / \mathrm{K})\left(5.0^{\circ} \mathrm{C}-0.0^{\circ} \mathrm{C}\right)=m_{\text {ice }}(21,000 \mathrm{~J} / \mathrm{kg})
$$

The sum of these four quantities must be zero:

$$
\begin{aligned}
Q_{\text {tea }}+Q_{\text {ice }}+Q_{\text {melt }}+Q_{\text {melted ice }}= & -754,000 \mathrm{~J}+m_{\text {ice }}(10,000 \mathrm{~J} / \mathrm{kg}) \\
& +m_{\text {ice }}(334,000 \mathrm{~J} / \mathrm{kg})+m_{\text {ice }}(21,000 \mathrm{~J} / \mathrm{kg})=0 .
\end{aligned}
$$

Therefore,

$$
m_{\mathrm{ice}}=\frac{754,000 \mathrm{~J}}{(10,000 \mathrm{~J} / \mathrm{kg})+(334,000 \mathrm{~J} / \mathrm{kg})+(21,000 \mathrm{~J} / \mathrm{kg})}=2.1 \mathrm{~kg}
$$

It takes a minimum of 2.1 kg of ice to cool the tea down.
evaluate Despite your best effort, the tea will be watery. Putting the ice in a bag will prevent the melted ice water from mixing with the tea. More importantly, we see how we must proceed stepwise through calorimetry problems.

## 5: Heat flow through three bars

A composite rod is made up of three equal lengths and cross sections of aluminum, brass, and copper. The free aluminum end is maintained at $100^{\circ} \mathrm{C}$, and the free end of the copper rod is maintained at $0^{\circ} \mathrm{C}$. If the surface of the rod is insulated to prevent radial heat flow, find the temperature at each junction.

## Solution

IDENTIFY We will use the heat current through the rods to find the temperatures at the rod junctionsthe target variables.

SET UP A sketch of the rod is shown in Figure 17.2. The heat current is the same through each segment of the rod, so we'll write the heat equations for each segment and set them equal to each other to solve the problem.


Figure 17.2 Problem 5 sketch.
EXECUTE The heat current through any rod is

$$
H=k A \frac{T_{H}-T_{C}}{L}
$$

We can write the heat current per unit area times the length through the individual rods as

$$
\begin{aligned}
\frac{H_{A B} L}{A} & =k_{\mathrm{Al}}\left(T_{H}-T_{C}\right)=(205 \mathrm{~W} / \mathrm{m} \cdot \mathrm{~K})\left(100^{\circ} \mathrm{C}-T_{B}\right) \\
\frac{H_{B C} L}{A} & =k_{\mathrm{Br}}\left(T_{H}-T_{C}\right)=(109 \mathrm{~W} / \mathrm{m} \cdot \mathrm{~K})\left(T_{B}-T_{C}\right) \\
\frac{H_{C D} L}{A} & =k_{\mathrm{Co}}\left(T_{H}-T_{C}\right)=(385 \mathrm{~W} / \mathrm{m} \cdot \mathrm{~K})\left(T_{C}-0^{\circ} \mathrm{C}\right)
\end{aligned}
$$

These expressions must all be equal, since they each reference the same heat current, area, and length of the rod segment. Setting the last two equations equal to each other results in

$$
(109 \mathrm{~W} / \mathrm{m} \cdot \mathrm{~K}) T_{B}-(109 \mathrm{~W} / \mathrm{m} \cdot \mathrm{~K}) T_{C}=(385 \mathrm{~W} / \mathrm{m} \cdot \mathrm{~K}) T_{C},
$$

or

$$
T_{B}=\frac{385+109}{109} T_{C}=4.53 T_{C}
$$

Setting the first and last equations equal to each other gives

$$
(205 \mathrm{~W} / \mathrm{m} \cdot \mathrm{~K})\left(100^{\circ} \mathrm{C}\right)-(205 \mathrm{~W} / \mathrm{m} \cdot \mathrm{~K}) T_{B}=(385 \mathrm{~W} / \mathrm{m} \cdot \mathrm{~K}) T_{C} .
$$

Replacing $T_{B}$ yields

$$
\begin{gathered}
205\left(100^{\circ} \mathrm{C}\right)-205\left(4.53 T_{C}\right)=385 T_{C}, \\
T_{C}=15.6^{\circ} \mathrm{C}
\end{gathered}
$$

Then

$$
T_{B}=70.7^{\circ} \mathrm{C}
$$

The aluminum-brass junction is at $70.7^{\circ} \mathrm{C}$ and the brass/copper junction is at $15.6^{\circ} \mathrm{C}$.
EVALUATE Although one might expect the three equal segments to have equal temperature differences, we see that the varying thermal conductivities of the segments caused a nonuniform temperature distribution. The segment with the highest thermal conductivity (copper) had the smallest temperature difference between its ends, and the segment with the lowest thermal conductivity (brass) had the greatest temperature difference between its ends.

## 6: Time required to melt a block of ice

A long steel rod that is insulated to prevent heat loss along its sides is in perfect thermal contact with a large container of boiling water at one end and a $3.0-\mathrm{kg}$ block of ice at the other. The steel rod is 1.2 m long with cross-sectional area $3.50 \mathrm{~cm}^{2}$. How long does it take for the block of ice to melt? The ice block is initially at $0^{\circ} \mathrm{C}$.

## Solution

IDENTIFY We'll combine our knowledge of heat conduction with our knowledge of heat of fusion to solve this problem. The target variable is the time required for the ice to melt.
SET UP A sketch of the problem is shown in Figure 17.3. We begin by determining the heat required to melt the ice. We then find the rate of heat flow into the ice. With that information, we can find the time it takes for the ice to melt.


Figure 17.3 Problem 6 sketch.
EXECUTE The heat required to melt the ice is the heat of fusion for ice:

$$
Q_{\mathrm{melt}}=m_{\mathrm{ice}} L_{f}=(3.0 \mathrm{~kg})\left(3.34 \times 10^{5} \mathrm{~J} / \mathrm{kg}\right)=1.0 \times 10^{6} \mathrm{~J}
$$

The rate of heat flow is given by

$$
H=\frac{\Delta Q}{\Delta t}=k A \frac{T_{H}-T_{C}}{L}
$$

where $k$ is the thermal conductivity, $A$ and $L$ are, respectively, the area and length of the bar, and $T_{H}$ and $T_{C}$ are, respectively, the temperatures of the hot and cold sides of the bar. We find that

$$
H=\frac{\Delta Q}{\Delta t}=k A \frac{T_{H}-T_{C}}{L}=(50.2 \mathrm{~W} /(\mathrm{m} \cdot \mathrm{~K}))\left(6.5 \times 10^{-4} \mathrm{~m}^{2}\right) \frac{\left(100^{\circ} \mathrm{C}\right)-\left(0^{\circ} \mathrm{C}\right)}{(1.2 \mathrm{~m})}=2.7 \mathrm{~W},
$$

where we used $50.2 \mathrm{~W} / \mathrm{m} / \mathrm{K}$ as the thermal conductivity of steel. The time required to melt the ice is

$$
\Delta t=\frac{Q_{\text {melt }}}{\mathrm{H}}=\frac{\left(10^{6} \mathrm{~J}\right)}{(2.7 \mathrm{~W})}=370,000 \mathrm{~s}
$$

The time required to melt the ice is $370,000 \mathrm{~s}$, or 103 hours.
EVALUATE We see that the thin steel bar is a relatively poor conductor of heat. Replacing the steel with a copper bar would increase the rate by almost a factor of 8 , due to the differences in thermal conductivity. Increasing the rod's diameter and shortening the rod would also increase the rate of melting.

## Try It Yourself!

## 1: Niagara Falls

Water flowing at a speed of $5.0-\mathrm{m} / \mathrm{s}$-falls over a $50-\mathrm{m}$-high waterfall into a still pool below. Calculate the approximate rise in water temperature due to the conversion of mechanical energy into thermal energy

## Solution Checkpoints

IDENTIFY AND SET UP Use energy conservation, equating the loss in mechanical energy to heat.
EXECUTE The water has kinetic energy and gravitational energy that together convert to heat. Equating the two forms of energy

$$
\frac{1}{2} m v^{2}+m g h=m c \Delta T .
$$

The rise in temperature is $0.12^{\circ} \mathrm{C}$.
EVALUATE Did you need to know the mass of the water?

## 2: Melting ice, again

A copper calorimeter of mass 2.0 kg initially contains 1.5 kg of ice at $-10^{\circ} \mathrm{C}$. (a) What will the final temperature be if the heat added is $5 \times 10^{5} \mathrm{~J}$ ? (b) What will the final temperature be if the heat added is $10^{6} \mathrm{~J}$ ?

## Solution Checkpoints

IDENTIFY AND SET UP Use heat capacity and heat of fusion to determine how the temperature of the ice increases with the heat provided. Follow a series of steps and determine whether the heat at each step exceeds the heat provided.

EXECUTE (a) Examining Problem 4, we see that $5 \times 10^{5} \mathrm{~J}$ would be exhausted during the melting phase of the problem. The final temperature is $0^{\circ} \mathrm{C}$.
(b) With $10^{6} \mathrm{~J}$ of heat, all of the ice melts but the temperature doesn't reach $100^{\circ} \mathrm{C}$. The temperature can be found by adding up the heat required in each step:

$$
\Delta Q=m_{\mathrm{ice}} c_{\mathrm{ice}}\left(10^{\circ} \mathrm{C}\right)+m_{\text {copper }} c_{\text {copper }} \Delta T+m_{\text {ice }} c_{\text {water }}\left(\Delta T-10^{\circ} \mathrm{C}\right)+m_{\text {ice }} L_{f}
$$

The final temperature is $64.8^{\circ} \mathrm{C}$.
EVALUATE Why is $-10^{\circ} \mathrm{C}$ required in the term expressing the heat capacity of water?

## Thermal Properties of Matter

## Summary

In this chapter, we extend our investigation into thermodynamics, viewing systems from both the macroscopic and microscopic perspectives and building links between the two perspectives. We will learn about equations of state for materials and examine the ideal-gas equation as one such equation. This investigation will allow us to build a model for the kinetic energy of individual molecules and predict the behavior of gases. We will define thermodynamic systems and examine the energy of those systems. This analysis will lead to the first law of thermodynamics and thermodynamic processes. Four common thermodynamics processes will be highlighted, and their implications for ideal gases examined.

## Objectives

After studying this chapter, you will understand

- How to define the mole and Avogadro's number.
- How to define equations of state and how to apply the ideal-gas equation.
- How to determine the kinetic energy of gases and how to apply that energy to individual particles.
- The origins of molar heat capacities for materials and gases.
- How to apply the first law of thermodynamics.
- The four common thermodynamic processes and how to apply them to find the changes in heat, work, and internal energy of thermodynamic systems.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Mole | One mole (mol) is the amount of substance that contains the same number of <br> elementary units as there are atoms in 0.012 kg of carbon 12. The number of <br> molecules in a mole is Avogadro's number, $N_{A}=6.022 \times 10^{23} \mathrm{molecules} \mathrm{per}$ <br> mole. The molar mass is the mass of 1 mole of a substance. |
| Equation of State | An equation of state expresses the relation among pressure, temperature, and <br> volume of a certain amount of a substance in equilibrium. The pressure $p$, <br> volume $V$, and absolute temperature $T$ are the state variables. |
| Ideal-Gas Equation | The ideal-gas equation is the equation of state for an ideal gas that approxi- <br> mates the behavior of a real gas at a low pressure and a high temperature. The <br> pressure $p$, temperature $T$, volume $V$, and number of moles, $n$, of the gas are <br> related by |
|  | $\quad p V=n R T$, |
| where $R$ is the ideal-gas constant. In SI units, when pressure is given in Pa |  |
| and volume is given in $\mathrm{m}^{3}, R=8.3145 \mathrm{~J} /($ mol $\cdot \mathrm{K})$. |  |

For a single molecule, the average translational kinetic energy is

$$
K_{\mathrm{av}}=\frac{3}{2} k T \text {, }
$$

where $k=R / N_{\mathrm{A}}=1.381 \times 10^{-23} \mathrm{~J} /($ molecule $\cdot \mathrm{K})$ is the Boltzmann constant. The mean free path of molecules in an ideal gas is given by

$$
\lambda=v t_{\text {mean }}=\frac{V}{4 \pi \sqrt{2} r^{2} N} .
$$

| Molar Heat Capacity | The amount of heat $Q$ needed for a temperature change $\Delta T$ is |
| :--- | :---: |
| $Q=n C \Delta T$, |  |

where $n$ is the number of moles of the substance and $C$ is the molar heat capacity. The molar heat capacity at constant volume is given in certain cases by

$$
\begin{array}{cl}
C_{V}=\frac{3}{2} R & \text { (monatomic gas) } \\
C_{V}=\frac{5}{2} R & \text { (diatomic gas) } \\
C_{V}=3 R & \text { (monatomic solid). }
\end{array}
$$

| Molecular Speeds | The speeds of molecules in an ideal gas are given by the Maxwell-Boltzmann |
| :--- | :--- | distribution:

$$
f(v)=4 \pi\left(\frac{m}{2 \pi k T}\right)^{3 / 2} v^{2} e^{-m v^{2} / 2 k T} .
$$

The quantity $f(v) d v$ describes the fraction of molecules with speeds between $v$ and $v+d v$.
Phases of Matter
Ordinary matter exists in solid, liquid, and gas phases. A phase diagram shows the conditions under which two phases can coexist in phase equilibrium. All three phases can coexist at the triple point.

## Conceptual Questions

## 1: Don't hold your breath

Explain why scuba divers are taught not to hold their breath as they ascend to the surface from depths under the water.

## Solution

IDENTIFY, SET UP, AND EXECUTE We know from fluid statics that pressure increases with depth in water. The ideal-gas equation states that pressure and volume are inversely proportional for a given temperature and quantity of gas. Ascending to the surface reduces the ambient pressure, causing an increase in volume. (We assume that the temperature is constant in the water). By holding her breath, a scuba diver traps a quantity of air inside her lungs. As the pressure decreases upon her ascent, her lungs expand, possibly damaging some lung tissue. If the diver exhales during the ascent, the pressure cannot build to dangerous levels.

EVALUATE This problem combines our knowledge of fluid statics and our knowledge of ideal gases and helps illustrate the relation between pressure and volume. High-altitude weather balloons also expand as they rise, so they are partially filled at the ground to prevent the balloons from bursting as they ascend.

## 2: Atmosphere on the earth and moon

Why does the earth, but not the moon, have an atmosphere?

## Solution

IDENTIFY, SET UP, AND EXECUTE The escape velocity of molecules on the earth is about $11 \mathrm{~km} / \mathrm{s}$, much higher than the average rms speed of molecules in the atmosphere. Without sufficient speed, the molecules remain near earth, thus creating an atmosphere.

The gravitational potential on the moon's surface is about 20 times weaker than that on the earth's surface, so the escape speed is about 20 times less, or $2400 \mathrm{~m} / \mathrm{s}$. This lesser escape speed greatly enhances the probability that molecules of whatever atmosphere the moon might have had have all escaped into space.

EVALUATE The problem illustrates how molecular motion can help explain common physics phenomena.

## Problems

## 1: Changing volume in a diving bell

A diving bell (a circular cylinder 3.0 m high, open at the bottom) is lowered into a lake. By how much does the water rise as the bell is lowered 75 m ? The surface temperature of the lake is $25^{\circ} \mathrm{C}$ and the temperature at the 75 m depth is $15^{\circ} \mathrm{C}$.

## Solution

IDENTIFY We assume that the gas is ideal, so we use the ideal-gas equation to relate the surface values of pressure, temperature, and volume to the values at depth. The target value is the height of the water in the diving bell at depth.

SET UP Figure 18.1 shows a diagram of the situation. We'll use fluid statics to relate the pressure at the surface to the pressure at depth. These two relations will be combined to find the final height of water in the bell.


Figure 18.1 Problem 1 sketch.
EXECUTE The same amount of gas is trapped inside the bell both at the surface and at depth; therefore,

$$
\frac{p_{S} V_{S}}{T_{S}}=\frac{p_{D} V_{D}}{T_{D}}=\text { constant }
$$

where the subscript $S$ indicates the value at the surface and subscript $D$ indicates the value at depth. Substituting the known values gives

$$
\frac{\left(1.01 \times 10^{5} \mathrm{~Pa}\right)(A)(3.0 \mathrm{~m})}{(298 \mathrm{~K})}=\frac{p_{D} A l}{(288 \mathrm{~K})}
$$

where we replaced the volume of the cylinder with its area times its height. There two unknowns in this equation:

$$
p_{D} l=(288 \mathrm{~K}) \frac{\left(1.01 \times 10^{5} \mathrm{~Pa}\right)(3.0 \mathrm{~m})}{(298 \mathrm{~K})}=2.93 \times 10^{5} \mathrm{~Pa} \mathrm{~m} .
$$

We can find the pressure at depth from fluid statics, using

$$
p_{D}=p_{S}+\rho g(D+l)
$$

Substituting the known values yields

$$
\begin{aligned}
& p_{D}=\left(1.01 \times 10^{5} \mathrm{~Pa}\right)+\left(1 \times 10^{3} \mathrm{~kg} / \mathrm{m}^{3}\right)\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(72 \mathrm{~m}+l) \\
& p_{D}=\left(8.07 \times 10^{5} \mathrm{~Pa}\right)+\left(7.06 \times 10^{5} \mathrm{~Pa} / \mathrm{m}\right) l
\end{aligned}
$$

This equation also has two unknowns. Combining the two equations to eliminate $p_{D}$ results in

$$
p_{\mathrm{D}}=\frac{\left(2.93 \times 10^{5} \mathrm{~Pa} \mathrm{~m}\right)}{l}=\left(8.07 \times 10^{5} \mathrm{~Pa}\right)+\left(7.06 \times 10^{5} \mathrm{~Pa} / \mathrm{m}\right) l .
$$

Rearranging terms gives

$$
\left(7.06 / \mathrm{m}^{2}\right) l^{2}+(8.07 / \mathrm{m}) l-2.93=0
$$

This is a quadratic equation with solutions $l=0.290,-1.43 \mathrm{~m}$. The negative root is nonphysical, so the correct $l$ is 0.29 m . The water rises $3.0 \mathrm{~m}-0.29 \mathrm{~m}$, or 2.7 m , as the bell descends.

EVALUATE This problem illustrates how the increased pressure at depth reduces the volume of gas in the diving bell. If you consider the reverse process, you can see how the volume would increase as the bell rises to the surface, as we discussed in Conceptual Question 1. You can also try both situations by using a bucket of water. Submerge an inverted glass in a bucket of water, and see how the water level in the glass rises as the glass is lowered. Then use a hose to add air to the bottom of an inverted glass at the bottom of the bucket. As you raise the glass, you should see air leaving it.

## 2: Spacing between hydrogen molecules

Find the average spacing between $\mathrm{H}_{2}$ molecules, assuming that the molecules are at the vertices of a fictitious cubic structure, in (a) gaseous $\mathrm{H}_{2}$ at STP, (b) liquid $\mathrm{H}_{2}$ at 20 K , where the density is $41,060 \mathrm{~mol} / \mathrm{m}^{3}$, and (c) solid $\mathrm{H}_{2}$ at 4.2 K , where the molar volume is $22.91 \times 10^{-6} \mathrm{~m}^{3} / \mathrm{mol}$.

## Solution

IDENTIFY We will use the definitions of density, Avogadro's number, and molar mass to find the average spacing.

SET UP We'll assign a spacing of $a$ to the distance between $\mathrm{H}_{2}$ molecules, giving a volume of $a^{3}$ per molecule. One mole of gas occupies 22.4 L at STP, and 1 L is $0.001 \mathrm{~m}^{3}$.

EXECUTE (a) For the gaseous $\mathrm{H}_{2}$, we multiply the volume of one molecule by $N_{A}$ to get the volume at STP:

$$
N_{A} a^{3}=22.4 \mathrm{~L} / \mathrm{mol}=22.4 \times 10^{-3} \mathrm{~m}^{3} / \mathrm{mol}
$$

Solving for $a$ gives

$$
a=\sqrt[3]{\frac{22.4 \times 10^{-3} \mathrm{~m}^{3} / \mathrm{mol}}{N_{A}}}=\sqrt[3]{\frac{22.4 \times 10^{-3} \mathrm{~m}^{3} / \mathrm{mol}}{6.203 \times 10^{23}}}=3.34 \times 10^{-9} \mathrm{~m}
$$

(b) We are given that each cubic meter contains 41,060 moles of liquid $\mathrm{H}_{2}$. In equation form, this is

$$
1 \mathrm{~m}^{3}=41,060 N_{A} a^{3}
$$

Solving for $a$ results in

$$
a=\sqrt[3]{\frac{1 \mathrm{~m}^{3}}{41,060 N_{A}}}=\sqrt[3]{\frac{1 \mathrm{~m}^{3}}{(41,060)\left(6.203 \times 10^{23}\right)}}=3.44 \times 10^{-10} \mathrm{~m}
$$

(c) We are given the molar volume:

$$
N_{A} a^{3}=22.91 \times 10^{-6} \mathrm{~m}^{3} / \mathrm{mol}
$$

Solving for $a$ produces

$$
a=\sqrt[3]{\frac{22.91 \times 10^{-6} \mathrm{~m}^{3} / \mathrm{mol}}{N_{A}}}=\sqrt[3]{\frac{22.91 \times 10^{-6} \mathrm{~m}^{3} / \mathrm{mol}}{6.203 \times 10^{23}}}=3.37 \times 10^{-10} \mathrm{~m}
$$

The average spacing is $3.34 \times 10^{-9} \mathrm{~m}$ for the gaseous $\mathrm{H}_{2}, 3.44 \times 10^{-10} \mathrm{~m}$ for the liquid $\mathrm{H}_{2}$, and $3.37 \times 10^{-10} \mathrm{~m}$ for the solid $\mathrm{H}_{2}$.

EVALUATE We see that the average spacing for the liquid and solid $\mathrm{H}_{2}$ is similar and about 10 times closer than the spacing for the gaseous $\mathrm{H}_{2}$.

## 3: Spacing in a vacuum

A common type of laboratory vacuum pump produces an ultimate pressure of 10 microns. How many molecules (in moles) of gas are present in a volume of $0.15 \mathrm{~m}^{3}$ reduced to that pressure at 300 K ?

## Solution

IDENTIFY We will use the ideal-gas law to solve for the number of molecules in the volume-the target variable.

SET UP To use the ideal-gas law, we need to know the pressure in standard units. One micron is a measure of pressure in terms of the height of mercury. One atmosphere is 760 mm of mercury; one micron is one-thousandth of a millimeter of mercury.

EXECUTE First convert the pressure to atmospheres, using the information provided:

$$
10 \text { microns }=\left(10 \times 10^{-6}\right)\left(\frac{1 \mathrm{~atm}}{0.760 \mathrm{~m}}\right)=1.32 \times 10^{-5} \mathrm{~atm}
$$

The ideal-gas law is

$$
p V=n R T
$$

Solving for $n$ and using the molar gas constant yields

$$
n=\frac{p V}{R T}=\frac{\left(1.32 \times 10^{-5} \mathrm{~atm}\right)\left(1.5 \times 10^{2} \mathrm{~m}^{3}\right)}{(0.08206 \mathrm{~L} \cdot \mathrm{~atm} / \mathrm{mol} \cdot \mathrm{~K})(300 \mathrm{~K})}=8.04 \times 10^{-5} \mathrm{~mol}
$$

There is $8.04 \times 10^{-5}$ of a mole of molecules in the container after evacuating.
EVALUATE This may seem like a small number, but there are almost $5 \times 10^{19}$ molecules remaining in the volume. One of the toughest challenges in physics research is creating ultrahigh vacuums to study the behavior of small numbers of particles. Physicists don't want collisions with remaining air molecules to interfere with the molecules whose behavior they are studying.

Did we assume that the gas was an ideal gas? Yes, we took the gas to be an ideal gas. This assumption is valid because the gas is at low pressure.

## 4: Mixing gases

A 1-liter flask at 293 K contains a mixture of 3 g of $\mathrm{N}_{2}$ and 3 g of $\mathrm{H}_{2}$ gas. Assuming that the gases behave as ideal gases, (a) calculate the partial pressure exerted by both gases and (b) calculate the rms speeds of the two gases.

## Solution

IDENTIFY We will use the ideal-gas law and kinetic theory to solve the problem. The target variables are the partial pressures and rms speeds of the two gases.

SET UP Partial pressure is the pressure exerted by each gas separately. We'll use the ideal-gas law to find the partial pressure of each gas. The rms speed may be calculated by an expression found in the text.

EXECUTE (a) The volume and temperature are given, so we need to convert the amount of each gas to moles. We have

$$
\begin{aligned}
& n_{\mathrm{N}_{2}}=\frac{3 \mathrm{~g}}{28 \mathrm{~g} / \mathrm{mol}}=\frac{3}{28} \mathrm{~mol}, \\
& n_{\mathrm{H}_{2}}=\frac{3 \mathrm{~g}}{2 \mathrm{~g} / \mathrm{mol}}=\frac{3}{2} \mathrm{~mol} .
\end{aligned}
$$

The ideal-gas law is

$$
p V=n R T
$$

Solving for $p$ for each gas yields

$$
\begin{aligned}
& p_{\mathrm{N}_{2}}=\frac{n R T}{V}=\frac{\left(\frac{3}{28} \mathrm{~mol}\right)(8.314 \mathrm{~J} / \mathrm{mol} \cdot \mathrm{~K})(293 \mathrm{~K})}{10^{-3} \mathrm{~m}^{3}}=2.61 \times 10^{5} \mathrm{~Pa}=2.6 \mathrm{~atm}, \\
& p_{\mathrm{H}_{2}}=\frac{n R T}{V}=\frac{\left(\frac{3}{2} \mathrm{~mol}\right)(8.314 \mathrm{~J} / \mathrm{mol} \cdot \mathrm{~K})(293 \mathrm{~K})}{10^{-3} \mathrm{~m}^{3}}=3.65 \times 10^{6} \mathrm{~Pa}=36 \mathrm{~atm}
\end{aligned}
$$

(b) The rms speed is found from

$$
v_{\mathrm{rms}}=\sqrt{\frac{3 R T}{M}}
$$

Solving for our gases gives

$$
\begin{aligned}
& v_{\mathrm{N}_{2}}=\sqrt{\frac{3 R T}{M}}=\sqrt{\frac{3(8.314 \mathrm{~J} / \mathrm{mol} \cdot \mathrm{~K})(293 \mathrm{~K})}{\left(28 \times 10^{-3} \mathrm{~kg} / \mathrm{mol}\right)}}=511 \mathrm{~m} / \mathrm{s} \\
& \mathrm{v}_{\mathrm{H}_{2}}=\sqrt{\frac{3 R T}{M}}=\sqrt{\frac{3(8.314 \mathrm{~J} / \mathrm{mol} \cdot \mathrm{~K})(293 \mathrm{~K})}{\left(2 \times 10^{-3} \mathrm{~kg} / \mathrm{mol}\right)}}=1910 \mathrm{~m} / \mathrm{s} .
\end{aligned}
$$

EVALUATE We see that both the pressure and the rms speed are higher for the $\mathrm{H}_{2}$ gas, consistent with its smaller mass.

## 5: Escape velocities of the sun and earth

Calculate the escape velocities at the surface of the sun and the earth, and determine whether any molecules have rms thermal speeds comparable with these escape velocities at 300 K .

## Solution

IDENTIFY We will use energy conservation to determine the escape velocities and kinetic theory to find the rms speeds. The target variables are the escape velocities at the surface of the sun and earth and the rms speeds for gases at 300 K . The mass and radius values are found in Appendix F of the text.

SET UP For a molecule to escape from the sun or the earth, all of the molecule's gravitational potential energy must convert to kinetic energy. We use that relation to solve for the minimum escape velocity. The rms speed is given by an expression found in the text.

EXECUTE The escape velocity is found from energy conservation:

$$
\frac{G M m}{R}=\frac{1}{2} m v^{2}
$$

Solving for the velocity gives

$$
v=\sqrt{\frac{2 G M}{R}}
$$

Solving for the escape velocities of the sun and the earth yields

$$
\begin{aligned}
& v_{\text {sun }}=\sqrt{\frac{2 G M_{\text {sun }}}{R_{\text {sun }}}}=\sqrt{\frac{2\left(6.67 \times 10^{-11} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{kg}^{2}\right)\left(2 \times 10^{30} \mathrm{~kg}\right)}{7.0 \times 10^{8} \mathrm{~m}}}=6.2 \times 10^{5} \mathrm{~m} / \mathrm{s} \\
& v_{\text {earth }}=\sqrt{\frac{2 G M_{\text {earth }}}{R_{\text {earth }}}}=\sqrt{\frac{2\left(6.67 \times 10^{-11} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{kg}^{2}\right)\left(6.0 \times 10^{24} \mathrm{~kg}\right)}{6.38 \times 10^{6} \mathrm{~m}}}=1.12 \times 10^{4} \mathrm{~m} / \mathrm{s}
\end{aligned}
$$

The rms speed is found from

$$
v_{\mathrm{rms}}=\sqrt{\frac{3 R T}{M}}
$$

The hydrogen atom has the highest rms speed of any gaseous atom, so we calculate the speed for hydrogen. For a hydrogen atom at 300 K ,

$$
v_{\mathrm{H}}=\sqrt{\frac{3 R T}{M}}=\sqrt{\frac{3(8.314 \mathrm{~J} / \mathrm{mol} \cdot \mathrm{~K})(300 \mathrm{~K})}{\left(1 \times 10^{-3} \mathrm{~kg} / \mathrm{mol}\right)}}=2700 \mathrm{~m} / \mathrm{s} .
$$

The rms speed is $2700 \mathrm{~m} / \mathrm{s}$, less than the escape velocity of the sun or the earth.
EVALUATE We see that it is unlikely for any gas to have sufficient rms speed to escape the sun or the earth. The sun's average surface temperature is roughly 6000 K , corresponding to $12,000 \mathrm{~m} / \mathrm{s}$, a speed still too low for escape. Why does the moon not have an atmosphere?

Practice Problem: Find the escape velocity of a particle on the moon. Answer: $2370 \mathrm{~m} / \mathrm{s}$, enough for most gases to escape.

## Try It Yourself!

## 1: Helium gas

Helium gas is admitted to a volume of $200 \mathrm{~cm}^{3}$ at a temperature of 77 K until the pressure is equal to 1 atm . (a) If the temperature of the container is raised to $20^{\circ} \mathrm{C}$, what will the pressure inside the container be? (b) If the system has a relief valve that will not permit the pressure to exceed 1 atm, what fraction of gas remains at $20^{\circ} \mathrm{C}$ ?

## Solution Checkpoints

IDENTIFY AND SET UP Use the ideal-gas law to solve the problem. Is that law a valid approximation?
EXECUTE (a) For the closed system, of the number of moles, pressure, temperature, and volume, which changes as the temperature rises? Only pressure and temperature change, so their ratio remains constant:

$$
\frac{p_{1}}{T_{1}}=\frac{p_{2}}{T_{2}} .
$$

The final pressure is 3.8 atm .
(b) For the valved system, what remains constant? Both the moles and temperature change, yielding

$$
n_{1} R T_{1}=n_{2} R T_{2}
$$

$26.3 \%$ of the gas remains.
EVALUATE How could the system be changed so that both the pressure and the number of moles remain constant?

## 2: Nitrogen gas

(a) Calculate the volume occupied by 1 mole of nitrogen gas at the critical temperature ( 162.2 K ) and critical pressure $\left(33.9 \times 10^{5} \mathrm{~N} / \mathrm{m}^{2}\right)$. Express your answer as a ratio of the volume to the known critical volume $\left(90.1 \times 10^{-6} \mathrm{~m}^{3}\right)$. (b) Calculate the pressure at the critical volume and temperature.

## Solution Checkpoints

IDENTIFY AND SET UP Treat the nitrogen as an ideal gas and use the ideal-gas law to solve the problem.

EXECUTE (a) The ideal-gas law can be used to find the volume:

$$
V=\frac{n R T}{p}
$$

The volume is $3.44 V_{C}$.
(b) The critical pressure is $3.44 p_{c}$.

EVALUATE What do the critical pressure, temperature, and volume refer to? Do you think the results are valid?

## 3: RMS speed values

In a gas at 300 K that is a mixture of the diatomic molecules $\mathrm{H}_{2}(2 \mathrm{~g} / \mathrm{mol})$ and $\mathrm{D}_{2}(4 \mathrm{~g} / \mathrm{mol})$, find the rms speeds of the molecules

## Solution Checkpoints

IDENTIFY AND SET UP Kinetic theory gives the rms speed of molecules.
EXECUTE The rms speed is found from

$$
v=\sqrt{\frac{3 R T}{M}}
$$

The rms speed is $1930 \mathrm{~m} / \mathrm{s}$ for hydrogen and $1370 \mathrm{~m} / \mathrm{s}$ for deuterium $\left(\mathrm{D}_{2}\right)$.
EVALUATE Could you devise a method for separating these two molecules by using the differences in their rms speeds? Explain.

## The First Law of Thermodynamics

## Summary

In this chapter, we investigate and quantify thermodynamic processesprocesses that exchange heat and do work. We will examine thermodynamic systems and energy in these systems. Our examination will lead us to the first law of thermodynamics and thermodynamic processes. Four common thermodynamics processes will be highlighted, and the implications of those processes for ideal gases will be examined.

## Objectives

After studying this chapter, you will understand

- How to define thermodynamic processes.
- How heat is transferred and work is done in a thermodynamic process.
- The definition of, and how to apply, the first law of thermodynamics.
- How a path between initial and final states affects a thermodynamic process.
- The four common thermodynamic processes (adiabatic, isochoric, isobaric, and isothermal).
- How to apply the common thermodynamic processes to find the changes in heat, work, and internal energy of thermodynamic systems.
- How ideal gases are described in the common thermodynamic processes.


## Concepts and Equations

| Term |
| :--- |
| Heat and Work in |
| Thermodynamic Processes |

## Description

A thermodynamic system may exchange energy with its surroundings by heat transfer or by mechanical work. The work done by the system is given by

$$
\begin{aligned}
W & =\int_{V_{1}}^{V_{2}} p d V \\
& =p\left(V_{2}-V_{1}\right) \quad(\text { constant pressure only }) .
\end{aligned}
$$

In any thermodynamic process, the heat added to the system and the work done by the system depend on the steps the system takes from its initial to its final states, as well as on the initial and final states themselves.

## First Law of Thermodynamics

The first law of thermodynamics states that when heat $Q$ is added to a system while work $W$ is performed by the system, the internal energy $U$ changes by

$$
\Delta U=Q-W
$$

For infinitesimal changes,

$$
\boldsymbol{d} U=\boldsymbol{d} Q-\boldsymbol{d} W
$$

The internal energy of any thermodynamic system depends only on its state. The change in internal energy in any process depends only on the initial and final states.

## Thermodynamic Processes

Common thermodynamic processes include the adiabatic process, in which no heat flows into or out of the system $(Q=0)$; the isochoric process, in which the volume remains constant $(W=0)$; the isobaric process, in which the pressure remains constant $\left[W=p\left(V_{2}-V_{1}\right)\right]$; and the isothermal process, in which the temperature remains constant.

## Properties of an Ideal Gas

pressure or volume. The molar heat capacity at constant volume $\left(C_{V}\right)$ and the molar heat capacity at constant pressure $\left(C_{p}\right)$ for an ideal gas are related by

$$
C_{p}=C_{V}+R
$$

For an adiabatic process in an ideal gas, both $T V^{\gamma-1}$ and $p V^{\gamma}$ are constant, where $\gamma=C_{p} / C_{V}$. The work done by an ideal gas during an adiabatic expansion is given by

$$
\begin{aligned}
W & =n C_{V}\left(T_{1}-T_{2}\right) \\
& =\frac{C_{V}}{R}\left(p_{1} V_{1}-p_{2} V_{2}\right) \\
& =\frac{1}{\gamma-1}\left(p_{1} V_{1}-p_{2} V_{2}\right) .
\end{aligned}
$$

## Conceptual Questions

## 1: pV diagrams

One mole of helium gas is placed in a sealed container and undergoes an isochoric process that results in a doubling of the helium's pressure. Next, the gas undergoes an adiabatic process until the volume of the container is tripled. It then undergoes an isobaric expansion which results in a volume that is four times its original volume. Finally, the helium undergoes an isothermal compression that leaves the container with the same volume that it had after the first process. Sketch a $p V$ diagram for this combined process.

## Solution

SET UP AND SOLVE Figure 19.1 shows the resulting $p V$ diagram. The diagram starts at point $a$ with initial pressure $p_{0}$ and volume $V_{0}$. Then comes the isochoric process, at constant volume, represented by a vertical line to point $b$, where the pressure has doubled. Next is the adiabatic process, in which no heat is exchanged. This process follows the path to point $c$, where the volume has tripled. Next is the isobaric process, carried out at constant pressure and represented by a horizontal line to point $d$, where the volume has increased by $V_{0}$ from point $c$. Finally, in the isothermic process, the segment follows the path to point $e$, where the pressure has increased to $2 p_{0}$.


Figure 19.1 Question 1.
REFLECT This problem helps clarify the differences in the four common thermodynamic processes. $p V$ diagrams are a valuable aid in solving problems involving these processes. The diagrams will help lead us through a particular problem, as well as provide a check on our results.

## 2: Internal energy in a thermodynamic process

A container of argon gas undergoes a multistep process. First, it undergoes an isobaric expansion that triples its volume. Next, it goes through an isochoric process that results in a doubling of the argon's pressure. Then it cools adiabatically by 50 K . After that, it undergoes a second isochoric process that doubles its volume. Finally, it undergoes isobaric compression that leaves it at its initial temperature. Find the total change in internal energy.

## Solution

SET UP AND SOLVE The change in the internal energy of an ideal gas depends only upon temperature. Argon is an ideal gas. Since the final temperature is the same as the initial temperature, the total change in internal energy is zero.

REFLECT This complicated scenario shows that we need to focus on the important parts of the process to interpret the results properly. Although it is trivial to find the change in internal energy, it would be cumbersome to find the heat added to the system.

## Problems

## 1: Adiabatic compression of helium

Helium gas is expanded adiabatically from a 12 -liter volume at STP to a 33 -liter volume. Find the final temperature and pressure of the gas and the work done on the gas.

## Solution

IDENTIFY Since the helium is expanded adiabatically, both $p V^{y}$ and $T V^{\gamma-1}$ are constant during the process. The target variables are the final temperature and pressure of the gas and the work done on the gas.

SET UP For helium, $\gamma=1.67$ (from Table 19.1 in the text). Standard temperature and pressure (STP) refers to a temperature of 273 K and a pressure of 1 atmosphere.

EXECUTE To find the final pressure, we use the relation

$$
p V^{\gamma}=\text { constant }=p_{1} V_{1}^{\gamma}=p_{2} V_{2}^{\gamma}
$$

where the subscripts 1 and 2 indicate "before" and "after," respectively. Rearranging terms to find the final pressure, we obtain

$$
p_{2}=\frac{p_{1} V_{1}^{\gamma}}{V_{2}^{\gamma}}=\frac{\left(1.01 \times 10^{5} \mathrm{~Pa}\right)(12 \ell)^{1.67}}{(33 \ell)^{1.67}}=1.86 \times 10^{4} \mathrm{~Pa} .
$$

To find the final temperature, we use the relation

$$
T V^{\gamma-1}=\mathrm{constant}=T_{1} V_{1}^{\gamma-1}=T_{2} V_{2}^{\gamma-1}
$$

Rearranging terms in the preceding equation gives a final temperature of

$$
T_{2}=\frac{T_{1} V_{1}^{\gamma-1}}{V_{2}^{\gamma-1}}=\frac{(273 \mathrm{~K})(12 \ell)^{0.67}}{(33 \ell)^{0.67}}=139 \mathrm{~K} .
$$

The work done by an ideal gas in an adiabatic process is

$$
W=n C_{V}\left(T_{1}-T_{2}\right)
$$

For helium, $C_{V}$ is 12.47 , from Table 19.1. The number of moles is

$$
n=\frac{12 \mathrm{~L}}{22.4 \mathrm{~L}}=0.536 \mathrm{~mol} .
$$

The work done by the gas is

$$
W=n C_{V}\left(T_{1}-T_{2}\right)=(0.536 \mathrm{~mol})(12.47 \mathrm{~J} / \mathrm{mol} \cdot \mathrm{~K})(273 \mathrm{~K}-139 \mathrm{~K})=895 \mathrm{~J}
$$

The work done on the gas is the opposite, or -895 J . The final pressure is $1.86 \times 10^{4} \mathrm{~Pa}$ and the final temperature is 139 K .

EVALUATE We see that both the temperature and the pressure decreased in this adiabatic expansion. That makes sense, since no heat was transferred into or out of the system, so having a larger volume required a lower pressure and temperature.

Would you find the same final temperature if you used the ideal-gas equation? If you check, you'll find that you indeed do find the same final temperature.

## 2: Isochoric and isobaric process with helium

Two moles of helium gas are taken from point $a$ to point $c$ in the diagram shown in Figure 19.2. Find the change in internal energy along path $a b c$.


Figure 19.2 Problem 2.

## Solution

IDENTIFY The target variable is the change in internal energy.
SET UP We break the process up into two segments, one from $a$ to $b$ and one from $b$ to $c$. The first segment is an isochoric process (carried out at constant volume) and the second is an isobaric process (carried out at constant pressure). We'll use the relations for those segments to determine the work, heat, and temperature changes. We'll combine the work and heat changes to find the change in internal energy.

EXECUTE We find the change in internal energy along $a b$ by first finding the change in temperature along $a b$ :

$$
\Delta T_{a b}=\frac{\left(p_{b}-p_{a}\right) V_{a}}{n R}=\frac{\left(3.03 \times 10^{5} \mathrm{~Pa}-1.01 \times 10^{5} \mathrm{~Pa}\right)(15 \ell)\left(10^{-3} \mathrm{~m}^{3} / \ell\right)}{(2 \mathrm{~mol})(8.31 \mathrm{~J} / \mathrm{mol} / \mathrm{K})}=182 \mathrm{~K}
$$

The heat transferred during $a b$ is then

$$
Q_{a b}=n C_{\mathrm{v}} \Delta T_{a b}=(2 \mathrm{~mol})(12.47 \mathrm{~J} / \mathrm{mol} / \mathrm{K})(182 \mathrm{~K})=4500 \mathrm{~J}
$$

where we used the molar heat capacity at constant volume for helium ( $C_{V}=12.47 \mathrm{~J} / \mathrm{mol} / \mathrm{K}$, from Table 19.1). The work done during segment $a b$ is zero, since that segment is isochoric. The change in internal energy for segment $a b$ is then

$$
\Delta U_{a b}=Q_{a b}-W_{a b}=4500 \mathrm{~J}-0=4500 \mathrm{~J} .
$$

For segment $b c$, we follow the same procedure. The change in temperature along $b c$ is

$$
\Delta T_{b c}=\frac{p_{b}\left(V_{c}-V_{b}\right)}{n R}=\frac{\left(3.03 \times 10^{5} \mathrm{~Pa}\right)\left(2.2 \times 10^{-2} \mathrm{~m}^{3}-1.5 \times 10^{-2} \mathrm{~m}^{3}\right)}{(2 \mathrm{~mol})(8.31 \mathrm{~J} / \mathrm{mol} / \mathrm{K})}=127 \mathrm{~K}
$$

The heat transferred during $b c$ is

$$
Q_{b c}=n C_{\mathrm{P}} \Delta T_{b c}=(2 \mathrm{~mol})(20.78 \mathrm{~J} / \mathrm{mol} / \mathrm{K})(127 \mathrm{~K})=5300 \mathrm{~J},
$$

where we used the molar heat capacity at constant pressure for helium ( $C_{P}=20.78 \mathrm{~J} / \mathrm{mol} / \mathrm{K}$, from Table 19.1). The work done during segment $b c$ is

$$
W_{b c}=p_{b} \Delta V_{b c}=\left(3.03 \times 10^{5} \mathrm{~Pa}\right)\left(\left(2.2 \times 10^{-3} \mathrm{~m}^{3}\right)-\left(1.5 \times 10^{-3} \mathrm{~m}^{3}\right)\right)=2100 \mathrm{~J}
$$

The change in internal energy during segment $b c$ is

$$
\Delta U_{b c}=Q_{b c}-W_{b c}=5300 \mathrm{~J}-2100 \mathrm{~J}=3200 \mathrm{~J}
$$

The total change in internal energy is

$$
\Delta U=\Delta U_{a b}+\Delta U_{b c}=4500 \mathrm{~J}+3200 \mathrm{~J}=7700 \mathrm{~J}
$$

The total change in internal energy in the system is 7700 J .
EVALUATE We see that by breaking up a process into segments, determining the type of process that takes place during each segment, and knowing which variables change and which remain constant during each segment, one can easily find the change in internal energy.

How should the change in internal energy along path $a d c$ compare with the change along path $a b c$ ? They should be the same for helium, an ideal gas.

Practice Problem: Find the change in internal energy along segments $a d$ and $d c$, and compare their sums with the change in internal energy along path abc. Answer: $\Delta U_{a d}=1100 \mathrm{~J}, \Delta U_{d_{c}}=6600 \mathrm{~J}$, $\Delta U_{a d c}=7700 \mathrm{~J}$.

## 3: Monatomic gas process

One mole of an ideal monatomic gas starts at point $A$ in Figure 19.3 ( $T=273 \mathrm{~K}, p=1 \mathrm{~atm}$ ) and undergoes an adiabatic expansion to point $B$, where the volume of the gas has doubled. These two processes are followed by an isothermal compression to the original volume at point $C$ and an isobaric increase to the original point $A$. Find (a) the temperature at point $B$, (b) the pressure at point $C$, and (c) the total work done for the entire cycle. Take $\gamma$ to be $5 / 3$.


Figure 19.3 Problem 3.

## Solution

IDENTIFY The target variables are the temperature at point $B$, the pressure at point $C$, and the total work done during the process.

SET UP We'll break the process up into the three segments shown in the figure. We'll use the relations for those segments to determine the temperature, pressure, and work.

EXECUTE (a) For the adiabatic process, we use the relation

$$
T V^{\gamma-1}=\mathrm{constant}=T_{1} V_{1}^{\gamma-1}=T_{2} V_{2}^{\gamma-1} .
$$

Rearranging terms to find the temperature at point $B$ gives

$$
T_{2}=\frac{T_{1} V_{1}^{\gamma-1}}{V_{2}^{\gamma-1}}=\frac{(273 \mathrm{~K})(V)^{0.67}}{(2 V)^{0.67}}=172 \mathrm{~K}
$$

(b) The temperature at point $C$ is the same as at point $B$, since the second step is isothermal. Also, the volume is the same at point $C$ as it was at point $A$. We use the standard ideal-gas law:

$$
\frac{p_{1} V_{1}}{T_{1}}=\frac{p_{3} V_{3}}{T_{3}} .
$$

Solving for $p_{3}$ gives

$$
p_{3}=p_{1} \frac{T_{3}}{T_{1}}=(1 \mathrm{~atm}) \frac{(172 \mathrm{~K})}{(273 \mathrm{~K})}=0.630 \mathrm{~atm} .
$$

(c) The total work done by the gas during the entire cycle is the sum of the separate amounts of work done during each cycle. The path $A B$ is adiabatic, so no heat is exchanged and the work is

$$
W_{A B}=-C_{V}\left(T_{2}-T_{1}\right)=-\left(\frac{3}{2}(8.314 \mathrm{~J} / \mathrm{mol} \cdot \mathrm{~K})\right)(172 \mathrm{~K}-273 \mathrm{~K})=1260 \mathrm{~J}
$$

For segment $B C$, the temperature is constant. The work is given by the integral

$$
\begin{aligned}
W & =\int_{V_{1}}^{V_{2}} p d V \\
& =\int_{V_{2}}^{V_{3}} \frac{R T}{V} d V=\left.R T \ln V\right|_{V_{2}} ^{V_{3}} \\
& =R T \ln \left(V_{3} / V_{2}\right)=(8.314 \mathrm{~J} / \mathrm{mol} \cdot \mathrm{~K})(172 \mathrm{~K}) \ln (V / 2 \mathrm{~V}) \\
& =-991 \mathrm{~J} .
\end{aligned}
$$

No work is done during segment $C A$, since the system changes isobarically. The total work done is the sum, 268 J .

EVALUATE We solved this problem by breaking the process into segments and working through each segment to find the final state variables. Can you use the area inside the cycle to find the amount of work? The area is greater than zero, indicating positive work.

## 4: Expansion process for argon

One mole of argon is initially at $25^{\circ} \mathrm{C}$ and occupies a volume of 35 liters. The argon is first expanded at constant pressure until the volume is doubled and then expanded adiabatically until the temperature returns to $25^{\circ} \mathrm{C}$. Find the total change in internal energy, the total work done by the argon, and the final volume and pressure of the argon.

## Solution

IDENTIFY Sketch the process and then use the definitions to solve. The target variables are the total change in internal energy, the total work done, and the final volume and pressure.

SET UP Figure 19.4 shows the $p V$ diagram for the process. We'll break the process up into two segments, one from $a$ to $b$ and one from $b$ to $c$. The first segment is an isobaric process (carried out under constant pressure) and the second is adiabatic (no heat exchanged). We'll use the relations for those segments to determine the work, heat, and temperature changes. We'll combine these results to find the quantities of interest.


Figure 19.4 Problem 4 sketch.
EXECUTE The total change in internal energy is zero, since the internal energy of an ideal gas depends only on temperature and the final temperature is equal to the initial temperature.
Next, we find the temperature at point $b$. Segment $a b$ is at constant pressure, so

$$
\frac{V_{a}}{T_{a}}=\frac{V_{b}}{T_{b}} .
$$

The temperature at $b$ is

$$
T_{b}=\frac{V_{b}}{V_{a}} T_{a}=\frac{(70 \ell)}{(35 \ell)}(273+25) \mathrm{K}=596 \mathrm{~K}
$$

The heat supplied during $a b$ is

$$
Q_{a b}=n C_{\mathrm{P}} \Delta T_{a b}=(1 \mathrm{~mol})(20.78 \mathrm{~J} / \mathrm{mol} / \mathrm{K})(596 \mathrm{~K}-298 \mathrm{~K})=6190 \mathrm{~J},
$$

where we used the molar heat capacity at constant pressure for $\operatorname{argon}\left(C_{P}=20.78 \mathrm{~J} / \mathrm{mol} / \mathrm{K}\right.$, from Table 19.1). The heat transferred in $b c$ is zero, since segment $b c$ is adiabatic. The total heat supplied in the complete process is 6190 J . Because the total internal energy change is zero, the heat supplied must be equal to the work done by the argon. The work done by the argon is 6190 J .

Next, we find the final volume and pressure. To find the final volume in the adiabatic process $(b c)$, we use the relation

$$
T V^{\gamma-1}=\mathrm{constant}=T_{b} V_{b}^{\gamma-1}=T_{c} V_{c}^{\gamma-1}
$$

where $\gamma=1.67$ for argon. Rearranging terms to find the final volume gives

$$
V_{c}=\sqrt[\gamma-1]{\frac{T_{b} V_{b}^{\gamma-1}}{T_{c}}}=\sqrt[0.67]{\frac{(596 \mathrm{~K})(70 \ell)^{0.67}}{(298 \mathrm{~K})}}=1971 .
$$

We can use the equation of state for an ideal gas to find the final pressure:

$$
p_{c}=\frac{n R T_{c}}{V_{c}}=\frac{(1 \mathrm{~mol})(8.31 \mathrm{~J} / \mathrm{mol} / \mathrm{K})(298 \mathrm{~K})}{\left(197 \times 10^{-3} \mathrm{~m}^{3}\right)}=1.26 \times 10^{4} \mathrm{~Pa}
$$

The final pressure is $12,600 \mathrm{~Pa}$ and the final volume is 197 liters.
EVALUATE We again see that we need to break the process up into segments and work through each segment to find the final state variables. We also see that the $p V$ graph is useful in solving problems involving thermodynamic processes.

## Try It Yourself!

## 1: Ideal-gas process

Consider $n$ moles of an ideal gas that undergo the constant-volume and constant-pressure processes along the paths shown in Figure 19.5 from the initial state $a$ to $b$, then from $b$ to $c$, then from $c$ to $d$, and then back from $a$ to $d$. For each of these processes, calculate (a) the work done by the system and (b) the heat taken in by the system.


Figure 19.5 Try It Yourself 1.

## Solution Checkpoints

IDENTIFY AND SET UP We break the process up into segments and use the relations for isochoric and isobaric processes.

EXECUTE (a) No work is done by the gas during the isochoric processes. For the isobaric processes, the work done by the gas is the pressure times the change in volume:

$$
\begin{aligned}
W_{b \rightarrow c} & =p_{b}\left(V_{c}-V_{a}\right), \\
W_{d \rightarrow a} & =p_{a}\left(V_{a}-V_{c}\right) .
\end{aligned}
$$

The work done from $b$ to $c$ is a positive quantity, and the work done from $d$ to $a$ is negative, resulting in negative total work. Work is done on the gas.
(b) The heat taken in by the ideal gas will be the number of moles, times the change in temperature, times the heat capacity at constant volume or constant pressure. This relationship gives

$$
\begin{aligned}
Q_{a \rightarrow b} & =n C_{V}\left(T_{b}-T_{a}\right), \\
Q_{b \rightarrow c} & =n C_{p}\left(T_{c}-T_{b}\right), \\
Q_{c \rightarrow d} & =n C_{V}\left(T_{d}-T_{c}\right), \\
Q_{d \rightarrow a} & =n C_{p}\left(T_{a}-T_{d}\right) .
\end{aligned}
$$

EVALUATE How does the heat taken in by the gas compare with the work done by the gas?

## 2: Ideal-gas process, version 2

A total of $n$ moles of an ideal monatomic gas is taken from point 1 on the $T_{1}$ isotherm to point 3 on the $T_{2}$ isotherm along the path $1 \rightarrow 2 \rightarrow 3$ as shown in Figure 19.6. (a) Calculate the change in internal energy of the gas and the heat that must be added to it in this process. (b) Suppose the path $1 \rightarrow 4 \rightarrow 3$ is followed instead. Calculate the heat added along this path.


Figure 19.6 Try It Yourself 2.

## Solution Checkpoints

IDENTIFY AND SET UP We break the process up into segments and use the relations for isothermic and isobaric processes.

EXECUTE (a) In going from 1 to 3 , the change in internal energy depends on the temperature differences and not the path taken. This gives

$$
\Delta U=n C_{V}\left(T_{3}-T_{1}\right)
$$

The heat taken in from 1 to 3 is

$$
\Delta Q=n C_{p}\left(T_{2}-T_{1}\right)+n R T \ln \left(\frac{V_{3}}{V_{2}}\right)
$$

(b) The change in internal energy is the same as in (a). The heat is

$$
\Delta Q=n C_{V}\left(T_{2}-T_{1}\right)+p_{3}\left(V_{3}-V_{1}\right)
$$

EVALUATE How does the heat exchanged on the two paths differ. Why?

## The Second Law of Thermodynamics

## Summary

In this chapter, we will complete our investigation of thermodynamics, examining thermodynamic processes and the second law of thermodynamics. Heat engines and refrigerators transform heat into work or energy in cyclic processes. Thermal efficiency and performance coefficients for engines and refrigerators will be defined. The second law of thermodynamics limits the efficiency of engines and has profound implications in many physical processes. The second law can be quantified in terms of entropy, a measure of disorder. We will examine several common cyclic processes to aid our understanding of thermodynamics.

## Objectives

After studying this chapter, you will understand

- How to define and identify reversible processes.
- How to analyze heat engines and refrigeration cycles.
- How to apply the second law of thermodynamics.
- How to calculate entropy for a variety of systems.
- How to apply thermodynamic principles to a variety of engine and refrigeration cycles.

Concepts and Equations

| Term |
| :--- |
| Directions of Thermodynamic Processes |


|  | which the system is always in, or very close to, thermal e <br> thermodynamic processes are irreversible. |
| :--- | :--- |
| Heat Engine | A heat engine takes heat $Q_{H}$ from a source, converts part <br> $W$, and discards the remaining heat $\left\|Q_{C}\right\|$ at a lower tempe <br> engine's thermal efficiency $e$ is |
| $\qquad e=\frac{W}{Q_{H}}=1+\frac{Q_{C}}{\mathrm{Q}_{H}}=1-\frac{\left\|Q_{C}\right\|}{\left\|Q_{H}\right\|}$. |  |


| Otto Cycle | A gasoline engine operating in the Otto cycle has a theoretical maximum thermal efficiency given by $e=1-\frac{1}{r^{\gamma-1}},$ <br> where $r$ is the compression ratio. |
| :---: | :---: |
| Refrigerator | A refrigerator takes heat $Q_{C}$ from a cold source, performs work $W$, and discards the heat $\left\|Q_{H}\right\|$ to a warmer source. The performance coefficient $K$ is $K=\frac{Q_{C}}{W}=\frac{\left\|Q_{C}\right\|}{\left\|Q_{H}\right\|-\left\|Q_{C}\right\|} .$ |
| The Second Law of Thermodynamics | The second law of thermodynamics states that it is impossible for any cyclic system to convert heat completely into work. It also states that no cyclic process can transfer heat from a cold place to a hot place without any input of work. |
| Carnot Cycle | The Carnot cycle operates between two heat reservoirs and represents the most efficient heat engine. The Carnot cycle combines the reversible adiabatic and isothermal expansion and contraction between two heat reservoirs at temperatures $T_{H}$ and $T_{C}$, respectively. The efficiency of the Carnot cycle is $e_{\text {Carroot }}=1-\frac{T_{C}}{T_{H}}=\frac{T_{H}-T_{C}}{T_{H}} .$ |

## Entropy

Entropy is a quantitative measure of the disorder of a system. The entropy change in a reversible thermodynamic system is

$$
\Delta S=\int_{1}^{2} \frac{d Q}{T} .
$$

The second law of thermodynamics can be stated as "The entropy of an isolated system may increase, but not decrease. The total entropy of a system interacting with its surroundings may never decrease."

## Conceptual Questions

## 1: Cleaning your room

Your parents are always nagging you about cleaning your room. After learning about the second law of thermodynamics, you explain to your parents that it is impossible to clean your room, since cleaning would reduce the entropy inside your room and violate the second law of thermodynamics. Your mother recalls her college physics course and convinces you that you can clean your room without violating the second law. How does she convince you?

## Solution

IDENTIFY, SET UP, AND EXECUTE Your mother agrees with you that the entropy of a closed system can never decrease. But she notes that when you clean your room, the system consists of you plus your belongings. You can decrease the entropy of your belongings in your room by increasing the entropy of your body, as long as the total entropy increases. You can certainly clean your room!

EVALUATE This problem shows how the entropy of isolated components in a system may decrease as long as the system's total entropy increases. It also shows that you shouldn't argue with your mother, although you may want to try the argument on your father, who doesn't remember his physics course.

## 2: Leaving a refrigerator door open to cool a room

When the air-conditioning system at your house fails, your younger brother suggests leaving the refrigerator door open to cool the house. Is this method effective?

## Solution

IDENTIFY, SET UP, AND EXECUTE A refrigerator cools its contents by taking heat away from the contents, performing work, and expelling heat to a warmer region. The heat expelled is always greater than the heat removed from the contents. The refrigerator must add net heat to its surroundings. Opening the refrigerator will result in a warmer room, so it is not an effective method of cooling the room.

EVALUATE Can opening the refrigerator warm the house on a cold day? Yes, since it must expel heat to operate. It wouldn't be a very efficient heat source, but it would provide some heat to the room.

## 3: Water as a fuel

Some people have suggested using water as a clean fuel. The idea is to break apart water molecules into hydrogen and oxygen. Then, when the hydrogen (combined with water) is burned, it produces energy without pollution. How does the second law of thermodynamics relate to this idea?

## Solution

IDENTIFY, SET UP, AND EXECUTE Because the breaking apart of the water and the burning of hydrogen constitutes a reversible cycle, the net entropy must increase. The process may actually create more pollution, since it takes more energy to dissociate the water than is recovered by burning the hydrogen. For example, if gasoline is used to generate the hydrogen, it would take more gasoline to generate an equivalent amount of hydrogen-based power than if the gasoline were used to operate the vehicle directly.

EVALUATE There is the possibility that pollution would be reduced. If the hydrogen-generating plants installed high-quality pollution filters, than there could be less pollution generated overall by the plant compared with the pollution generated by many cars. However, a hydrogen-burning car will always require more total energy to operate. Hydrogen should probably be considered an alternative energy storage method.

## Problems

## 1: Work in a heat engine

A heat engine carries 0.2 mol of argon through the cyclic process shown in Figure 20.1. Process $a b$ is isochoric, process $b c$ is adiabatic, and process $c a$ is isobaric at a pressure of 2.0 atm . Find the net work done by the gas in the complete cycle. The temperatures at the of endpoints of the process are $T_{a}=290 \mathrm{~K}, T_{b}=650 \mathrm{~K}$, and $T_{c}=440 \mathrm{~K}$.


Figure 20.1 Problem 1.

## Solution

IDENTIFY Use the principles of heat engines to find the work done.
SET UP We'll break the cycle up into processes and find the work done during each process. We'll need to find the pressure and volume at the three points before finding the work. Argon is an ideal gas, so we'll use the ideal-gas relations.

EXECUTE Starting at point $a$, we find the volume $V_{a}$ from the ideal-gas equation:

$$
V_{a}=\frac{n R T_{a}}{p_{a}}=\frac{(0.2 \mathrm{~mol})(8.31 \mathrm{~J} / \mathrm{mol} / \mathrm{K})(290 \mathrm{~K})}{\left(2.02 \times 10^{5} \mathrm{~Pa}\right)}=2.39 \times 10^{-3} \mathrm{~m}^{3}
$$

At point $b$, the volume is the same as at point $a$. We find the pressure at $b$ :

$$
p_{b}=\frac{n R T_{b}}{V_{b}}=\frac{(0.2 \mathrm{~mol})(8.31 \mathrm{~J} / \mathrm{mol} / \mathrm{K})(650 \mathrm{~K})}{\left(2.39 \times 10^{-3} \mathrm{~m}^{3}\right)}=4.52 \times 10^{5} \mathrm{~Pa}
$$

We find the volume $V_{c}$ at $c$ :

$$
V_{c}=\frac{n R T_{c}}{p_{c}}=\frac{(0.2 \mathrm{~mol})(8.31 \mathrm{~J} / \mathrm{mol} / \mathrm{K})(440 \mathrm{~K})}{\left(2.02 \times 10^{5} \mathrm{~Pa}\right)}=3.62 \times 10^{-3} \mathrm{~m}^{3}
$$

With these values, we find the work done during each process. There is no work done during the process $a b$, since it is an isochoric process. Process $b c$ is adiabatic and there is no heat exchanged. The work is opposite the change in internal energy, or

$$
W_{b c}=-\Delta U_{b c}=-n C_{\mathrm{V}}\left(T_{c}-T_{b}\right)=-(0.2 \mathrm{~mol})(12.47 \mathrm{~J} / \mathrm{mol} / \mathrm{K})(440 \mathrm{~K}-650 \mathrm{~K})=524 \mathrm{~J},
$$

where we used the molar heat capacity at constant volume for $\operatorname{argon}\left(C_{V}=12.47 \mathrm{~J} / \mathrm{mol} / \mathrm{K}\right.$, from Table 19.1). Process $c a$ is isobaric and the work is

$$
W_{c a}=p_{c} \Delta V_{c a}=\left(2.02 \times 10^{5} \mathrm{~Pa}\right)\left(\left(2.39 \times 10^{-3} \mathrm{~m}^{3}\right)-\left(3.62 \times 10^{-3} \mathrm{~m}^{3}\right)\right)=-248 \mathrm{~J} .
$$

Note that the work done by the gas is negative, since it is compressed in $c a$. The work done during the complete cycle is the sum of the separate amounts of work done during each of the three processes:

$$
W=W_{a b}+W_{b c}+W_{c a}=0+524 \mathrm{~J}-248 \mathrm{~J}=276 \mathrm{~J} .
$$

The gas does 276 J of work in one cycle.
EVALUATE Since the area inside the $p V$ cycle diagram is equal to the work, our positive result is in agreement with the area shown in the diagram. Much of this problem is based on what we learned in Chapter 19. We are now combining the processes of Chapter 19 into a complete cycle.

## 2: Efficiency of a heat engine

Find the thermal efficiency of an engine that operates in accordance with the cycle shown in Figure 20.2, in which 2 moles of helium stored at 2.0 atm in a 10 -liter vessel starts at point $a$, undergoes an isochoric process to quadruple its pressure at point $b$, triples in volume in an isobaric expansion to point $c$, reduces its pressure to one-fourth its pressure at point $c$ through an isochoric process at point $d$, and goes through an isobaric compression reducing its volume by one-third to return to point $a$.


Figure 20.2 Problem 2.

## Solution

IDENTIFY The target variable is the thermal efficiency of the engine.
SET UP We need to know the work and heat of the cycle to find the efficiency of the engine. We break the cycle up into processes and use our knowledge of isochoric (constant-volume) and isobaric (constant-pressure) processes. We are given the changes in pressure and volume, so we can proceed immediately to calculating the work done during each of the four processes. Helium is an ideal gas, so we can use the ideal-gas relations.

EXECUTE We start by determining the pressure and volume at the points $b, c$, and $d$. The pressure at points $b$ and $c$ is $8 \mathrm{~atm}\left(4 P_{a}\right)$ and the pressure at $d$ is $2 \mathrm{~atm}\left(P_{a}\right)$. The volume at $b$ is 10 liters $\left(V_{a}\right)$ and the volume at $c$ and $d$ is 30 liters $\left(3 V_{a}\right)$. Next, we find the work for each process. No work is done during processes $a b$ and $c d$, since they are isochoric. Process $b c$ is isobaric and the work done is

$$
W_{b c}=p_{b} \Delta V_{b c}=\left(8.08 \times 10^{5} \mathrm{~Pa}\right)\left(\left(30 \times 10^{-3} \mathrm{~m}^{3}\right)-\left(10 \times 10^{-3} \mathrm{~m}^{3}\right)\right)=16,200 \mathrm{~J} .
$$

Process $d a$ is also isobaric and the work done is

$$
W_{d a}=p_{d} \Delta V_{d a}=\left(2.02 \times 10^{5} \mathrm{~Pa}\right)\left(\left(10 \times 10^{-3} \mathrm{~m}^{3}\right)-\left(30 \times 10^{-3} \mathrm{~m}^{3}\right)\right)=-4040 \mathrm{~J} .
$$

The total work is the sum of the separate amounts of work done during each of the four processes:

$$
W=W_{a b}+W_{b c}+W_{c d}+W_{d a}=0+16,200 \mathrm{~J}+0-4040 \mathrm{~J}=12,100 \mathrm{~J} .
$$

We need to find the heat flowing into the engine. Heat flows into the engine during processes $a b$ and $b c$. To find the heat flow into the engine, we need to know the temperature at points $a, b$, and $c$. The ideal-gas equation gives

$$
\begin{aligned}
& T_{a}=\frac{p_{a} V_{a}}{n R}=\frac{\left(2.02 \times 10^{5} \mathrm{~Pa}\right)\left(10 \times 10^{-3} \mathrm{~m}^{3}\right)}{(2 \mathrm{~mol})(8.31 \mathrm{~J} / \mathrm{mol} / \mathrm{K})}=122 \mathrm{~K}, \\
& T_{b}=\frac{p_{b} V_{b}}{n R}=\frac{\left(8.08 \times 10^{5} \mathrm{~Pa}\right)\left(10 \times 10^{-3} \mathrm{~m}^{3}\right)}{(2 \mathrm{~mol})(8.31 \mathrm{~J} / \mathrm{mol} / \mathrm{K})}=486 \mathrm{~K}, \\
& T_{c}=\frac{p_{c} V_{c}}{n R}=\frac{\left(8.08 \times 10^{5} \mathrm{~Pa}\right)\left(30 \times 10^{-3} \mathrm{~m}^{3}\right)}{(2 \mathrm{~mol})(8.31 \mathrm{~J} / \mathrm{mol} / \mathrm{K})}=1460 \mathrm{~K} .
\end{aligned}
$$

The heat flow in process $a b$ (carried out at constant volume) is

$$
Q_{a b}=n C_{\mathrm{V}} \Delta T_{a b}=(2 \mathrm{~mol})(12.47 \mathrm{~J} / \mathrm{mol} / \mathrm{K})(486 \mathrm{~K}-122 \mathrm{~K})=9080 \mathrm{~J}
$$

where we used the molar heat capacity at constant volume for helium ( $C_{V}=12.47 \mathrm{~J} / \mathrm{mol} / \mathrm{K}$, from Table 19.1). The heat flow in process $b c$ (carried out at constant pressure) is

$$
Q_{b c}=n C_{\mathrm{P}} \Delta T_{b c}=(2 \mathrm{~mol})(20.78 \mathrm{~J} / \mathrm{mol} / \mathrm{K})(1460 \mathrm{~K}-486 \mathrm{~K})=40,500 \mathrm{~J},
$$

where we used the molar heat capacity at constant pressure for helium ( $C_{P}=20.78 \mathrm{~J} / \mathrm{mol} / \mathrm{K}$, from Table 19.1). The total heat flowing into the engine in one cycle is therefore

$$
Q_{H}=Q_{a b}+Q_{b c}=40,500 \mathrm{~J}+9080 \mathrm{~J}=49,600 \mathrm{~J}
$$

The efficiency of the engine is

$$
e=\frac{W}{Q_{H}}=\frac{12,100 \mathrm{~J}}{49,600 \mathrm{~J}}=24.4 \% .
$$

EVALUATE We see that the engine is $24.4 \%$ efficient. Note that, to find the efficiency, we started by determining the state variables for the points on the $p V$ diagram. Then we found the work and heat flow in the cycle and combined these two pieces of information to solve the problem.

## 3: Efficiency of a diesel engine

Find the thermal efficiency of the diesel cycle shown in Figure 20.3 for an engine having a compression ratio $R=18$, an expansion ratio $E=6=R V / V_{C}$, and $C_{P} / C_{V}=1.4$.


Figure 20.3 Problem 3.

## Solution

IDENTIFY The target variable is the thermal efficiency of the engine. The efficiency involves the heat into and out of the system.

SET UP We start by finding the heat into and out of the system. Along the adiabatic lines, no heat is transferred. We will be able to find the heat in terms of the changes in temperature and then solve for the temperatures. We will substitute the compression and expansion factors where appropriate.
EXECUTE The efficiency is given by

$$
e=1+\frac{Q_{C}}{Q_{H}} .
$$

The heat entering the cycle between points $b$ and $c$ is given by

$$
Q_{b c}=n C_{p}\left(T_{c}-T_{b}\right)
$$

The heat leaving the cycle between points $d$ and $a$ is given by

$$
Q_{d a}=n C_{V}\left(T_{a}-T_{d}\right)
$$

No heat is exchanged between any other pairs of points, as they lie upon adiabatic lines. The efficiency is then

$$
e=1+\frac{T_{a}-T_{d}}{\gamma\left(T_{c}-T_{b}\right)},
$$

where we replaced the ratio of the specific heats with $\gamma$. To find the efficiency, we need to find the temperatures. We seek relations between the temperatures that should cancel out. Points $b$ and $c$ are at the same pressure, so we have

$$
\frac{V}{T_{b}}=\frac{V_{c}}{T_{c}}
$$

The ratio of the volumes can be substituted for the volumes:

$$
T_{c}=T_{b} \frac{R}{E}
$$

To relate these temperatures to the temperature at $a$, we use the adiabatic relation

$$
T_{b} V^{\gamma-1}=T_{a} V_{a}^{\gamma-1}=T_{a}(R V)^{\gamma-1}
$$

So we have

$$
T_{b}=T_{a} R^{\gamma-1}, \quad T_{c}=T_{a} R^{\gamma-1} \frac{R}{E}
$$

We find the temperature at point $d$ by using the other adiabatic line:

$$
T_{c} V_{c}^{\gamma-1}=T_{d} V_{d}^{\gamma-1}=T_{d}(R V)^{\gamma-1}
$$

Solving for $T_{d}$ yields

$$
T_{d}=\left(\frac{R}{E}\right)^{\gamma} T_{a}
$$

We now replace the temperatures with their expressions in terms of $T_{a}$ to find the efficiency:

$$
\begin{aligned}
e & =1+\frac{T_{a}-T_{d}}{\gamma\left(T_{c}-T_{b}\right)} \\
& =1+\frac{\left(1-(R / E)^{\gamma}\right)}{\gamma\left(R^{\gamma} / E-R^{\gamma-1}\right)} \\
& =1+\frac{\left(1-(18 / 6)^{1.4}\right)}{1.4\left(18^{1.4} / 6-18^{0.4}\right)}=0.59 .
\end{aligned}
$$

The efficiency is $59 \%$.
EVALUATE We see that the diesel engine is 59\% efficient. Diesel engines are generally more efficient than gasoline engines.

## 4: Entropy change in melting ice

A heat reservoir at $50^{\circ} \mathrm{C}$ is used to melt 25 kg of ice at $0^{\circ} \mathrm{C}$. What is the entropy change in the melted ice? What is the entropy change in the reservoir? What is the total entropy change in the system?

## Solution

IDENTIFY The target variables are the entropy changes for the ice, reservoir, and total system.
SET UP Entropy change in a reversible process is equal to the heat transferred divided by the temperature of the material. We can find the heat required to melt the ice, which must be equal to the heat provided by the reservoir.

EXECUTE The heat required to melt the ice is given by the heat of fusion relation and is

$$
Q=m L_{\mathrm{f}}=(25 \mathrm{~kg})\left(335 \times 10^{3} \mathrm{~J} / \mathrm{kg}\right)=8.375 \times 10^{6} \mathrm{~J}
$$

where we used the latent heat of fusion for ice $\left(335 \times 10^{3} \mathrm{~J} / \mathrm{kg}\right)$. The change in entropy for the water is

$$
\Delta S_{\text {ice }}=\frac{Q}{T}=\frac{8.375 \times 10^{6} \mathrm{~J}}{273 \mathrm{~K}}=30,700 \mathrm{~J} / \mathrm{K} .
$$

The change in entropy for the reservoir is equal to the heat leaving the reservoir divided by the temperature of the reservoir. The reservoir loses as much heat as the ice gains. The entropy change is

$$
\Delta S_{\text {reservoir }}=\frac{-Q}{T}=\frac{-8.375 \times 10^{6} \mathrm{~J}}{323 \mathrm{~K}}=-25,900 \mathrm{~J} / \mathrm{K}
$$

The total change in entropy is the sum of the entropies for the ice and reservoir:

$$
\Delta S_{\text {total }}=\Delta S_{\text {ice }}+\Delta S_{\text {reservoir }}=30,700 \mathrm{~J} / \mathrm{K}-25,900 \mathrm{~J} / \mathrm{K}=4800 \mathrm{~J} / \mathrm{K}
$$

The entropy of the ice increases by $30,700 \mathrm{~J} / \mathrm{K}$, the entropy of the reservoir decreases by $25,900 \mathrm{~J} / \mathrm{K}$, and the total entropy of the system increases by $4800 \mathrm{~J} / \mathrm{K}$.

EVALUATE The entropy of the reservoir decreased, but this does not violate the second law, since the system is the combination of the ice, the water, and the reservoir. The total change in entropy is positive, as expected.

## 5: Entropy change in isothermal expansion

Find the change in entropy for an ideal gas that undergoes an isothermal expansion from an initial volume to a final volume that is twice the initial volume.

## Solution

IDENTIFY The target variable is the entropy change for the gas.
SET UP Entropy change is related to the work and the temperature of the material. For an ideal gas, the internal energy doesn't change. We'll use the standard definitions to find the entropy change in the process.

EXECUTE Writing the change in entropy in terms of work and temperature gives

$$
T d S=d Q=0+d W
$$

since the change in internal energy is zero. The work done by an ideal gas undergoing an isothermal process is given by

$$
d W=p d V=\frac{n R T}{V} d V
$$

Combining the two equations yields an expression for the change in entropy:

$$
d S=\frac{n R}{V} d V
$$

We integrate to find the change:

$$
S=\int_{V_{1}}^{V_{2}} \frac{n R}{V} d V=\left.n R \ln V\right|_{V} ^{2 V}=n R \ln 2
$$

The change in entropy is $R \ln 2$ for each mole of gas.
EVALUATE The entropy of the gas increased, as expected.

## Try It Yourself!

## 1: Efficiency of a heat engine

Following the cycle shown in Figure 20.4, find the efficiency of a heat engine using an ideal monatomic gas as its working substance. Take $\gamma=1.4$ and $r=V_{b} / V_{a}=2.5$.


Figure 20.4 Try It Yourself 1.

## Solution Checkpoints

IDENTIFY AND SET UP Break the process up into segments and the heat exchanged during each segment.

EXECUTE No heat is exchanged in segment $a b$. (Why?) Along segment $b c$, heat is removed from the system and is equivalent to

$$
Q=n R T_{c} \ln \left(V_{a} / V_{b}\right)
$$

Along segment $c a$, heat is put into the system and is equivalent to

$$
Q=n C_{V}\left(T_{a}-T_{c}\right)
$$

The temperatures are written in terms of $T_{c}$ to find the efficiency, which is

$$
e=1-\frac{((\gamma-1) \ln r)}{r^{\gamma-1}-1}
$$

EVALUATE What effect does increasing $r$ have on the efficiency?

## 2: Entropy in mixed water

Equal volumes of water at $80^{\circ} \mathrm{C}$ and $20^{\circ} \mathrm{C}$ are mixed together. Find the increase in entropy for a total volume of $1.0 \mathrm{~m}^{3}$.

## Solution Checkpoints

IDENTIFY AND SET UP The final temperature of the mixture should be halfway between the initial temperatures of the two volumes of water if the heat capacities are constant and independent of temperature. No process is noted, so we can choose any reversible process. Assume that the process keeps the volume constant.

EXECUTE The change in entropy for the system is

$$
\Delta S=\int_{1}^{2} \frac{d Q}{T}=C_{V} \int_{T_{1}}^{T_{2}} \frac{d T}{T}=C_{V}\left(\frac{\mathrm{~T}_{2}}{\mathrm{~T}_{1}}\right)
$$

where the temperatures are in kelvins. The entropy of the hot water decreases, while the entropy of the cold water increases. The total change in entropy is $3.62 \times 10^{4} \mathrm{~J} / \mathrm{K}$.

EVALUATE How did you determine $C_{V}$ ?

## 21 <br> Electric Charge and Electric Field

## Summary

In this chapter, we begin our investigation into electricity and magnetism with a study of interactions between static electric charges. We will see that the electric interaction is governed by electric charges and that the electric charge is based on the structure of atoms and matter. Materials will be classified as conductors or insulators, depending on how charge moves within them. The force between two charges will be defined, and we will learn how to find the force on a charge due to many charges. We will also find an alternative description of the electric interaction through the definition of the electric field. We will learn to calculate electric fields for a variety of collections of charge and how to represent electric fields graphically.

## Objectives

After studying this chapter, you will understand

- The electric interaction, electric charge, insulators, and conductors.
- How to apply Coulomb's law to systems of two or more charges.
- How to find the resultant electric force by summing several electric forces.
- The electric field representation of the electric interaction.
- How to calculate electric fields for various distributions of charges.
- The movement of charges through electric fields.

Concepts and Equations

| Term | Description |
| :--- | :--- |
| Electric Charge | Electric charge is a fundamental property of particles that is responsible for <br> electrical interactions. Particles may be positively charged, negatively <br> charged, or neutral. Like charges repel and unlike charges attract. Electric <br> charge is conserved: It may be transferred between objects, but cannot be <br> created or destroyed. The SI unit of electric charge is the coulomb (C). The <br> fundamental unit of electric charge is the magnitude of the charge of one <br> electron or one proton, denoted $e:$ |
| $\qquad e=1.60217653 \times 10^{-19} \mathrm{C}$. |  |

where the constant $k=8.987551789 \times 10^{9} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{C}^{2}$ and $\epsilon_{0}=8.854 \times$ $10^{-12} \mathrm{C}^{2} / \mathrm{N} \cdot \mathrm{m}^{2}$ is the permittivity of free space. The force on each charge acts along the line joining the two charges and is repulsive if the two charges have the same sign and attractive if the two charges have opposite signs.

## Electric Field

An electric field, denoted $\vec{E}$, is a vector quantity that transmits the electric force to a particle. A charge or collection of charges creates an electric field at a point $P$. A test charge $q_{0}$ placed at point $P$ will be acted upon by a force

$$
\vec{F}=q_{0} \vec{E} .
$$

The magnitude of the electric field a distance $r$ from a point charge $q$ is

$$
\vec{E}=\frac{1}{4 \pi \epsilon_{0}} \frac{q}{r^{2}} \hat{r} .
$$

The direction of the electric field is along the line between the charge and $P$, pointing away from positive charges and toward negative charges. The electric field for a collection of charges is the vector sum of the electric fields for the individual charges:

$$
\vec{E}_{\text {total }}=\vec{E}_{1}+\vec{E}_{2}+\vec{E}_{3}+\cdots
$$

## Electric Field Lines

Electric field lines are a graphical representation of electric fields. The direction of the electric field at any point in space is tangent to the field line, and the magnitude of the electric field is proportional to the density or number of lines per unit area perpendicular to their direction. Electric field lines point away from positive charges and toward negative charges.

## Electric Dipoles

An electric dipole is a pair of equal and opposite charges of magnitude $q$ separated by a distance $d$. The electric dipole moment $\vec{q}$ has magnitude $q d$ and points from negative to positive charge. In an electric field, the electric dipole is acted upon by a torque

$$
\tau=p E \sin \phi,
$$

or

$$
\vec{\tau}=\vec{p} \times \vec{E}
$$

where $\phi$ is the angle between the electric dipole and the electric field. The potential energy for an electric dipole in an electric field is given by

$$
U=-\vec{p} \cdot \vec{E}
$$

## Conceptual Questions

## 1: Charges on pith balls

Three conducting pith balls are suspended from thin threads. Charged rods are then brought near or in contact with the pith balls. When you bring two pith balls $A$ and $B$ near a third pith ball $C$, you find that $A$ and $B$ repel each other whereas $A$ and $C$ attract each other. What can you conclude about the charges on the three pith balls?

## Solution

IDENTIFY, SET UP, AND EXECUTE The charge on pith balls $A$ and $B$ must be of the same sign, since the two objects repel each other. By contrast, because pith ball $C$ is attracted to pith ball $A$, either $A$ and $C$ are oppositely charged or $C$ is neutral. We cannot be sure that pith ball $C$ picked up a charge when the rods were brought near to it. For example, if pith ball $C$ was neutral and the rods were brought nearby but did not touch, pith ball $C$ would not acquire a charge. A neutral object will be inductively charged when brought near a charged object, causing like charges to be displaced away from the charge and unlike charges to be displaced toward the charge, resulting in an overall attraction. Oppositely charged objects will also attract. We cannot determine whether pith ball $C$ is neutral or oppositely charged from the information provided. Nor can we determine the sign of the charge on pith balls $A$ and $B$; we can be sure only that the charges on $A$ and $B$ are of the same sign. More experiments would have to be done to draw any additional conclusions.

EVALUATE We see how we must carefully interpret the data presented in problems involving charge. We must keep in mind that charge may or may not be transferred when the charged rods are brought near or in contact with the pith balls. An attraction between two objects does not indicate that the objects are oppositely charged, only that at least one object is charged. Solely with repulsion can we conclude that the charges on both objects are of the same sign, although we cannot determine whether that sign is positive or negative.

## 2: Where is the net force zero?

Charged balls $A$ and $B$ have charges $-Q$ and $+4 Q$, respectively, and are fixed at a separation distance of $R$, as shown in Figure 21.1. Is it possible to place another charged ball ( $C$ ) with charge $Q_{0}$ on the dashed line such that ball $C$ will be in electrostatic equilibrium? If so, indicate where to place the ball $C$. If not, explain why not.


Figure 21.1 Question 2.

## Solution

IDENTIFY, SET UP, AND EXECUTE The net force on ball $C$ is the sum of the forces due to the charges on balls $A$ and $B$. For equilibrium, the net force on $C$ must be zero. The forces due on $C$ to $A$ and $B$ will have to be in opposite directions for the net force to be zero. Let's examine the net force on $C$ in each of the three regions (between the balls, to the right of ball $B$, and to the left of ball $A$ ) to find whether an equilibrium position exists in any region. For simplicity, we'll take the charge on ball $C$ to be positive. We'll check our results to find the impact of having a negatively charged ball $C$.

If we place ball $C$ between balls $A$ and $B, C$ will be attracted to $A$ and repelled from $B$, resulting in both forces acting in the same direction, a situation that could not result in equilibrium. The same holds for a negatively charged ball $C$, but with the net force directed to the right. We conclude that there is no position between the balls that would result in equilibrium.

If we place ball $C$ to the right of ball $B, C$ will be attracted to ball $A$ and repelled from $B$, resulting in forces acting in opposite directions, a situation that could lead to equilibrium. However, $B$ has more charge than $A$, so $C$ would have to be placed closer to $A$ than to $B$ in order for the magnitudes of the forces to be equal (since Coulomb's law states that the force decreases with the square of the separation). For all points to the right of $B$, the force between $B$ and $C$ is larger than the force between $A$ and $C$. The same holds for a negatively charged $C$. We conclude that there is no position to the right of $B$ that would result in equilibrium.

If we place ball $C$ to the left of ball $A, C$ will be attracted to $A$ and repelled from ball $B$, resulting in forces acting in opposite directions, a situation that could lead to equilibrium. In the region to the left of ball $A$, we could place $C$ closer to $A$ than to $B$, giving balanced forces. If we place $C$ a distance $R$ to the left of $A$, then the magnitude of the force between $A$ and $C$ will be

$$
\left|F_{A C}\right|=\frac{1}{4 \pi \epsilon_{0}} \frac{Q Q_{0}}{R^{2}}
$$

The magnitude of the force between $B$ and $C$ will be

$$
\left|F_{B C}\right|=\frac{1}{4 \pi \epsilon_{0}} \frac{(4 Q) Q_{0}}{(2 R)^{2}}=\frac{1}{4 \pi \epsilon_{0}} \frac{Q Q_{0}}{R^{2}}
$$

The magnitudes of these forces are the same, and their directions are opposite each other, resulting in no net force. Therefore, this is the equilibrium position. It will also be an equilibrium position if charge $Q_{0}$ is negative.

EVALUATE This question combined concepts of vector addition with the radial dependence of Coulomb's law. Often, a quick guess will place the third charged ball between the other two balls without considering the direction of the net force on the third ball. As you can see, carefully considering the possibilities leads to the correct answer.

## 3: Direction of electric field for a collection of charges

Two thin, straight glass rods of equal length $L$ are placed perpendicular to each other with their ends almost touching, as shown in Figure 21.2. A charge $+Q$ is distributed uniformly along the top rod, and a charge of $-Q$ is distributed uniformly along the left rod. What is the direction of the electric field at a point $P$ located a distance $L / 2$ from each rod?


Figure 21.2 Question 3.

## Solution

IDENTIFY, SET UP, AND EXECUTE The electric field is a vector, so we'll find the electric field for each rod separately and sum the two electric fields. The electric field points away from positive charges and toward negative charges. The electric field for the top, positively charged rod points vertically downward at $P$. This is because $P$ is located halfway along the rod and the symmetric arrangement of the charges gives an electric field that is perpendicular to the rod. The electric field for the left, negatively charged rod points horizontally to the left at point $P$ for the same reason. The magnitude of the electric field for each rod is the same, since $P$ is located the same distance from each rod and the magnitude of the charge on each rod is the same. To find the resulting direction, we sketch the two components and their sum in Figure 21.3.


Figure 21.3 Question 3 sketch.
We see that the resulting electric field points downward and to the left in the figure, $45^{\circ}$ below the negative $x$-axis.

EVALUATE This problem illustrates where the electric field for a symmetric collection of charges points and how to combine electric fields. You should always consider what you expect for a resulting electric field before calculating the field. Later, you can check your results.

## 4: Symmetric Charges

Identical positive charges $Q$ are placed in the four corners of a square of side $L$. A fifth positive charge $q$ is placed at the center of the square. What is the force on the fifth charge?

## Solution

IDENTIFY, SET UP, AND EXECUTE Each of the four charges on the corner exert a force on the center charge. All four charges are equidistant from the charge at the center, so all exert forces of equal magnitude on the center charge. Charges in opposite corners exert forces in opposite directions, thus canceling the forces. The net force on the charge at the center is zero.

EVALUATE Symmetry is an important attribute to consider in electric force and field problems. We'll use symmetry throughout to simplify problem solving.

## Problems

## 1: Electric force with three charges

Three charges are arranged as shown in Figure 21.4: charge $A(+4.0 \mu \mathrm{C})$ is located on the $y$-axis at $y=25 \mathrm{~cm}$, charge $B(-6.0 \mu \mathrm{C})$ is located on the $y$-axis at $y=-25 \mathrm{~cm}$, and charge $C(+5.0 \mu \mathrm{C})$ is located on the $x$-axis at $x=15 \mathrm{~cm}$. Find the resulting force acting on charge $C$.


Figure 21.4 Problem 1.

## Solution

IDENTIFY We will compute the force on charge $C$ due to the other two charges and then find the vector sum of the two forces. The target variable is the force acting on charge $C$.

SET UP A free-body diagram showing the forces due to the two charges acting on charge $C$ is given in Figure 21.5 a . The force due to charge $A$ is repulsive and the force due to charge $B$ is attractive. Coulomb's law gives the magnitudes of the forces. Because the forces are vectors, we will add components to find the resultant. We have added an $x-y$ coordinate system to the free-body diagram and resolved the forces into their $x$ and $y$ components. Figure 21.5 b shows the resulting $x$ and $y$ components of the force and the resultant force vector.


Figure 21.5 Problem 1 free-body diagram.
EXECUTE We start by finding the distance between the charges so that we can the find the magnitudes of the two electric forces. The distance between charge $C$ and charge $A$ is the same as that between charge $C$ and charge $B$ and is

$$
r_{A C}=r_{B C}=\sqrt{(0.25 \mathrm{~m})^{2}+(0.15 \mathrm{~m})^{2}}=0.291 \mathrm{~m} .
$$

The magnitude of the force between charges $C$ and $A$ is given by Coulomb's law:

$$
F_{A C}=\frac{1}{4 \pi \epsilon_{0}} \frac{\left|q_{A} q_{C}\right|}{r_{A C}^{2}}=\left(8.99 \times 10^{9} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{C}^{2}\right) \frac{\left|\left(4.0 \times 10^{-6} \mathrm{C}\right)\left(5.0 \times 10^{-6} \mathrm{C}\right)\right|}{(0.291 \mathrm{~m})^{2}}=2.12 \mathrm{~N} .
$$

The magnitude of the force between charges $C$ and $B$ is

$$
F_{B C}=\frac{1}{4 \pi \epsilon_{0}} \frac{\left|q_{B} q_{C}\right|}{r_{B C}^{2}}=\left(8.99 \times 10^{9} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{C}^{2}\right) \frac{\left|\left(-6.0 \times 10^{-6} \mathrm{C}\right)\left(5.0 \times 10^{-6} \mathrm{C}\right)\right|}{(0.291 \mathrm{~m})^{2}}=3.18 \mathrm{~N} .
$$

To add the two vectors, we need the angles between the $x$-axis and the lines between charges $A$ and $C$ and charges $B$ and $C$. Examining the triangles formed by the $x$-axis, the $y$-axis, and the lines between the charges, we see that the triangles are similar, since they have sides of the same length. We find the angle between the $x$-axis and the lines between the charges by taking the inverse tangent of the ratio of the two sides of the triangles:

$$
\theta=\tan ^{-1} \frac{0.25 \mathrm{~m}}{0.15 \mathrm{~m}}=59.0^{\circ}
$$

We now add the $x$ and $y$ components of the two forces. The $x$ components sum to

$$
F_{x}=F_{A C} \cos 59.0^{\circ}-F_{B C} \cos 59.0^{\circ}=(2.12 \mathrm{~N}) \cos 59.0^{\circ}-(3.18 \mathrm{~N}) \cos 59.0^{\circ}=0.545 \mathrm{~N} .
$$

The $y$ components sum to

$$
F_{y}=-F_{A C} \sin 59.0^{\circ}-F_{B C} \sin 59.0^{\circ}=-(2.12 \mathrm{~N}) \sin 59.0^{\circ}-(3.18 \mathrm{~N}) \sin 59.0^{\circ}=-4.54 \mathrm{~N} .
$$

The magnitude of the net force is found by summing the squares of the components and taking the square root of the result:

$$
F=\sqrt{F_{x}^{2}+F_{y}^{2}}=\sqrt{(0.545 \mathrm{~N})^{2}+(-4.54 \mathrm{~N})^{2}}=4.57 \mathrm{~N} .
$$

The net force points downward and to the right. The angle between the positive $x$-axis and the net force is

$$
\phi=\tan ^{-1} \frac{F_{y}}{F_{x}}=\tan ^{-1} \frac{(-4.54 \mathrm{~N})}{(0.545 \mathrm{~N})}=-83.2^{\circ} .
$$

The resulting force acting on charge $C$ has magnitude 4.57 N and points $83.2^{\circ}$ below the positive $x$ axis, as shown in Figure 21.5b.
EVALUATE This problem shows how we can calculate the net force on one charge due to several charges. Once we identified and quantified the Coulomb forces, we followed the standard procedure for summing vectors, a skill we learned in the first half of the course. If the process is a bit unfamiliar to you, you should review and practice your vector-adding skills, since we'll be utilizing vectors throughout our exploration of electricity and magnetism.
Practice Problem: The problem could have been made simpler with symmetric charges. What charge should be substituted for charge $B$ to make the problem symmetric? Answer: $+4.0 \mu \mathrm{C}$, the same as charge $A$, giving a force with components only along the $x$-axis.

## 2: Charge on suspended pith balls

Two pith balls, shown in Figure 21.6, are charged, and each has a mass of 12.0 g . Each pith ball is suspended by an insulating thin thread of length $L=20.0 \mathrm{~cm}$. The balls are in equilibrium at $\theta=24.0^{\circ}$. If charge $A$ is twice charge $B$, find the magnitude of charge $A$.


Figure 21.6 Problem 2.

## Solution

IDENTIFY We'll apply Newton's first law to solve the problem. Coulomb's law will lead to the charge on $A$, the target variable.

SET UP Since the two pith balls are in equilibrium, the net force acting on either ball must be zero. A free-body diagram of charge $A$ (including an $x-y$ coordinate system) is shown in Figure 21.7. Three forces act on charge $A$ : the Coulomb force, the tension due to the string, and gravity. We'll ignore the mass of the string, since the string is thin and has mass much less than that of the pith ball.


Figure 21.7 Problem 2 free-body diagram.

EVALUATE We begin by quantifying the forces and then applying Newton's first law in two dimensions. Coulomb's law gives the magnitude of the electric force between the two charges:

$$
F_{E}=\frac{1}{4 \pi \epsilon_{0}} \frac{\left|q_{A} q_{B}\right|}{r_{A B}^{2}} .
$$

Since charge $A$ is twice charge $B$, we can replace $q_{B}$ with $\frac{1}{2} q_{A}$. The distance $r_{A B}$ between the charges is found to be

$$
r_{A B}=2 L \sin 24^{\circ}
$$

by examining Figure 21.6. Substituting into the electric force equation then yields

$$
F_{E}=\frac{1}{4 \pi \epsilon_{0}} \frac{\left|q_{A} q_{B}\right|}{r_{A B}^{2}}=\frac{1}{4 \pi \epsilon_{0}} \frac{\frac{1}{2} q_{\mathrm{A}}^{2}}{\left(2 L \sin 24^{\circ}\right)^{2}}=\frac{1}{4 \pi \epsilon_{0}} \frac{q_{\mathrm{A}}^{2}}{8 L^{2} \sin ^{2} 24^{\circ}}
$$

The absolute-value sign is removed, since $q_{A}$ is squared. We now apply Newton's first law to the charges in equilibrium. In the $x$ direction, the forces in the horizontal component are tension and the electric force. These must sum to zero:

$$
\sum F_{x}=0=F_{E}-T \sin 24^{\circ}
$$

Solving for tension, we obtain

$$
T=\frac{F_{E}}{\sin 24^{\circ}} .
$$

In the $y$ direction, the forces in the vertical component are tension and gravity. These also must sum to zero:

$$
\sum F_{y}=0=T \cos 24^{\circ}-m g .
$$

Substituting for the tension yields

$$
F_{E} \cot 24^{\circ}=m g
$$

Rearranging terms then gives

$$
F_{E}=\frac{1}{4 \pi \epsilon_{0}} \frac{q_{\mathrm{A}}^{2}}{8 L^{2} \sin ^{2} 24^{\circ}}=m g \tan 24^{\circ} .
$$

Solving for $q_{A}$, we get

$$
\begin{aligned}
q_{A} & =2 L \sin 24^{\circ} \sqrt{\frac{2 m g \tan 24^{\circ}}{k}} \\
& =2(0.20 \mathrm{~m}) \sin 24^{\circ} \sqrt{\frac{2(0.012 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \tan 24^{\circ}}{\left(8.99 \times 10^{9} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{C}^{2}\right)}}=5.55 \times 10^{-7} \mathrm{C}
\end{aligned}
$$

The charge on $A$ is $0.555 \mu \mathrm{C}$.
EVALUATE This second problem should also be a reminder of problems encountered in the first half of the course. The only new component is the electrostatic force; all other steps are similar to those we've used before.

Although we can find the magnitudes of the charges, we cannot find their signs. We only know that both charges must have the same sign, since they repel.

## 3: Finding charges, given an electric field

Two charges placed on the $y$-axis create a resultant electric field of $145,000 \mathrm{~N} / \mathrm{C}$ pointing along the $x$ axis at a point $P$ on the $x$-axis, as shown in Figure 21.8. Both charges are placed 50.0 cm from the origin and the electric field is measured at $x=40.0 \mathrm{~cm}$. Find the magnitudes and signs of the two charges.


Figure 21.8 Problem 3.

## Solution

IDENTIFY The target variables are the signs and magnitudes of the two charges creating the electric field.

SET UP The electric field due to each charge is shown in Figure 21.9. The triangles formed by the origin, the charge, and point $P$ are similar for each charge. Therefore, both charges are equidistant from point $P$, and the angle between the $x$-axis and the line connecting point $P$ and each charge is the same. Since the resulting electric field is along the $x$-axis, the vertical components of the electric field due to each charge must be equal and the charges must be of the same magnitude. Since the field points away from the charges along the $x$-axis, both charges must be positive.

With this information, we can proceed to find the magnitude of the charge. We'll call the charge $q$ and find the electric field at $P$. Knowing the magnitude of the field will allow us to solve for the magnitude of the charge.


Figure 21.9 Problem 3 sketch of the electric fields.
EXECUTE The magnitude of the electric field due to a point charge is given by

$$
E=\frac{1}{4 \pi \epsilon_{0}} \frac{|q|}{r^{2}}=k \frac{|q|}{r^{2}} .
$$

The magnitude of the electric field at $P$ due to the upper charge is then

$$
E_{\text {upper }}=k \frac{|q|}{r^{2}}=k \frac{q}{\left(x^{2}+y^{2}\right)},
$$

where we removed the absolute-value designation from $q$ ( $q$ is positively charged) and replaced the squared distance to the charge by the sum of the squared components along the $x$ - and $y$-axis. The $x$ component of the electric field due to the upper charge is

$$
E_{\text {upper }, x}=E_{\text {upper }} \cos \theta=k \frac{q}{x^{2}+y^{2}} \frac{x}{\sqrt{x^{2}+y^{2}}}=k \frac{q x}{\left(x^{2}+y^{2}\right)^{3 / 2}},
$$

where we substituted for the cosine. The $x$ component of the electric field due to the lower charge is identical, since the lower charge has the same magnitude, is located the same distance away, and is positioned symmetrically to the upper charge. The total $x$ component of the electric field due to both charges is twice the electric field due to the upper charge:

$$
E_{\text {total }, x}=2 E_{\text {upper }, x}=2 k \frac{q x}{\left(x^{2}+y^{2}\right)^{3 / 2}} .
$$

Solving for the charge yields

$$
q=E_{\text {total }, x} \frac{\left(x^{2}+y^{2}\right)^{3 / 2}}{2 k x}=(145,000 \mathrm{~N} / \mathrm{C}) \frac{\left((0.50 \mathrm{~m})^{2}+(0.40 \mathrm{~m})^{2}\right)^{3 / 2}}{2\left(8.99 \times 10^{9} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{C}^{2}\right)(0.40 \mathrm{~m})}=+5.29 \times 10^{-6} \mathrm{C}
$$

Each charge is $+5.29 \mu \mathrm{C}$.
EVALUATE This problem shows how to find the magnitudes of charges given the locations of the charges and the net electric field they produce. Most of the effort expended in solving the problem was determining that both charges have the same sign and magnitude. Once that was established, completing the problem was easy.

## 4: Electric field on the $\boldsymbol{x}$-axis

Two charges, $Q$ and $2 Q$, are placed on the $y$-axis as shown in Figure 21.10. Each charge is located a distance $a$ above or below the $x$-axis. Find the magnitude of the electric field everywhere on the $x$-axis.


Figure 21.10 Problem 4.

## Solution

IDENTIFY The target variable is the electric field on the $x$-axis.
SET UP The electric fields due to each charge are shown in Figure 21.11. The triangles formed by the origin, the charge, and the point on the $x$-axis are similar for each charge. Therefore, both charges are equidistant from the point on the axis, and the angle between the axis and the line to the charges is the same. The magnitude of the field due to each charge at the point on the axis is different, so the field due to each charge is different. We'll add the two fields to find the solution.


Figure 21.11 Problem 4 sketch of the electric fields.

EXECUTE The magnitude of the electric field due to a point charge is given by

$$
E=\frac{1}{4 \pi \epsilon_{0}} \frac{|q|}{r^{2}} .
$$

The magnitude of the electric field due to the upper charge is then

$$
E_{\text {upper }}=\frac{1}{4 \pi \epsilon_{0}} \frac{Q}{r^{2}}=\frac{1}{4 \pi \epsilon_{0}} \frac{Q}{\left(x^{2}+a^{2}\right)} .
$$

The magnitude of the electric field due to the lower charge is

$$
E_{\text {lower }}=\frac{1}{4 \pi \epsilon_{0}} \frac{2 Q}{r^{2}}=\frac{1}{4 \pi \epsilon_{0}} \frac{2 Q}{\left(x^{2}+a^{2}\right)} .
$$

We must add these electric fields as vectors, so we need their components. To find the components, we need the sine and cosine of the angles, given respectively by

$$
\begin{aligned}
\sin \theta & =\frac{a}{\sqrt{x^{2}+a^{2}}} \\
\cos \theta & =\frac{x}{\sqrt{x^{2}+a^{2}}}
\end{aligned}
$$

The $x$ component of the electric field due to the charges is

$$
\begin{aligned}
E_{x} & =E_{\text {upper }} \cos \theta-E_{\text {lower }} \cos \theta \\
& =\frac{1}{4 \pi \epsilon_{0}} \frac{Q}{x^{2}+a^{2}} \frac{x}{\sqrt{x^{2}+a^{2}}}-\frac{1}{4 \pi \epsilon_{0}} \frac{2 Q}{x^{2}+a^{2}} \frac{x}{\sqrt{x^{2}+a^{2}}} \\
& =-\frac{1}{4 \pi \epsilon_{0}} \frac{Q x}{\left(x^{2}+a^{2}\right)^{3 / 2}} .
\end{aligned}
$$

Note how the $x$ components are in opposite directions. Similarly, the $y$ component of the electric field due to the charges is

$$
\begin{aligned}
E_{y} & =-E_{\text {upper }} \sin \theta-E_{\text {lower }} \sin \theta \\
& =-\frac{1}{4 \pi \epsilon_{0}} \frac{Q}{x^{2}+a^{2}} \frac{a}{\sqrt{x^{2}+a^{2}}}-\frac{1}{4 \pi \epsilon_{0}} \frac{2 Q}{x^{2}+a^{2}} \frac{a}{\sqrt{x^{2}+a^{2}}} \\
& =-\frac{1}{4 \pi \epsilon_{0}} \frac{3 Q a}{\left(x^{2}+a^{2}\right)^{3 / 2}} .
\end{aligned}
$$

To find the magnitude, we sum the squares of the components and take the square root, yielding

$$
\begin{aligned}
E & =\sqrt{E_{x}^{2}+E_{y}^{2}} \\
& =\sqrt{\left(-\frac{1}{4 \pi \epsilon_{0}} \frac{Q x}{\left(x^{2}+a^{2}\right)^{3 / 2}}\right)^{2}+\left(\frac{1}{4 \pi \epsilon_{0}} \frac{3 Q a}{\left(x^{2}+a^{2}\right)^{3 / 2}}\right)^{2}} \\
& =\frac{1}{4 \pi \epsilon_{0}} \frac{Q}{\left(x^{2}+a^{2}\right)^{3 / 2}} \sqrt{x^{2}+9 a^{2}} .
\end{aligned}
$$

EVALUATE This problem shows how to find the magnitude of the electric field, given the locations and magnitudes of the charges that produce the field. Most of our effort was spent combining vectors; the electric field calculation was straightforward.

## 5: Motion of a charge in a uniform electric field

An electron with velocity $2.00 \times 10^{7} \mathrm{~m} / \mathrm{s}$ is projected into a uniform electric field between parallel plates as shown in Figure 21.12. The electric field points vertically upward and is zero outside the plates. If the electron enters the field near the top plate as shown and emerges from the other side after missing the bottom plate, find the maximum electric field between the plates.


Figure 21.12 Problem 5.

## Solution

IDENTIFY We will combine our knowledge of electric forces and fields, Newton's laws, and constantacceleration kinematics to solve this problem. The target variable is the maximum electric field.

SET UP The free-body diagram of the charge in the electric field is shown in Figure 21.13. Only one force acts on the electron. (We'll ignore gravity, since the electron's mass is minuscule.) The force is due to the electric field, which is constant; therefore, the electric force and acceleration are constant in the region between the plates. Since the acceleration is constant, we can apply the kinematics relations of Chapter 3 to the problem. The force acts in the vertical direction, and there is no acceleration in the horizontal direction.

The maximum electric field corresponds to the electron just missing the edge of the bottom plate as it emerges. A larger electric field would result in a larger force and acceleration, causing the electron to strike the bottom plate.


Figure 21.13 Problem 5 free-body diagram.

EXECUTE The time required for the electron to pass between the plates is found by examining the horizontal motion. Since there is no acceleration in the $x$ direction, the time required is equal to the horizontal distance the electron travels divided by its initial horizontal velocity:

$$
\Delta t=\frac{\Delta d}{v_{0, x}}=\frac{(0.030 \mathrm{~m})}{\left(2.00 \times 10^{7} \mathrm{~m} / \mathrm{s}\right)}=1.50 \times 10^{-9} \mathrm{~s}
$$

We need to find the acceleration in the vertical direction that corresponds to the electron's moving from the top plate to the bottom plate in the time it takes to pass between the plates. The initial vertical velocity is zero, so

$$
y=y_{0}+v_{0, y} t+\frac{1}{2} a_{y} t^{2}=\frac{1}{2} a_{y} t^{2}
$$

where we have taken the origin to be the left edge of the top plate $\left(y_{0}=0\right)$ and we take positive $y$ downward. The vertical acceleration is then

$$
a_{y}=\frac{2 y}{t^{2}}=\frac{2(0.020 \mathrm{~m})}{\left(1.50 \times 10^{-9} \mathrm{~s}\right)^{2}}=1.78 \times 10^{16} \mathrm{~m} / \mathrm{s}^{2}
$$

The net force is the electron's mass times its acceleration and is also the electric field times the electron's charge. The magnitude of the electric field is then

$$
E=\frac{F_{E}}{q}=\frac{m a}{q}=\frac{m a}{|e|}=\frac{\left(9.11 \times 10^{-31} \mathrm{~kg}\right)\left(1.78 \times 10^{16} \mathrm{~m} / \mathrm{s}^{2}\right)}{\left(1.60 \times 10^{-19} \mathrm{C}\right)}=1.01 \times 10^{5} \mathrm{~N} / \mathrm{C}
$$

The maximum electric field between the plates is $101,000 \mathrm{~N} / \mathrm{C}$.
EVALUATE This problem resembles the problems involving motion under constant acceleration that we initially encountered in Chapter 3, but with the acceleration now resulting from the electric force. As we' ve seen, many problems are similar to problems we've completed earlier, but modified here to include the electric force. You may want to review the earlier material before continuing with problems in the current chapter.

## 6: Electric field of a charge distribution

Find the electric field along the axis at a distance $a$ from the end of a straight insulating rod of length $L$. Charge $+Q$ is evenly distributed along the rod's length.

## Solution

IDENTIFY Charge is distributed throughout the rod, so the rod must be split into infinitesimal segments and the resulting equations must be integrated to find the field. The target variable is the electric field on the axis.

SET UP Figure 21.14 shows a sketch of the situation, including the electric field $d E$ for a small segment of charge $d q$. We will find the charge in the segment and integrate from $x=a$ to $x=a+L$ to solve. The electric field points away from the rod and along the $x$-axis.


Figure 21.14 Problem 6 sketch.
EXECUTE The magnitude of the electric field due to a point charge is given by

$$
E=\frac{1}{4 \pi \epsilon_{0}} \frac{|q|}{r^{2}}
$$

The magnitude of the electric field due to the segment of charge $d q$ on the $x$-axis is then

$$
d E=\frac{1}{4 \pi \epsilon_{0}} \frac{d q}{x^{2}}
$$

The charge in segment $d q$ is the total charge of the rod, divided by the length, multiplied by the length of the segment. This is equivalent to the linear charge density (charge per unit length) times the length of $d q$. The segment has length $d x$, so

$$
d q=\frac{Q}{L} d x
$$

Replacing $d q$ yields

$$
d E=\frac{Q}{4 \pi \epsilon_{0}} L \frac{d x}{x^{2}} .
$$

We integrate this electric field expression between the two endpoints of the rod to find the magnitude of the electric field:

$$
E=\int d E=\int_{a}^{a+L} \frac{Q}{4 \pi \epsilon_{0} L} \frac{d x}{x^{2}}=\left.\frac{Q}{4 \pi \epsilon_{0} L} \frac{-1}{x^{1}}\right|_{a} ^{a+L}=\frac{Q}{4 \pi \epsilon_{0} L}\left(\frac{1}{a}-\frac{1}{a+L}\right)
$$

The electric field is directed away from the end of the $\operatorname{rod}$ (toward negative $x$ in the sketch in Figure 21.14).

EVALUATE This problem illustrates how to find the electric field by integrating contributions to the field. You must carefully draw a sketch and set up the variables when you encounter this type of problem.

Why was the rod made of insulating material in this problem? Because charges don't move on an insulator, thus allowing the charge to be distributed uniformly along the length of the rod. Charge would redistribute itself in a conducting rod.

## Try It Yourself!

## 1: Moving electrons

Two small, neutral spheres are separated by 20.0 cm . Electrons are removed from one sphere and placed on the second sphere. After the electrons are transferred, the force is found to be $2.0 \times 10^{-5} \mathrm{~N}$. How many electrons were transferred?

## Solution Checkpoints

IDENTIFY AND SET UP: Use Coulomb's law to find the magnitude of the charge on the spheres. Both spheres have equal charges. (Why?)

EXECUTE The magnitude of the force is

$$
F=\frac{1}{4 \pi \epsilon_{0}} \frac{q q}{r^{2}} .
$$

This gives a charge of $9.43 \times 10^{-9} \mathrm{C}$, or $5.9 \times 10^{10}$, electrons.
EVALUATE Many electrons are required to create a relatively small force between two spheres. How many electrons must be transferred to create a force of 1 N ?

## 2: Electric field due to two charges

Two charges, each of magnitude 0.10 nC , are placed on the $y$-axis as shown in Figure 21.15. Find the electric field on the $x$-axis at $x=+10.0 \mathrm{~cm}$.


Figure 21.15 Try It Yourself 2.

## Solution Checkpoints

IDENTIFY AND SET UP: Use Coulomb's law to find the electric field due to each charge, and then sum the fields. Start with a free-body diagram. Can you use symmetry to simplify the solution?

EXECUTE The magnitude of the electric field due to a point charge is given by

$$
E=\frac{1}{4 \pi \epsilon_{0}} \frac{|q|}{r^{2}}
$$

The net electric field is along the $x$-axis at $x=+10.0 \mathrm{~cm}$. (Why?) The cosine of the angle is given by

$$
\cos \theta=\frac{x}{\sqrt{x^{2}+a^{2}}}
$$

The $x$ component of the electric field due to the charges is

$$
\begin{aligned}
E_{x} & =E_{\text {upper }} \cos \theta+E_{\text {lower }} \cos \theta \\
& =\frac{1}{4 \pi \epsilon_{0}} \frac{2 Q x}{\left(x^{2}+a^{2}\right)^{3 / 2}}
\end{aligned}
$$

Evaluating the result gives $E=1.58 \times 10^{2} \mathrm{~N} / \mathrm{C}$. In which direction does the electric field point?
EVALUATE This problem is similar to Problem 4, although you can apply symmetry here.

## 3: Electric field due to a charge distribution

Find the electric field along the axis of, and a distance $z$ away from, a circular insulating disk of radius $a$ having a total charge $Q$ distributed uniformly across its surface.

## Solution Checkpoints

IDENTIFY AND SET UP: Charge is distributed throughout the disk, so it must be split into infinitesimal segments and the resulting equations integrated to find the field. Sketch the situation and decide how to integrate.

The disk can be split into thin concentric rings of width $d r$ and length $2 \pi r$. The charge on the ring is

$$
d q=\frac{Q}{\pi a^{2}} 2 \pi r d r
$$

EXECUTE The magnitude of the electric field due to the ring of charge $d q$ on the axis is

$$
d E=\frac{1}{4 \pi \epsilon_{0}} \frac{Q}{\pi a^{2}} \frac{2 \pi r d r}{r^{2}+z^{2}} \cos \theta
$$

Find the cosine of the angle and write it in terms of the given variables. Integrating then gives

$$
E=\int_{0}^{a} \frac{1}{4 \pi \epsilon_{0}} \frac{Q}{\pi a^{2}} \frac{z 2 \pi r d r}{\left(r^{2}+z^{2}\right)^{3 / 2}}=\frac{z Q}{2 \pi \epsilon_{0} A^{2}}\left(\frac{1}{z}-\frac{1}{\sqrt{z^{2}+a^{2}}}\right)
$$

The electric field is directed away from the disk.
EVALUATE This problem illustrates how to find the electric field due to a continuous charge distribution.

## Problem Summary

Chapter 21 adds the electric interaction to our physics knowledge base-a new force to our inventory of concepts. Electricity builds on our foundation developed in the previous chapters, allowing us to integrate this new force into our well-developed problem-solving techniques, which remain consistent:

- Identifying the general procedure to find the solution.
- Sketching the situation when no figure is provided.
- Identifying the forces and torques acting on the bodies.
- Identifying the forms of energy included in the problem.
- Drawing free-body diagrams of the bodies.
- Applying appropriate coordinate systems to the diagrams.
- Applying the equations of motion to find relations among the forces, masses, and accelerations.
- Applying conservation of energy or conservation of momentum as appropriate.
- Solving the equations through algebra and substitutions.
- Reflecting on the results and checking for inconsistencies.

You should take advantage of your success in the first half of the course as you ponder the new forces we will examine in the second half.

## 22 <br> Gauss's Law

## Summary

In this chapter, we investigate Gauss's law, an elegant tool that helps us calculate electric fields for symmetric systems while affording an insight into those very electric fields. We begin by examining electric flux, or the "flow" of electric field lines through a surface. We'll see how Gauss's law relates the electric flux through a surface to the charge enclosed within that surface. We will then apply Gauss's law to symmetric systems to find the electric field inside and around those systems. We'll also look at the implications of Gauss's law and learn how charges distribute themselves on conductors. We'll find that the electric field inside any conductor is zero.

## Objectives

After studying this chapter, you will understand

- How to calculate the electric flux passing through a surface.
- The definition and conceptual application of Gauss's law.
- How to apply Gauss's law to spherically and cylindrically symmetric charge distributions.
- How charges are arranged on conductors.
- Why the electric field inside a conductor is zero for static charges.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Electric Flux | Electric flux is a measure of the "flow" of electric field through a surface. <br> Flux is found by integrating the electric field over the surface: |
| $\qquad$The electric flux is proportional to the number of electric field lines crossing <br> the surface. |  |
| Gauss's Law | Gauss's law states that the total electric flux through any closed volume is <br> proportional to the total charge $Q_{\text {encl }}$ enclosed inside the volume: |
| $\qquad \Phi_{E}=\oint \vec{E} \cdot d \vec{A}=\int E_{\perp} d A=\int E \cos \phi d A$. |  |
| $Q_{\text {encl }}$. |  |

## Conceptual Questions

## 1: Electric flux through a cylinder

A positive charge is placed just above a cylinder as shown in Figure 22.1. What is the sign of the electric flux through the top of the cylinder, the sides of the cylinder, and the bottom of the cylinder due to the charge? What is the total electric flux through the cylinder due to the charge?


Figure 22.1 Question 1.

## Solution

IDENTIFY, SET UP, AND EXECUTE Electric flux is proportional to the number of electric field lines passing through a surface. If we consider how the electric field lines cross the surfaces of the cylinder, we can find the electric flux through the surfaces.

Examining the top surface, we see that electric field lines point into the surface, resulting in a negative electric flux. Looking at the bottom surface, we see that electric field lines point out of the surface, resulting in a positive electric flux. The sides of the cylinder are a bit more complicated, but if you examine them carefully, you'll see that the electric field lines point out of the cylinder's sides. The electric flux through the side of the cylinder is positive.

The net electric flux through the whole cylinder is zero. You can see this by examining how all the field lines that enter the cylinder exit the cylinder, resulting in zero net flux. This finding agrees with Gauss's law: Since there is no charge enclosed inside the cylinder, the net electric flux is zero.

EVALUATE This problem shows how we can interpret electric flux as the number of electric field lines passing through a surface. Such an interpretation can be useful in calculating the electric flux and reflecting on your results.

## 2: Electric flux through several surfaces

Three small spheres, depicted in Figure 22.2, carry charges $q_{1}=5.0 \mu \mathrm{C}, q_{2}=-3.0 \mu \mathrm{C}$, and $q_{3}=2.0 \mu \mathrm{C}$, respectively. Find the net electric flux through the three surfaces $S_{1}, S_{2}$, and $S_{3}$, shown in the figure.


Figure 22.2 Question 2.

## Solution

IDENTIFY AND SET UP: Gauss's law gives the net flux through a surface as the charge enclosed by the surface divided by $\epsilon_{0}$. Our task is to identify the net charge enclosed by each surface and divide that number by $\epsilon_{0}$.

EXECUTE Surface $S_{1}$ encloses charges $q_{1}$ and $q_{2}$, so the electric flux is

$$
\Phi_{E}=\frac{Q_{\mathrm{encl}}}{\epsilon_{0}}=\frac{q_{1}+q_{2}}{\epsilon_{0}}=\frac{(5.0 \mu \mathrm{C})+(-3.0 \mu \mathrm{C})}{\epsilon_{0}}=\frac{2.0 \mu \mathrm{C}}{\epsilon_{0}}=2.26 \times 10^{5} \mathrm{~N} / \mathrm{C}
$$

Surface $S_{2}$ encloses charges $q_{2}$ and $q_{3}$, so the electric flux is

$$
\Phi_{E}=\frac{Q_{\mathrm{encl}}}{\epsilon_{0}}=\frac{q_{2}+q_{3}}{\epsilon_{0}}=\frac{(-3.0 \mu \mathrm{C})+(2.0 \mu \mathrm{C})}{\epsilon_{0}}=\frac{-1.0 \mu \mathrm{C}}{\epsilon_{0}}=-1.13 \times 10^{5} \mathrm{~N} / \mathrm{C}
$$

Surface $S_{3}$ encloses all three charges, so the electric flux is

$$
\Phi_{E}=\frac{Q_{\mathrm{encl}}}{\epsilon_{0}}=\frac{q_{1}+q_{2}+q_{3}}{\epsilon_{0}}=\frac{(5.0 \mu \mathrm{C})+(-3.0 \mu \mathrm{C})+(2.0 \mu \mathrm{C})}{\epsilon_{0}}=\frac{0.0 \mu \mathrm{C}}{\epsilon_{0}}=0
$$

EVALUATE We see that Gauss's law provides a straightforward method of evaluating the flux through a surface: Simple addition and division produces the results. We see also that the flux can be positive, negative, or zero.

## 3: Electric field where electric flux is zero

Is the electric field zero along a surface where the electric flux is zero?

## Solution

IDENTIFY, SET UP, AND EXECUTE The total electric flux is zero whenever the net charge inside a surface is zero. The electric field may or may not be zero if the total electric flux is zero. Consider the surface $S_{3}$ from the last question; the net flux was zero, but the field would not be zero along the surface. Near charge $q_{3}$, the electric field would be pointing outward, and near charge $q_{2}$, the electric field would be pointing inward.

EVALUATE Another interpretation of the electric flux is that it is equal to the number of electric field lines out of a surface minus the number of electric field lines into the surface. If the number of electric field lines into the surface is equal to the number of electric field lines out of the surface, then the electric flux is zero.

## Problems

## 1: Electric field for a spherical shell

A conducting spherical shell with inner radius $a$ and outer radius $b$ contains a charge $+2 Q$. Find the electric field in all regions of space.

## Solution

IDENTIFY The system is spherically symmetric, so we will apply Gauss's law to a spherical surface. The target variable is the electric field in all space.
SET UP Figure 22.3 shows a cutaway sketch of the spherical shell. The charge is located on the outer surface of the sphere, since it is a conductor. To find the electric field in any region, we place a spherical Gaussian surface in that region. Inside radius $b$, there is no field, since the charge is on the surface at that radius. We only need to calculate the electric field outside the radius $b$. Outside the shell, the electric field points radially outward and is constant at a given radius.


Figure 22.3 Problem 1 sketch.
EXECUTE We start by imagining a spherical Gaussian surface at a radius $r$ outside of the shell, as shown in the figure. We evaluate Gauss's law at that surface. Gauss's law states that

$$
\oint \vec{E} \cdot d \vec{A}=\frac{Q_{\mathrm{encl}}}{\epsilon_{0}}
$$

The charge enclosed is the charge on the shell, $+2 Q$. We evaluate the surface integral by examining the electric field outside the shell. The electric field is directed radially outward, as is the surface area element $d A$, so their dot product is just the product of their magnitudes:

$$
\oint \vec{E} \cdot d \vec{A}=\oint E d A
$$

The electric field is constant along the Gaussian surface, so we may move it outside of the integral:

$$
\oint \vec{E} \cdot d \vec{A}=\oint E d A=E \oint d A
$$

The integral of $d A$ is just the area of the spherical surface, $4 \pi r^{2}$ :

$$
\oint \vec{E} \cdot d \vec{A}=\oint E d A=E \oint d A=E 4 \pi r^{2}
$$

Summing up these results, we have

$$
\oint \vec{E} \cdot d \vec{A}=E 4 \pi r^{2}=\frac{+2 Q}{\epsilon_{0}}
$$

Solving for the electric field yields

$$
E=\frac{Q}{2 \pi \epsilon_{0} r^{2}}
$$

The electric field is directed radially outward.
EVALUATE We see how, outside the conducting shell, the electric field is the same as that of a $+2 Q$ point charge located at the center of the shell.

## 2: Electric field for a cylindrical insulator

A long, nonconducting cylinder with radius $R$ contains a uniform volume charge density $\rho$. Find the electric field inside and outside of the cylinder.

## Solution

IDENTIFY The system is cylindrically symmetric, so we will apply Gauss's law to a cylindrical surface. The target variable is the electric field.
SET UP Figure 22.4 shows a sketch of the cylinder with two cylindrically symmetric Gaussian surfaces, one inside the cylinder and one outside. The charge is distributed throughout the cylinder, so we will sum the charge inside the surface to find the field. Inside the cylinder, the Gaussian surface contains only a fraction of the charge. Outside the cylinder, the Gaussian surface contains the total charge.


Figure 22.4 Problem 2 sketch.

EXECUTE We start by imagining a cylindrical Gaussian surface of length $L$ at a radius $r$ inside the cylinder, as shown in the figure. We evaluate Gauss's law at that surface. Gauss's law states that

$$
\oint \vec{E} \cdot d \vec{A}=\frac{Q_{\mathrm{encl}}}{\epsilon_{0}} .
$$

The charge enclosed is the charge within the cylinder and is equal to the volume charge density times the volume of the cylinder enclosed by the surface, or

$$
Q_{\mathrm{encl}}=\rho V_{\mathrm{encl}}=\rho \pi r^{2} L
$$

We evaluate the surface integral by first examining the electric field inside the cylinder. The electric field points outward, perpendicular to the axis of the cylinder. The surface integral can be broken into three pieces: the curved sidewall and the two flat ends. Mathematically,

$$
\oint \vec{E} \cdot d \vec{A}=\int_{\text {curved side }} \vec{E} \cdot d \vec{A}+\int_{\text {left end }} \vec{E} \cdot d \vec{A}+\int_{\text {right end }} \vec{E} \cdot d \vec{A}
$$

At the two flat ends, the electric field is perpendicular to the area vector, so the dot product is zero. On the curved side, the electric field is directed outward and is parallel to the area vector. This gives

$$
\oint \vec{E} \cdot d \vec{A}=\int_{\text {curved side }} E d A+0+0
$$

The electric field is constant along the curved side and so may be taken outside of the integral. The integral becomes the integral of the area of the curved side, $2 \pi r L$. Combining gives the net electric flux:

$$
\oint \vec{E} \cdot d \vec{A}=E 2 \pi r L
$$

Summing up our results, we have

$$
\oint \vec{E} \cdot d \vec{A}=E 2 \pi r L=\frac{\rho \pi r^{2} L}{\epsilon_{0}} .
$$

Solving for the electric field inside the cylinder yields

$$
E=\frac{\rho r}{2 \epsilon_{0}} .
$$

The electric field is directed outward.
We follow the same procedure to find the field outside of the cylinder. The field points in the same direction, so the flux integral is the same as that inside the cylinder. The charge enclosed is now the full charge inside the cylinder, $\rho 2 \pi R^{2} L$. So we have

$$
\oint \vec{E} \cdot d \vec{A}=E 2 \pi r L=\frac{\rho \pi R^{2} L}{\epsilon_{0}},
$$

which, after we solve for the electric field, becomes

$$
E=\frac{\rho R^{2}}{2 r \epsilon_{0}} .
$$

The electric field is directed outward.
EVALUATE Inside the cylinder, the electric field increases with the radius. Outside the cylinder, the electric field decreases inversely with the radius.

## 3: Electric field for a spherical insulating shell

A hollow insulating spherical shell with inner radius $a$ and outer radius $b$ contains a charge $+Q$ uniformly distributed throughout its volume. Find the electric field in all regions of space.

## Solution

IDENTIFY The system is spherically symmetric, so we will apply Gauss's law to a spherical surface. The target variable is the electric field.

SET UP Figure 22.5 shows a cutaway sketch of the spherical shell. The charge is distributed uniformly throughout the shell, since the shell is an insulator. To find the electric field in any region, we place a spherical Gaussian surface in that region. Inside radius $a$, there is no field, because no charge is located inside the hollow shell. We only need to calculate the electric field inside and outside of the shell. The electric field points radially outward and is constant at a given radius.


Figure 22.5 Problem 3 sketch.
EXECUTE We start by imagining a spherical Gaussian surface at a radius $r$ for $a<r<b$, as shown in the figure. We evaluate Gauss's law at that surface. Gauss's law states that

$$
\oint \vec{E} \cdot d \vec{A}=\frac{Q_{\mathrm{encl}}}{\epsilon_{0}}
$$

The charge enclosed is the charge within the Gaussian surface and is equal to the volume charge density times the volume of the cylinder enclosed by the surface. The volume charge density is the total charge divided by the volume of the shell. The volume of the shell is the volume at radius $b$ minus the volume at radius $a$. The charge density is then

$$
\rho=\frac{Q}{\frac{4}{3} \pi b^{3}-\frac{4}{3} \pi a^{3}} .
$$

The volume enclosed by the Gaussian surface is the volume enclosed by the surface minus the volume at the radius $a$. The charge enclosed is given by

$$
Q_{\mathrm{encl}}=\rho V_{\mathrm{encl}}=\left(\frac{Q}{\frac{4}{3} \pi b^{3}-\frac{4}{3} \pi a^{3}}\right)\left(\frac{4}{3} \pi r^{3}-\frac{4}{3} \pi a^{3}\right)=Q\left(\frac{r^{3}-a^{3}}{b^{3}-a^{3}}\right)
$$

We evaluate the surface integral by examining the electric field outside the shell. The electric field is directed radially outward, as is the surface area element $d A$, so their dot product is just the product of their magnitudes:

$$
\oint \vec{E} \cdot d \vec{A}=\oint E d A .
$$

The electric field is constant along the Gaussian surface and so may be moved outside of the integral:

$$
\oint \vec{E} \cdot d \vec{A}=\oint E d A=E \oint d A
$$

The integral of $d A$ is just the surface area of the spherical surface, $4 \pi r^{2}$ :

$$
\oint \vec{E} \cdot d \vec{A}=\oint E d A=E \oint d A=E 4 \pi r^{2}
$$

Summing up these results, we have

$$
\oint \vec{E} \cdot d \vec{A}=E 4 \pi r^{2}=\frac{Q}{\epsilon_{0}}\left(\frac{r^{3}-a^{3}}{b^{3}-a^{3}}\right) .
$$

Solving for electric field inside the shell yields

$$
E=\frac{Q}{4 \pi \epsilon_{0} r^{2}}\left(\frac{r^{3}-a^{3}}{b^{3}-a^{3}}\right)
$$

The electric field is directed radially outward.
We follow the same procedure to find the field outside of the shell. The field points in the same direction, so the flux integral is the same as that inside the shell. The charge enclosed is now the full charge $Q$ inside the cylinder. So we have

$$
\oint \vec{E} \cdot d \vec{A}=E 4 \pi r^{2}=\frac{Q}{\epsilon_{0}}
$$

which, after we solve for the electric field outside of the shell, becomes

$$
E=\frac{1}{4 \pi \epsilon_{0}} \frac{Q}{r^{2}} .
$$

The electric field is directed outward.
EVALUATE By this point, you should see a common theme developing: Problems involving Gauss's law are solved by finding the flux through the surface and the charge enclosed inside the surface. The flux integral reduces to one of three choices, given the symmetry we have encountered. The challenge is often in determining the charge enclosed. Frequently, we can find the enclosed charge by taking ratios of volumes.

## 4: Electric field for a combination of thin sheets

Three very large nonconducting thin sheets respectively carry uniform charge densities $\sigma_{1}, \sigma_{2}$, and $\sigma_{3}$ on their surfaces, as shown in Figure 22.6. If $\sigma_{1}=+4.00 \mu \mathrm{C} / \mathrm{m}^{2}, \sigma_{2}=-6.00 \mu \mathrm{C} / \mathrm{m}^{2}$, and $\sigma_{3}=+7.00 \mu \mathrm{C} / \mathrm{m}^{2}$, find the electric field at points $A, B$, and $C$. You may assume that the sheets are infinitely thin.


Figure 22.6 Problem 4.

## Solution

IDENTIFY The electric field for each sheet is constant. We will sum the electric fields due to each sheet to find the electric field, the target variable.

SET UP The electric field for an infinite sheet is given by

$$
E=\frac{\sigma}{2 \epsilon_{0}}
$$

We will calculate the magnitude and direction of the field due to each sheet at each point and then sum the results to find the total field at each point. We will use Figure 22.7 to guide us through the summation.


Figure 22.7 Problem 4 sketch.

EXECUTE At point $A$, the electric field due to the sheets carrying $\sigma_{1}$ and $\sigma_{3}$ point to the left, while the sheet carrying $\sigma_{2}$ points to the right. The field at point $A$ is given by

$$
\begin{aligned}
E_{A} & =-E_{1}+E_{2}-E_{3} \\
& =-\frac{\left|\sigma_{1}\right|}{2 \epsilon_{0}}+\frac{\left|\sigma_{2}\right|}{2 \epsilon_{0}}-\frac{\left|\sigma_{3}\right|}{2 \epsilon_{0}} \\
& =-\frac{\left|4.00 \mu \mathrm{C} / \mathrm{m}^{2}\right|}{2 \epsilon_{0}}+\frac{\left|-6.00 \mu \mathrm{C} / \mathrm{m}^{2}\right|}{2 \epsilon_{0}}-\frac{\left|7.00 \mu \mathrm{C} / \mathrm{m}^{2}\right|}{2 \epsilon_{0}} \\
& =\frac{-17.00 \mu \mathrm{C} / \mathrm{m}^{2}}{2 \epsilon_{0}}=-960,000 \mathrm{~N} / \mathrm{C} .
\end{aligned}
$$

The negative value indicates that the field points to the left.
At point $B$, the electric field due to the sheets carrying $\sigma_{2}$ and $\sigma_{3}$ point to the left, while the sheet carrying $\sigma_{1}$ points to the right. The field at point $B$ is given by

$$
\begin{aligned}
E_{B} & =E_{1}-E_{2}-E_{3} \\
& =\frac{\left|\sigma_{1}\right|}{2 \epsilon_{0}}-\frac{\left|\sigma_{2}\right|}{2 \epsilon_{0}}-\frac{\left|\sigma_{3}\right|}{2 \epsilon_{0}} \\
& =\frac{\left|4.00 \mu \mathrm{C} / \mathrm{m}^{2}\right|}{2 \epsilon_{0}}-\frac{\left|-6.00 \mu \mathrm{C} / \mathrm{m}^{2}\right|}{2 \epsilon_{0}}-\frac{\left|7.00 \mu \mathrm{C} / \mathrm{m}^{2}\right|}{2 \epsilon_{0}} \\
& =\frac{-9.00 \mu \mathrm{C} / \mathrm{m}^{2}}{2 \epsilon_{0}}=-508,000 \mathrm{~N} / \mathrm{C} .
\end{aligned}
$$

The negative value indicates that the field points to the left.

At point $C$, the electric field due to the sheets carrying $\sigma_{1}$ and $\sigma_{3}$ point to the right, while the sheet carrying $\sigma_{2}$ points to the left. The field at point $C$ is given by

$$
\begin{aligned}
E_{A} & =E_{1}-E_{2}+E_{3} \\
& =\frac{\left|\sigma_{1}\right|}{2 \epsilon_{0}}-\frac{\left|\sigma_{2}\right|}{2 \epsilon_{0}}+\frac{\left|\sigma_{3}\right|}{2 \epsilon_{0}} \\
& =\frac{\left|4.00 \mu \mathrm{C} / \mathrm{m}^{2}\right|}{2 \epsilon_{0}}-\frac{\left|-6.00 \mu \mathrm{C} / \mathrm{m}^{2}\right|}{2 \epsilon_{0}}+\frac{\left|7.00 \mu \mathrm{C} / \mathrm{m}^{2}\right|}{2 \epsilon_{0}} \\
& =\frac{5.00 \mu \mathrm{C} / \mathrm{m}^{2}}{2 \epsilon_{0}}=282,000 \mathrm{~N} / \mathrm{C} .
\end{aligned}
$$

The positive value indicates that the field points to the right.
EVALUATE We see that we can easily add electric fields based on several geometries.

## Try It Yourself!

## 1: Conducting cylindrical shells

Two long conducting cylindrical shells of radii $a$ and $b$ are placed such that one is inside the other and they share a common axis. The inner shell carries a charge per unit length $\lambda$ and the outer shell carries a charge per unit length of $-2 \lambda$. Find the electric field in all regions of space.

## Solution Checkpoints

IDENTIFY AND SET UP The system is cylindrically symmetric, so apply Gauss's law to cylindrical surfaces. Sketch the cylinders and place Gaussian surfaces at the appropriate radii.
EXECUTE There is no field inside radius $a$. (Why?) Between $a$ and $b$, Gauss's law yields

$$
\oint \vec{E} \cdot d \vec{A}=E 2 \pi r L=\frac{\lambda L}{\epsilon_{0}}
$$

So the field between the cylindrical shells is

$$
E=\frac{\lambda}{2 \pi \epsilon_{0} r}
$$

and points radially outward from the axis of the cylinder.
Outside radius $b$, the Gaussian surface encloses both cylindrical shells, so Gauss's law gives

$$
\oint \vec{E} \cdot d \vec{A}=E 2 \pi r L=\frac{(\lambda-2 \lambda) L}{\epsilon_{0}}
$$

The field between the cylinders is then

$$
E=\frac{\lambda}{2 \pi \epsilon_{0} r}
$$

with the field pointing inwards toward the cylinder's axis.

EVALUATE Why do the magnitudes of the field in both regions seem to agree? The total charge enclosed for the region outside the shells is negative; how does that enter into the solution?

## 2: Nonconducting cylindrical shells

A long nonconducting cylindrical tube of inner radius $a$ and outer radius $b$ contains a uniform volume charge density $\rho$ throughout its volume. Find the electric field in all regions of space.

## Solution Checkpoints

IDENTIFY AND SET UP The system is cylindrically symmetric, so apply Gauss's law to cylindrical surfaces. Sketch the cylinders and place Gaussian surfaces at the appropriate radii. Inside the cylinder, the Gaussian surface will enclose only a portion of the total charge.

EXECUTE There is no field inside radius $a$. (Why?) Between $a$ and $b$, Gauss's law yields

$$
\oint \vec{E} \cdot d \vec{A}=E 2 \pi r L=\frac{Q_{\mathrm{encl}}}{\epsilon_{0}} .
$$

The charge enclosed is the volume charge density times the volume of the cylinder enclosed by the surface:

$$
Q_{\mathrm{encl}}=\rho\left(\pi r^{2} L-\pi a^{2} L\right)
$$

So the field between the cylindrical shells is

$$
E=\frac{\rho\left(r^{2}-a^{2}\right)}{2 \epsilon_{0} r}
$$

and points radially outward from the axis of the cylinder.
Outside radius $b$, the Gaussian surface encloses the entire tube, so Gauss's law gives

$$
\oint \vec{E} \cdot d \vec{A}=E 2 \pi r L=\frac{\rho\left(\pi b^{2} L-\pi a^{2} L\right)}{\epsilon_{0}}
$$

The field between the cylinders is then

$$
E=\frac{\rho\left(b^{2}-a^{2}\right)}{2 \epsilon_{0} r},
$$

with the field pointing inwards toward the cylinder's axis.
EVALUATE Sketch the graph of the electric field as a function of the radius of the cylinder. Does the expression for the electric field inside the tube agree with the expression for the electric field outside the tube at radius $b$ ?

## 23 <br> Electric Potential

## Summary

In this chapter, we will examine the energy associated with electric interactions. We will begin by defining the electric potential energy and incorporate that concept into our general energy problem-solving methods. Electric potential energy will then be recast in terms of electric potential, which is often a more useful quantity. We will learn to find the electric potential from our knowledge about collections of charges, as well as from our knowledge of the electric field. We will also see how to find the electric field from the electric potential.

## Objectives

After studying this chapter, you will understand

- The definition of electric potential energy and electric potential.
- How to find the electric potential energy and electric potential for a system of charges.
- How to find the electric potential for distributions of charge.
- How equipotential surfaces are used to visualize the electric potential.
- The connection between electric fields and electric potentials and how to calculate one from knowledge of the other.


## Concepts and Equations

Term $\quad$ Description

## Electric Potential Energy

## Description

The electric potential energy $U$ is a scalar quantity equivalent to the work done by the electric force on a charge when it is moved:

$$
W_{a \rightarrow b}=U_{a}-U_{b} .
$$

The electric potential energy for a point charge $q_{0}$ in the electric field of a stationary point charge $q$, with a distance $r$ separating the charges, is

$$
U=\frac{1}{4 \pi \epsilon_{0}} \frac{q_{0} q}{r} .
$$

The unit of electric potential energy is the joule.

## Electric Potential

The electric potential, or potential, $V$ is a scalar quantity that is equal to the potential energy per unit charge:

$$
V=\frac{U}{q_{0}} .
$$

The potential for a point charge $q$ at any point is

$$
V=\frac{1}{4 \pi \epsilon_{0}} \frac{q}{r} .
$$

The potential due to a collection of charges is the sum of the potentials due to each charge:

$$
V=\frac{1}{4 \pi \epsilon_{0}} \sum_{i} \frac{q_{i}}{r_{i}} .
$$

For a continuous distribution of charge, the potential is found by integrating

$$
V=\frac{1}{4 \pi \epsilon_{0}} \int \frac{d q}{r} .
$$

The potential difference between two points $a$ and $b$ is given by the line integral over those points:

$$
V_{a}-V_{b}=\int_{a}^{b} \vec{E} \cdot d \vec{l}=\int_{a}^{b} E \cos \phi d l
$$

The SI unit of potential is the volt $(\mathrm{V})$ equal to 1 joule per coulomb.
An equipotential surface is the graphical representation of potential given by a surface on which every point has the same potential. Electric field lines are perpendicular to equipotential surfaces. The surface of a conductor is an equipotential surface when charges are stationary and all points within the surface have the same potential.
Electric Field from Electric Potential
The electric field can be found from the electric potential if the potential function is known in a region. The components of the electric field are given by the partial derivatives

$$
E_{x}=-\frac{\partial V}{\partial x} \quad E_{y}=-\frac{\partial V}{\partial y} \quad E_{z}=-\frac{\partial V}{\partial z}
$$

where

$$
\vec{E}=-\left(\frac{\partial V}{\partial x} \hat{i}+\frac{\partial V}{\partial y} \hat{j}+\frac{\partial V}{\partial z} \hat{k}\right) .
$$

## Conceptual Questions

## 1: Electric potential of a collection of charges

The electric potential at the center of a square is -8.0 V when one $-2 Q$ charge is placed in a corner of the square. What is the electric potential at the center of the square when a $-Q$ charge is placed in each of the remaining three corners of the square?

## Solution

IDENTIFY, SET UP, AND EXECUTE The electric potential for a collection of charges is the scalar sum of the potentials for each individual charge. The electric potential for a charge depends on the distance from the charge. In this case, all four charges are the same distance from the center of the square, so we do not need to adjust the potentials for distance.

Since the $-2 Q$ charge creates a potential of -8.0 V at the center of the square, we conclude that a single $-Q$ charge would create a potential of -4.0 V at the center of the square. When the three new $-Q$ charges are added, each adds -4.0 V to the potential, for a total of -12.0 V . The total potential due to all four charges is therefore -20.0 V .

EVALUATE This problem illustrates how we add electric potentials. Unlike, the electric field, electric potential is a scalar quantity and does not depend on direction.

## 2: Electric potential and electric potential energy

Two charges are separately brought near a charge $+Q$. The first charge, $+q$, is brought to a distance $r$ from $+Q$. Later, the second charge, $+3 q$, is brought to the same distance $r$ from $+Q$. (a) Compare the electric potentials due to the charge $+Q$ when the two charges are brought to a distance $r$ from $+Q$. (b) Compare the electric potential energies when the two charges are brought to a distance $r$ from $+Q$.

## Solution

IDENTIFY, SET UP, AND EXECUTE (a) The electric potential due to charge $+Q$ depends only on that charge and the distance from it. The potential is thus the same for both charges when they are placed a distance $r$ from the charge $+Q$.
(b) The electric potential energy for a pair of charges depends on the distance between the charges and the magnitudes of the two charges. The electric potential energy between the charges $+Q$ and $+q$ is greater than the electric potential energy between the charges $+Q$ and $+q$.

EVALUATE Electric potential and electric potential energy are different quantities. They are similar to the electric field and electric force in that the members of each pair of quantities are related to each other, but the two pairs have different physical meanings and interpretations.

## 3: Zero electric field and zero electric potential

Is it possible to place four charges of the same magnitude in the corners of a square such that the electric field and the electric potential are zero at the center of the square?

## Solution

IDENTIFY, SET UP, AND EXECUTE Both the electric field and electric potential depend on distance. The center of the square is equidistant from the four corners, so the distance dependence can be satisfied. The net electric potential is the sum of the potentials due to the four charges. To cancel, two would need to be positive and two would need to be negative. The result would not depend on the locations of the charges, because potential is a scalar quantity. However, the electric field is a vector quantity, so the locations of the charges would matter. To get a zero electric field at the center, charges of like signs would be placed in opposite corners.

Both electric potential and the electric field will be zero at the center of the square if two negative and two positive charges, all of equal magnitude, are placed in opposite corners of the square.

EVALUATE Does a zero electric potential always imply a zero electric field or vice versa? We see from this problem that we could have a zero electric potential and a nonzero electric field by placing opposite charges in adjacent corners of the square. And halfway between two equal and opposite charges the electric field is zero, but the electric potential is not.

A common mistake occurs when students assume that the electric field (electric potential) is zero because the electric potential (electric field) is zero.

## Problems

## 1: Closest approach of two protons

Two protons are accelerated toward one another with initial speeds of $1200 \mathrm{~km} / \mathrm{s}$ relative to the earth. The protons are initially very far apart. What is their distance of closest approach?

## Solution

IDENTIFY Energy is conserved, so we will use it to find the distance of closest approach-the target variable.

SET UP The only force between the protons is the electric force, which is a conservative force. The two forms of energy are kinetic energy and elastic potential energy. Initially, the protons have only kinetic energy. (They are far apart, so their initial electric potential energy can be ignored.) As they get closer together, the electric force is repulsive, so the protons slow until they momentarily stop and then return to their original directions. When they are at their closest approach, their kinetic energy is zero and there is only electric potential energy. We'll equate their initial energy to their final energy to solve the problem.

EXECUTE Energy conservation relates the initial to the final energy:

$$
K_{i}+U_{i}=K_{f}+U_{f}
$$

Initially, the two protons have only kinetic energy. At the final position, the protons have only electric potential energy. Thus,

$$
K_{1, i}+K_{2, i}=U_{f},
$$

where we have included the kinetic energy of each proton explicitly. Substituting for the energies, we have

$$
\frac{1}{2} m_{\mathrm{p}} v^{2}+\frac{1}{2} m_{\mathrm{p}} v^{2}=m_{\mathrm{p}} v^{2}=k \frac{q q}{r^{2}}=k \frac{e e}{r_{\text {close }}^{2}}
$$

Solving for the closest approach, $r_{\text {close }}$, yields

$$
r_{\text {close }}=\frac{e}{v} \sqrt{\frac{k}{m_{\mathrm{p}}}}=\frac{\left(1.6 \times 10^{-19} \mathrm{C}\right)}{\left(1.2 \times 10^{6} \mathrm{~m} / \mathrm{s}\right)} \sqrt{\frac{\left(8.99 \times 10^{9} \mathrm{~N} \cdot \mathrm{~m}^{2} / \mathrm{C}^{2}\right)}{\left(1.67 \times 10^{-27} \mathrm{~kg}\right)}}=3.09 \times 10^{-7} \mathrm{~m}
$$

The closest approach of the two protons is 309 nm .
EVALUATE This problem adds a new type of energy to our energy conservation repertoire. As with many energy conservation problems, the key point was to determine the initial and final positions. We could have tried solving the problem by using forces; however, it would have been quite challenging, since the force (and acceleration) increases as the protons move closer together.

## 2: Proton or electron?

A proton or an electron is placed between two equal and oppositely charged parallel plates and is released near plate $A$. The particle accelerates toward plate $B$. If the particle attains a speed of $1.87 \times 10^{6} \mathrm{~m} / \mathrm{s}$ just before striking the other plate, which plate is positively charged? The potential difference between the two plates is 10.0 V .

## Solution

IDENTIFY To find the sign of the charge on the plates, we need to know whether the unknown particle is a proton or an electron. We will use energy conservation to determine which particle is accelerated.

SET UP Protons and electrons have equal and opposite charges and different masses. To find the mass, we'll use energy conservation. Initially, the particle has only electric potential energy at plate $A$. This energy transforms to kinetic energy as the particle accelerates to plate $B$. Once we identify the particle between the plates, we can ascertain the sign of the charge on the plates.

EXECUTE Energy conservation relates the initial to the final energy:

$$
K_{1}+U_{1}=K_{2}+U_{2}
$$

Initially, the particle has only electric potential energy. At plate $B$, it has only kinetic energy. Thus,

$$
U_{1}=K_{2} .
$$

Substituting for the energies, we have

$$
q \Delta V=\frac{1}{2} m v^{2} .
$$

Solving for the mass yields

$$
m=\frac{2 q V}{v^{2}}=\frac{2 e V}{v^{2}}=\frac{2\left(1.6 \times 10^{-19} \mathrm{C}\right)(10.0 \mathrm{~V})}{\left(1.87 \times 10^{6} \mathrm{~m} / \mathrm{s}\right)^{2}}=9.11 \times 10^{-31} \mathrm{~kg}
$$

The mass is that of an electron. The electron is negatively charged, so it accelerates toward positive charge. Plate $B$ is positively charged and plate $A$ is negatively charged.

EVALUATE Can we also learn how the electric potential energy and electric potential change as the electron moves between the plates? The answer is yes. The electron moves toward lower electric potential energy. Due to the negative charge of the electrons, lower electric potential energy corresponds to higher electric potential. Thus, the electron moves toward higher electric potential.

## 3: Charges in a square?

Three charges of $+5.00 \mu \mathrm{C}$ are placed in three corners of a square with sides of length 1.60 m . A fourth charge of $+3.00 \mu \mathrm{C}$ is placed in the center of the square and is released. What is the velocity of the fourth charge as it passes the corner. The fourth charge has a mass of 0.0050 kg .

## Solution

IDENTIFY Energy conservation will be used to calculate the velocity of the charge as it passes the corner-the target variable.

SET UP The fourth charge initially has electric potential energy at the center that converts to kinetic plus electric potential energy at the corner. We must find the electric potential energy of the system when the fourth charge is at the center and at the corner. With four charges, the expression for the electric potential energy includes six pairs of charges. We simplify our calculation by finding the electric potential at the center and at the fourth corner due to the three corner charges, finding the change in potential, and then calculating the change in potential energy for the fourth charge.

EXECUTE The electric potential at the center of the square due to the three charges is the sum of the potentials due to each charge. Each charge is the same in magnitude and is equidistant from the center. The potential is then

$$
V_{\text {center }}=\frac{1}{4 \pi \epsilon_{0}} \sum_{i} \frac{q_{i}}{r_{i}}=3\left(\frac{1}{4 \pi \epsilon_{0}} \frac{q}{r}\right)=3\left(\frac{1}{4 \pi \epsilon_{0}} \frac{5.0 \mu \mathrm{C}}{\frac{\sqrt{2}}{2}(1.6 \mathrm{~m})}\right)=119,000 \mathrm{~V} .
$$

The electric potential at the fourth corner of the square due to the three charges is again the sum of the potentials due to each charge. Each charge is the same magnitude, but the charges all vary in distance from the corner. The potential is then

$$
\begin{aligned}
V_{\text {corner }} & =\frac{1}{4 \pi \epsilon_{0}} \sum_{i} \frac{q_{i}}{r_{i}}=2\left(\frac{1}{4 \pi \epsilon_{0}} \frac{q}{r_{\text {side }}}\right)+\left(\frac{1}{4 \pi \epsilon_{0}} \frac{q}{r_{\text {diag }}}\right) \\
& =2\left(\frac{1}{4 \pi \epsilon_{0}} \frac{5.0 \mu \mathrm{C}}{(1.6 \mathrm{~m})}\right)+\left(\frac{1}{4 \pi \epsilon_{0}} \frac{5.0 \mu \mathrm{C}}{\sqrt{2}(1.6 \mathrm{~m})}\right)=76,000 \mathrm{~V} .
\end{aligned}
$$

The change in potential undergone by the fourth charge is the difference of the two potentials we have just found. The change in potential energy of the fourth charge is this change in potential times that charge.

Energy conservation for the fourth charge relates the initial to the final energy:

$$
K_{1}+U_{1}=K_{2}+U_{2} .
$$

Initially, the fourth charge has only electric potential energy. At the corner, it has electric potential energy plus kinetic energy. Thus,

$$
U_{1}-U_{2}=K_{2} .
$$

Substituting for the energies, we have

$$
U_{1}-U_{2}=q \Delta V=\frac{1}{2} m v^{2} .
$$

Solving for the velocity yields

$$
v=\sqrt{\frac{2 q \Delta V}{m}}=\sqrt{\frac{2(3.0 \mu \mathrm{C})(119,000 \mathrm{~V}-76,000 \mathrm{~V})}{(0.005 \mathrm{~kg})}}=7.2 \mathrm{~m} / \mathrm{s}
$$

The fourth charge has a velocity of $7.2 \mathrm{~m} / \mathrm{s}$ when it passes the empty corner.

EVALUATE How does the force vary as the fourth charge moves to the corner? The force decreases, so we could not have used the constant-acceleration kinematics relations we developed in Chapter 3. This problem illustrates how electric potential simplifies problems by removing the challenges of varying forces and thereby simplifying the calculation of potential energy for a system of charges.

## 4: Potential of two charges

Two $+Q$ charges are placed on the $y$-axis a distance $a$ above and below the $x$-axis, as shown in Figure 23.1. Find the electric potential on the $x$-axis at a distance $x$ from the origin.


Figure 23.1 Problem 4.

## Solution

IDENTIFY AND SET UP Electric potential is a scalar, so the net electric potential is the sum of the potentials due to the two charges. We need to find the distance to the two charges in terms of $a$ and $x$.

EXECUTE The charges are equidistant from the point on the axis. The distance to each charge is then

$$
r=\sqrt{x^{2}+a^{2}} .
$$

The net electric potential is the sum of the potentials due to each charge:

$$
\begin{aligned}
V & =\frac{1}{4 \pi \epsilon_{0}} \sum_{i} \frac{q_{i}}{r_{i}} \\
& =\frac{1}{4 \pi \epsilon_{0}}\left(\frac{Q}{\sqrt{x^{2}+a^{2}}}+\frac{Q}{\sqrt{x^{2}+a^{2}}}\right) \\
& =\frac{2}{4 \pi \epsilon_{0}}\left(\frac{Q}{\sqrt{x^{2}+a^{2}}}\right) .
\end{aligned}
$$

EVALUATE Finding the electric potential can be relatively simple, as in this problem. One use of the potential is in calculating the electric field, as we will see in Problem 5.

## 5: Electric field due to two charges

Two $+Q$ charges are placed on the $y$-axis a distance $a$ above and below the $x$-axis, as shown in Figure 23.1. Find the electric field on the $x$-axis at a distance $x$ from the origin.

## Solution

IDENTIFY We can calculate the electric field in two ways: We can find the electric field due to each charge and add their vectors, or we can use the partial derivative of the electric potential to find the electric field. We will use the second method in this problem, since we found the electric potential in the previous problem.
SET UP The electric field is the negative partial derivative of the electric potential. From the expression for the electric potential in Problem 4, we see that the potential varies only as a function of $x$, so the electric field has only an $x$ component.
EXECUTE The $x$ component of the electric field is the negative partial derivative of the electric potential with respect to $x$ :

$$
\begin{aligned}
E_{x} & =-\frac{\partial V}{\partial x} \\
& =-\frac{\partial}{\partial x}\left(\frac{2}{4 \pi \epsilon_{0}}\left(\frac{Q}{\sqrt{x^{2}+a^{2}}}\right)\right) \\
& =-\frac{2 Q}{4 \pi \epsilon_{0}} \frac{\partial}{\partial x}\left(\frac{1}{\sqrt{x^{2}+a^{2}}}\right) \\
& =-\frac{2 Q}{4 \pi \epsilon_{0}}\left(\frac{\left(-\frac{1}{2}\right) 2 x}{\left(x^{2}+a^{2}\right)^{3 / 2}}\right) \\
& =\frac{1}{4 \pi \epsilon_{0}} \frac{2 Q x}{\left(x^{2}+a^{2}\right)^{3 / 2}} .
\end{aligned}
$$

EVALUATE This result agrees with Chapter 21, Do It Yourself Problem 2, illustrating how finding the electric field from an electric potential can be simpler than finding the electric field directly from a vector sum.

## Try It Yourself!

## 1: Potential of two charges

Two $1.0 \times 10^{-9} \mathrm{C}$ charges are placed on the $y$-axis a distance $a$ above and below the $x$-axis, as shown in Figure 23.2. An electron located at the origin is given a slight push. Find the velocity of the electron when it is at $y=a / 2$.


Figure 23.2 Try It Yourself 1.

## Solution Checkpoints

IDENTIFY AND SET UP Use energy conservation to solve. What forms of energy are present?
EXECUTE The net electric potential is the sum of the potentials due to each charge:

$$
\begin{aligned}
V & =\frac{1}{4 \pi \epsilon_{0}}\left(\frac{Q}{a-y}+\frac{Q}{a+y}\right) \\
& =\frac{2}{4 \pi \epsilon_{0}}\left(\frac{Q a}{a^{2}-y^{2}}\right) .
\end{aligned}
$$

The change in energy is related to the change in potential. The kinetic energy of the electron increases as the electron moves along the axis. At any point, the kinetic energy is given by

$$
\Delta K=\frac{1}{2} m v^{2}=\frac{2}{4 \pi \epsilon_{0}}\left(\frac{Q a}{a^{2}-y^{2}}\right)-\frac{2}{4 \pi \epsilon_{0}}\left(\frac{Q}{a}\right) .
$$

The final velocity is $1.45 \times 10^{7} \mathrm{~m} / \mathrm{s}$.
EVALUATE How could you use forces to solve this problem?

## 2: Conducting sphere

A solid conducting sphere of radius $a$ is concentric with a hollow spherical shell of inner radius $b$ and outer radius $c$. The potential difference between the spheres is $V_{a b}$. Find the electric field between the spheres and find the charge on the inner sphere.

## Solution Checkpoints

IDENTIFY AND SET UP Apply Gauss's law between the conductors. Sketch the spheres. $Q$ is not given, but can be found from the potential difference.

EXECUTE Assuming a charge $Q$ on the inner conductor, we find that the field is given by

$$
E=\frac{Q}{4 \pi \epsilon_{0} r^{2}} .
$$

The potential difference between the spheres is given by

$$
\begin{aligned}
V_{a b} & =\int_{a}^{b} \vec{E} \cdot d \vec{r} \\
& =\frac{Q}{4 \pi \epsilon_{0}}\left(\frac{1}{a}-\frac{1}{b}\right) .
\end{aligned}
$$

This equation can be rearranged to find the charge and the electric field:

$$
\begin{aligned}
Q & =\frac{4 \pi \epsilon_{0} V_{a b}}{\left(\frac{1}{a}-\frac{1}{b}\right)} \\
E & =\frac{V_{a b}}{\left(\frac{1}{a}-\frac{1}{b}\right) r^{2}}
\end{aligned}
$$

EVALUATE We see that if the charge isn't explicitly given, we can assign it a value and later evaluate the charge in terms of the known values.

## 3: The hydrogen atom

(a) Estimate the energy necessary to separate the electron from the proton in a hydrogen atom. Assume that the two particles are bound by the coulomb force and are an average of $10^{-8} \mathrm{~m}$ apart. (b) Through what voltage must an electron be accelerated to gain this energy?

## Solution Checkpoints

IDENTIFY AND SET UP Can energy conservation be used to solve the problem?
EXECUTE (a) The energy necessary to separate the electron from the proton is equal to the electric potential energy of the pair:

$$
U=\frac{e^{2}}{4 \pi \epsilon_{0} r}=2.3 \times 10^{-18} \mathrm{~J}
$$

(b) The potential energy is equal to the electron's charge times the change in potential:

$$
\Delta V=\frac{U}{e}=14.4 \mathrm{~V}
$$

EVALUATE This result is actually quite close to the energy needed to remove an electron from a hydrogen atom. We'll learn about nuclear forces later in the text.

## Capacitance and Dielectrics

## Summary

In this chapter, we will examine capacitance and the capacitor, a device for storing energy and charge. We will begin to analyze circuits and consider the energy and flow of charge in a circuit. We will find that the familiar term voltage is equivalent to electric potential. We will learn how to combine capacitors in a circuit and find the charge on a capacitor. We will also see how capacitors store energy and how to quantify that energy. Finally, we will examine dielectrics-materials that change the capacitance of capacitors.

## Objectives

After studying this chapter, you will understand

- Capacitors and capacitance and how to define them.
- How to recognize, differentiate, and find the equivalent capacitance of capacitors connected in series and in parallel.
- How to find the energy stored in a capacitor and in the electric field of a capacitor.
- Dielectrics and their effect on capacitors.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Capacitors | A capacitor is any two conductors separated by vacuum or a material. When <br> charge is added to the conductors of a capacitor, a potential exists between <br> the conductors. The capacitance $C$ is the ratio of the charge to the potential <br> difference: <br> $C C=\frac{Q}{V}$. |

The SI unit of capacitance is the farad; $1 \mathrm{~F}=1 \mathrm{C} / \mathrm{V}$. A parallel-plate capacitor consists of two parallel plates of area $A$ separated by a distance $d$. The capacitance of a parallel-plate capacitor is

$$
C=\epsilon_{0} \frac{A}{d}
$$

## Combinations of Capacitors

Electric Field Energy

When capacitors are combined in series, the equivalent capacitance is given by

$$
\frac{1}{C_{\mathrm{eq}}}=\frac{1}{C_{1}}+\frac{1}{C_{2}}+\frac{1}{C_{3}}+\cdots,
$$

and all capacitors have the same magnitude of charge on their plates. When capacitors are combined in parallel, the equivalent capacitance is given by

$$
C_{\mathrm{eq}}=C_{1}+C_{2}+C_{3}+\cdots,
$$

and all capacitors have the same potential difference across their plates.
The electric field energy $U$ is the energy required to charge a capacitor $C$ to a potential difference $V$ with charge $Q$ :

$$
U=\frac{1}{2} Q V=\frac{Q^{2}}{2 C}=\frac{1}{2} C V^{2} .
$$

This energy is stored in the electric field, with an energy density of

$$
u=\frac{1}{2} \epsilon_{0} E^{2} .
$$

The nonconducting material placed between the plates of a capacitor is a dielectric. For a constant charge, dielectrics decrease the electric field and potential between the plates of a capacitor, thus increasing the capacitance by the dielectric constant $K$ of the material. A parallel-plate capacitor filled with a dielectric has capacitance

$$
C=\kappa C_{0}=\kappa \epsilon_{0} \frac{A}{d}=\epsilon \frac{A}{d} .
$$

Gauss's law in a dielectric becomes

$$
\oint_{\kappa} \vec{E} \cdot d \vec{A}=\frac{Q_{\text {encl free }}}{\epsilon_{0}},
$$

where $Q_{\text {encl free }}$ is the total free charge enclosed by the Gaussian surface. The maximum electric field strength a material can withstand without ionizing (i.e., without the dielectric breaking down) is the dielectric strength of the material.

## Conceptual Questions

## 1: Changes in a disconnected capacitor

A parallel-plate capacitor is charged to a potential difference of 100 V and is disconnected from the power source. The plates are then pushed closer together by means of insulated rods. How do the charge, potential difference, and capacitance change when the plates are pushed closer together?

## Solution

IDENTIFY, SET UP, AND EXECUTE When the capacitor is disconnected from the power source and the plates are pushed closer together with the insulated rods, no charge can leave the plates and the charge remains constant. As the plates come closer together, the capacitance will increase, since the capacitance is inversely proportional to the separation of the plates. Because the capacitance increases, the potential difference must decrease in order for the charge to remain constant. Summarizing, we say that the charge remains constant, the potential difference decreases, and the capacitance increases.

EVALUATE Capacitance, charge, and potential difference are related in a simple manner; understanding how these quantities vary under different circumstances develops a better understanding of capacitors.

Practice Problem: How would the charge on the plates, the potential difference, and the capacitance change if the capacitor remained connected to the power source as the plates were pushed together? Answer: The capacitance would still increase, the potential difference would remain constant, and the charge on the plates would increase.

## 2: Combinations of capacitors

How would you combine the two arrangements of capacitors shown in Figure 24.1 to find the equivalent capacitance between $a$ and $b$ ?


Figure 24.1 Question 2.

## Solution

IDENTIFY, SET UP, AND EXECUTE Capacitors $C_{1}$ and $C_{2}$ are positioned parallel to each other. However, closer inspection reveals that there is only one path leading between $a$ and $b$ through the two capacitors. Capacitors $C_{1}$ and $C_{2}$ are connected in series between $a$ and $b$. Their equivalent capacitance would be found by adding reciprocals.

Capacitors $C_{3}$ and $C_{4}$ are positioned somewhat diagonally from each other. This time, closer inspection reveals that there are two paths between $a$ and $b$, one leading through capacitor $C_{3}$ and the other through capacitor $C_{4}$. Capacitors $C_{3}$ and $C_{4}$ are connected in parallel between $a$ and $b$. Their equivalent capacitance would be found by adding their capacitances.

EVALUATE Recognizing and differentiating series and parallel connections in circuits is an important aspect of this and the next chapter. Be aware that series and parallel connections refer to the electrical connections and not any geometric arrangment. Once you determine how components are connected, finding the equivalent capacitance is straightforward. You should practice combining capacitors until you've developed a good eye for recognizing the nature of the connections.

## Problems

## 1: Changing parallel-plate capacitor

A parallel-plate capacitor has a charge of 1.0 nC when the plates are separated by 1.0 cm and it is connected to a $12.0-\mathrm{V}$ battery. (a) If the plates are pulled to a $2.0-\mathrm{cm}$ separation while keeping the battery connected, what is the new charge on the plates? (b) If the plates are disconnected from the battery and then pulled to a $2.0-\mathrm{cm}$ separation, what is the new potential between the plates?

## Solution

IDENTIFY We'll use the definition of capacitance and the capacitance of a parallel-plate capacitor to solve the problem.

SET UP We must recognize what remains constant and what changes in the two parts of the problem. The potential remains constant in part (a) and the charge remains constant in part (b). The capacitance changes in both cases, so the charge changes in part (a) and the potential changes in part (b).

EXECUTE The capacitance drops by a factor of two in both parts of the problem, since capacitance for a parallel-plate capacitor is inversely proportional to the distance between the plates:

$$
C=\frac{\epsilon_{0} A}{d} .
$$

The charge, potential, and capacitance are related by

$$
q=C V .
$$

(a) The charge must decrease by a factor of two in part (a), since the potential remains constant and the capacitance drops by a factor of two. The final charge is 0.5 nC .
(b) The potential must increase by a factor of two in part (b), since the charge remains constant and the capacitance drops by a factor of two. The final potential is 24.0 V .

EVALUATE This problem illustrates how we will approach many problems involving capacitors. The key is to recognize what parameters change and what parameters remain constant through the course of the problem. When capacitors are disconnected from a battery, their charge remains constant since the charge has no path for leaving the capacitor. When capacitors are connected to batteries, the charge may flow onto or off of the plates of the capacitor as the plates are manipulated.

## 2: Charge on a capacitor

The circuit of Figure 24.2 is attached to a $12-\mathrm{V}$ battery across terminal $a b$. Find the charge on the $21-\mu \mathrm{F}$ capacitor.


Figure 24.2 Problem 2.

## Solution

IDENTIFY To find the charge on the capacitor, we need to know the potential difference across the capacitor. We will combine the capacitors to find the potential on the capacitor.

SET UP In this problem, we know only the voltage across the terminal, so we will need to combine the capacitors initially in order to find the charges and potentials in the system. Then we will separate the system into its components to isolate the $21-\mu \mathrm{F}$ capacitor.

EXECUTE We start by finding the equivalent capacitance of the system. To do that, we use our series and parallel rules to find two capacitors that we can combine. Examining the figure, we see that the $21-\mu \mathrm{F}$ and $16-\mu \mathrm{F}$ capacitors are in series. Combining those two, we find that

$$
\frac{1}{C_{21+16}}=\frac{1}{21 \mu \mathrm{~F}}+\frac{1}{16 \mu \mathrm{~F}}, C_{21+16}=9.08 \mu \mathrm{~F}
$$

where the subscript refers to the combined capacitors. Combining the $41-\mu \mathrm{F}$ capacitor in parallel with the $21-\mu \mathrm{F}$ and $16-\mu \mathrm{F}$ capacitors gives

$$
C_{21+16+41}=9.08 \mu \mathrm{~F}+41 \mu \mathrm{~F}=50.1 \mu \mathrm{~F}
$$

The $33-\mu \mathrm{F}$ capacitor is in series with the rest of the capacitors, so the total equivalent capacitance becomes

$$
\frac{1}{C_{21+16+41+33}}=\frac{1}{33 \mu \mathrm{~F}}+\frac{1}{50.1 \mu \mathrm{~F}}, C_{21+16+41+33}=19.9 \mu \mathrm{~F} .
$$

With the total equivalent capacitance found, we turn to calculating the charges and potentials in the circuit. The total charge on the capacitor combination is

$$
Q_{\text {total }}=C V=(19.9 \mu \mathrm{~F})(12 \mathrm{~V})=239 \mu \mathrm{C}
$$

Since the $33-\mu \mathrm{F}$ capacitor is in series with the $50.1-\mu \mathrm{F}$ combination, the total charge on both capacitors is $239 \mu \mathrm{C}$. The potential difference across the $50.1-\mu \mathrm{F}$ combination is

$$
V_{21+16+41}=\frac{Q_{\text {total }}}{C_{21+16+41}}=\frac{239 \mu \mathrm{C}}{50.1 \mu \mathrm{~F}}=4.77 \mathrm{~V}
$$

Since these capacitors are in parallel, the potential difference across the $21-\mu \mathrm{F}$ and $16-\mu \mathrm{F}$ combination is 4.77 V . We now find the charge on the $21-\mu \mathrm{F}$ and $16-\mu \mathrm{F}$ combination:

$$
\mathrm{Q}_{21+16}=C_{21+16} V_{21+16}=(9.08 \mu \mathrm{~F})(4.77 \mathrm{~V})=43.3 \mu \mathrm{C}
$$

The charge on the $21-\mu \mathrm{F}$ and $16-\mu \mathrm{F}$ combination is $43.3 \mu \mathrm{C}$. Since the capacitors are in series, both have the same charge. The charge on the $21-\mu \mathrm{F}$ capacitor is $43.3 \mu \mathrm{C}$.

EVALUATE This problem has shown us how to find the equivalent capacitance of a circuit and then determine the charges and voltages on various capacitors. In combining the capacitors, we saw that the concepts of series and parallel connections are not geometric concepts, but electrical concepts. In this circuit, the $21-\mu \mathrm{F}$ and $16-\mu \mathrm{F}$ capacitors are oriented perpendicular to each other, but are in series electrically. The $41-\mu \mathrm{F}$ capacitor is diagonal to the $21-\mu \mathrm{F}$ and $16-\mu \mathrm{F}$ capacitors, but is electrically parallel to them. Learning to identify series and parallel circuits will greatly benefit your problem-solving skills in this and the next chapter.

## 3: Connecting capacitors

A $30.0-\mu \mathrm{F}$ capacitor is charged to a potential of 500 V , is disconnected from the source, and is connected to an uncharged $10.0-\mu \mathrm{F}$ capacitor. Find the final charge on each capacitor.

## Solution

IDENTIFY We will use the properties of capacitors and potentials to solve the problem. The target variables are the charges on the two capacitors after they are connected.

SET UP The $30.0-\mu \mathrm{F}$ capacitor is initially charged to 500 V , so we find the initial charge on it. This charge remains on the capacitor after it is disconnected, as it has no path for leaving the plates. When the capacitors are connected, the charge is shared between the two capacitors. The amount of charge on the two capacitors after reconnection is found by realizing that the potential differences across each capacitor must be equal.

EXECUTE We start by finding the initial charge on the $30.00-\mu \mathrm{F}$ capacitor. The charge is

$$
Q_{30}=C_{30} V=(30.00 \mu \mathrm{~F})(500 \mathrm{~V})=15,000 \mu \mathrm{C}
$$

where the subscript refers to the capacitor. After connection, the charge is shared between the two capacitors; that is,

$$
Q_{30}=Q_{30}^{\prime}+Q_{10}^{\prime},
$$

where the primes indicate the values after connection. When the capacitors are connected, the potential is the same across each one, or

$$
\begin{aligned}
V_{30} & =V_{10} \\
\frac{Q_{30}^{\prime}}{C_{30}} & =\frac{Q_{10}^{\prime}}{C_{10}}
\end{aligned}
$$

We can substitute and solve for the charge on the $30.0-\mu \mathrm{F}$ capacitor:

$$
\begin{aligned}
& \frac{Q_{30}^{\prime}}{C_{30}}=\frac{Q_{30}-Q_{30}^{\prime}}{C_{10}} \\
& Q_{30}^{\prime}=\frac{Q_{30}}{C_{10}}\left(\frac{1}{C_{30}}+\frac{1}{C_{10}}\right)^{-1}=\frac{(15,000 \mu \mathrm{C})}{10.0 \mu \mathrm{~F}}\left(\frac{1}{(30.0 \mu \mathrm{~F})}+\frac{1}{10.0 \mu \mathrm{~F}}\right)^{-1}=11,250 \mu \mathrm{C} .
\end{aligned}
$$

The charge on the other capacitor is then

$$
Q_{10}^{\prime}=Q_{30}-Q_{30}^{\prime}=15,000 \mu \mathrm{C}-11,250 \mu \mathrm{C}=3750 \mu \mathrm{C} .
$$

The charge on the $30.00-\mu \mathrm{F}$ capacitor is 11.3 mC , and the charge on the $10.00-\mu \mathrm{F}$ capacitor is 3.8 mC .

EVALUATE This problem illustrates how charge moves on capacitors after the capacitors are connected. One method to better understand the movement of charges on capacitors is to draw diagrams showing the distribution of charges before and after connection. You may also want to use arrows to indicate how charges move between capacitors. Try several drawings to illustrate the movement of charges in this problem.

Practice Problem: Find the potential difference across the capacitors after they are connected. Answer: 375 V.

## 4: Charge on rearranged capacitors

A $3.00-\mu \mathrm{F}$ capacitor and a $7.00-\mu \mathrm{F}$ capacitor are connected in parallel, and the combination is connected to a battery that provides a potential difference. After the capacitors are charged, they are disconnected from the battery and each other and then are reconnected to each other with their terminals reversed. If the final potential difference across the $3.00-\mu \mathrm{F}$ capacitor is 175 V , find the voltage of the battery.

## Solution

IDENTIFY We will use the properties of capacitors and potentials to solve the problem. The target variable is the voltage of the battery.

SET UP The capacitors are initially charged to the same voltage, with the $7.00-\mu \mathrm{F}$ capacitor acquiring more charge than the $3.00-\mu \mathrm{F}$ capacitor. When they are disconnected, the charge remains on the capacitors. When the capacitors are reconnected with their terminals reversed, the positive charge on the $7.00-\mu \mathrm{F}$ capacitor combines with the negative charge on the $3.00-\mu \mathrm{F}$ capacitor, neutralizing part of the total charge, with the remaining positive charge split between the two capacitors. The same occurs on the other terminals, resulting in an equal amount of negative charge split between the two capacitors. The amount of charge on the two capacitors after reconnection is found by realizing that the potential differences across each capacitor must be equal. We are told that the potential across the $3.00-\mu \mathrm{F}$ capacitor is 175 V , so the potential across the $7.00-\mu \mathrm{F}$ capacitor is also 175 V .

Knowing the final voltages across the capacitors allows us to find the final charges on the capacitors. We can then find the initial charge on the capacitors and the voltage of the battery.

EXECUTE We start by finding the charge on the two capacitors after they have been reconnected. The charges are

$$
\begin{aligned}
& Q_{3}^{\prime}=C_{3} V^{\prime}=(3.00 \mu \mathrm{~F})(175 \mathrm{~V})=525 \mu \mathrm{C} \\
& Q_{7}^{\prime}=C_{7} V^{\prime}=(7.00 \mu \mathrm{~F})(175 \mathrm{~V})=1225 \mu \mathrm{C}
\end{aligned}
$$

where the subscripts refer to the capacitors and the primes indicate the values after reconnection. The total charge on the capacitors after reconnection is the sum of the charges on the capacitors:

$$
Q_{\text {total }}^{\prime}=Q_{3}^{\prime}+Q_{7}^{\prime}=(525 \mu \mathrm{C})+(1225 \mu \mathrm{C})=1750 \mu \mathrm{C}
$$

The initial charges on the capacitors (while connected to the battery) are found in the same manner and are

$$
\begin{aligned}
& Q_{3}=C_{3} V, \\
& Q_{7}=C_{7} V,
\end{aligned}
$$

where $V$ is the battery's voltage. When the capacitors are reconnected, the positive charge on the $7.00-\mu \mathrm{F}$ capacitor combines with the negative charge on the $3.00-\mu \mathrm{F}$ capacitor, giving a total charge of

$$
Q_{\text {total }}=Q_{7}-Q_{3}=C_{7} V-C_{3} V=\left(C_{7}-C_{3}\right) V
$$

This total charge is the same as the total charge after the reconnection. We solve for the battery's voltage:

$$
V=\frac{Q_{\text {total }}}{\mathrm{C}_{7}-C_{3}}=\frac{Q_{\text {total }}^{\prime}}{\mathrm{C}_{7}-C_{3}}=\frac{(1750 \mu \mathrm{C})}{(7.00 \mu \mathrm{~F})-(3.00 \mu \mathrm{~F})}=438 \mathrm{~V}
$$

The battery's voltage is 438 V .
EVALUATE This problem illustrates how charge remains on capacitors after they are disconnected from a battery and how charge rearranges itself after capacitors are reconnected. The relation among charge, voltage, and capacitance is straightforward; however, mastery of capacitance problems comes only after considering a variety of problems in which either the charge or the voltage remains constant while other parameters change. The next problem illustrates what happens when the voltage is kept constant.

## 5: Charge on a dielectric

A capacitor is made of two square sheets of aluminum foil with sides of length 50.0 cm placed 2.00 mm apart in air. The capacitor is connected to a power supply that keeps a constant potential difference of 24.0 V across the plates. A piece of neoprene is placed between the pieces of foil, completely filling the space. How much charge flows onto or off of the plates when the neoprene is added? What is the change in energy stored in the capacitor when the neoprene is added?

## Solution

IDENTIFY The target variable is the amount of charge moving on or off of the capacitor plates and the change in energy stored.

SET UP To find the charge flow, we calculate the initial and final charges on the capacitor. We'll need to calculate the capacitance of the foil, using the parallel-plate capacitor formula, and we have to look up the dielectric strength of the neoprene. Once we have the charges, we will calculate the initial and final energy stored on the capacitor. We'll take the dielectric constant of air to be equal to vacuum to simplify the analysis.

EXECUTE The capacitance of a parallel-plate capacitor is

$$
C=\epsilon_{0} \frac{A}{d}
$$

where $A$ is the area of the plates and $d$ is the separation distance. For the capacitor with air between the plates, the capacitance is

$$
C_{\mathrm{air}}=\epsilon_{0} \frac{A}{d}=\left(8.85 \times 10^{-12} \mathrm{~F} / \mathrm{m}\right) \frac{(0.500 \mathrm{~m})(0.500 \mathrm{~m})}{(0.00200 \mathrm{~m})}=1.106 \mathrm{nF}
$$

When the neoprene is added, the capacitance increases by the dielectric constant. The dielectric constant for neoprene is 6.70 , from Table 18.1. The capacitance is

$$
C_{\text {neoprene }}=\kappa C_{\text {air }}=(6.70)(1.106 \mathrm{nF})=7.412 \mathrm{nF} .
$$

The charge on the capacitor with and without the neoprene is, respectively,

$$
\begin{aligned}
Q_{\text {air }} & =C_{\text {air }} V=(1.106 \mathrm{nF})(24 \mathrm{~V})=26.5 \mathrm{nC}, \\
Q_{\text {neoprene }} & =C_{\text {neeprene }} V=(7.412 \mathrm{nF})(24 \mathrm{~V})=177.9 \mathrm{nC} .
\end{aligned}
$$

The amount of charge flowing onto the plates when the neoprene is added is

$$
\Delta Q=Q_{\text {neoprene }}-Q_{\text {air }}=177.9 \mathrm{nC}-26.5 \mathrm{nC}=151.4 \mathrm{nC} .
$$

The energy stored in a capacitor is given by

$$
E=\frac{1}{2} C V^{2} .
$$

The energy stored in the capacitor with and without the neoprene is, respectively,

$$
\begin{aligned}
E_{\text {air }} & =\frac{1}{2} C_{\text {air }} V^{2}=\frac{1}{2}(1.106 \mathrm{nF})(24 \mathrm{~V})^{2}=318 \mathrm{~nJ}, \\
E_{\text {neoprene }} & =\frac{1}{2} C_{\text {neoprene }} V^{2}=\frac{1}{2}(7.412 \mathrm{nF})(24 \mathrm{~V})^{2}=2135 \mathrm{~nJ} .
\end{aligned}
$$

When the neoprene is added, the energy increases by an amount

$$
\Delta E=E_{\text {neoprene }}-E_{\text {air }}=2135 \mathrm{~nJ}-318 \mathrm{~nJ}=1817 \mathrm{~nJ} .
$$

When the neoprene is inserted, 151 nC of charge flows onto the plates, increasing the energy stored in the capacitor by 1820 nJ .

EVALUATE This problem shows us how to find the capacitance and energy of a parallel-plate capacitor with and without a dielectric. We see that, since the potential remains constant, the charge on the capacitor must rearrange itself to compensate for the change in capacitance when the neoprene is added.

Practice Problem: Would the neoprene have to be pushed between the plates, or would it be drawn into the plates? Answer: Since the energy increases when the neoprene is added, the neoprene would have to be pushed between the plates.

## 6: Adding a dielectric

A parallel-plate capacitor of area $A$ and plate separation $L$ is filled with a removable dielectric slab of dielectric constant $K$ ? The capacitor is given a charge $Q$ with the slab removed, and the capacitor is disconnected from the battery. Then the slab is inserted. Find the change in potential difference when the slab is inserted.

## Solution

IDENTIFY The target variable is the change in potential difference when the slab is added.
SET UP We will calculate the initial and final voltage on the capacitor. When the dielectric is added, the charge remains constant.

EXECUTE The initial potential difference is given by

$$
V=\frac{Q}{C}=\frac{Q L}{\epsilon_{0} A} .
$$

When the dielectric is added, the potential decreases, since the capacitance increases. We then have

$$
V^{\prime}=\frac{Q}{C^{\prime}}=\frac{Q L}{\kappa \epsilon_{0} A} .
$$

We need to find the difference of these two potentials. The difference is

$$
\Delta V=V^{\prime}-V=\frac{Q L}{\kappa \epsilon_{0} A}-\frac{Q L}{\epsilon_{0} A}=\frac{Q L}{\epsilon_{0} A}\left(\frac{1}{\kappa}-1\right)
$$

The potential difference is negative, indicating that the potential decreases when the dielectric is added.

EVALUATE Does the dielectric have to be pushed into place, or is it pulled into place? The energy after adding the dielectric is less; therefore, the dielectric is pulled into place by the field.

## Try It Yourself!

## 1: Charge on a capacitor

Find the charges on, and the potential drops across, each capacitor in the circuit of Figure 24.3. The potential across the terminal $a b$ is 1000.0 V .


Figure 24.3 Try It Yourself 1.

## Solution Checkpoints

IDENTIFY AND SET UP To find the charge on the capacitors, we need to know the potential difference across the capacitors. We will initially combine the capacitors to find the charges and potentials in the system.

EXECUTE Capacitors $C_{1}$ and $C_{3}$ are combined first, and then capacitors $C_{2}$ and $C_{4}$ are added to the combination:

$$
\frac{1}{C_{\text {total }}}=\frac{1}{2.0 \mu \mathrm{~F}}+\frac{1}{1.0 \mu \mathrm{~F}+3.0 \mu \mathrm{~F}}+\frac{1}{4.0 \mu \mathrm{~F}}, C_{\text {total }}=1.0 \mu \mathrm{~F} .
$$

The charge on the equivalent capacitor is

$$
Q_{\text {total }}=C V=1000 \mu \mathrm{C}
$$

Capacitors $C_{2}$ and $C_{4}$ have this same charge, since they are in series. The potential differences are

$$
\begin{aligned}
& V_{2}=\frac{Q_{2}}{C_{2}}=500 \mathrm{~V} \\
& V_{4}=\frac{Q_{4}}{C_{4}}=250 \mathrm{~V} \\
& \mathrm{~V}_{1}=V_{3}=1000 \mathrm{~V}-V_{2}-V_{4}=250 \mathrm{~V} .
\end{aligned}
$$

The charges on capacitors $C_{1}$ and $C_{3}$ are

$$
\begin{aligned}
& Q_{1}=C_{1} V_{1}=250 \mu \mathrm{C} \\
& Q_{3}=C_{3} V_{3}=750 \mu \mathrm{C}
\end{aligned}
$$

EVALUATE How would you check these results?

## 2: Changing capacitors

Imagine that capacitors $C_{1}$ and $C_{3}$ are removed from the circuit of Figure 24.3 and are reconnected to each other with terminals of opposite signs together. Find the new charge on, and potential across, each capacitor.

## Solution Checkpoints

IDENTIFY AND SET UP How does the initial charge on the capacitors distribute after they reconnect? How do the potentials across the capacitors compare after the reconnection?

EXECUTE The total charge on the two capacitors is

$$
Q^{\prime}=Q_{1}-Q_{3}=Q_{1}^{\prime}+Q_{3}^{\prime}=500 \mu \mathrm{C}
$$

The potentials across the capacitors are the same, giving

$$
\begin{aligned}
V_{1}^{\prime} & =V_{3}^{\prime} \\
\frac{Q_{1}^{\prime}}{C_{1}} & =\frac{Q_{3}^{\prime}}{C_{3}}
\end{aligned}
$$

Solving these two equations gives

$$
\begin{aligned}
& Q_{1}^{\prime}=125 \mu \mathrm{C} \\
& Q_{3}^{\prime}=375 \mu \mathrm{C}
\end{aligned}
$$

EVALUATE How would you check these results? Did the overall energy increase or decrease?

## 3: Adding a dielectric

A parallel-plate capacitor of area $A$ and plate separation $L$ is filled with a removable dielectric slab of dielectric constant $K$ ? The capacitor is connected to a battery with the slab removed, and then the slab is inserted while the capacitor remains connected to the battery. Find the change in charge on the capacitor when the slab is inserted.

## Solution Checkpoints

IDENTIFY AND SET UP How does the initial charge on the capacitors distribute after the dielectric is inserted? How do the potentials across the capacitor compare after the slab is inserted?

EXECUTE The initial charge is given by

$$
Q=V C=V \frac{\epsilon_{0} A}{L}
$$

When the dielectric is added, the potential remains constant. The new charge is

$$
Q^{\prime}=V C^{\prime}=V \frac{\kappa \epsilon_{0} A}{L}
$$

We need to find the difference of the two charges. The difference is

$$
\Delta Q=Q^{\prime}-Q=V \frac{\kappa \epsilon_{0} A}{L}-V \frac{\epsilon_{0} A}{L}=V \frac{\epsilon_{0} A}{L}(\kappa-1) .
$$

EVALUATE Does the charge on the capacitor increase or decrease when the dielectric is inserted?

## Current, Resistance, and Electromotive Force

## Summary

In this chapter, we will study the movement of charge through electric circuits. We begin by defining current (the flow of charge through a conductor), resistance (a quantity related to how easily charge flows through a conductor), and voltage (electric potential difference, which causes charge to move). We'll then analyze electric circuits and see how energy and power are provided to and dissipated by the devices in a circuit. By the end of the chapter, we will have built a foundation from which we can begin understanding circuits in electronic devices that we use every day.

## Objectives

After studying this chapter, you will understand

- The flow of current and charges through an electric circuit.
- The definition of resistivity and conductivity.
- The resistance of a circuit and the implications of Ohm's law.
- How an electromotive force (emf) causes charges to flow in a circuit.
- How to determine energy and power for a circuit.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Current | Current is the movement of electrical charge from one region to another. A net <br> charge $d Q$ flowing through a cross-sectional area $A$ in time $d t$ is a current |
| $\quad I=\frac{d Q}{d t}=n\|q\| v_{d} A$, |  |

where $n$ is the concentration of charge carriers, $q$ is the charge per charge carrier, and $v_{d}$ is the magnitude of the charge carrier's drift velocity. The SI unit of current is the ampere $(\mathrm{A}): 1 \mathrm{~A}=1 \mathrm{C} / \mathrm{s}$. Current is conventionally described in terms of the flow of positive charge. The current density is the current per unit cross-sectional area, given by

$$
\vec{J}=n q \vec{v}_{d} .
$$

## Resistivity

The resistivity $\rho$ of a material is the ratio of the magnitude of the electric field to the magnitude of the current density:

$$
\rho=\frac{E}{J}
$$

Resistivity usually increases with temperature in accordance with the relationship

$$
\rho(T)=\rho_{0}\left[1+\alpha\left(T-T_{0}\right)\right]
$$

where $\alpha$ is the temperature coefficient of resistivity.

| Resistance | The resistance $R$ of a conductor is the ratio of the potential difference across the <br> conductor to the current through the conductor: |
| :--- | :--- | conductor to the current through the conductor:

$$
R=\frac{V}{I}
$$

The SI unit of resistance is the ohm $(\Omega): 1 \Omega=1 \mathrm{~V} / \mathrm{A}$. Many materials obey Ohm's law, which asserts that the potential difference across a conductor is proportional to the current through the conductor, where the proportionality constant is the conductor's resistance. The resistance of a cylindrical conductor of resistivity $\rho$, length $L$, and cross-sectional area $A$ is

$$
R=\rho \frac{L}{A}
$$

| Electric Circuits |
| :--- |
| Power in an Electric Circuit |

A complete electric circuit is a conducting loop that provides a continuous current-carrying path. An electromotive force (emf, $\mathcal{E}$ ) establishes a potential difference in the circuit that causes the charges in the circuit to move.
The power $P$ of an electric circuit is the rate of energy transferred into or out of the circuit and is equivalent to

$$
P=V I
$$

Resistors always remove energy from a circuit, converting the electrical energy to thermal energy at a rate equal to

$$
P=V I=I^{2} R=V^{2} / R
$$

## Conceptual Questions

## 1: Comparing light bulbs

Consider two light bulbs, one rated at 50 W and the other rated at 75 W , both designed to operate at 120 V. (a) Which bulb has the higher resistance? (b) Which bulb carries the greater current?

## Solution

IDENTIFY, SET UP, EXECUTE Both bulbs operate at the same voltage, so we'll examine the power equations that include voltage as a variable. (a) Power is inversely proportional to resistance at a fixed voltage ( $P=V^{2} / R$ ), so the $50-\mathrm{W}$ bulb has a higher resistance.
(b) Power is proportional to current at a fixed voltage $(P=I V)$, so the 75 -W bulb operates at a higher current.

The $75-\mathrm{W}$ bulb has less resistance than the $50-\mathrm{W}$ bulb, allowing more current to pass through the bulb and thus generating more power.

EVALUATE In comparing the relations among power, current, resistance, and voltage in electrical devices, it is important to select the proper relation to understand the problem. Using $P=I^{2} R$ for this problem would lead to a confusing analysis, since both current and resistance vary between the two bulbs.

## 2: Taking appliances to Europe

What would happen to an electric appliance made for use in North America if it were used in Europe? What would happen to an electric appliance made for use in Europe if it were used in North America? The electrical system in North America operates at 120 V , and that in Europe operates at 220 V .

## Solution

IDENTIFY, SET UP, EXECUTE The electric appliance has a fixed resistance regardless of its operating voltage. Taking the appliance made for use in North America to Europe would result in higher power consumption when the appliance was in use. Power is proportional to voltage squared $\left(P=V^{2} / R\right)$, so the power used would almost quadruple and most likely damage the appliance. When the European appliance is used in North America, the power used would decrease by almost a factor of four. The appliance would probably not function correctly, but damage should not occur as a result.

EVALUATE Inexpensive adapter plugs are often found in airport gift shops and are bought by international travelers. Many North American travelers accidentally damage their hair dryers when they use these adapters in Europe, as the adapters don't convert voltage. More expensive adapters that convert voltage are available and can be used to prevent such damage. Many small appliances (such as cellphone chargers, laptop power supplies, etc.) are made to work at multiple operating voltages.

## Problems

## 1: Heating a wire

A copper wire has a resistance of $10^{-2} \Omega$ at $20^{\circ} \mathrm{C}$. What is its resistance at $100^{\circ} \mathrm{C}$ ?

## Solution

IDENTIFY We will use the temperature dependence of resistivity to find the resistance at $100^{\circ} \mathrm{C}$, the target variable.

SET UP We will use the given resistance and temperature as the reference resistance and temperature in the resistivity equation. The thermal coefficient of expansion for copper is small, and we will ignore effects due to thermal changes in the size of the wire.
EXECUTE The resistivity varies with temperature according to

$$
\rho(T)=\rho_{0}\left[1+\alpha\left(T-T_{0}\right)\right]
$$

We find the temperature dependence of resistance by multiplying both sides by $L / A$, giving

$$
\frac{\rho(T) L}{A}=R(T)=R_{0}\left[1+\alpha\left(T-T_{0}\right)\right]
$$

The temperature coefficient of resistivity $\alpha$ is $0.00393 /{ }^{\circ} \mathrm{C}$, according to Table 25.2 in the text. The resistance at $100^{\circ} \mathrm{C}$ is then

$$
R\left(100^{\circ} \mathrm{C}\right)=\left(10^{-2} \Omega\right)\left[1+\left(0.00393 /{ }^{\circ} \mathrm{C}\right)\left(100^{\circ} \mathrm{C}-20^{\circ} \mathrm{C}\right)\right]=1.31 \times 10^{-2} \Omega
$$

EVALUATE We see that the resistance increased by $31 \%$ when the copper wire was heated by $80^{\circ} \mathrm{C}$.
We could use copper wire to make a thermometer if we have an accurate ohmmeter.

## 2: How much charge?

The current in a wire varies with time according to

$$
I=I_{0} e^{-(t / \tau)}
$$

where $\tau=1.0 \times 10^{-6} \mathrm{~s}$ and $I_{0}=2.0 \mathrm{~A}$. Find the total charge that passes through the wire (a) between $t=0$ and $t=\tau$ and (b) between $t=0$ and $t \gg \tau$.

## Solution

IDENTIFY AND SET UP Current is the rate of change of charge with respect to time, so we will integrate the current to find the charge passing through the wire in each of the two time intervals. The charge is the target variable.

EXECUTE The current is the derivative of the charge, or

$$
I=\frac{d Q}{d t}
$$

The charge is the integral of the current and is given by

$$
Q=\int I d t=\int I_{0} e^{-(t / \tau)} d t
$$

The total charge passing through the wire in time $t$ is

$$
Q=\int_{0}^{t} I_{0} e^{-\left(t^{\prime} / \tau\right)} d t^{\prime}=\left.I_{0}(-\tau) e^{-\left(t^{\prime} \mid \tau\right)}\right|_{o} ^{t}=I_{0} \tau\left(1-e^{-(t / \tau)}\right)
$$

(a) For the first time interval, $t=\tau$ and

$$
Q=I_{0} \tau\left(1-e^{-(\tau / \tau)}\right)=I_{0} \tau\left(1-e^{-1}\right)=(2.0 \mathrm{~A})\left(1.0 \times 10^{-6} \mathrm{~s}\right)\left(1-e^{-1}\right)=1.26 \times 10^{-6} \mathrm{C}
$$

(b) For the second time interval, we take $t$ to be very large (essentially infinite), so the exponential term is zero, giving

$$
Q=I_{0} \tau\left(1-e^{-(\infty / \tau)}\right)=I_{0} \tau(1)=(2.0 \mathrm{~A})\left(1.0 \times 10^{-6} \mathrm{~s}\right)=2.0 \times 10^{-6} \mathrm{C}
$$

EVALUATE This problem illustrates the nature of exponential functions as it explores the definition of current and charge. The current decreases exponentially and becomes zero after a long time. After one time interval equal to $\tau$ (called the time constant), $63 \%$ of the total charge has passed through the wire.

## 3: Current and potential in a circuit

(a) Find the current in the circuit shown in Figure 25.1. (b) Find the potential differences $V_{a b}, V_{b c}, V_{c d}$, $V_{d e}$, and $V_{e a}$. What is the sum of these potential differences?


FIGURE 25.1 Problem 3.

## Solution

IDENTIFY We will use energy conservation to find the current in the circuit. Once we have the current in the circuit, we will be able to find the various potential differences.

SET UP There are two types of components in the circuit: resistors and emfs. Energy conservation tells us that the energy produced in the circuit is equal to the energy dissipated in the circuit. Electric potential is the energy per unit charge, so we will equate the potential provided to the circuit by the emfs with the potential drops in the resistors. We will assume that the current travels counterclockwise in the circuit.

EXECUTE (a) The potential provided by the two emfs is the sum of the potentials across the batteries. In this circuit, the $18-\mathrm{V}$ emf is oriented in the direction of the current and the $3-\mathrm{V}$ emf is oriented opposite to the direction of the current. The net potential provided is then

$$
\sum \mathcal{E}=18 \mathrm{~V}-3 \mathrm{~V}=15 \mathrm{~V}
$$

The potential drops across the three resistors are equal to the current multiplied by the resistance. The total decrease in potential is

$$
\sum V_{\text {resistors }}=I(8 \Omega+3 \Omega+9 \Omega)=I(20 \Omega)
$$

Equating the increases to the decreases gives

$$
15 \mathrm{~V}=I(20 \Omega), \quad I=0.75 \mathrm{~A}
$$

(b) The potential differences are defined by

$$
V_{a b}=V_{a}-V_{b} .
$$

We use the same procedure as in part (a) to find the potential differences, keeping the preceding definition in mind. We have

$$
\begin{aligned}
V_{a b} & =I(8 \Omega)=6 \mathrm{~V}, \\
V_{b c} & =I(3 \Omega)=2.25 \mathrm{~V}, \\
V_{c d} & =-18 \mathrm{~V}, \\
V_{d e} & =I(9 \Omega)=6.75 \mathrm{~V}, \\
V_{e a} & =3 \mathrm{~V} .
\end{aligned}
$$

Note that for all of the resistors, the potential decreases as the current passes through the resistor, giving us positive values in our definition. Also, for the emfs, the potential at the first point is less than that at the second point for the $18-\mathrm{V}$ emf and greater at the first point than that at the second point for the 3 V emf. The total of the potential differences in the circuit is zero.

EVALUATE The techniques we used to solve this problem will be developed further in the next chapter. We will combine charge conservation with energy conservation and apply both to solve a multitude of electric circuit problems.

## 4: Investigating a light bulb

(a) What is the resistance of a $60-\mathrm{W}$ light bulb designed for use on a $120-\mathrm{V}$ outlet? (b) What power would the light bulb draw if connected to a $240-\mathrm{V}$ outlet?

## Solution

IDENTIFY AND SET UP We will use the relations among power, voltage, and resistance to solve the problem.

EXECUTE (a) The resistance is found by dividing the square of the voltage by the power. We have

$$
P=\frac{V^{2}}{R}, \quad R=\frac{V^{2}}{P}=\frac{(120 \mathrm{~V})^{2}}{(60 \mathrm{~W})}=240 \Omega
$$

(b) The resistance remains constant when connected to the higher voltage. The power is then

$$
P=\frac{V^{2}}{R}=\frac{(240 \mathrm{~V})^{2}}{(240 \Omega)}=240 \mathrm{~W}
$$

EVALUATE This problem confirms what we found in Question 2: The $60-\mathrm{W}$ light bulb made for $120-\mathrm{V}$ electric service would be destroyed if it were connected to $240-\mathrm{V}$ electric service.

## Try It Yourself!

## 1: A resistive thermometer

A tungsten wire is used as a thermometer in a physics research lab. At room temperature $\left(20^{\circ} \mathrm{C}\right)$, the resistance of the wire is found to be $2.0 \times 10^{-2} \Omega$. When the tungsten wire is in equilibrium with the device, the resistance is found to be $3.8 \times 10^{-3} \Omega$. What is the temperature of the device?

## Solution Checkpoints

IDENTIFY AND SET UP Use the temperature dependence of resistivity to find the temperature of the device. Ignore effects due to thermal changes in the size of the wire.

EXECUTE The temperature dependence of resistance is given by

$$
R(T)=R_{0}\left[1+\alpha\left(T-T_{0}\right)\right]
$$

The temperature coefficient of resistivity is $0.0045 /{ }^{\circ} \mathrm{C}$ for tungsten, according to Table 25.2 in the text. This gives a final temperature of $-160^{\circ} \mathrm{C}$.
EVALUATE Are you convinced that resistance can be used to measure temperature?

## 2: Internal resistance

A battery has an open-circuit potential difference of 3.0 V and a short-circuit current of 10.0 A . (a) Find its internal resistance. (b) If the battery is connected to an external resistor and then provides 1.0 A of current, what is value of the external resistance?

## Solution Checkpoints

IDENTIFY AND SET UP The internal resistance of a battery is modeled by a resistor placed in series with an emf. Use the relations among current, resistance, and voltage to find the target variables.

EXECUTE (a) When the circuit is shorted, its only resistance is due to the internal resistor. The resistance is

$$
r=\frac{\mathcal{E}}{I}=0.3 \Omega
$$

(b) Adding the external resistor adds a second resistor to the circuit. The new resistor's resistance is

$$
R=\frac{\mathcal{E}-I r}{I}=2.7 \Omega
$$

EVALUATE How much of the power provided by the emf is used by the internal resistor in part (b)?

## 3: Power in a circuit

A $12-\mathrm{V}$ battery with an internal resistance of $0.3 \Omega$ is connected to an $11.7-\Omega$ resistor. (a) What power is dissipated by the resistor? (b) If the battery is charged by connecting it to a $24-\mathrm{V}$ source, how much power does the source provide?

## Solution Checkpoints

IDENTIFY AND SET UP Use the relations among power, current, resistance, and voltage to find the target variables.

EXECUTE (a) Model the circuit with one emf and two resistances. When the circuit is connected to the $11.7-\Omega$ resistor, the current in the circuit is 1.0 A . The power dissipated by the resistor is 11.7 W .
(b) Model this circuit with two emfs and one internal resistor. The internal resistor must have a $12-\mathrm{V}$ potential difference across it (why?), so the current in the circuit is 40.0 A . The power source provides

$$
P=I V=960 \mathrm{~W}
$$

EVALUATE Can the battery be charged to 12 V with a $12-\mathrm{V}$ power supply?

## 26

## Direct-Current Circuits

## Summary

In this chapter, we will further our investigation of the movement of charge through electric circuits. We'll build on our circuit analysis foundation to analyze electric circuits constructed out of combinations of resistors, batteries, and capacitors. Conservation of charge and conservation of energy will form the basis of a powerful technique used to analyze a wide variety of circuits. Kirchhoff's rules will allow us to investigate time-varying circuits near the end of the chapter, by which time we will have built a foundation that can be applied to investigate many circuits, including those in electronic devices that we use every day.

## Objectives

After studying this chapter, you will understand

- How to recognize series and parallel combinations of resistors and how to determine their equivalent resistance.
- How to apply Kirchhoff's rules to a variety of circuits in order to learn about the flow of current in circuits.
- How to apply Kirchhoff's rules to resistor--capacitor circuits in order to bring out the dependence on time of charge, current, and voltage in these circuits.
- Several examples of electric circuits in use in everyday life.


## Concepts and Equations

Term
Resistors in Series and Parallel

## Kirchhoff's Rules

Circuits can be systematically analyzed with the use of Kirchhoff's rules. Kirchhoff's junction rule states that the algebraic sum of the currents moving into any junction is zero:

$$
\sum_{\text {at junction }} I=0 \text {. }
$$

Kirchhoff's loop rule states that the algebraic sum of the potential differences around any loop of a circuit is zero:

$$
\sum_{\text {around loop }} V=0 .
$$

A junction in a circuit is a point at which three or more conductors meet. A loop in a circuit is any closed conducting path.

Resistance-Capacitance Circuits
The charge on a capacitor and the current passing through a resistor for a series resistor-capacitor circuit being charged by a battery are given as functions of time by

$$
q=Q_{\mathrm{final}}\left(1-e^{-t / R C}\right), \quad i=I_{0} e^{-t / R C},
$$

where $Q_{\text {final }}$ is the final charge $(C \mathcal{E})$ on the capacitor and $I_{0}$ is the initial current $(\mathcal{E} / R)$ in the circuit. The charge on a capacitor and the current through a circuit for a capacitor discharged through a resistor are given as functions of time by

$$
q=Q_{0} e^{-t / R C}, \quad i=I_{0} e^{-t / R C}
$$

where $Q_{0}$ and $I_{0}$ are, respectively, the initial charge and the current on the capacitor. The time required for a significant change in the amount of charge is given by the time constant $\tau=R C$.

## Conceptual Questions

## 1: Where to start?

Your task is to find the equivalent resistance for the three circuits shown in Figure 26.1. Identify which resistors you'd combine first and the type of combination in which they are configured.

## Solution

IDENTIFY, SET UP, AND EXECUTE For problems of this type, we'll need to find two resistors that are connected either in parallel or in series. These two resistors will be our starting point for calculating the equivalent resistance.


(c)

Figure 26.1 Question 1.
In Figure 26.1a, we see that most of the resistors are not in series or parallel with adjoining resistors. However, in the upper half of the circuit, resistors $R_{3}$ and $R_{4}$ are connected in series with each other. Recall that a series connection implies that the same current must pass through both resistors. There is no junction between resistors $R_{3}$ and $R_{4}$ so the same current passes through both. We should start, then, by combining resistors $R_{3}$ and $R_{4}$ in series.

Figure 26.1 b is a bit more complicated. All of the resistors are positioned in a straight line on the page, suggesting that they are all in series. However, the connections around the middle two resistors are not in a straight electrical line. The current leaving $R_{1}$ can go into either $R_{2}$ or $R_{3}$. The current passing through both $R_{2}$ and $R_{3}$ combines and passes through $R_{4}$. Therefore, $R_{2}$ and $R_{3}$ are connected in parallel. We should start by combining resistors $R_{2}$ and $R_{3}$ in parallel.

Figure 26.1c is also complicated. A quick scan shows that there are no resistors in series. Resistors $R_{1}, R_{2}$, and $R_{3}$ are each connected to junctions with more than one component, indicating that these resistors are not connected in parallel. However, looking at the bottom half of the figure, we see that the current leaving $R_{2}$ splits, goes through $R_{4}$ and $R_{5}$, and then rejoins before entering $R_{3} . R_{4}$ and $R_{5}$ are connected in parallel. Accordingly, we should start by combining resistors $R_{4}$ and $R_{5}$ in parallel.

EVALUATE Identifying the starting point for calculating the equivalent resistance is of ten the most difficult part of a problem. Many circuits include a variety of combinations. Developing an eye for recognizing series and parallel circuits helps simplify what might appear to be a complicated circuit. One way to identify the starting point is to imagine cutting out two resistors. If you can cut them out without disrupting the rest of the circuit, these resistors probably represent a good starting point.

## 2: Checking equations

Examine each of the following equations derived from the circuit shown in Figure 26.2, and determine whether they are valid or erroneous:
(a) $-I_{2}-I_{3}+I_{1}=0$
(b) $-I_{1}(14 \Omega)-I_{1}(32 \Omega)+12 \mathrm{~V}+24 \mathrm{~V}=0$
(c) $28 \mathrm{~V}+I_{2}(32 \Omega)=I_{3}(23 \Omega)$
(d) $-I_{1}(14 \Omega)-I_{3}(23 \Omega)+16 \mathrm{~V}-I_{2}(32 \Omega)+24 \mathrm{~V}=0$


Figure 26.2 Question 2.

## Solution

IDENTIFY, SET UP, AND EXECUTE The equations are derived from Kirchhoff's rules, so we will use the junction and loop rules to guide our analysis.

The first equation appears to come from the junction rule. If we were to apply the junction rule to the bottom junction and assign positive values to currents entering that junction, we would get equation (a). We conclude that equation (a) is valid.

The remaining equations appear to come from the loop rule applied to different loops. We'll use the values of the resistors to determine the loop and confirm that the equation is consistent with traversing the loop. Equation (b) involves the $14-\Omega$ and $32-\Omega$ resistors, so we'll examine the left loop of the circuit. The first term is negative; thus, the $14-\Omega$ resistor was traversed in the direction of the current. Continuing around the loop clockwise, we confirm that the next term is correct. Passing over the $12-\mathrm{V}$ emf should give a negative sign, but the equation has a positive sign. This is incorrect and should be fixed. The last term corresponds to the $24-\mathrm{V}$ emf and has the correct sign. We conclude that equation (b) is incorrect.

Equation (c) involves the $32-\Omega$ and $23-\Omega$ resistors, so we'll examine the right loop of the circuit. The first term is positive; hence, an emf was traversed from the negative to the positive terminal. The value 28 V does not correspond to any emf, but it does correspond to the sum of the $12-\mathrm{V}$ and $16-\mathrm{V}$ emfs in the right loop. Continuing around the loop counterclockwise, we confirm that the next term is correct as we traverse the $32-\Omega$ resistor against the current. Next, we travel across the $23-\Omega$ resistor, which should give a negative value. However, the $23-\Omega$ resistor's term has been moved to the other side of the equation, which is correct. We conclude that equation (c) is valid.

Equation (d) involves all three resistors. There is no single loop that includes all three currents, so we should be suspicious. The equation could correspond to several equations combined, so we need to continue checking. The first two terms are negative, corresponding to traversing the $14-\Omega$ and $23-\Omega$ resistors from left to right. Continuing clockwise around the outer loop, we travel across the $16-\mathrm{V}$ emf, consistent with the next term in the equation. A term corresponding to the potential difference across the $32-\Omega$ resistor is next and appears out of place. Then comes a term that corresponds to the $24-\mathrm{V} \mathrm{emf}$, also consistent with a clockwise trip around the outer loop. Except for the fourth term, the equation is valid. We conclude that equation (d) is incorrect, since it has the extra term for the potential difference across the $32-\Omega$ resistor.

We have found that equations (a) and (c) are correct and (b) and (d) are incorrect.
EVALUATE Generating equations from Kirchhoff's rules is straightforward, but may involve many terms with various signs. Double-checking your equations before solving for the unknowns can save aggravation and wasted effort.

CAUTION Each Branch has a Different Current! In this problem we see that the current in each branch was labeled separately, indicating that the currents in the branches may be different. A common error is to label all branches of a circuit with the same current, resulting in an inconsistent solution.

## 3: Blown fuses

In an introductory physics laboratory, students often blow fuses when using ammeters, but rarely blow fuses when using voltmeters. Why? (In each case, the fuse is placed in series with the measuring device.)

## Solution

IDENTIFY, SET UP, AND EXECUTE To solve this problem, we need to consider the design and operation of the two measuring devices. An ammeter measures current in a circuit and is placed in series with the branch being measured. To minimize the effect on the circuit, ammeters have very little resistance. A voltmeter measures the potential difference between two points in a circuit and is placed in parallel with the components being measured. To minimize the effect on the circuit, voltmeters have very high resistance.

If an ammeter is placed in series in a circuit, it should perform as expected. If an ammeter is placed in parallel in a circuit, then a large current may flow through the meter, since it has little resistance. Depending on where the ammeter is placed in parallel, the current passing through the meter could exceed the design limit and blow the fuse. For example, if a student tries to measure the total current provided to a circuit by placing the ammeter leads across the emf supply, then the maximum current available from the supply will pass through the ammeter.

If a voltmeter is placed in series in a circuit, its high resistance would prevent excessive current. The reading of the voltmeter would not be accurate, but no damage to the voltmeter would occur.

EVALUATE Understanding how electrical measuring devices operate will improve your use of those devices, as well as help you interpret the differences between series and parallel combinations.

## 4: Short- and long-term behavior in an RC circuit

An initially uncharged capacitor (with capacitance $C$ ) and a resistor (with resistance $R$ ) are connected in series with an emf (with voltage V ). Just after the connection is made, what are the potentials across the resistor and capacitor, the charge on the capacitor, and the current in the circuit? A very long time after the connection is made, what are the potentials across the resistor and capacitor, the charge on the capacitor, and the current in the circuit?

## Solution

IDENTIFY, SET UP, AND EXECUTE This problem asks us to consider how charge, current, and potential difference vary in an RC circuit. We'll work through each part, using the techniques we've learned in the chapter.

Just after the connection is made, there is essentially no charge on the capacitor. Therefore, the potential difference across the capacitor is zero. (Recall that $V=Q / C$.) If the potential across the capacitor is zero, then the potential across the resistor must be equal to the voltage $V$ across the emf. The current through the resistor is then $I=V / R$, the maximum current supplied to the circuit.

A long time after the connection is made, the capacitor is fully charged. Therefore, the current in the circuit is zero, because no charges are moving. The potential difference across the resistor must also be zero, since there is no current. The potential difference across the capacitor is equal to the voltage $V$ of the emf, since the voltage across the resistor is zero.

EVALUATE We see that the circuit behaves as if there is only a resistor just after the connection is made and as if there is only a capacitor after a long time. Keeping this long-term and short-term behavior in mind will help you interpret RC circuits.

## Problems

## 1: Combining resistances

Find the equivalent resistance between points $a$ and $b$ in the circuit shown in Figure 26.3.


Figure 26.3 Problem 1.

## Solution

IDENTIFY The network is a combination of series and parallel resistors, so we will identify and replace combinations that are purely in series or purely in parallel. The target variable is the equivalent resistance.

SET UP We redraw the circuit as shown in Figure 26.4a, without the resistance values shown for convenience. Examining the circuit, we find that resistors $R_{2}$ and $R_{3}$ are in series, since there is only one path for the current to flow through them. We replace these resistors with their equivalent resistance $R_{23}$, shown in Figure 26.4b. $R_{4}$ is combined in parallel with $R_{23}$, since the current could flow through $R_{4}$ or $R_{23}$ (Figure 26.4c). Next, we add $R_{5}$ in series with $R_{234}$ (Figure 26.4d). Finally, we add $R_{1}$ in parallel with $R_{2345}$ to find the equivalent resistance of the circuit (Figure 26.4).


Figure 26.4 Problem 1.

EXECUTE We have determined how to combine the resistors. We now find the numeric values of the resistances. We have

$$
R_{23}=R_{2}+R_{3}=21 \Omega+14 \Omega=35 \Omega
$$

Next, we combine resistors $R_{23}$ and $R_{4}$ in parallel:

$$
\frac{1}{R_{234}}=\frac{1}{R_{23}}+\frac{1}{R_{4}}=\frac{1}{35 \Omega}+\frac{1}{26 \Omega}, R_{234}=14.9 \Omega
$$

Combining resistors $R_{234}$ and $R_{5}$ in series yields

$$
R_{2345}=R_{234}+R_{5}=14.9 \Omega+33 \Omega=47.9 \Omega
$$

Finally, combining resistors $R_{2345}$ and $R_{1}$ in parallel gives

$$
\frac{1}{R_{12345}}=\frac{1}{R_{2345}}+\frac{1}{R_{1}}=\frac{1}{47.9 \Omega}+\frac{1}{57 \Omega}, R_{12345}=26.0 \Omega
$$

The equivalent resistance of the circuit is $26.0 \Omega$.
EVALUATE The key to finding the equivalent resistance is identifying where to begin. Once we found which two resistors to combine first, the remainder of the problem followed directly. You can see two patterns that developed in this problem. First, after we combined the first pair of resistors, each additional combination included the resistor combination from the previous step. Second, the types of combinations alternated between series and parallel. These two patterns show up in many resistor combination problems, so keep an eye out for them.

Developing an eye for identifying series and parallel connections results from examining and solving circuit problems. The procedure is much like the one we developed in the previous chapter for combining capacitors with resistors and for current replacing capacitors and charge.

CAUTION Watch Your Inverse Sums! The most common mistake found in adding resistors in parallel is not inverting the inverse sum of the resistances, much like the common mistake that is made when capacitors are combined in parallel. Carefully watching your units also helps identify mistakes.

## 2: Practicing Kirchhoff's rules

Find the values of $\mathcal{E}_{1}, \mathcal{E}_{2}$, and $I$ in the circuit shown in Figure 26.5.


Figure 26.5 Problem 2.

## Solution

IDENTIFY We will apply Kirchhoff's rules to the circuit to solve for the target varialbles.
SET UP We'll take current $I$ to be directed upwards, as shown in the diagram. We will need three equations for the three unknowns.

EXECUTE We find the current $I$ from the junction rule. Examining the lower junction and taking currents entering the junctions as positive, we have

$$
\begin{aligned}
& 3.0 \mathrm{~A}-I-1.2 \mathrm{~A}=0 \\
& I=1.8 \mathrm{~A}
\end{aligned}
$$

To find $\mathcal{E}_{1}$, we use the loop rule and apply it to the left loop. Starting at the top left corner and traveling around the loop clockwise, we obtain

$$
-\mathcal{E}_{1}+I(17 \Omega)-(1.2 \mathrm{~A})(23 \Omega)=0
$$

Substituting the value for the current and solving gives

$$
\mathcal{E}_{1}=(1.8 \mathrm{~A})(17 \Omega)-(1.2 \mathrm{~A})(23 \Omega)=3.0 \mathrm{~V}
$$

To find $\mathcal{E}_{2}$, we use the loop rule and apply it to the outer loop. Starting at the top left corner and traveling around the loop clockwise yields

$$
\begin{aligned}
& \mathcal{E}_{2}-(1.2 \mathrm{~A})(23 \Omega)=0 \\
& \mathcal{E}_{2}=(1.2 \mathrm{~A})(23 \Omega)=27.6 \mathrm{~V}
\end{aligned}
$$

We have found that $\mathcal{E}_{1}=3.0 \mathrm{~V}, \mathcal{E}_{2}=27.6 \mathrm{~V}$, and $I=1.8 \mathrm{~A}$.
EVALUATE This problem has shown us how to apply Kirchhoff's rules to a circuit. In the next problem, we'll apply Kirchhoff's rules to a similar circuit. The procedure for solving that circuit will be similar, but the algebra is a bit more complicated.

## 3: Two-loop circuit

Find the currents $I_{1}, I_{2}$, and $I_{3}$ in the circuit shown in Figure 26.6.


Figure 26.6 Problem 3.

## Solution

IDENTIFY This is a two-loop circuit, and we will need to apply both of Kirchhoff's rules.
SET UP With the three emf's in the three branches of the circuit, we will need to apply Kirchhoff's rules to find the currents. We will need to apply the junction rule once and the loop rule twice to find three equations that will lead to the solution. The three currents and their directions are indicated in the figure. We'll use the two loops shown as well.

EXECUTE We will find three equations and then manipulate them to find the three currents. One equation comes from the junction rule. Examining the top junction, where the three currents meet, we see that current $I_{3}$ enters the junction and currents $I_{1}$ and $I_{2}$ exit the junction. Taking currents entering the junction as positive and currents leaving the junction as negative, we have

$$
-I_{1}-I_{2}+I_{3}=0
$$

Two more equations come from applying the loop rule to the two loops. Starting in the upper left corner of loop l, we proceed clockwise and find that

$$
12 \mathrm{~V}-I_{2}(3 \Omega)-6 \mathrm{~V}+I_{1}(4 \Omega)=0
$$

Note that the $12-\mathrm{V}$ emf produces an increase in potential, since we traveled from the negative to the positive terminal, the $3-\Omega$ resistor gives a decrease in potential, since we traveled with the current across the resistor, the $6-\mathrm{V}$ emf generates a decrease in potential, since we traveled from the positive to the negative terminal, and the $4-\Omega$ resistor produces an increase in potential, since we traveled against the current across the resistor. We'll also start in the upper left corner of loop 2 and proceed clockwise to find that

$$
I_{3}(6 \Omega)+20 \mathrm{~V}+I_{2}(3 \Omega)-12 \mathrm{~V}=0
$$

Here, the $6-\Omega$ resistor gives an increase in potential, since we traveled against the current across the resistor; the $20-\mathrm{V}$ emf results in an increase in potential, since we traveled from the negative to the positive terminal; the $4-\Omega$ resistor produces an increase in potential, since we traveled against the current across the resistor; and the $12-\mathrm{V}$ emf gives a decrease in potential, since we traveled from the positive to the negative terminal.

At this point, we can proceed to solve for the currents in several ways. Let's use the junction rule equation to substitute $I_{1}+I_{2}$ for $I_{3}$ in the equation for loop 2 :

$$
\left(I_{1}+I_{2}\right)(6 \Omega)+20 \mathrm{~V}+I_{2}(3 \Omega)-12 \mathrm{~V}=0
$$

Simplifying, we obtain

$$
I_{1}(6 \Omega)+I_{2}(9 \Omega)+8 \mathrm{~V}=0
$$

We rewrite this equation to find $I_{2}$ in terms of $I_{1}$ :

$$
I_{2}=\frac{-I_{1}(6 \Omega)-8 \mathrm{~V}}{(9 \Omega)}
$$

Using this expression for $I_{2}$ in the equation for loop 1 gives

$$
12 \mathrm{~V}-\left(\frac{-I_{1}(6 \Omega)-8 \mathrm{~V}}{(9 \Omega)}\right)(3 \Omega)-6 \mathrm{~V}+I_{1}(4 \Omega)=0
$$

Solving, we get

$$
I_{1}=-\frac{26 \mathrm{~V}}{18 \Omega}=-1.44 \mathrm{~A}
$$

We now replace $I_{1}$ in the previous expressions to solve for the other currents. We find $I_{2}$ from

$$
I_{2}=\frac{-I_{1}(6 \Omega)-8 \mathrm{~V}}{(9 \Omega)}=\frac{-(-1.44 \mathrm{~A})(6 \Omega)-8 \mathrm{~V}}{(9 \Omega)}=+0.0711 \mathrm{~A}
$$

Current $I_{3}$ is the sum $I_{1}+I_{2}$ :

$$
I_{3}=I_{1}+I_{2}=-1.44 \mathrm{~A}+0.0711 \mathrm{~A}=-1.37 \mathrm{~A}
$$

We found current $I_{1}$ to be $-1.44 \mathrm{~A}, I_{2}$ to be 0.0711 A , and $I_{3}$ to be -1.37 A ; the negative signs indicate that the currents are opposite the directions indicated.

EVALUATE We check our results by replacing the currents with their values in each of the three equations. If the three equations are satisfied (i.e., if they each sum to zero), then we conclude that our numeric results match the equations. We should also double-check the three equations to prevent any sign errors.

Kirchhoff's rules offer multiple paths, all of which lead to the correct result. Here, you could have started at any point in the loop circuit, proceeded either clockwise or counterclockwise, and found equivalent equations. You could also have written an equation for the outer loop to replace one of the other loop equations. You can proceed to apply the algebra in a variety of ways, solving for currents in different orders. Experience will show that all methods result in the same answers when applied consistently and correctly.

CAUTION Watch Signs in Loop Problems Careful interpretation of signs in multiloop problems is critical to obtaining an accurate result. Common mistakes occur in forming the equations, as well as in interpreting negative current solutions. Label your currents carefully, check and recheck your equations before solving for unknowns, and interpret negative currents as moving opposite to the initial direction.

## 4: Investigating an RC circuit

A $5.00-\mu \mathrm{F}$ capacitor that is initially uncharged is connected in series with a $4.75-\mathrm{M} \Omega$ resistor and a $350-\mathrm{V}$ emf supply. How long after the circuit is completed does the capacitor reach $90 \%$ of its maximum charge? What is the potential difference across the resistor at that time? How much power is the emf providing at that time?

## Solution

IDENTIFY We will use the equations developed for charging RC circuits to solve the problem.
SET UP The charge on a capacitor and the current in the circuit are not constant in an RC circuit; equation 21.12 in the text gives the time dependence of the charge. We'll use this equation and the relation among charge, current, voltage, and power to solve the problem.

EXECUTE The charge on a capacitor in series with a resistor and an emf is given by

$$
q=Q_{f}\left(1-e^{-t / R C}\right)
$$

In this problem, we want to find the time when the charge $q$ is $90 \%$ of $Q_{f}$. We are given the values of resistance and capacitance, so we substitute to find the time:

$$
q=(90 \%) Q_{f}=0.90 Q_{f}=Q_{f}\left(1-e^{-t / R C}\right)
$$

Canceling $Q_{f}$ and rearranging terms gives

$$
0.10=e^{-t / R C}
$$

To find the time, we take the natural logarithm of both sides:

$$
\begin{aligned}
& \ln (0.10)=\ln \left(e^{-t / R C}\right)=-\frac{t}{R C} \\
& t=-R C \ln (0.10)=-(4.75 \mathrm{M} \Omega)(5.00 \mu \mathrm{~F})(-2.30)=54.6 \mathrm{~s}
\end{aligned}
$$

The potential difference across the resistor is the product of the current and the resistance. The current as a function of time is

$$
i=\frac{d q}{d t}=I_{0} e^{-t / R C}
$$

where $I_{0}$ is the maximum current $(\mathcal{E} / R)$. The potential difference across the resistor is

$$
V=i R=\left(I_{0} e^{-t / R C}\right) R=\left(\frac{\mathcal{E}}{R} e^{-t / R C}\right) R=\mathcal{E} e^{-t / R C}=(350 \mathrm{~V}) e^{-(54.6 \mathrm{~s}) /(4.75 \mathrm{M} \Omega)(5.00 \mu \mathrm{~F})}=35.1 \mathrm{~V}
$$

The power provided by the emf is product of the current and the voltage:

$$
P=i V=\left(I_{0} e^{-t / R C}\right) \mathcal{E}=I_{0} \mathcal{E} e^{-t / R C}=\frac{\mathcal{E}^{2}}{R} e^{-t / R C}=\frac{(350 \mathrm{~V})^{2}}{(4.75 \mathrm{M} \Omega)} e^{-(54.6 \mathrm{~s})((4.75 \mathrm{M} \Omega)(5.00 \mu \mathrm{~F})}=2.59 \mathrm{~mW}
$$

The capacitor reaches $90 \%$ of its maximum charge 54.6 s after the circuit is connected; the resistor has a 35.1-V potential difference across it at that time; and the emf provides 2.59 mW of power at that time.

EVALUATE We see that it takes almost 1 minute for the capacitor to reach $90 \%$ of its maximum charge. That it takes that long is due to the resistor and capacitor and does not depend on the value of the emf. We also see that, at this time, only $10 \%$ of the resistor's initial potential difference remains and the emf is supplying little power to the circuit.

## Try It Yourself!

## 1: Equivalent resistance

Find the equivalent resistance of the resistor network shown in Figure 26.7.


Figure 26.7 Try It Yourself Problem 1.

## Solution Checkpoints

IDENTIFY AND SET UP Identify series and parallel groups of resistors, and then replace each group with an equivalent resistance. Continue until all groups are combined. Draw new sketches as each group is replaced.

EXECUTE The two top resistors are combined in series. The bottom left two resistors (the $3-\Omega$ and $6-\Omega$ resistors) are combined in parallel. This leaves a network of five resistors, shown in Figure 26.8.


Figure 26.8 Try It Yourself Problem 1.

At this point, it should be clear that the top three resistors are connected in parallel and the bottom two are connected in series. When these two groups are combined, the two equivalent resistors ( $1.5 \Omega$ on top and $3 \Omega$ on the bottom) are combined in parallel, giving a total combined equivalent resistance of $1 \Omega$.

EVALUATE Can you find another order in which to combine the resistors? Do you find the same result?

## 2: Two-loop circuit

Find the currents $I_{1}, I_{2}$, and $I_{3}$ in the circuit shown in Figure 26.9.


Figure 26.9 Try It Yourself Problem 2.

## Solution Checkpoints

IDENTIFY AND SET UP Apply both of Kirchhoff's rules to the circuit to find three equations that will lead to the three unknown currents. Use the directions of the currents given in the figure.

EXECUTE Taking currents entering the left junction as positive and currents leaving the left junction as negative gives

$$
I_{1}-I_{2}+I_{3}=0
$$

Three equations can be found from the loop rule. Starting at the upper left corner of the top loop and proceeding clockwise gives

$$
-12 \mathrm{~V}+I_{1}(6 \Omega)+6 \mathrm{~V}+I_{2}(2 \Omega)=0
$$

Repeating for the lower loop, starting at the left junction and proceeding clockwise yields

$$
-I_{2}(2 \Omega)-6 \mathrm{~V}-I_{3}(6 \Omega)+4 \mathrm{~V}=0
$$

Finally, for the outer loop, starting at the top left corner and proceeding clockwise gives

$$
-12 \mathrm{~V}+I_{1}(6 \Omega)-I_{3}(6 \Omega)+4 \mathrm{~V}=0
$$

Two of the three equations can be combined with the junction rule equation to solve for the three currents. Solving, we get

$$
I_{1}=\frac{13}{15} \mathrm{~A}, \quad I_{2}=\frac{2}{5} \mathrm{~A}, \quad I_{3}=-\frac{7}{15} \mathrm{~A}
$$

The negative value for $I_{3}$ indicates that $I_{3}$ is opposite to the direction shown in the figure.
EVALUATE Check your results by replacing the currents with their values in each of the three equations.

## 3: An RC circuit

Two capacitors and a resistor are connected to a battery by a switch, as shown in Figure 26.10. (a) Find the final charge on each capacitor after the switch has been closed for a long time. (b) Find the time it takes for the charges and potential differences on the capacitors to reach half their final values.


Figure 26.10 Try It Yourself Problem 3.

## Solution Checkpoints

IDENTIFY AND SET UP After a long time, both capacitors are fully charged and there is no current in the circuit. The capacitors can then be combined and replaced with their equivalent capacitance. Use the relation for a charging RC circuit.

EXECUTE (a) The equivalent capacitance of the two capacitors is $2 / 3 \mu \mathrm{~F}$. After a long time, the voltage across the capacitors is 6 V , so there is $4 \mu \mathrm{C}$ of charge on each capacitor. The potentials across the two capacitors are 4 V and 2 V .
(b) Both the charge and the voltage reach half their final value at the same time, since the charge and voltage are proportional to each other. Both capacitors charge at the same rate, so we can find the time required to reach half the charge on one capacitor, which will be the same as the time required to reach half the charge on the other capacitor. This time can be found from

$$
\frac{Q}{Q_{\mathrm{f}}}=\left(1-e^{-t / R C}\right)=\frac{1}{2} .
$$

Taking the natural logarithm of both sides gives

$$
t=R C \ln 2 .
$$

The time constant is $12 \mu \mathrm{~s}$, so $t=8.32 \times 10^{-6} \mathrm{~s}$.
EVALUATE What is the initial charge on the capacitors? What is the initial current in the circuit?

## Magnetic Field and Magnetic Forces

## Summary

We will investigate magnetism and magnetic forces in this chapter. The magnetic interaction can be either attractive or repulsive, much like the electric interaction. But, as we shall see, the magnetic force is more complicated than the electric force, as it depends on two vectors: the magnetic field and the velocity of a charge moving through the field. We'll use magnetic field lines to get a visual representation of the magnetic field. We will investigate the motion of charged particles in magnetic fields and apply our knowledge to several common applications. We'll also investigate the net force, torque, and energy associated with a current-carrying loop. By the end of the chapter, we'll have a good understanding of magnetic fields and the magnetic force. In the next chapter, we'll learn how magnetic fields are generated.

## Objectives

After studying this chapter, you will understand

- How to calculate the magnetic force between permanent magnets.
- How to calculate the direction and magnitude of the magnetic force for moving charges and currents.
- How magnetic field lines are used to represent the magnetic field.
- How to analyze the motion of charged particles in a magnetic field.
- How to apply magnetism to the several practical applications, including the velocity selector and mass spectrometer.
- How to calculate the force and torque on a current-carrying loop.

| Term |
| :--- |
| Magnetic Forces |
| Magnetic Fields |

## Magnetic Force on a CurrentCarrying Conductor

## Description

Bar magnets have north ( N ) and south ( S ) poles. Opposite poles attract and like poles repel. Both moving charges and currents are acted upon by a force in the presence of a magnetic field. The direction of the magnetic force is perpendicular to both the magnetic field and the direction of the moving charge and is given by the right-hand rule. The magnitude of the magnetic force is given by

$$
\vec{F}=q \vec{v} \times \vec{B}
$$

where $q$ is the charge, $v$ is the velocity, and $B$ is the magnetic field. The SI unit of magnetic field is the tesla $(T): 1 \mathrm{~T}=1 \mathrm{~N} /(\mathrm{A} \cdot \mathrm{m})$. Also commonly used to measure the magnetic field is the gauss $(\mathrm{G})$, where $1 \mathrm{G}=10^{-4} \mathrm{~T}$.

Magnetic field lines are used to represent the magnetic field graphically. The direction of the magnetic field is tangent to the magnetic field line, and the magnitude is proportional to the density of magnetic field lines. Magnetic flux is defined analogously to the electric flux as

$$
\Phi_{B}=\int \vec{B} \cdot d \vec{A}
$$

and has units of webers $\left(1 \mathrm{~Wb}=1 \mathrm{~T} \mathrm{~m}^{2}\right)$. The net magnetic flux through any closed surface is zero:

$$
\oint \vec{B} \cdot d \vec{A}=0
$$

Gauss's law for magnetism indicates that magnetic field lines always close in on themselves.

The magnetic force acts in a direction perpendicular to the velocity and therefore influences only the direction, and not the magnitude, of the particle's velocity. In a uniform magnetic field, a particle with initial velocity perpendicular to the field will move in a circle of radius $R$ given by

$$
R=\frac{m v}{|\boldsymbol{q}| B} .
$$

The magnetic force on a current-carrying conductor is due to the magnetic force on the individual charges moving within the conductor. The force is given by

$$
\vec{F}=\vec{l} \vec{l} \times \vec{B}
$$

For a small segment of wire, the contribution to the force is

$$
d \vec{F}=I d \vec{l} \times \vec{B}
$$

In a uniform magnetic field, the total magnetic force acting on a current-carrying loop is zero. The magnetic force creates a torque on the loop of magnitude

$$
\tau=I A B \sin \phi
$$

where $I$ is the current in the loop, $A$ is the cross-sectional area of the loop, and $\phi$ is the angle between the normal to the loop and the direction of the magnetic field. The torque tends to rotate the loop toward decreasing $\phi$. The torque on the loop can be expressed in terms of the magnetic moment of the loop, given by

$$
\vec{\mu}=I \vec{A}
$$

The torque on the current loop is

$$
\vec{\tau}=\vec{\mu} \times \vec{B}
$$

The potential energy of the current loop is

$$
U=-\vec{\mu} \cdot \vec{B}
$$

## Conceptual Questions

## 1: Signs of charges

Four particles enter a region of uniform magnetic field. Their trajectories are shown in Figure 27.1. What are the signs of the charge of all four particles?


Figure 27.1. Question 1.

## Solution

IDENTIFY, SET UP, AND EXECUTE The magnetic force acts perpendicular to a particle's velocity and to the direction of the magnetic field. By examining the figure, we can determine the direction in which the force acts and then find the sign of the particle's charge. The right-hand rule relates the force to the direction of the magnetic field and the particle's velocity. With a magnetic force directed into the page, a positive particle will deflect to the left.

Particle A deflects to the right as it passes through the magnetic field, so we conclude that it is negatively charged. Particle B also deflects to the right, so it also must be negatively charged. Note that it may appear that particle B deflects opposite to particle A , but it is also traveling in the opposite direction. Particle $C$ deflects to the left, so it is positively charged. Particle D travels along a straight line and so must be neutrally charged.

EVALUATE This problem gives us practice applying the right-hand rule to find the magnetic force acting on a moving charge. We will be applying the right-hand rule to a variety of problems in the next two chapters.

The technique set forth here is used by elementary-particle physicists to determine the signs of particles they discover at accelerator facilities.

CAUTION Practice the Right-Hand Rule We will use the right-hand rule in a variety of applications in the next two chapters. You must practice to become proficient at its use, so you should check and confirm every problem that involves the right-hand rule in this guide.

## 2: The odd magnetic force

We've seen that a charged particle moving through a magnetic field is acted upon by a force. In addition, we've seen that a net force causes acceleration. But we also learned that a charged particle
moving through a uniform magnetic field moves with constant speed. How can this charged particle accelerate without changing speed?

## Solution

IDENTIFY, SET UP, AND EXECUTE Acceleration and velocity are vectors and thus have both magnitude and direction. Here, the force-and therefore the acceleration-is perpendicular to the direction of motion. Since the acceleration is perpendicular to the motion, the magnitude of the velocity (i.e., the speed) does not change. Instead, the acceleration causes the direction of the particle's velocity to change.

We can also consider the work done on the charged particle as it moves through the field. Since the force acts perpendicular to the displacement, no work is done and the speed remains constant.

EVALUATE Where have we seen motion similar to this before? When we learned about circular motion, we saw that a centripetal force could change the direction of motion of an object while keeping the object's speed constant.

## 3: Protons with differing velocities in a magnetic field

Two protons with different velocities enter a region having a uniform magnetic field that is perpendicular to their velocities. The region is large enough that the protons can execute complete circular trajectories. How do the radii of their circular paths compare? Which particle takes longer to complete one revolution?

## Solution

IDENTIFY, SET UP, AND EXECUTE A charged particle moving in a magnetic field follows a trajectory with radius

$$
r=\frac{m v}{|q| B} .
$$

The two protons have the same charge and mass and are in the same magnetic field. They do have different velocities, however, and the proton with the higher speed will move in a circular path with a greater radius.

The time to complete one revolution $(T)$ can be found from the velocity $(v)$ and the circumference $(2 \pi r)$ :

$$
v=\frac{2 \pi r}{T}, \quad T=\frac{2 \pi r}{v}
$$

Substituting for the radius, we find that

$$
T=\frac{2 \pi r}{v}=\frac{2 \pi}{v} \frac{m v}{|q| B}=\frac{2 \pi m}{|q| B}
$$

This expression shows us that the time required to complete one revolution (the period) is independent of the velocity. Both protons complete one revolution in the same amount of time.

EVALUATE This result is the basis for one type of particle accelerator: the cyclotron. Cyclotrons use varying electric fields to accelerate particles inside a uniform magnetic field. Inside the cyclotron, particles with different velocities move in unison, allowing for all particles to be accelerated at the same time.

## Problems

## 1: Flux through a wedge

Calculate the flux through the surfaces $A B C D$ and $A E F D$ shown in Figure 27.2 for a constant and uniform magnetic field $B=0.8 \mathrm{~T}$ directed along the positive $y$-axis.


Figure 27.2 Problem 1.

## Solution

IDENTIFY The target variable is the magnetic flux through the two surfaces. The flux is related to the magnetic field lines passing through the area.

SET UP The flux is the integral of the scalar product of the magnetic field and the area vector. The magnetic field is constant and uniform, so it is taken out of the integral. The area vector has magnitude equal to the area and is directed along the outward normal to the surface.

EXECUTE For surface $A B C D$, the area vector has magnitude $17.5 \mathrm{~m}^{2}$ and is rotated slightly from the magnetic field. The angle between the magnetic field and the area vector is $90^{\circ}-\theta$ and is found by examining the sides of the wedge:

$$
\tan \left(90^{\circ}-\theta\right)=\frac{3}{4}, \quad 90^{\circ}-\theta=36.9^{\circ}
$$

The flux through $A B C D$ is then

$$
\phi=\int \vec{B} \cdot d \vec{A}=B A \cos \left(90^{\circ}-\theta\right)=(0.8 \mathrm{~T})\left(17.5 \mathrm{~m}^{2}\right)(0.80)=11.2 \mathrm{~Wb}
$$

For surface $A E F D$, the area vector has magnitude $14.0 \mathrm{~m}^{2}$ and is directed opposite and parallel to the magnetic field. The flux through $A E F D$ is then

$$
\phi=\int \vec{B} \cdot d \vec{A}=B A \cos \left(180^{\circ}\right)=(0.8 \mathrm{~T})\left(14.0 \mathrm{~m}^{2}\right)(-1)=-11.2 \mathrm{~Wb}
$$

EVALUATE We see that the fluxes through the two surfaces are equal and opposite. If we examine the other three surfaces, we see that the area vector is perpendicular to the magnetic field, giving no flux
for any of the three surfaces. When we sum up the net flux through all surfaces, we must get zero, and we do.

As with electric flux, we can associate the magnetic flux with the number of field lines exiting a surface minus the number of field lines entering the surface. In this problem, any field line entering the surface on side $A E F D$ exits on side $A B C D$, indicating that the fluxes through the two surface are equal and opposite.

## 2: Mass spectrometer

Protons enter a mass spectrometer that consists of a velocity selector followed by a region with a 1.40-T uniform magnetic field, as shown in Figure 27.3. The velocity selector consists of two plates separated by 1.5 mm , with a potential difference of 125 V between the plates and a uniform magnetic field directed through the plates and perpendicular to the path of the protons. The protons strike the detector after traveling a distance of 22.0 mm following their exit from the velocity selector. Find the magnetic field inside the velocity selector.


Figure 27.3 Problem 2.

## Solution

IDENTIFY We will find the magnetic field inside the velocity selector by finding the electric field and the velocity of the protons.

SET UP To find the electric field, we use the potential difference and the separation between the plates in the velocity selector. We find the velocity of the protons from their path in the uniform magnetic field.

EXECUTE The velocity of the protons is determined from their path in the uniform magnetic field. The protons strike 22.0 mm from where they exit the velocity selector. The radius of their trajectory is half of this value. Ions in a magnetic field follow a circular path of radius

$$
R=\frac{m v}{|q| B} .
$$

We can rearrange this equation to solve for the protons' velocity. The protons each have a charge of $+e$ and a mass of $1.67 \times 10^{-27} \mathrm{~kg}$. Their velocity is

$$
v=\frac{|q| R B}{m}=\frac{+e(D / 2) B}{m}=\frac{\left(1.60 \times 10^{-19} \mathrm{C}\right)(0.022 \mathrm{~m}) / 2(1.40 \mathrm{~T})}{\left(1.67 \times 10^{-27} \mathrm{~kg}\right)}=1.475 \times 10^{6} \mathrm{~m} / \mathrm{s}
$$

The electric field between the plates of the velocity selector is

$$
E_{\mathrm{vs}}=\frac{\Delta V}{d}=\frac{(125 \mathrm{~V})}{(0.0015 \mathrm{~m})}=8.33 \times 10^{4} \mathrm{~V} / \mathrm{m}
$$

For the protons not to deflect as they pass through the velocity selector, the net force due to the magnetic and electric fields must be zero (i.e., the forces must be equal and opposite). Therefore,

$$
F_{\mathrm{E}}=F_{\mathrm{B}}, \quad q E_{\mathrm{vs}}=q v B_{\mathrm{vs}}
$$

The magnetic field in the velocity selector must then be

$$
B_{\mathrm{vs}}=\frac{E_{\mathrm{vs}}}{v}=\frac{\left(8.33 \times 10^{4} \mathrm{~V} / \mathrm{m}\right)}{\left(1.475 \times 10^{6} \mathrm{~m} / \mathrm{s}\right)}=5.65 \times 10^{-2} \mathrm{~T}
$$

The magnetic field inside the velocity selector is 0.0565 T .
EVALUATE Mass spectrometers are scientific instruments that are based on fundamental principles of electricity and magnetism. Problems involving mass spectrometers allow us to review these principles.

PRACTICE PROBLEM What potential is needed to accelerate the protons to the velocity given in the problem? Answer: 11,300 V.

## 3: Torque on a wire

A long wire is bent to form a rectangular section and is placed in a region with a uniform magnetic field, as shown in Figure 27.4. The magnetic field is directed downward and has a magnitude of 2.5 T . The wire carries a 25 -A current. Find the net torque on the wire about the $z$-axis.


Figure 27.4 Problem 3.

## Solution

IDENTIFY We will find the torque by finding the magnetic force on each segment of wire and then finding the torque due to each segment. The torques will sum to the net torque.

SET UP There are five segments of wire: two lying along the $z$-axis, two vertical segments, and the $1.0-\mathrm{m}$-long horizontal segment that crosses the $y$-axis. The two segments lying along the $z$-axis have a
net force acting on them, but no torque, since the moment arm about them is zero. The two vertical segments have no force acting on them, because the segments are parallel to the magnetic field. The remaining segment is the horizontal segment at the center. There is a net force, as well as a torque, acting on this segment. Figure 27.5 provides a view from the $z$-axis. We see that the force is directed to the right, due to the right-hand rule. The torque is along the negative $z$-axis and will tend to rotate the segment clockwise.


Figure 27.5 Problem 3.

EXECUTE We find the magnitude of the torque on the horizontal segment by first finding the magnitude of the force on that segment. The magnetic force on a segment of current-carrying wire is

$$
F=I l B \sin \phi
$$

In this case, the current is perpendicular to the magnetic field, so $\phi=90^{\circ}$ and $\sin \phi=1$. The distance between the $z$-axis and the force is 0.25 m , and the force is perpendicular to the moment arm, so the torque is

$$
\tau=r F=r l l B=(0.25 \mathrm{~m})(25 \mathrm{~A})(1.0 \mathrm{~m})(2.5 \mathrm{~T})=15.6 \mathrm{~N} \cdot \mathrm{~m}
$$

The net torque on the wire is 15.6 Nm , directed along the negative $z$-axis.
EVALUATE This problem illustrates how to find the torque on a current-carrying wire. We had to use first principles to find the torque on each segment. Problems involving the torque on a current-carrying loop can be solved with the use of the formula included in the "Concepts and Equations" section in this chapter.

## 4: A hanging loop

A long wire carries a current $I$ and produces a magnetic moment $\mu$, as in Figure 27.6. The coil is pivoted in a frictionless manner about the $z$-axis. The loop has mass $m$, and its center of gravity is located a distance $L$ from the $z$-axis. A constant magnetic field $B$ is directed along the negative $x$-axis. If $\theta$ is $37^{\circ}$, what value of $B$ is needed for equilibrium? Would this value of $B$ give equilibrium for any $\theta$ ?


Figure 27.6 Problem 4.

## Solution

IDENTIFY For the loop to be in equilibrium, the net torque must be zero. We'll use the equilibrium condition to find the magnetic field, the target variable.

SET UP Gravity and the magnetic force create torques on the loop. We will set these torques equal to each other to find the magnetic field.

EXECUTE The magnetic field makes an angle $\theta$ with the magnetic moment of the loop. The torque due to the magnetic field is then

$$
\tau_{B}=B M \sin \theta
$$

This torque tends to rotate the loop counterclockwise.
The torque due to gravity is given by

$$
\tau_{g}=m g L \sin \theta
$$

The torque due to gravity tends to rotate the loop clockwise. The net torque is zero, given by

$$
\sum \tau=0=B M \sin \theta-m g L \sin \theta
$$

Solving for $B$, we see that the loop will remain in equilibrium if

$$
B=\frac{m g L}{M} .
$$

Note that the angle cancels and the result does not depend on the angle $\theta$. The loop would remain in equilibrium at any angle.

EVALUATE This problem is similar to earlier torque equilibrium problems, but with the addition of the magnetic force.

## Try It Yourself!

## 1: Force on a charged particle

A particle with charge $e$, mass $2.0 \times 10^{-27} \mathrm{~kg}$, and velocity given by $\vec{v}=\left(2.0 \times 10^{5} \mathrm{~m} / \mathrm{s}\right) \times$ $(\hat{\imath}+2 \hat{\jmath}+2 \hat{k})$ enters a region with a constant uniform magnetic field $\vec{B}=(2.5 \mathrm{~T}) \hat{\jmath}$. (a) Calculate the force on this particle due to the magnetic field. (b) Describe the path of the particle.

## Solution Checkpoints

IDENTIFY AND SET UP Use the definition of the vector product to find the force.
EXECUTE (a) After applying the vector product, we obtain the force from

$$
\vec{F}=e\left(2.0 \times 10^{5} \mathrm{~m} / \mathrm{s}\right)(2.5 \mathrm{~T})(\hat{k}-2 \hat{\imath}) .
$$

The magnitude of the force is found by summing the squares of the components of the force and taking the square root. The force has magnitude of $1.79 \times 10^{-13} \mathrm{~N}$.
(b) There is no force component along the $y$ direction, so the $y$ component of the particle's motion is constant. In the $x z$ plane, the orbit is circular. Combining the two motions, we can describe the path of the particle as a helix.

EVALUATE What is the radius of the helix?

## 2: Investigating a current loop

For the square loop shown in Figure 27.7, (a) calculate the magnetic flux through the loop and (b) calculate the torque about an axis parallel to the $x$-axis and through the center of the coil.


Figure 27.7 Try It Yourself Problem 2.

## Solution Checkpoints

IDENTIFY AND SET UP Use the definitions of flux and torque to solve.
EXECUTE (a) The normal is rotated by $30^{\circ}$, so the flux is

$$
\Phi_{B}=B L^{2} \cos 30^{\circ}
$$

(b) The loop creates a magnetic moment. The magnitude of the torque is then

$$
\tau=I A B \sin \theta=I L^{2} B \sin 30^{\circ}
$$

EVALUATE Do you get the same value for the torque if you calculate the torques about the axis for the segments $a b$ and $c d$ ?

## 3: Magnetic force on power lines

Let's investigate the magnetic force on power transmission lines due to the earth's magnetic field. Suppose the magnetic field due to the earth is represented by

$$
\vec{B}=B\left(\cos 70^{\circ}(-\hat{\imath})+\sin 70^{\circ}(-\hat{k})\right),
$$

where $B=2.5 \times 10^{-5} \mathrm{~T}$. A straight wire 1.0 m long carrying 500 A of current runs parallel to the ground in the $x y$ plane. (See Figure 27.8.) (a) Find the force on the wire when the current is in the $-x$ direction. (b) Find the force on the wire when the current is in the $+y$ direction.


Figure 27.8 Try It Yourself Problem 3.

## Solution Checkpoints

IDENTIFY AND SET UP Use the definition of magnetic force to solve.
EXECUTE (a) When the current in the $-x$ direction,

$$
\vec{L}=L(-\hat{\imath}) .
$$

Taking the vector product gives a force of magnitude $1.18 \times 10^{-2} \mathrm{~N}$ in the $-y$ direction.
(b) When the current is in the $+y$ direction, the force becomes

$$
\vec{F}=I L B\left(-\sin 70^{\circ} \hat{\imath}+\cos 70^{\circ} \hat{k}\right)
$$

This equation gives a force that makes an angle of $20^{\circ}$ with the $-x$-axis in the $x z$ plane and that has a magnitude of $2.5 \times 10^{-2} \mathrm{~N}$.

EVALUATE If this represents the force on a small section of power transmission line, how much force acts on a 1-km segment of line? Do the power lines oscillate, since the transmission lines carry alternating current?

## 28 <br> Sources of Magnetic Fields

## Summary

We will investigate the sources of magnetic fields in this chapter. We will see how moving charges and currents generate magnetic fields and will use the law of Biot and Savart to find the magnetic field for a variety of current distributions. We will then look at the force between conductors by examining the magnetic force on one current due to the field generated by a second current. We will learn about Ampere's law and understand how it can be applied to symmetric current distributions to find the magnetic field. Finally, we will reexamine permanent magnets to learn how magnetic fields are generated by those devices.

## Objectives

After studying this chapter, you will understand

- How magnetic fields are generated by moving charges and currents.
- How to apply the law of Biot and Savart to a variety of current distributions.
- How to calculate the force between two current-carrying conductors.
- How to apply Ampere's law to symmetric current distributions.
- How the magnetic field surrounding permanent magnets originates.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Magnetic Field of a Current- | The law of Biot and Savart gives the magnetic field contribution due to an <br> element of current-carrying conductor: |

$$
d \vec{B}=\frac{\mu_{0}}{4 \pi} \frac{I d \vec{l} \times \hat{r}}{r^{2}}
$$

Here, $\mu_{0}$ is the permeability of vacuum and has a value of

$$
\mu_{0}=4 \pi \times 10^{-7} \mathrm{~T} \cdot \mathrm{~m} / \mathrm{A} .
$$

The total field is found by integrating over the length of the conductor.
The strength of the magnetic field generated by a long, straight currentcarrying conductor is given by

$$
B=\frac{\mu_{0} I}{2 \pi r},
$$

where $r$ is the distance from the center axis of the conductor.

| Magnetic Force between Current- | Two current carrying conductors can exert magnetic forces on each other. For <br> two long, straight parallel wires, the force per unit length is |
| :--- | :--- |
| Carrying Conductors |  | Carrying Conductors two long, straight parallel wires, the force per unit length is

$$
\frac{F}{l}=\frac{\mu_{0} I I^{\prime}}{2 \pi r} .
$$

The two wires attract each other if their currents are in the same direction and repel each other if their currents are in opposite directions.

| Magnetic Field of a Current Loop | The magnetic field at the center of a coil of $N$ loops of wire with radius $R$ is <br> given by |
| :--- | :--- |


|  | $B=\frac{\mu_{0} N I}{2 R}$ |
| :--- | :--- |
| Ampère's Law | $\oint \vec{B} \cdot d \vec{l}=\mu_{0} I_{\text {encl }}$. |
| Solenoids | Ampère's law states that the line integral of the magnetic field around any <br> closed path is proportional to the current though the area enclosed by the <br> path; algebraically, <br> Ampère's law can be used to find the magnetic field for symmetric current <br> distributions. |
| The magnetic field inside a solenoid (a collection of many closely spaced <br> windings) is |  |
|  | where $n$ is the number of turns per unit length. The field outside the solenoid <br> is approximately zero. |

## Conceptual Questions

## 1: Magnetic field due to two currents

Two insulated wires perpendicular to each other in the same plane carry equal currents as shown in Figure 28.1. Is there a region where the magnetic field is zero? If so, where is this region? If not, explain why the field is never zero.


Figure 28.1 Question 1.

## Solution

IDENTIFY, SET UP, AND EXECUTE A current in a long, straight wire creates a field everywhere around the current. To find a region with zero magnetic field requires that the fields due to each wire be equal and opposite, canceling the magnetic field. In the plane of the page, the magnetic field due to the currents points either out of or into the page. The field due to the horizontal current points into the page above the wire and out of the page below the wire. The field due to the vertical current points into the page to the right of the wire and out of the page to the left of the wire.

In the upper left and lower right quadrants, the fields due to the two currents are in opposite directions. (The fields from both currents in the lower left and upper right quadrants are in the same direction.) It is possible for the fields to cancel in the upper left and lower right quadrants.

For the fields to cancel, the magnitudes must be equal. The magnetic field due to a long, straight current depends inversely on the distance from the current. Therefore, the magnetic field will cancel where the distance to both currents is the same. A diagonal line from lower right to upper left will be equidistant from both currents. The magnetic field along this diagonal line is zero.

EVALUATE This problem shows that magnetic fields are vector quantities and can be added to yield a resulting magnetic field. It also shows that the field due to a long, straight wire varies inversely with distance.

## 2: Using Ampère's law

Can you use Ampère's law to find the magnetic field at the center of a ring of radius $R$ and carrying current $I$ ?

## Solution

IDENTIFY, SET UP, AND EXECUTE Ampère's law is useful when you can evaluate the line integral of the magnetic field around a loop. It can be also be useful when the magnetic field is zero and when the magnetic field is perpendicular to the integration path. The field around a circular loop of current varies as a function of position and distance from the loop. There is no path along which Ampère's law can be applied to determine the magnetic field at the center of the ring.

EVALUATE Ampère's law is useful only in certain symmetric situations in which the integral can be evaluated at constant magnetic field. It can be applied to any problem, but the results may not be usefuland the integral may prove to be exceedingly challenging to evaluate.

## Problems

## 1: Hanging parallel wires

Two long, straight wires are suspended from light threads so that the wires are parallel and in equilibrium, as shown in Figure 28.2. The threads make an angle of $25^{\circ}$ with each other. The wires carry equal currents and are separated by 15 cm . If the mass per unit length of the wires is $0.75 \mathrm{~kg} / \mathrm{m}$, find the currents in the wires. In what directions do the currents flow?


Figure 28.2 Problem 1.

## Solution

IDENTIFY The wires are in equilibrium, so we will use Newton's first law to find the currents in them.
SET UP For the wires to be separated, there must be a repulsive magnetic force. Therefore, the currents must be in opposite directions. To find the currents in the wires, we'll use a free-body diagram of the right-hand wire in Figure 28.3. This wire is in equilibrium, so the net force acting on the wire must be zero. Three forces act on the wire: tension, gravity, and a magnetic force. We'll follow the familiar procedure of setting the $x$ and $y$ force components to zero to solve the problem. An $x y$-coordinate system has been included in the free-body diagram, and the tension force has been broken down into components.


Figure 28.3 Problem 1.

EXECUTE Two components of force act along each axis. In the vertical direction, the forces include the vertical component of tension and the force of gravity, acting in opposite directions:

$$
\sum F_{y}=T \cos 12.5^{\circ}-m g=0
$$

We write the tension as

$$
T=\frac{m g}{\cos 12.5^{\circ}}
$$

In the horizontal direction, the forces include the horizontal component of tension and the magnetic force, acting in opposite directions:

$$
\sum F_{x}=F_{\mathrm{B}}-T \sin 12.5^{\circ}=0
$$

Replacing the tension gives

$$
F_{\mathrm{B}}=m g \tan 12.5^{\circ} .
$$

The magnetic force per unit length between two parallel conductors is

$$
\frac{F_{\mathrm{B}}}{l}=\frac{\mu_{0} I I^{\prime}}{2 \pi r} .
$$

In this case, the currents are the same: $I=I^{\prime}$. Recall that we are given the mass per unit length, so dividing both sides of the force equation by length will allow us to solve the problem:

$$
\frac{F_{\mathrm{B}}}{l}=\frac{\mu_{0} I^{2}}{2 \pi r}=\left(\frac{m}{l}\right) g \tan 12.5^{\circ} .
$$

Solving for the current gives

$$
I=\sqrt{\frac{\left(\frac{m}{l}\right) 2 \pi r g \tan 12.5^{\circ}}{\mu_{0}}}=\sqrt{\frac{(0.75 \mathrm{~kg} / \mathrm{m}) 2 \pi(0.15 \mathrm{~m})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \tan 12.5^{\circ}}{\left(4 \pi \times 10^{-7} \mathrm{~T} \cdot \mathrm{~m} / \mathrm{A}\right)}}=1100 \mathrm{~A} .
$$

The two parallel wires each carry 1100-A currents, flowing in opposite directions.
EVALUATE This problem is reminiscent of earlier equilibrium problems, but with the addition of the magnetic force. By this stage of the course, you should be quite familiar with forces and should even look forward to problems that include them.

## 2: Magnetic field due to two loops

A piece of wire is formed into two loops as shown in Figure 28.4. The wire carries a 25 A current. Find the magnetic field at the center of the loops.


Figure 28.4 Problem 2.

## Solution

IDENTIFY The net magnetic field is the sum of the field due to the two individual loops and the field due to the straight segments.
SET UP Examining the straight segments, we see that there are equal currents to the right and to the left, resulting in a zero net magnetic field at the center of the loops. This leaves the inner and outer loops. Each loop contributes to the field at the center, but the contributions of the two loops are in opposite directions, since the current travels in opposite directions through the loops.
EXECUTE The magnitude of the magnetic field at the center of a loop is given by

$$
B=\frac{\mu_{0} I}{2 r}
$$

The field due to the outer loop is directed into the page at the center and has magnitude

$$
B_{\text {outer }}=\frac{\mu_{0} I}{2 r_{b}}
$$

The field due to the inner loop is directed out of the page at the center and has magnitude

$$
B_{\text {inner }}=\frac{\mu_{0} I}{2 r_{a}}
$$

Let's assign positive fields as pointing out of the page. The total magnetic field at the center is then

$$
\begin{aligned}
B_{\text {total }} & =B_{\text {inner }}-B_{\text {outer }} \\
& =\frac{\mu_{0} I}{2 r_{a}}-\frac{\mu_{0} I}{2 r_{b}} \\
& =\frac{\mu_{0} I}{2}\left(\frac{1}{r_{a}}-\frac{1}{r_{b}}\right) \\
& =\frac{\left(4 \pi \times 10^{-7} \mathrm{~T} \cdot \mathrm{~m} / \mathrm{A}\right)(25 \mathrm{~A})}{2}\left(\frac{1}{(0.10 \mathrm{~m})}-\frac{1}{(0.25 \mathrm{~m})}\right) \\
& =9.43 \times 10^{-5} \mathrm{~T} .
\end{aligned}
$$

The magnetic field at the center of the loops is 0.943 G , directed out of the page.
EVALUATE This problem illustrates how we can find the magnetic field at a point by adding the magnetic field contributions due to many segments of current. Here, we see that the contributions due to the straight segments are cancelled at the center. Remember that magnetic field contributions are vectors, and you've added vectors many times throughout this course.

## 3: Magnetic field from the Biot and Savart law

A piece of wire is formed into the shape shown in Figure 28.5. The wire consists of two long, straight sections and a three-eighths semicircle. The radius of the semicircle is 15.0 cm , and the wire carries a $75-$ A current. Find the magnetic field at the center $P$ of the semicircle.


Figure 28.5 Problem 3.

## Solution

IDENTIFY The net magnetic field is the sum of the fields due to the three segments. We will use the law of Biot and Savart to find the field for each segment and add to find the net field.

SET UP We first consider the vector product for each of the segments. For the straight segments, the vector product is zero, since $d \vec{l}$ is parallel to $\hat{r}$. For the semicircular segment, the vector product has magnitude $d l$, since $\overrightarrow{d l}$ is perpendicular to $\hat{r}$. All points along the semicircle are equidistant from the center, so the integral will be straightforward to evaluate.

EXECUTE The magnitude of the magnetic field at the center of the semicircle is found from the law of Biot and Savart, given by

$$
d \vec{B}=\frac{\mu_{0}}{4 \pi} \frac{I d \vec{l} \times \hat{r}}{r^{2}}
$$

As we have determined, the magnitude of the vector product is $d l$ and the radius is constant. We find the magnitude of the field by integrating. We have

$$
\vec{B}=\int d \vec{B}=\int \frac{\mu_{0}}{4 \pi} \frac{I d \vec{l} \times \hat{r}}{r^{2}}=\frac{\mu_{0} I}{4 \pi R^{2}} \int d l
$$

where the constant terms were removed, leaving the integral of $d l$. This integral gives the length of the semicircular portion of the wire. We can integrate $r d \theta$ around the arc, or we can use our knowledge of a circle's circumference. The arc is three-eighths of a full circle, so it has a length three-eighths of the circumference of a circle with radius $R$. The magnetic field strength is then

$$
\vec{B}=\frac{\mu_{0} I}{4 \pi R^{2}} \int d l=\frac{\mu_{0} I}{4 \pi R^{2}}\left(\frac{3}{8} 2 \pi R\right)=\frac{3 \mu_{0} I}{16 R} .
$$

Substituting the given values results in

$$
\vec{B}=\frac{3 \mu_{0} I}{16 R}=\frac{3\left(4 \pi \times 10^{-7} \mathrm{~T} \cdot \mathrm{~m} / \mathrm{A}\right)(75 \mathrm{~A})}{16(0.15 \mathrm{~m})}=1.18 \times 10^{-4} \mathrm{~T}=1.18 \mathrm{G}
$$

The direction of the magnetic field is found from the right-hand rule. The magnetic field is directed outward at point $P$.

EVALUATE This problem illustrates that evaluating an equation generated by the law of Biot and Savart can be relatively easy. By splitting the problem up into multiple segments and carefully evaluating the current in each segment, we found the solution.

## 4: Evaluating a line integral around a wire

Use the results for the magnetic field of a long, straight wire to evaluate the integral $\oint \vec{B} \cdot d \vec{l}$ around the closed contour $a b c d a$ that lies in a plane perpendicular to the wire shown in Figure 28.6 on the next page.


Figure 28.6 Problem 4.

## Solution

IDENTIFY We will break the integral up into four segments, evaluate the integral for each segment, and sum the results.

SET UP The magnetic field lines for a long, straight wire form concentric rings around the wire; the magnetic field is directed tangent to the field lines. We use the right-hand rule and find that the magnetic field is directed counterclockwise when viewed from above, or from $b$ to $c$ along segment $b c$. The magnitude of the field depends on the distance from the wire and is constant at a given radius. We will keep this in mind as we evaluate each segment.

EXECUTE The magnitude of the magnetic field a distance $r$ from the wire is

$$
B=\frac{\mu_{0} I}{2 \pi r} .
$$

Along the segment $a b$, the magnetic field is perpendicular to $d \vec{l}$, so the scalar product is zero and doesn't contribute to the integral. Along segment $b c$, the magnetic field is parallel to $d \vec{l}$ at every point, so

$$
\vec{B} \cdot d \vec{l}=B d l .
$$

Evaluating the integral for segment $b c$, we obtain

$$
\int_{b c} \vec{B} \cdot d \vec{l}=\int \frac{\mu_{0} I}{2 \pi r} d l=\frac{\mu_{0} I}{2 \pi r_{2}} \int d l=\frac{\mu_{0} I}{2 \pi r_{2}} r_{2} \theta=\frac{\mu_{0} I}{2 \pi} \theta .
$$

Here, the line integral is evaluated and found to be the arc length of segment $b c$.
Along segment $c d$, the magnetic field is perpendicular to $d \vec{l}$, so the scalar product is zero and doesn't contribute to the integral. Along segment $d a$, the magnetic field is antiparallel to $d \vec{l}$ at every point, so

$$
\vec{B} \cdot d \vec{l}=-B d l
$$

Evaluating the integral for segment $d a$, we find that

$$
\int_{d a} \vec{B} \cdot d \vec{l}=-\int \frac{\mu_{0} I}{2 \pi r} d l=-\frac{\mu_{0} I}{2 \pi r_{1}} \int d l=-\frac{\mu_{0} I}{2 \pi r_{1}} r_{1} \theta=-\frac{\mu_{0} I}{2 \pi} \theta .
$$

As with segment $b c$, the line integral is evaluated and found to be the arc length of segment $d a$. We combine the results to find the integral around the complete path. This leads to

$$
\oint \vec{B} \cdot d \vec{l}=\int_{a b} \vec{B} \cdot d \vec{l}+\int_{b c} \vec{B} \cdot d \vec{l}+\int_{c d} \vec{B} \cdot d \vec{l}+\int_{d a} \vec{B} \cdot d \vec{l}=0+\frac{\mu_{0} I}{2 \pi} \theta+0-\frac{\mu_{0} I}{2 \pi} \theta=0 .
$$

The integral around the complete path is zero, consistent with Ampère's law.
EVALUATE Why is the result consistent with Ampère's law? Ampère's law states that the integral along a closed path is equal to the net current enclosed by the path. In this case, the path encloses no current; hence, the integral must be zero.

This problem also gives us practice evaluating the line integral of the magnetic field. We'll capitalize on this practice in the next problem as we apply Ampère's law.

## 5: Magnetic field for a distributed current

A long, straight conductor of radius $R$ carries a current $I$ uniformly distributed over its cross section. Find the magnetic field both inside and outside the conductor.

## Solution

IDENTIFY We will apply Ampere's law both inside and outside the conductor to solve for the magnetic field in both regions-the target variables.

SET UP To find the magnetic field inside the conductor, we place an Amperian loop inside the conductor. This loop will enclose only a portion of the total current, and we will determine that portion. We will follow the same procedure to find the field outside the conductor, but in this case all of the current will be enclosed. A sketch of the conductor with Amperian loops is shown in Figure 28.7.


Figure 28.7 Problem 5.
EXECUTE We place a circular Amperian loop inside the conductor to find the field in this region. The magnetic field is tangent to the cylinder, so the integrand is

$$
\vec{B} \cdot d \vec{l}=B d l .
$$

The magnetic field is constant at the radius $r$ and so can be taken out of the integral. The integral is then

$$
\oint \vec{B} \cdot d \vec{l}=\oint B d l=B \oint d l=B 2 \pi r
$$

To evaluate Ampère's law, we need to find the current enclosed by the loop. The current is evenly distributed, so the current enclosed is the total current, multiplied by the ratio of the enclosed area to the total cross-sectional area. Algebraically,

$$
I_{\mathrm{encl}}=I \frac{\pi r^{2}}{\pi R^{2}}
$$

Ampère's law states that

$$
\oint B d l=\mu_{0} I_{\mathrm{encl}}
$$

Evaluating Ampère's law leads to

$$
B 2 \pi r=\mu_{0} I \frac{r^{2}}{R^{2}}
$$

Solving for the magnetic field inside the conductor gives

$$
B=\frac{\mu_{0} I r}{2 \pi R^{2}}
$$

The result is the same outside of the conductor, except that the total current is enclosed. Ampere's law gives

$$
B 2 \pi r=\mu_{0} I
$$

which leads to

$$
B=\frac{\mu_{0} I}{2 \pi r}
$$

In both regions, the magnetic field is tangent to the radius.
EVALUATE We see that outside the conductor the field is the same as that for a long, straight wire carrying a current $I$, a situation reminiscent of how the electric field outside of a spherical charge distribution reduces to the field of a point charge. Inside the conductor, the field increases linearly from zero at the center to a maximum at the outer radius of the conductor.

## Try It Yourself!

## 1: Force on a loop

Find the force on each of the segments of the rectangular loop shown in Figure 28.8. The long, straight wire on the left carries a current $I_{1}$ and the loop carries a current $I_{2}$.


Figure 28.8 Try It Yourself Problem 1.

## Solution Checkpoints

IDENTIFY AND SET UP Use the magnetic force due to a current and the magnetic field for a long, straight wire to find the force on each segment of the loop. Two segments are at constant distance from the left wire, and two segments vary in distance from the left wire.

EXECUTE The force varies along segment $a b$, so integration is required to find the net force. The contribution to the force anywhere along the segment is

$$
d F=\frac{\mu_{0} I_{1} I_{2}}{2 \pi r} d r
$$

Integrating from $r_{a}$ to $r_{b}$ gives

$$
F_{a b}=\frac{\mu_{0} I_{1} I_{2}}{2 \pi} \ln \left(\frac{r_{b}}{r_{a}}\right) .
$$

The force is directed upwards, along the page.
The force is constant along segment $b c$, so we can write down the force on the segment as

$$
F_{b c}=\frac{\mu_{0} I_{1} I_{2} L}{2 \pi r_{b}}
$$

The force is directed away from the left wire.
The force varies along segment $c d$, as it did for segment $a b$. The integration is the same, except that the limits of integration are reversed, resulting in an equal, but opposite, force.

$$
F_{c d}=-F_{a b}=-\frac{\mu_{0} I_{1} I_{2}}{2 \pi} \ln \left(\frac{r_{b}}{r_{a}}\right) .
$$

This force is directed downwards, along the page.
The force is also constant along segment $d a$, so we can write down the force on that segment as

$$
F_{d a}=\frac{\mu_{0} I_{1} I_{2} L}{2 \pi r_{a}}
$$

The force is directed towards the left wire.
The two vertical forces cancel, leaving only the two horizontal forces. The net force is the sum of the two horizontal forces, directed towards the left wire and with a magnitude

$$
F_{\text {net }}=\frac{\mu_{0} I_{1} I_{2} L}{2 \pi}\left(\frac{1}{r_{b}}-\frac{1}{r_{a}}\right) .
$$

EVALUATE How can you check these results?

## 2: Magnetic field due to concentric cylinders

Two equal and opposite currents $I$ are carried in two long concentric thin cylinders of radii $R_{1}$ and $R_{2}$, as shown in Figure 28.9 on the next page. Find the magnetic field in all regions of space.


Figure 28.9 Try It Yourse lf Problem 2.

## Solution Checkpoints

IDENTIFY AND SET UP Apply Ampère's law to the three regions in the figure: inside the cylinders, between the cylinders, and outside the cylinders.

EXECUTE Inside the cylinders, no current passes through the Amperian loop, so the magnetic field is zero. Between the cylinders, the Amperian loop will enclose the current in the inner cylinder, resulting in

$$
B 2 \pi r=\mu_{0} I,
$$

which leads to

$$
B=\frac{\mu_{0} I}{2 \pi r} .
$$

Outside the cylinders, the net current enclosed by the Amperian loop is zero $(I-I=0)$, so the magnetic field is zero.

EVALUATE What is the direction of the magnetic field between the conductors? Does this problem help illustrate why many audio and video cables are made of concentric conductors?

## Electromagnetic Induction

## Summary

In this chapter, we will study the electromotive force (emf) and the current induced by magnetic interactions. We will learn to use Faraday's and Lenz's laws to analyze the behavior of time-dependent magnetic fields and magnetic fluxes. This investigation will shed light on a variety of electrical applications, including generators, motors, transformers, speakers, and microphones. The chapter will link magnetic and electric phenomena through changing magnetic fields that generate electric fields, thereby laying the foundation for the study of alternating-current circuits, covered in Chapter 31, and for the discovery that light is an electromagnetic wave (Chapter 32).

## Objectives

After studying this chapter, you will understand

- The concept of magnetic flux.
- The nature of the source of induced emfs and currents.
- How to use Faraday's and Lenz's laws to quantify induced emfs and currents.
- How to solve a variety of electromagnetic induction problems.


## Concepts and Equations

| Term | Description |
| :---: | :---: |
| Faraday's Law | A changing magnetic flux through a closed loop induces an emf in the loop. Faraday's law states that the magnitude of the induced emf in a closed loop equals the negative of the rate of change with respect to time of the magnetic flux through the loop: $\mathcal{E}=-\frac{d \Phi_{B}}{d t} .$ |
| Lenz's Law | Lenz's law states that an induced current or emf always acts to oppose the change that caused it. Lenz's law is used to find the direction of the induced effect in induction problems. |
| Motional Electromotive Force | The charges in a conductor moving in a magnetic field will be acted upon by a force and will induce a current in the conductor. A conductor of length $L$ moving with speed $v$ perpendicular to a magnetic field with magnitude $B$ will induce an emf $\mathcal{E}=v B L .$ |
| Induced Electric Fields | When an emf is induced by a changing magnetic flux, there is an induced electric field of nonelectrostatic origin given by $\oint \vec{E} \cdot d \vec{l}=-\frac{d \Phi_{B}}{d t} .$ <br> This field is not conservative and does not have an associated potential. |
| Displacement Current | A time-varying electric field generates a displacement current given by $i_{\mathrm{D}}=\mathcal{E} \frac{d \Phi_{E}}{d t} .$ <br> This displacement current generates magnetic fields in the same way that a conduction current generates fields. The displacement current is added to the conduction current in Ampere's law. |
| Maxwell's Equations | The relation between electric and magnetic fields and their sources are summarized in Maxwell's equations: $\begin{aligned} & \oint \vec{E} \cdot d \vec{A}=\frac{Q_{\mathrm{encl}}}{\mathcal{E}_{0}}, \\ & \oint \vec{B} \cdot d \vec{A}=0, \\ & \oint \vec{B} \cdot d \vec{l}=\mu_{0}\left(i_{\mathrm{C}}+\mathcal{E}_{0} \frac{d \Phi_{E}}{d t}\right), \\ & \oint \vec{E} \cdot d \vec{l}=-\frac{d \Phi_{B}}{d t} . \end{aligned}$ |

## Conceptual Questions

## 1: Induced current in a moving current loop

Figure 29.1 shows a current loop next to a long, straight wire carrying a current $I$. Find the direction of the induced current in the loop when the loop moves towards the long wire, when the loop moves away from the long wire, when the loop is stationary, and when the loop moves upwards along the wire.


Figure 29.1 Question 1.

## Solution

IDENTIFY, SET UP, AND EXECUTE We'll use Lenz's law to find the direction of the induced current. The right-hand rule tells us that the magnetic field due to the long wire is directed into the current loop.

When the loop moves towards the long wire, the magnetic flux increases, since the magnetic field increases near the wire. The induced current will oppose the change and create an induced field directed out of the page. The induced current must be counterclockwise as the loop moves towards the long wire.

When the loop moves away from the long wire, the magnetic flux decreases, since the magnetic field decreases away from the wire. The induced current will oppose the change and create an induced field directed into the page. The induced current must be clockwise as the loop moves away from the long wire.

When the loop is stationary, the magnetic flux remains constant. No current is induced in this case.
When the loop moves upwards along the long wire, the magnetic flux remains constant, since the loop remains a constant distance from the long wire. No current is induced in the loop in this case either.

EVALUATE Lenz's law can be applied to a variety of situations, resulting in a variety of solutions. Each situation may be subtly different, and you must apply the law carefully each time. Guessing can also lead to a solution, but only in about one-third of the cases will you be correct.

You may be confused about the last case, in which the loop moved along the long wire. You can imagine that as the loop moves, magnetic field lines exit and enter the loop. But since the number of field lines leaving the loop is equal to the number of lines entering the loop, the total magnetic flux remains constant.

## 2: Induced current in a stationary current loop

Figure 29.2 shows a stationary current loop next to a long, straight wire carrying a current $I$. Find the direction of the induced current in the loop when the current in the long wire is increasing with time.


Figure 29.2 Question 2.

## Solution

IDENTIFY, SET UP, AND EXECUTE We'll use Lenz's law to find the direction of the induced current. The right-hand rule tells us that the magnetic field due to the long wire is directed into the current loop.

When the current increases in the long wire, the magnetic field inside the loop increases, increasing the magnetic flux as well. There will be an induced current in the loop that opposes the change in flux. The induced current will create an induced field directed out of the loop. The induced current must be counterclockwise as the current increases in the long wire.

EVALUATE This problem presents another application of Lenz's law: that of a stationary loop and a changing magnetic field. We can also consider what will happen if the current in the long wire changes as the loop is moving.

## 3: Induced electric field in a capacitor

A capacitor's plates are connected through a loop of wire. A uniform magnetic field is directed outward through the loop, as shown in Figure 29.3. If the magnetic field decreases with time, in which direction will the electric field between the plates of the capacitor point?


Figure 29.3 Question 3.

## Solution

IDENTIFY, SET UP, AND EXECUTE The magnetic flux through the loop will decrease with time, so the loop will induce a magnetic field in the same direction as the existing field, or out of the page. The induced current in the loop will flow counterclockwise to create the induced field directed out of the page. Since the current is counterclockwise, positive charges leave the top plate of the capacitor and move to the bottom plate, leaving excess negative charges on the top plate and excess positive charges on the bottom plate. These charges will create an electric field that points from the capacitor's bottom plate to its top plate.

EVALUATE This problem combines our newly acquired knowledge of electromagnetic induction with our older knowledge of capacitors. We used our knowledge of both subjects to solve the problem.

## Problems

## 1: Horizontal rod-and-rail system

A conducting rod moves without friction to the left on a horizontal rail system, as shown in Figure 29.4. The rail system is placed inside a uniform 1.2-T magnetic field directed out of the plane of the
system. The rod has a resistance of $1.85 \Omega$, and the rest of the rail system has negligible resistance. If the ammeter reads a constant 12 A , find the velocity of the conducting rod. In which direction is the current flowing?


Figure 29.4 Problem 1.

## Solution

IDENTIFY AND SET UP As the bar moves to the left, an induced emf and current are generated in the rod-and-rail system due to the changing magnetic flux. The induced emf can be found from the motional emf formula, and the direction of the induced current can be found from Lenz's law.

EXECUTE The emf induced in the system is

$$
\mathcal{E}=v B L
$$

We are given the amount of induced current in the system. We relate that to the induced emf through the formula

$$
\mathcal{E}=I R=v B L .
$$

The only resistance in the circuit is the resistance of the bar. Rearranging terms to solve for the for velocity, we have

$$
v=\frac{I R}{B L}=\frac{(12 \mathrm{~A})(1.85 \Omega)}{(1.2 \mathrm{~T})(0.50 \mathrm{~m})}=37.0 \mathrm{~m} / \mathrm{s}
$$

The rod moves to the left at $37.0 \mathrm{~m} / \mathrm{s}$.
The magnetic flux is increasing, so the induced current will oppose the increase. The induced current will induce a magnetic field directed into the page, so the induced current must be clockwise.

EVALUATE What causes the rod to move to the left? There must be a force directed to the left, in addition to the magnetic force on the rod. This other force must be equal in magnitude to the magnetic force, since the rod moves with constant velocity. (The current remains constant.)

## 2: Vertical rail system

A conducting bar is connected between two vertical conducting rails in turn connected at the bottom through a resistor, as shown in Figure 29.5 on the next page. A horizontal, uniform magnetic field exists perpendicular to the plane of the rail system and is directed into the rail system. The conducting bar is allowed to fall, and it reaches its terminal velocity after several seconds. What is the terminal velocity of the bar? In which direction does the current flow through the rail? The bar has a mass of 0.75 kg , and you may ignore friction between the bar and the rail system.


Figure 29.5 Problem 2.

## Solution

IDENTIFY We will use the fact that the bar is in equilibrium when it is at terminal velocity, and we will use our knowledge of motional emf to find the terminal velocity, the target variable.

SET UP Recall that when an object reaches its terminal velocity, the net force acting on it is zero. Gravity and the magnetic force act on the rod, as shown in Figure 29.6. These two forces must be equal and opposite. The magnetic force is between the induced current in the rod and the uniform magnetic field. We will need to find the induced current to solve the problem.

As the bar falls, an induced emf and current are generated in the rod-and-rail system due to the changing magnetic flux. The induced emf can be found from the motional emf formula. The magnetic flux decreases as the rod falls, so the induced current tends to sustain the flux. The induced current will travel clockwise around the circuit in accordance with Lenz's law.


Figure 29.6 Problem 2.

EXECUTE The net force on the rod must be zero at terminal velocity. Newton's first law gives

$$
\sum F_{y}=0=F_{B}-m g .
$$

The magnetic force is

$$
F_{B}=I l B \sin \phi,
$$

where the current is the induced current and the angle between the current and the magnetic field is $90^{\circ}$. The induced emf will lead to the induced current. The induced emf is related to the velocity of the falling rod by the formula

$$
\mathcal{E}=v B l .
$$

The induced current is

$$
I=\frac{\mathcal{E}}{R}=\frac{v B l}{R}
$$

Combining these results gives

$$
F_{B}=I l B=\frac{v B l}{R} l B=m g .
$$

Solving for the velocity yields

$$
v=\frac{m g R}{l^{2} B^{2}}=\frac{(0.75 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(1.4 \Omega)}{(0.62 \mathrm{~m})^{2}(0.86 \mathrm{~T})^{2}}=36.2 \mathrm{~m} / \mathrm{s}
$$

The terminal velocity of the bar is $36.2 \mathrm{~m} / \mathrm{s}$, and the induced current in the bar flows from left to right.
EVALUATE This problem combines terminal velocity, net force, magnetic force, and electromagnetic induction. It may appear intimidating, but each step has been seen in previous problems.

## 3: Emf induced by a time-varying current

A power line carrying a current that varies with time according to the formula $I(t)=I_{0} \cos (2 \pi f t)$ is located near a rectangular coil, as shown in Figure 29.7. If $I_{0}=250 \mathrm{~A}$ and $f=60 \mathrm{~Hz}$, find the emf induced in the coil.


Figure 29.7 Problem 3.

## Solution

IDENTIFY We will apply Faraday's law to find the induced emf. We will have to integrate the flux in the coil, since the magnetic field varies across the coil.

SET UP To find the magnetic flux, we break the area of the coil into thin, vertical strips and integrate across the coil, as illustrated in Figure 29.8. The magnetic field due to the power line is constant in the thin strips. Once we have determined the magnetic flux, we take the derivative to find the induced emf.


Figure 29.8 Problem 3.
EXECUTE The flux through the shaded strip in Figure 29.8 is given by

$$
d \Phi=B l d x=\frac{\mu_{0} I(t)}{2 \pi x} l d x .
$$

Note that the area vector and the magnetic field are parallel, so the cosine term is unity. We find the total flux $\Phi$ for the coil by integrating across the coil:

$$
\begin{aligned}
\Phi & =\int d \Phi \\
& =\int_{x_{1}}^{x_{2}}\left(\frac{\mu_{0} I(t)}{2 \pi x} l d x\right) \\
& =\frac{\mu_{0} I(t) l}{2 \pi} \ln \frac{x_{2}}{x_{1}} \\
& =\frac{\mu_{0} I(t) l}{2 \pi} \ln \frac{(10 \mathrm{~cm})}{(5 \mathrm{~cm})} \\
& =\frac{\mu_{0} I(t) l}{2 \pi} \ln 2 .
\end{aligned}
$$

Faraday's law is used to determine the induced emf. Integrating, we have

$$
\begin{aligned}
\mathcal{E} & =\frac{d \Phi}{d t} \\
& =\frac{d}{d t}\left(\frac{\mu_{0} I(t) l}{2 \pi} \ln 2\right) \\
& =\left(\frac{\mu_{0} l}{2 \pi} \ln 2\right) \frac{d}{d t}(I(t)) \\
& =\left(\frac{\mu_{0} l}{2 \pi} \ln 2\right)(2 \pi f) I_{0} \sin (2 \pi f t) \\
& =\left(\mu_{0} l f I_{0} \ln 2\right) \sin (2 \pi f t)
\end{aligned}
$$

Substituting the given values results in

$$
\mathcal{E}=\left(1.31 \times 10^{-3} \mathrm{~V}\right) \sin ((377 \mathrm{~Hz}) t)
$$

EVALUATE Do you think that this problem illustrates a good way to "borrow" power from a power line? We see that the coil picks up on a very small emf. You could replace the coil with a coil having many loops, but you would need about $10^{5}$ loops to get to 120 V . Adding loops would add to the resistance, making this approach unfeasible - not to mention that you would have to place the coil near a power line, which is not a wise idea!

## 4: A dropping coil

A rectangular coil is dropped from an initial height $h$ into a region of constant magnetic field, as shown in Figure 29.9. The coil has a length of 12.0 cm , a width of 6.0 cm , a mass per unit length of $0.010 \mathrm{~kg} / \mathrm{m}$, and a resistance of $0.1 \Omega$. The magnetic field is 0.5 T . What should the initial height be in order for the coil to enter the magnetic field at constant velocity?


Figure 29.9 Problem 4.

## Solution

IDENTIFY For the coil to pass into the field at constant velocity, the net force must be zero. We will use this fact, along with our knowledge of motional emf and gravity, to find the initial height.

SET UP As the loop falls into the magnetic field, the flux through the coil changes, because the area of the coil inside the field changes. We will use the change in flux to find the induced emf and the current in the coil. There will be a magnetic force acting upon the induced current due to the magnetic field. The only other force acting on the coil is gravity. We will need to find the mass of the coil to find the gravitational force.

After we set the forces equal to each other, we will solve for the velocity required for equilibrium. Once we have the velocity, we will use energy conservation to solve for the initial height.

EXECUTE When a length $y$ of the coil is inside the magnetic field, the flux through the coil is given by

$$
\Phi=B l y
$$

The time rate of change of the length $y$ is the velocity of the coil, $d y / d t$. The induced emf is then

$$
\mathcal{E}=\frac{d \Phi}{d t}=B l \frac{d y}{d t}=B l v
$$

This emf in turn induces a current of magnitude

$$
I=\frac{\mathcal{E}}{R}=\frac{B l v}{R}
$$

The force on this current due to the magnetic field is

$$
F_{B}=I l B=\frac{(B l)^{2} v}{R}
$$

For the velocity to remain constant, there must be no net force acting on the coil: The weight must be equal to the magnetic force. To find the weight, we first find the mass. The mass is the total length of the coil, multiplied by the linear density:

$$
\begin{aligned}
m & =(2 l+2 w)(0.010 \mathrm{~kg} / \mathrm{m}) \\
& =(2(0.12 \mathrm{~m})+2(0.06 \mathrm{~m}))(0.010 \mathrm{~kg} / \mathrm{m}) \\
& =3.6 \times 10^{-3} \mathrm{~kg}
\end{aligned}
$$

Setting the weight equal to the magnetic force yields

$$
m g=\frac{(B l)^{2} v}{R}
$$

As the coil falls, gravitational potential energy is converted into kinetic energy. Energy conservation gives

$$
\begin{aligned}
m g h & =\frac{1}{2} m v^{2} \\
v & =\sqrt{2 g h}
\end{aligned}
$$

Combining these results yields

$$
m g=\frac{(B l)^{2} \sqrt{2 g h}}{R}
$$

Finally solving for $h$ produces

$$
h=\frac{R^{2} m^{2} g}{2(B l)^{4}}=\frac{(0.10 \Omega)^{2}\left(3.6 \times 10^{-3} \mathrm{~kg}\right)^{2}\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{2((0.5 \mathrm{~T})(0.12 \mathrm{~m}))^{4}}=4.9 \mathrm{~cm}
$$

EVALUATE This problem combines techniques, forces, energy conservation, circuits, magnetic force, magnetic fields, and induced emf, thus serving as a review of much of your physics course up to now.

What happens to the coil after the top edge enters the magnetic field? The flux no longer changes, so there is no induced current or magnetic force. The coil will accelerate at $g$ after fully entering the magnetic field.

## Try It Yourself!

## 1 : Rod on rail

A conducting rod of length 0.25 m , shown in Figure 29.10, moves to the right with a velocity of $3.0 \mathrm{~m} / \mathrm{s}$ in a constant magnetic field of magnitude 0.8 T . What external force must be applied to the rod for it to maintain constant velocity? Assume that the circuit has a total resistance of $0.2 \Omega$.


Figure 29.10 Try It Yourself Problem 1.

## Solution Checkpoints

IDENTIFY AND SET UP What must the net force be on the rod for it to maintain constant velocity? What forces act on the rod? What is the induced emf in the rod?

EXECUTE The moving rod induces an emf of magnitude Blv. The induced current in the loop is

$$
I=\frac{B l v}{R}
$$

This current is acted upon by a magnetic force. The force that must be applied to the rod to keep it at constant velocity must equal this magnetic force, which is given by

$$
F_{B}=I l B=\frac{B^{2} l^{2} v}{R}
$$

The magnitude of this force is 0.6 N . The force must be applied to the right.
EVALUATE Why must the force be applied to the right?

## 2: Induced emf due to a solenoid

A square coil with sides of length 0.25 m is placed around a solenoid of diameter 0.10 m and length 0.2 m wrapped with 1000 turns of wire. The coil has its normal parallel to the axis of the solenoid. The current through the solenoid when it energizes is given by

$$
I(t)=I_{0}\left(1-e^{-t / \tau}\right)
$$

where $I_{0}$ is 100.0 A and $\tau$ is 5.0 s . Calculate the induced emf in the square coil.

## Solution Checkpoints

IDENTIFY AND SET UP Find the flux through the square coil and the change in flux per unit time to determine the emf. Does a magnetic field exist in the entire space within the square coil?

EXECUTE There is a magnetic field only inside the solenoid, so the flux is given by

$$
\Phi=\frac{\pi d_{\text {solenoid }}^{2}}{4} B=\frac{\pi d_{\text {solenoid }}^{2}}{4} \mu_{0} n I(t)
$$

The induced emf is the derivative of the flux:

$$
\begin{aligned}
\mathcal{E} & =\frac{d \Phi}{d t} \\
& =\frac{d}{d t}\left(\frac{\pi d_{\text {solenoid }}^{2}}{4} \mu_{0} n I(t)\right) \\
& =\frac{\pi d_{\text {solenoid }}^{2}}{4} \mu_{0} n \frac{d}{d t}(I(t)) \\
& =\frac{\pi d_{\text {solenoid }}^{2}}{4} \mu_{0} n\left(\frac{I_{0}}{\tau} e^{-t / \tau}\right) .
\end{aligned}
$$

The peak value of the emf is $9.87 \times 10^{-4} \mathrm{~V}$.
EVALUATE Why was the area of the solenoid used to find the magnetic flux? Is the magnetic field outside a solenoid approximately zero?

## 30

## Inductance

## Summary

In this chapter, we will investigate the manifestations of induced emf and induced current in electric circuits. We will learn about the inductor, an electronic component that stores energy, much as a capacitor stores energy. The magnetic energy stored in an inductor leads to the study of magnetic energy density-the energy stored in a magnetic field. We will then look at circuits made out of combinations of inductors, resistors, and capacitors. We lay the foundation for the study of alternating-current circuits, covered in the next chapter.

## Objectives

After studying this chapter, you will understand

- The definition of mutual inductance, self-inductance, and inductors.
- How to find the energy stored in an inductor and in the magnetic field.
- How to analyze time-dependent RL, LC, and LRC circuits.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Mutual Inductance | A changing magnetic flux in one circuit linked to a second circuit can induce <br> an emf in the second circuit and vice versa. The induced emfs in each circuit <br> are related by |

$$
\mathcal{E}_{2}=-M \frac{d i_{1}}{d t}, \quad \mathcal{E}_{1}=-M \frac{d i_{2}}{d t},
$$

where $M$ is a constant called the mutual inductance. The SI unit of inductance is the henry $(\mathrm{H}): 1 \mathrm{H}=1 \mathrm{~Wb} / \mathrm{A}$.

| Self-Inductance | A changing current in any circuit causes a self-induced emf in the same cir- |
| :--- | :--- | cuit, given by

$$
\mathcal{E}=-L \frac{d i}{d t},
$$

where $L$ is the self-inductance or, simply, the inductance. The inductance of a coil with $N$ turns and average flux $\Phi$ through each turn caused by current $i$ is

$$
L=\frac{N \Phi_{B}}{i} .
$$

An inductor is a circuit device usually consisting of a coil of wire and having substantial inductance.

| Magnetic Field Energy | The magnetic field energy $U$ in an inductor with inductance $L$ carrying a cur- |
| :--- | :--- | rent $I$ is

$$
U=\frac{1}{2} L I^{2} .
$$

This energy is stored in the electric field with an energy density of

$$
u=\frac{B^{2}}{2 \mu_{0}} .
$$

## RL Circuits

The current through a charging circuit containing an inductor, a resistor, and an emf source connected in series increases and decreases exponentially. The time constant for this circuit is $\tau=L / R$, the time required for the current to approach the fraction $1 / e$ of its final value.
LC Circuits
A circuit containing a capacitor and an inductor undergoes electrical oscillations with angular frequency

$$
\omega=\sqrt{\frac{1}{L C}} .
$$

This is the circuit analog to mechanical harmonic motion.

## LRC Series Circuits

A circuit containing capacitance, inductance, and resistance undergoes damped oscillations for small resistances. The frequency of the damped oscillations is given by

$$
\omega^{\prime}=\sqrt{\frac{1}{L C}-\frac{R^{2}}{4 L^{2}}} .
$$

Critical damping occurs when

$$
R^{2}=\frac{4 L}{C} .
$$

## Conceptual Questions

## 1: Short- and long-term behavior of an RL circuit

An inductor (with inductance $L$ ) and a resistor (with resistance $R$ ) are connected in series with an emf source (with voltage $V$ ). Just after the connection is made, what is the potential across the resistor and inductor, and what is the current in the circuit? A very long time after the connection is made, what is the potential across the resistor and inductor, and what is the current in the circuit?

## Solution

IDENTIFY, SET UP, AND EXECUTE This problem asks us to consider how current and potential difference vary in an RL circuit. We'll work through each part, using the techniques we've learned in this chapter.

Just after the connection is made, there is essentially no current in the circuit. Therefore, the potential difference across the resistor is zero $(V=I R)$. But then the potential across the inductor must be equal to the voltage across the emf $(V)$. The rate of change of the current in the circuit is therefore maximal.

A long time after the connection is made, the current has reached its maximum and is no longer changing. The potential difference across the inductor must then be zero, since the current is constant. The potential difference across the resistor is therefore equal to the voltage of the emf $(V)$, because the voltage across the inductor is zero.

EVALUATE We see that the circuit behaves as if there is only an inductor just after the connection is made and as if there is only a resistor after a long time. This analysis is similar to the analysis we carried out in Question 4 of Chapter 19, which you may want to review at this point. Keeping long-term/ short-term behavior in mind will help your interpretation of RL circuits, just as it did with RC circuits.

## Problems

## 1: Mutual inductance of two coils

Consider a short coil (length $L_{1}$ ) of radius $R_{1}$ having $N_{1}$ turns inside a second long (length $L_{2} \gg L_{1}$ ) solenoid of radius $R_{2}$ having $N_{2}$ turns. Calculate the mutual inductance of the pair (a) if the axes of the two coils are parallel and (b) if the axes of the two coils are perpendicular.

## Solution

IDENTIFY Our target variable is the mutual inductance for the two orientations.
SET UP We will find the mutual inductance by finding the flux through the short coil due to the long solenoid. We will assign a current to the solenoid that will cancel when we calculate the mutual inductance.

EXECUTE (a) When the two coils have parallel axes, we start by finding the field due to the solenoid. Given the current $I$, the solenoid creates a magnetic field of magnitude

$$
B_{2}=\mu_{0} \frac{N_{2}}{L_{2}} I .
$$

The flux through one turn of the inner coil is the area of the inner coil times the field, since the field and the area vector are parallel:

$$
\Phi_{1}=\mu_{0} \frac{N_{2}}{L_{2}} I\left(\pi R_{1}^{2}\right)
$$

The inner coil has $N_{1}$ turns, so the total flux is

$$
N_{1} \Phi_{1}=N_{1} \mu_{0} \frac{N_{2}}{L_{2}} I\left(\pi R_{1}^{2}\right)
$$

The total flux is the mutual inductance times the current. The mutual inductance is therefore given by

$$
M_{21}=\frac{N_{1} \Phi_{1}}{I}=\mu_{0} \frac{N_{1} N_{2}}{L_{2}}\left(\pi R_{1}^{2}\right)
$$

(b) When the axes of the two coils are perpendicular, the area vector and the magnetic field are perpendicular, giving zero mutual inductance.

EVALUATE The mutual inductance of the coils is always proportional to the product of their number of turns. The mutual inductance also depends only on the geometry of the coils and not on the current.

## 2: Potential and energy in a toroidal solenoid

At a certain time, an inductor has a 500.0-A current passing through it that is decreasing at a rate of $320 \mathrm{~A} / \mathrm{s}$. The inductor is in the shape of a toroidal solenoid with a cross-sectional area of $1.44 \mathrm{~cm}^{2}$ and a mean radius of 8.00 cm . The inductor has 300 turns. Find the potential drop across the inductor and the energy stored in the inductor at the given time.

## Solution

IDENTIFY AND SET UP The potential across an inductor is the inductance times the rate of change of the current. We will need to find the inductance of the toroid. The energy stored in the inductor involves the inductance and the current.

EXECUTE The inductance of a toroidal solenoid is discussed in Example 30.3 in the textbook and is found to be

$$
L=\frac{\mu_{0} N^{2} A}{2 \pi r},
$$

where $A$ is the cross-sectional area, $N$ is the number of turns, and $r$ is the mean radius of the solenoid. The potential drop across the inductor is

$$
\mathcal{E}=L \frac{d i}{d t}
$$

Combining the results to find the potential drop gives

$$
\mathcal{E}=L \frac{d i}{d t}=\frac{\mu_{0} N^{2} A}{2 \pi r} \frac{d i}{d t}=\frac{\left(4 \pi \times 10^{-7} \mathrm{~T} \cdot \mathrm{~m} / \mathrm{A}\right)(300)^{2}\left(1.44 \times 10^{-4} \mathrm{~m}^{2}\right)}{2 \pi(0.0800 \mathrm{~m})}(320 \mathrm{~A} / \mathrm{s})=0.0104 \mathrm{~V}
$$

The energy stored in the inductor is

$$
U=\frac{1}{2} L I^{2}=\frac{1}{2} \frac{\mu_{0} N^{2} A}{2 \pi r} I^{2}=\frac{1}{2} \frac{\left(4 \pi \times 10^{-7} \mathrm{~T} \cdot \mathrm{~m} / \mathrm{A}\right)(300)^{2}\left(1.44 \times 10^{-4} \mathrm{~m}^{2}\right)}{2 \pi(0.0800 \mathrm{~m})}(500 \mathrm{~A})^{2}=4.05 \mathrm{~J}
$$

The potential drop across the inductor is 10.4 mV and the energy stored in the inductor is 4.05 J .
EVALUATE This problem introduces us to analyzing inductors in RL and LC circuits.

## 3: Exploring an RL Circuit

In the circuit shown in Figure 30.1, switch $S_{1}$ is closed until a constant current is established. Switch $S_{2}$ is then closed and switch $S_{1}$ is opened. How long after the switches are closed do the current and energy stored in the inductor reach half their initial values?


Figure 30.1 Problem 3.

## Solution

IDENTIFY We will use energy and exponential decay relations to solve the problem.
SET UP After the switches are thrown, the current in the circuit will decrease. We will need to find the initial steady-state current and then use the exponential decay of the RL circuit to find the time at which the current is half its initial value. We use the same procedure to find the time at which the energy drops to half its initial value.

EXECUTE The current in a decaying RL circuit is

$$
i=I_{0} e^{-(R / L) t}
$$

The initial current can be found from

$$
I_{0}=\frac{\mathcal{E}}{R}
$$

In this case, we seek the time at which the current becomes half the initial current, or

$$
i=\frac{1}{2} I_{0}=I_{0} e^{-(R / L) t}
$$

so that

$$
\frac{1}{2}=e^{-(R / L) t}
$$

To solve for the time, we take the natural logarithm of both sides of the last equation:

$$
\ln \frac{1}{2}=\ln \left(e^{-(R / L) t}\right)=-\frac{R}{L} t
$$

Solving for the time gives

$$
t=-\ln \frac{1}{2} \frac{L}{R}=-(-0.6931) \frac{(0.500 \mathrm{H})}{(30.0 \Omega)}=0.0116 \mathrm{~s}
$$

The current reaches half its initial value 11.6 ms after the switches are thrown.
The energy stored in the inductor is

$$
U=\frac{1}{2} L I^{2} .
$$

For the energy to decrease by one-half, the current must decrease by $\sqrt{\frac{1}{2}}$, since the energy depends on the square of the current. Following the same procedure we just applied, we find the time corresponding to this current:

$$
\begin{aligned}
& \sqrt{\frac{1}{2}} I_{0}=I_{0} e^{-(R / L) t}, \\
& t=-\ln \sqrt{\frac{1}{2}} \frac{L}{R}=-(-0.3466) \frac{(0.500 \mathrm{H})}{(30.0 \Omega)}=0.00578 \mathrm{~s} .
\end{aligned}
$$

The energy stored in the inductor reaches half its initial value 5.78 ms after the switches are thrown.
EVALUATE We see that the energy stored in the inductor drops twice as fast as the current decreases. If you are familiar with the properties of logarithms, you could have skipped the algebra in the second half of the solution and simply divided the value of the current by 2 .

## 4: Self-inductance of concentric cylinders

Calculate the self-inductance per unit length of a pair of concentric thin cylinders of radii $R_{1}$ and $R_{2}$, where $R_{2}>R_{1}$. Figure 30.2 illustrates the cylinders.


Figure 30.2 Problem 4.

## Solution

IDENTIFY The target variable is the self-inductance per unit length of the system.
SET UP We can find the self-inductance either by finding the flux through the space between the cylinders and using the definition of inductance or by finding the energy stored in the volume between the cylinders and using the energy relation for inductors. We will follow the second option. We will use
the magnetic field to find the energy density. Since the energy density varies with the radius of the cylinder (as does the field), we will have to integrate the energy density through the region in between the cylinders. A slice through the cylinders shown in Figure 30.3 illustrates our integration procedure.


Figure 30.3 Problem 4.
EXECUTE We find the magnetic field between the cylinders by a quick application of Ampère's law or by reviewing Chapter 28, Try It Yourself Problem 2. The field between the cylinders is

$$
B=\frac{\mu_{0} I}{2 \pi r}
$$

The energy per unit volume stored in the magnetic field is given by

$$
u=\frac{1}{2 \mu_{0}} B^{2}=\frac{U}{V}
$$

Since the magnetic field is not constant between the cylinders, we must integrate between them. The energy stored in the magnetic field in the vertical slice $d U$, shown in Figure 30.3, is

$$
d U=\frac{1}{2 \mu_{0}} B^{2} d V
$$

The volume of the slice is the circumference $(2 \pi r)$, times the thickness $(d r)$, times the height $(Y)$ of the cylinder. Combining these results gives

$$
d U=\frac{1}{2 \mu_{0}}\left(\frac{\mu_{0} I}{2 \pi r}\right)^{2}(2 \pi r Y d r)=\frac{\mu_{0} I^{2} Y}{4 \pi} \frac{d r}{r}
$$

Integrating from the inner to the outer radius yields

$$
U=\int_{R_{1}}^{R_{2}} \frac{\mu_{0} I^{2} Y}{4 \pi} \frac{d r}{r}=\frac{\mu_{0} I^{2} Y}{4 \pi} \int_{R_{1}}^{R_{2}} \frac{d r}{r}=\frac{\mu_{0} I^{2} Y}{4 \pi} \ln \left(\frac{R_{2}}{R_{1}}\right)
$$

This expression must be equivalent to

$$
U=\frac{1}{2} L I^{2}
$$

so we find the self-inductance by comparing terms:

$$
L=\frac{\mu_{0} Y}{2 \pi} \ln \left(\frac{R_{2}}{R_{1}}\right)
$$

EVALUATE You can check this result by using the first method mentioned in "Set Up": Find the flux through the space between the cylinders, and use the definition of inductance to find the selfinductance. This procedure, however, has similar difficulty and results in a similar integration. Both procedures lead to the same result for the self-inductance.

## Try It Yourself!

## 1: Mutual inductance between a wire and a loop

Calculate the mutual inductance between a long, straight wire and a rectangular coil, as shown in Figure 30.4.


Figure 30.4 Try It Yourself Problem 1.

## Solution Checkpoints

IDENTIFY AND SET UP Find the mutual inductance by finding the flux through the coil due to the long wire. You will need to use integration. Why?

EXECUTE What is the magnetic field due to the long wire? The flux through a thin vertical strip is

$$
d \Phi_{2}=\frac{\mu_{0} I_{1}}{2 \pi r} L d r
$$

Integrating across the coil gives

$$
\Phi_{2}=\frac{\mu_{0} I_{1} L}{2 \pi} \ln \left(\frac{r_{2}}{r_{1}}\right)
$$

From the flux, the mutual inductance is found to be $2.77 \times 10^{-8 \mathrm{H}}$.
EVALUATE What is the mutual inductance when the loop is turned such that the normal to its area is directed radially outward?

## 2: Equivalent inductances

Calculate the equivalent inductance required to replace two inductors $L_{1}$ and $L_{2}$ when they are (a) connected in series and (b) connected in parallel.

## Solution Checkpoints

IDENTIFY AND SET UP Equivalent inductance is like equivalent resistance and equivalent capacitance. Combine the inductors and find the proportionality between the potential across the combination to the change in current through the combination.

EXECUTE (a) In series, the potentials across the two inductors add to the total potential across the combination. The current through both is the same. The potential is

$$
V=-L_{1} \frac{d i}{d t}-L_{2} \frac{d i}{d t}=-\left(L_{1}+L_{2}\right) \frac{d i}{d t}
$$

The equivalent inductance is $L_{1}+L_{2}$.
(b) In parallel, the potentials across the two inductors are the same and the currents through each vary. The potential in terms of $L_{1}$ is

$$
V=-L_{1} \frac{d i_{1}}{d t}
$$

and similarly for $L_{2}$. The total current is split between the two inductors, so we write

$$
V=-L_{1} \frac{d i_{1}}{d t}=-L_{2} \frac{d\left(i-i_{1}\right)}{d t}
$$

Collecting the terms involving $i_{1}$ and $i$ on either side, rearranging terms, and applying some algebra yields

$$
V=-L_{1} \frac{d i_{1}}{d t}=-\left(\frac{L_{1} L_{2}}{L_{1}+L_{2}}\right) \frac{d i}{d t}
$$

The factor in front of the rightmost term is the equivalent inductance. Rearranging terms gives

$$
\frac{1}{L_{\mathrm{Eq}}}=\frac{1}{L_{1}}+\frac{1}{L_{2}}
$$

EVALUATE Are you surprised that inductors combine in series and in parallel in the same way that resistors combine in series and in parallel? Does this knowledge shed more light on the behavior of inductors?

## 31 Alternating Current

## Summary

In this chapter, we will learn how resistors, capacitors, and inductors behave in circuits with sinusoidally varying voltages and currents. To work with alternating-current circuits, we will expand our analysis toolbox with phasors and reactance. We will examine circuits containing resistors, capacitors, and inductors and see how impedance brings together the resistance, capacitance, and inductance of those circuits. We'll examine the special case of resonance, in which the maximum current occurs at a particular frequency.

## Objectives

After studying this chapter, you will understand

- How to analyze circuits with time-varying currents and emfs.
- How to use phasor diagrams to analyze ac circuits.
- How to apply capacitive and inductive reactances to circuit analysis.
- How to analyze LRC circuits and find their impedance.
- Resonance and the conditions for resonance in a circuit.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Alternating Current and Phasors | A source of alternating current produces an emf that varies sinusoidally with <br> time. Sinusoidal currents and voltages can be represented with phasors- <br> vectors whose length represents a given quantity's amplitude and that rotate <br> counterclockwise about the origin with constant angular velocity $\omega$. The pha- <br> sor's projection on the horizontal axis at any instant represents the instanta- <br> neous value of the quantity. Two useful quantities in ac circuits are the <br> root-mean-square values of the voltage and current: |
| $\qquad I_{\text {rms }}=\frac{I}{\sqrt{2}}, \quad V_{\text {rms }}=\frac{V}{\sqrt{2}}$. |  |

is the impedance of the circuit. The phase angle of the voltage relative to the current is given by

$$
\phi=\arctan \frac{\omega L-1 / \omega C}{R} .
$$

## Power in ac Circuits

## Resonance

## Transformers

The average power in an ac circuit is

$$
P=\frac{1}{2} V I \cos \phi=V_{\mathrm{mms}} I_{\mathrm{rms}} \cos \phi,
$$

where $\phi$ is the phase angle of the voltage relative to the current. Power is dissipated only through resistors. The quantify $\cos \phi$ is called the power factor.
Resonance occurs when the current in a series LRC circuit reaches a maximum and the impedance reaches a minimum at an angular frequency

$$
\omega_{0}=\frac{1}{\sqrt{L C}}
$$

called the resonance frequency. The current and voltage are in phase at resonance, and the impedance is equal to the resistance.
Transformers are used to change the voltage and current in an alternating-current circuit. The ratio of the primary to the secondary voltage for an ideal
transformer with no energy loss is

$$
\frac{V_{2}}{V_{1}}=\frac{N_{2}}{N_{1}},
$$

## Conceptual Questions

## 1: Changing the power factor

A series LRC circuit has an impedance of $50 \Omega$ and a power factor of 0.700 at 60 Hz , and the source voltage leads the current. Should an inductor or a capacitor be added in series to the circuit to decrease the power factor to 0.400 ?

## Solution

IDENTIFY, SET UP, AND EXECUTE The impedance diagram for this problem is shown in Figure 31.1. Since the voltage leads the current, the impedance must be located counterclockwise of the resistance. The vertical projection of the impedance is the difference between the inductive and capacitive reactances. For the power factor to decrease, the cosine of the phase angle must decrease, rotating the impedance counterclockwise. To rotate the impedance counterclockwise, the difference between the inductive and capacitive reactances increases. To increase the difference by adding an element, you must add inductance. We conclude that we must add an inductor in series to the circuit in order for the power factor to decrease.


Figure 31.1 Question 1.
EVALUATE The key to solving this problem was the use of the impedance diagram. Once we created the diagram, we proceeded directly to the solution.

## 2: What's in the circuit?

You are given a box with electrical components connected in series inside. You cannot see inside the box, but are allowed to connect an ac emf source to the box. You observe that the voltage and current are in phase. You also determine that the resistance inside the box is $50 \Omega$. What can you conclude about the contents of the box?

## Solution

IDENTIFY, SET UP, AND EXECUTE When the voltage and current are in phase, the phase angle is zero and the impedance equals the resistance. The box could contain a single resistor. However, it could also contain a capacitor and an inductor in addition to the resistor. For there to be a capacitor and an inductor inside the box, their reactances must be equal and the circuit must be in resonance.

EVALUATE With ac circuits, our analysis of circuits becomes more challenging than with dc circuits: We must now include capacitive reactance and inductive reactance in the analysis.

## 3: Potentials in an LRC circuit

A series LRC circuit is analyzed and is found to have potential differences of 12 V across the capacitor, 21 V across the resistor, and 17 V across the inductor at a certain instant in time. Is the circuit operating above, below, or at its resonant angular frequency?

## Solution

IDENTIFY, SET UP, AND EXECUTE The analysis indicates that the inductive reactance is greater than the capacitive reactance, since the potential difference is proportional to the reactance. When the inductive reactance is greater than the capacitive reactance, the frequency is above the resonant frequency. The circuit is operating above the resonant angular frequency.

EVALUATE This problem elucidates the relationships among potential differences across circuit elements, reactances, and resonance.

## Problems

## 1: Adding an inductor to an LRC circuit

An LRC series circuit has an impedance of $45 \Omega$ and a power factor of 0.600 at 60 Hz , and the voltage lags the current. What value of inductor should be added in series to the circuit so that the power factor will be raised to unity?

## Solution

IDENTIFY The target variable is the inductance to be added to the circuit.
SET UP For a power factor of unity, the phase angle must be $0^{\circ}$. Since the voltage lags the current, the capacitive reactance must be larger than the inductive reactance, so we will need to add inductance for a phase angle of $0^{\circ}$. To guide us, Figure 31.2 shows the impedance diagram for the situation. We see that adding an inductor with the vertical projection of the impedance $\left(X_{L}-X_{C}\right)$ will cancel the reactance.


Figure 3 1.2 Problem 1.

EXECUTE The original impedance and phase angle are given by

$$
Z=\sqrt{R^{2}+\left(X_{L}-X_{C}\right)^{2}}, \quad \tan \phi=\frac{X_{L}-X_{C}}{R}
$$

We first eliminate $R$ from the equations by squaring the impedance and phase angle and substituting:

$$
\begin{aligned}
& Z^{2}=R^{2}+\left(X_{L}-X_{C}\right)^{2}, \quad R^{2}=\frac{\left(X_{L}-X_{C}\right)^{2}}{\tan ^{2} \phi} \\
& Z^{2}=\frac{\left(X_{L}-X_{C}\right)^{2}}{\tan ^{2} \phi}+\left(X_{L}-X_{C}\right)^{2}=\left(X_{L}-X_{C}\right)^{2}\left(\frac{1}{\tan ^{2} \phi}+1\right) \\
& Z^{2}=\left(X_{L}-X_{C}\right)^{2} \csc ^{2} \phi \\
& \left(X_{L}-X_{C}\right)=Z \sin \phi
\end{aligned}
$$

The impedance we need to add is simply $X_{L}-X_{C}$, so we solve for that difference. If $\cos \phi=0.600$, then $\sin \phi=0.800$, and

$$
X=X_{L}-X_{C}=Z \sin \phi=(45 \Omega)(0.800)=36.0 \Omega
$$

The inductance must be

$$
L=\frac{X}{\omega}=\frac{36.0 \Omega}{2 \pi(60 \mathrm{rad} / \mathrm{s})}=0.0955 \mathrm{H}
$$

We will need to add an inductor with 95.5 mH of inductance to the circuit.
EVALUATE Our first ac circuit problem has explored impedance, the power factor, and reactance. Using the impedance diagram guided us to the answer, as sketches and figures have guided us in most problems.

## 2: Adding a capacitor to an LRC circuit

A series LRC circuit uses 350 W of power from a $120-\mathrm{V}(\mathrm{rms}), 60-\mathrm{Hz}$ ac power source. The power factor is 0.400 and the voltage leads the current. What size of capacitor should be placed in series with the circuit to raise the power factor to unity?

## Solution

IDENTIFY The target variable is the capacitance to be added to the circuit.
SET UP For a power factor of unity, the phase angle must be $0^{\circ}$. Since the voltage leads the current, the inductive reactance must be larger than the capacitive reactance, so we will need to add capacitance for a phase angle of $0^{\circ}$. To guide us, Figure 31.3 on the next page shows the impedance diagram for the situation. We see that adding a capacitor with reactance equal to the vertical projection of the impedance $\left(X_{L}-X_{C}\right)$ will cancel the reactance.


Figure 31.3 Problem 2.
EXECUTE The power can be written

$$
P=\frac{1}{2} I V \cos \phi
$$

Substituting for current gives

$$
P=\frac{1}{2} \frac{V^{2}}{Z} \cos \phi
$$

Examining the phasor diagram, we see that

$$
\sin \phi=\frac{X}{Z}
$$

where $X$ is the capacitive reactance we need to add to the circuit. Substituting yields

$$
P=\frac{1}{2} \frac{V^{2}}{X / \sin \phi} \cos \phi=\frac{1}{2} \frac{V^{2} \sin \phi \cos \phi}{X} .
$$

Solving for $X$ gives

$$
X=\frac{V^{2} \sin \phi \cos \phi}{2 P}=\frac{(\sqrt{2} \times 120 \mathrm{~V})^{2}(0.917)(0.400)}{2(350 \mathrm{~W})}=15.1 \Omega
$$

We were given the rms voltage, so we multiplied $V$ by $\sqrt{2}$ to find the maximum voltage. The capacitance is then

$$
C=\frac{1}{\omega X}=\frac{1}{2 \pi(60 \mathrm{rad} / \mathrm{s})(15.1 \Omega)}=176 \mu \mathrm{~F}
$$

You will need to add a capacitor with $176 \mu \mathrm{~F}$ of capacitance to the circuit.
EVALUATE This ac circuit problem explored power and reactance and utilized an impedance diagram as a guide.

## 3: Finding the resonance frequency

An LRC circuit is made by placing a parallel-plate capacitor with plates of area $22.6 \mathrm{~m}^{2}$, separated by 2.0 mm and completely filled with Mylar®, in series with a $27-\Omega$ resistor and a toroidal solenoid inductor with 1000 turns of wire, a cross-sectional area of $2.75 \mathrm{~cm}^{2}$, and a mean radius of 5.4 cm . The components are connected to an emf source operating with a $50-\mathrm{V}$ amplitude at 120 Hz . What is the resonance angular frequency of this circuit?

## Solution

IDENTIFY The target variable is the resonance frequency of the circuit.
SET UP The resonance angular frequency is

$$
\omega_{0}=\frac{1}{\sqrt{L C}} .
$$

We will need to find the inductance and capacitance of the components to determine the frequency.
EXECUTE Recall that a parallel-plate capacitor with a dielectric has capacitance

$$
C=\frac{\kappa \mathcal{E}_{0} A}{d} .
$$

In this case, the dielectric strength of Mylar® is 3.1 (from Table 18.1). The capacitance is

$$
C=\frac{\kappa \mathcal{E}_{0} A}{d}=\frac{(3.1)\left(8.85 \times 10^{-12} \mathrm{~F} / \mathrm{m}\right)\left(22.6 \mathrm{~m}^{2}\right)}{(0.0002 \mathrm{~m})}=3.10 \mu \mathrm{~F}
$$

The inductance of a toroidal solenoid is given in Example 21.8 in the textbook:

$$
L=\frac{\mu_{0} N^{2} A}{2 \pi r} .
$$

For the inductor in this problem,

$$
L=\frac{\mu_{0} N^{2} A}{2 \pi r}=\frac{\left(4 \pi \times 10^{-7} \mathrm{~T} \cdot \mathrm{~m} / \mathrm{A}\right)(1000)^{2}\left(2.75 \times 10^{-4} \mathrm{~m}^{2}\right)}{2 \pi(0.054 \mathrm{~m})}=1.02 \mathrm{mH}
$$

We can now calculate the resonance angular frequency:

$$
\omega_{0}=\frac{1}{\sqrt{L C}}=\frac{1}{\sqrt{(1.02 \mathrm{mH})(3.10 \mu \mathrm{~F})}}=17,800 \mathrm{rad} / \mathrm{s}
$$

The resonance angular frequency of the circuit is $17,800 \mathrm{rad} / \mathrm{s}$, or 2.83 kHz .
EVALUATE This problem recalled our previous work with capacitors and inductors and combined it with our knowledge of resonance. We see that the resonance angular frequency does not depend on the resistance or the ac source; rather, it depends only on the inductance and the capacitance.

## 4: Transformer example

A transformer connected to a $110-\mathrm{V}$ line delivers 10 V to the secondary circuit. If the power drawn from the primary circuit is 220 W , what is the equivalent resistance of the secondary circuit?

## Solution

IDENTIFY We will use the relations among current, voltage, and number of turns to solve the problem. The target variable is the equivalent resistance of the secondary circuit.

SET UP We start by finding the current in the secondary circuit. We will then find the equivalent resistance.

EXECUTE Assuming that there are no power losses in the transformer, we know that the power in the primary and secondary are the same:

$$
V_{1} I_{1}=V_{2} I_{2}
$$

The current in the secondary is found by substitution, yielding

$$
I_{2}=\frac{V_{1} I_{1}}{V_{2}}=\frac{(220 \mathrm{~W})}{(10 \mathrm{~V})}=22 \mathrm{~A}
$$

With the current and voltage known, we find the resistance:

$$
R=\frac{V_{2}}{I_{2}}=\frac{(10 \mathrm{~V})}{(22 \mathrm{~A})}=0.46 \Omega
$$

EVALUATE What is the ratio of the primary resistance to the secondary resistance? It is the square of the ratio of the number of turns in the primary circuit to the number of turns in the secondary circuit.

PRACTICE PROBLEM What is the ratio of the number of turns in the two circuits? Answer: 11.

## Try It Yourself!

## 1: An LR circuit

A resistor and an inductor are connected in series to a $60-\mathrm{Hz}$ voltage source that produces a maximum voltage of 155.6 V . If the resistor has a resistance of $200 \Omega$ and the inductor has an inductance of 0.3 H , find (a) the impedance of the circuit, (b) the maximum current, and (c) the phase angle between the current and voltage.

## Solution Checkpoints

IDENTIFY AND SET UP Use the definitions to find the target variables. A phasor diagram will guide you to the solution.

EXECUTE (a) The impedance is given by

$$
Z=\sqrt{R^{2}+X_{L}^{2}}=230 \Omega
$$

Note that the angular frequency is $377 \mathrm{rad} / \mathrm{s}$.
(b) The maximum current is 0.676 A .
(c) The phase angle is found by examining the phasor diagram, yielding

$$
\phi=\cos ^{-1}\left(\frac{R}{Z}\right)
$$

The source voltage leads the current by $29.6^{\circ}$.
EVALUATE Can you find the time average provided by the source? Would you find the power used by the resistor and inductor and add those values?

## 2: An LRC circuit

A series LRC circuit is made from a $100-\Omega$ resistor, an $11-\mathrm{mH}$ inductor, and a $0.21-\mu \mathrm{F}$ capacitor. The elements are connected to an alternating-voltage source with a maximum voltage of 50 V operating at a frequency of 5000 Hz . Find the maximum current and the phase difference between the current and voltage.

## Solution Checkpoints

IDENTIFY AND SET UP Use the definitions to find the target variables. A phasor diagram will guide you to the solution.

EXECUTE The angular frequency is $3.14 \times 10^{4} \mathrm{rad} / \mathrm{s}$. The reactances are found to be

$$
\begin{aligned}
& X_{L}=346 \Omega, \\
& X_{C}=152 \Omega .
\end{aligned}
$$

The impedance is then

$$
Z=\sqrt{R^{2}+\left(X_{L}-X_{C}\right)^{2}}=218 \Omega
$$

The maximum current is

$$
I=\frac{V}{Z}=0.229 \mathrm{~A}
$$

The phase angle is found to be

$$
\phi=\cos ^{-1}\left(\frac{R}{Z}\right)=62.7^{\circ}
$$

Since the inductive reactance is greater than the capacitive reactance, the voltage leads the current by $62.7^{\circ}$.

EVALUATE Would the voltage lag or lead the current if the inductive reactance were less than the capacitive reactance?

## Electromagnetic Waves

## Summary

In this chapter, we'll study electromagnetic waves. Light, televisions, radios, cellular phones, microwave ovens, and radioactive nuclei have their origins in electromagnetic waves; their study is fundamental to our understanding of such waves. Maxwell's equations predict the existence of electromagnetic waves and describe them as timevarying electric and magnetic fields propagating through space. Electromagnetic waves carry energy and momentum, and we'll learn to describe the flow of waves, energy, and momentum through space.

## Objectives

After studying this chapter, you will understand

- The nature of electromagnetic waves as predicted by Maxwell's equations.
- The propagation of electromagnetic waves through space and matter.
- How to describe the propagation of sinusoidal waves through space.
- How electromagnetic waves transmit energy and momentum through space.

Concepts and Equations

| Term | Description |
| :---: | :---: |
| Maxwell's Equations and Electromagnetic Waves | Maxwell's equations predict the existence of electromagnetic waveselectromagnetic disturbances composed of time-varying electric and magnetic fields that propagate at the speed of light, $c$, in vacuum. The speed of light is given by the expression $c=\frac{1}{\sqrt{\mathcal{E}_{0} \mu_{0}}}=3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}$ <br> The direction of propagation is the direction of $\vec{E} \times \vec{B}$. Electromagnetic waves are transverse: Their electric and magnetic fields are perpendicular to the direction of propagation and to each other. The electromagnetic spectrum is the full range of electromagnetic waves, covering frequencies from at least 1 to $10^{24} \mathrm{~Hz}$. |
| Sinusoidal Electromagnetic Waves | A sinusoidal plane electromagnetic wave traveling in vacuum in the $+x$ direction is described by the equations $\begin{aligned} & \vec{E}(x, t)=\hat{\jmath} E_{\max } \cos (k x-\omega t), \\ & \vec{B}(x, t)=\hat{k} B_{\max } \cos (k x-\omega t), \\ & E_{\max }=c B_{\max } . \end{aligned}$ |

## Electromagnetic Waves in Matter

## Energy and Momentum in Electromagnetic Waves

When an electromagnetic wave travels through a dielectric, the wave speed $v$ is less than the speed of light, $c$, in vacuum. The speed of light in a dielectric is given by

$$
c=\frac{1}{\sqrt{\mathcal{E}_{\mu}}}=\frac{c}{\sqrt{K K_{\mathrm{m}}}} .
$$

The Poynting vector $\vec{S}$ gives the energy flow rate of an electromagnetic wave in a vacuum:

$$
\vec{S}=\frac{1}{\mu_{0}} \vec{E} \times \vec{B} .
$$

The magnitude of the time-averaged value of the Poynting vector is the intensity $I$ of the wave:

$$
I=S_{\mathrm{av}}=\frac{E_{\max } B_{\max }}{2 \mu_{0}}=\frac{E_{\max }{ }^{2}}{2 \mu_{0} c}=\frac{1}{2} \sqrt{\frac{\mathcal{E}_{0}}{\mu_{0}}} E_{\max }^{2}=\frac{1}{2} \mathcal{E}_{0} c E_{\max }{ }^{2} .
$$

Electromagnetic waves carry momentum and exert radiation pressure $p_{\text {rad }}$ when striking a surface. For a surface perpendicular to the wave that is totally absorbing,

$$
p_{\mathrm{rad}}=\frac{I}{c} .
$$

For a surface perpendicular to the wave that is a perfect reflector,

$$
p_{\mathrm{rad}}=\frac{2 I}{c}
$$

The flow rate of electromagnetic radiation is given by

$$
\frac{1}{A} \frac{d p}{d t}=\frac{S}{c}=\frac{E B}{\mu_{0} c} .
$$

## Conceptual Questions

## 1: Solar sail

Is it possible to build a space vehicle that "sails" on the light from the sun?

## Solution

IDENTIFY, SET UP, AND EXECUTE Electromagnetic waves exert a radiation pressure on a surface. Light is an electromagnetic wave, so the light from the sun can be used to exert pressure on a surface. A space vehicle can be designed to operate off of a solar sail.

EVALUATE Solar sails have been designed for space vehicles, but are often supplemented by other forms of energy (such as gravitational energy). Among the drawbacks of solar sails are their inability to guide a vehicle toward the sun and their inability to provide large amounts of force. They are, however, inexpensive.

## 2: Designing a solar sail

Given that you are now convinced that you can create a space vehicle that sails on the sun's light, from what type of material should the sail be manufactured, reflective or absorbing? Or does it matter?

## Solution

IDENTIFY, SET UP, AND EXECUTE The radiation pressure of a perfectly reflecting surface perpendicular to the incident wave is twice that of a completely absorbing surface. To maximize the radiation pressure, which maximizes the force on the sail, a perfect reflector should be chosen as the material for the sail.

EVALUATE Even with twice the radiation pressure, the solar sail does not provide much power.

## 3: Reflection versus absorption

Why does a perfect reflector exert twice the radiation pressure on a surface compared with a total absorber?

## Solution

IDENTIFY, SET UP, AND EXECUTE Light reflecting from a perfect reflector will leave the surface with its initial momentum, thus changing its momentum by twice its initial momentum. Light being totally absorbed by a surface will have no momentum after absorption, thus changing its momentum by its initial momentum. The change in the light's momentum is momentum imparted onto the surface. Since light from the perfect reflector changes its momentum by twice that of light absorbed by the surface, light exerts a greater radiation pressure on the perfect reflector.

EVALUATE Recall that a ball bouncing off of a surface imparts more momentum to the surface than a piece of clay that is thrown at the surface and sticks. Since light has a similar property, can light, which we know is an electromagnetic wave, exhibit particle properties, as the ball or piece of clay does? We'll see that it does in Chapter 38.

## Problems

## 1: Radiating waves

A source of electromagnetic waves with $10^{7} \mathrm{~W}$ of power radiates uniformly in all directions. Calculate the amplitude of the electric field vector for waves at a distance of (a) 100.0 m from the source and (b) 1.0 km from the source.

## Solution

IDENTIFY We solve for the quantities in the problem by using the definition of intensity in terms of field magnitudes.

SET UP The waves radiate uniformly in all directions, so we start by finding the intensity at a spherical surface located an arbitrary distance $r$ from the source. From the intensity, we solve for the amplitude of the electric field vector.

EXECUTE Light radiates outward into a spherical surface. The area of that surface is

$$
A=4 \pi r^{2}
$$

The intensity at a distance $r$ is given by

$$
I=\frac{P}{A}=\frac{P}{4 \pi r^{2}} .
$$

Intensity is related to the amplitude of the electric field vector by

$$
I=\frac{1}{2} \mathcal{E}_{0} c E^{2}
$$

Solving for the amplitude gives

$$
E=\sqrt{\frac{P}{2 \pi r^{2} \mathcal{E}_{0} c}}
$$

Substituting, we find the amplitude of the electric field at 100.0 m :

$$
\begin{aligned}
E & =\sqrt{\frac{P}{2 \pi r^{2} \mathcal{E}_{0} c}} \\
& =\sqrt{\frac{\left(10^{7} \mathrm{~W}\right)}{2 \pi(100.0 \mathrm{~m})^{2}\left(8.85 \times 10^{-12} \mathrm{C}^{2} / \mathrm{N} \cdot \mathrm{~m}^{2}\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}} \\
& =245 \mathrm{~V} / \mathrm{m} .
\end{aligned}
$$

Similarly, at 1.0 km , the amplitude of the electric field is $24.5 \mathrm{~V} / \mathrm{m}$.
EVALUATE We see that the amplitude of the electric field is proportional to the inverse of the distance. Thus, the amplitude of the electric field drops by a factor of 10 when the distance to the source increases by a factor of 10 .

## 2: From Neptune to Pluto

What are the minimum and maximum times required for light to travel from Neptune to Pluto?

## Solution

IDENTIFY We will find the time taken for light to travel between the two planets by using the speed of light and the distance between the planets.

SET UP The maximum time required for light to travel between the two planets occurs when they are on opposite sides of the sun, completely opposite each other in their orbits. The distance between them in this case is the sum of their orbital radii. The minimum time required for light to travel between the two planets occurs when both planets are aligned on the same side of the sun. The distance between them in this case is the difference of their orbital radii.

Appendix F, on the inside of the book's cover, lists the orbital radius of Neptune as $4.50 \times 10^{12} \mathrm{~m}$ and that of Pluto as $5.91 \times 10^{12} \mathrm{~m}$. Space approximates a vacuum, so we use the speed of light in a vacuum in this problem.

EXECUTE The minimum distance between the two planets is the difference of their orbital radii. Light traveling between them in this configuration will take a time given by the distance the light travels divided by the speed of light, or

$$
\begin{aligned}
t_{\text {minimum }} & =\frac{\text { distance }}{c} \\
& =\frac{r_{\text {Pluto }}-r_{\text {Neptune }}}{c} \\
& =\frac{\left(5.91 \times 10^{12} \mathrm{~m}\right)-\left(4.50 \times 10^{12} \mathrm{~m}\right)}{\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}=4700 \mathrm{~s}=1.3 \mathrm{~h}
\end{aligned}
$$

The maximum distance between the two planets is the sum of their orbital radii. Light traveling between them in this configuration will take a time given by the distance the light travels divided by the speed of light, or

$$
\begin{aligned}
t_{\text {maximum }} & =\frac{\text { distance }}{c} \\
& =\frac{r_{\text {Pluto }}+r_{\text {Neptune }}}{c} \\
& =\frac{\left(5.91 \times 10^{12} \mathrm{~m}\right)+\left(4.50 \times 10^{12} \mathrm{~m}\right)}{\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}=34,700 \mathrm{~s}=9.6 \mathrm{~h}
\end{aligned}
$$

EVALUATE We can get an idea of the vastness of our solar system by realizing how much time it takes for light to travel across it. This problem shows that it takes almost 10 hours for light to travel between the outermost two bodies of the solar system.

## 3: Energy from the sun

The intensity of radiation from the sun is $1.4 \mathrm{~kW} / \mathrm{m}^{2}$. (a) Find the maximum values of the electric and magnetic fields. If a beam of sunlight falls on a perfectly reflecting surface of area $1 \mathrm{~m}^{2}$ for 1 minute, find (b) the energy reflected from the mirror, (c) the momentum delivered to the mirror during that time, and (d) the force acting on the mirror.

## Solution

IDENTIFY AND SET UP We solve for the quantities in the problem by using the definition of intensity in terms of field magnitudes and in terms of radiation pressure.
EXECUTE The intensity of an electromagnetic wave given in terms of the maximum electric field is

$$
I=\frac{1}{2} \mathcal{E}_{0} c E_{\max }^{2} .
$$

Rearranging terms to find the maximum electric field yields

$$
E_{\max }=\sqrt{\frac{2 I}{\mathcal{E}_{0} c}}=\sqrt{\frac{2\left(1400 \mathrm{~W} / \mathrm{m}^{2}\right)}{\left(8.85 \times 10^{-12} \mathrm{C}^{2} /\left(\mathrm{N} \cdot \mathrm{~m}^{2}\right)\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}}=1030 \mathrm{~V} / \mathrm{m}
$$

The maximum magnetic field is then

$$
B_{\max }=\frac{E_{\max }}{c}=\frac{(1030 \mathrm{~V} / \mathrm{m})}{\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}=3.4 \times 10^{-6} \mathrm{~T}
$$

The intensity is the energy per area per unit time, so we find the energy by multiplying the intensity by area and time:

$$
\Delta U=I A \Delta t=\left(1400 \mathrm{~W} / \mathrm{m}^{2}\right)\left(1 \mathrm{~m}^{2}\right)(60 \mathrm{~s})=84,000 \mathrm{~J}
$$

The average momentum transferred per unit time per unit area is given by

$$
\frac{1}{A} \frac{\Delta p}{\Delta \mathrm{t}}=\frac{I}{c}
$$

Our surface is reflecting, so the momentum imparted is twice this value. The momentum imparted is then

$$
\Delta p=2 \frac{I A \Delta t}{c}=2 \frac{\left(1400 \mathrm{~W} / \mathrm{m}^{2}\right)\left(1 \mathrm{~m}^{2}\right)(60 \mathrm{~s})}{\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}=5.60 \times 10^{-4} \mathrm{~kg} \cdot \mathrm{~m} / \mathrm{s}
$$

The impulse of the force is the change in momentum, so the force is the change in momentum divided by the time:

$$
F=\frac{\Delta p}{\Delta t}=\frac{\left(5.60 \times 10^{-4} \mathrm{~kg} \cdot \mathrm{~m} / \mathrm{s}\right)}{(60 \mathrm{~s})}=9.33 \times 10^{-6} \mathrm{~N}
$$

EVALUATE This problem illustrates how electromagnetic waves transmit energy, momentum, and force through their electric and magnetic waves. You can feel the transmission of electromagnetic wave energy when you open an oven door or sit outside in the sun.

## 4: Force on the earth

Find the average force exerted by the sun's light on the earth. Take the average intensity of radiation from the sun to be $1.4 \mathrm{~kW} / \mathrm{m}^{2}$.

## Solution

IDENTIFY We will use the relations among intensity, power, radiation pressure, and force to solve the problem.

SET UP We model the earth as a circular, flat surface with a radius equal to that of the earth $\left(6.38 \times 10^{6} \mathrm{~m}\right)$. We will find the force by combining the radiation pressure and the area of the circle.

We will find the radiation pressure from the intensity of the sun's light that is incident upon the earth. We will take the earth as a total absorber.

EXECUTE The radiation pressure for a totally absorbing surface oriented perpendicular to the incident wave is given by

$$
p_{\mathrm{rad}}=\frac{I}{c}
$$

The radiation pressure for the light from the sun is therefore

$$
p_{\mathrm{rad}}=\frac{I}{c}=\frac{\left(1.4 \times 10^{3} \mathrm{~W} / \mathrm{m}^{2}\right)}{\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}=4.7 \times 10^{-6} \mathrm{~Pa}
$$

The average force is equal to the radiation pressure times the area of the circle:

$$
\begin{aligned}
F & =p_{\mathrm{rad}} A \\
& =p_{\mathrm{rad}} \pi r_{\text {earth }}{ }^{2} \\
& =\left(4.7 \times 10^{-6} \mathrm{~Pa}\right) \pi\left(6.38 \times 10^{6} \mathrm{~m}\right)^{2}=6.01 \times 10^{8} \mathrm{~N}
\end{aligned}
$$

The sun's light exerts an average force on the earth of $6.01 \times 10^{8} \mathrm{~N}$.
EVALUATE With such a large force, why doesn't the earth move away from the sun? The earth has a very large mass, so the acceleration due to this force on the earth alone is tiny.

PRACTICE PROBLEM What is the acceleration of the earth resulting from the average force on the earth due to the sun's light, assuming that no other forces are present. The earth's mass is $5.97 \times$ $10^{24} \mathrm{~kg}$. Answer: $1.01 \times 10^{-16} \mathrm{~m} / \mathrm{s}^{2}$.

## Try It Yourself!

## 1: Light traveling from the sun to the earth

How much time does it take for light to travel from the sun to the earth?

## Solution Checkpoints

IDENTIFY AND SET UP Light travels at a constant rate. Use the average distance between the sun and the earth.

EXECUTE The time taken for light to travel between the sun and the earth is

$$
\begin{aligned}
t & =\frac{\text { distance }}{c} \\
& =\frac{r_{\text {earth }}}{c} \\
& =\frac{\left(1.50 \times 10^{11} \mathrm{~m}\right)}{\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}=500 \mathrm{~s}=8.3 \mathrm{~m}
\end{aligned}
$$

EVALUATE We see that the earth is relatively close to the sun, considering that it take less than 10 minutes for light to travel from the sun to the earth. Compare that travel time with the 10 hours that is needed for light to travel from Pluto to Neptune in Problem 2.

## 2: Electric field due to a source

At a distance of 15.0 km from a source that radiates electromagnetic waves uniformly in all directions, the electric field amplitude is found to be $125 \mathrm{~V} / \mathrm{m}$. Find (a) the time-averaged value of the Poynting vector and (b) the time-averaged power radiated by the source.

## Solution Checkpoints

IDENTIFY AND SET UP Use the relations among intensity, power, and electric field amplitude to solve the problem.

EXECUTE (a) The average Poynting vector $S$ is the intensity of an electromagnetic wave. Given in terms of the maximum electric field,

$$
S_{\mathrm{av}}=I=\frac{1}{2} \mathcal{E}_{0} c E_{\max }^{2}
$$

Rearranging terms and solving leads to a value of $20.7 \mathrm{~W} / \mathrm{m}^{2}$.
(b) The intensity is the power per unit area, so find the power radiated by multiplying the intensity by the area of a sphere with radius 15.0 km :

$$
P=I A=\left(20.7 \mathrm{~W} / \mathrm{m}^{2}\right) \pi(15.0 \mathrm{~km})^{2}=5.86 \times 10^{10} \mathrm{~W}
$$

EVALUATE If you measured the electric field amplitude for an electromagnetic wave at several locations, could you determine the location of the radiation source and its power?

## The Nature and Propagation of Light

## Summary

Light is an electromagnetic wave, and we will study the properties of light in this chapter. Light allows us to see, and is used to provide insight, into physical processes at the atomic scale. We will focus on visible light to understand how light is reflected from and propagates through materials. That study will lay the foundation for our examination of optics in the next chapter.

## Objectives

After studying this chapter, you will understand

- How light rays are used to represent the propagation of light through space.
- How to apply the law of reflection.
- How to apply the law of refraction.
- The meaning and application of total internal reflection.
- Polarization and how to apply it to a variety of problems.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Light | Light is an electromagnetic wave that also shows particle properties as it is <br> emitted or absorbed. A wave front is a surface of constant phase. Wave fronts <br> move at the wave's propagation speed. A ray represents light as a line along <br> the direction of propagation, which is perpendicular to the wave front. |
| When light is transmitted from one material to another, its frequency remains <br> constant, but its wavelength changes. The index of refraction, $n$, of a material <br> is the ratio of the speed of light in vacuum, $c$, to the speed $v$ in the material: <br> $n=c / v$. The variation of the index of refraction with wavelength is called <br> dispersion. |  |

## Reflection and Refraction

At a smooth interface between two optical materials, the incident, reflected, and refracted rays are related. The law of reflection states that the angles of incidence and reflection are the same:

$$
\theta_{r}=\theta_{a} .
$$

The law of reflection also relates the angle of incidence to the angle of refraction as

$$
n_{a} \sin \theta_{a}=n_{b} \sin \theta_{b},
$$

where $n$ is the index of refraction of the material. Angles of incidence, reflection, and refraction are always measured with respect to the normal to the interface.

## Total Internal Reflection

## Polarization

Total internal reflection occurs when light travels within a material with a greater index of refraction, $n_{a}$, towards an interface with a smaller index of refraction, $n_{b}$, and the angle of incidence exceeds a critical angle given by

$$
\theta_{\text {crit }}=\sin ^{-1} \frac{n_{b}}{n_{a}} .
$$

Light rays with angles greater than $\theta_{\text {crit }}$ are totally reflected; no light is refracted.

Polarized electromagnetic waves fluctuate along a single axis, with the direction of polarization given by the direction the electric field. Polarization filters allow the passage of radiation that is linearly polarized along the polarization axis and block radiation that is polarized perpendicular to the axis. Malus's law relates the incident intensity $I_{\max }$ to the transmitted intensity $I$ through a polarizer with its axis oriented at an angle $\phi$ to the incident wave's polarization axis:

$$
I=I_{\max } \cos ^{2} \phi
$$

Brewster's law states that, for unpolarized light that strikes an interface between two materials, the reflected light is completely polarized perpendicular to the plane of incidence at an angle of

$$
\theta_{\mathrm{p}}=\tan ^{-1} \frac{n_{b}}{n_{a}},
$$

where $n_{a}$ is the index of refraction of the incident material and $n_{b}$ is the index of refraction of the reflected surface.

## Conceptual Questions

## 1: Fishing with a Spear

You are standing on the side of a river and spot a delicious fish below the surface of the water. To spear the fish, should you aim above, below, or directly at its image?

## Solution

IDENTIFY, SET UP, AND EXECUTE The light rays emanating from the fish in the water are refracted as they exit the water and enter the air. Since the index of refraction of water is greater than that of air, the image of the fish appears above the actual fish. (See Figure 33.1.) You will therefore want to aim below the image of the fish to spear it.


Figure 33.1 Question 1.
EVALUATE Refraction plays an important role in the formation of images, as we will see in the next chapter. Lenses make use of refraction to magnify images of faraway objects.

Practice Problem Where should you aim a laser beam to stun the fish? Answer: Aim the laser beam directly at the image of the fish. The light from the laser beam will refract in the same way as the light coming from the fish.

## 2: Light through crossed linear polarizers

Two linearly polarizing filters are arranged with their axes perpendicular to each other so that no light passes through the filters. Can a third linearly polarizing filter be placed somewhere so that light passes through the system?

## Solution

IDENTIFY, SET UP, AND EXECUTE Placing the third linearly polarizing filter outside of the original two will not allow light to pass through. However, placing it between the original two and orienting its axis at an angle that does not coincide with either of the two filters' axes will allow light to pass through. With this configuration, the axes of any two successive filters will not be perpendicular, and light will pass through according to Malus's law.

EVALUATE Crossed linear polarizers prevent the transmission of light only when they are placed next to each other. If any elements are placed between the polarizers, transmission may occur.

## 3: Polarized glasses on the beach

If you sit on a beach wearing polarized sunglasses, glare from the sunlight reflected off the ocean is reduced. If you lay down with your head sideways and view the ocean, the glare from the sunlight reflecting off the ocean is reduced only slightly. Why?

## Solution

IDENTIFY, SET UP, AND EXECUTE Sunlight that reflects off the ocean is partially polarized, mostly in the horizontal direction. Polarized sunglasses reduce glare by orienting the polarization axis vertically, thus reducing the horizontally polarized light (and therefore the glare). When you are sitting, your sunglasses are oriented correctly and the glare is reduced. When you lie down, the axis of the glasses tilts $90^{\circ}$ and the polarization axis is no longer perpendicular to the horizontal direction, so the sunglasses don't reduce the glare significantly.

EVALUATE The next time you're wearing polarized sunglasses, tilt your head to see how the glare increases. This exercise illustrates how reflection polarizes light.

## Problems

## 1: Light reflecting off of plane mirrors

A light ray is incident on two plane mirrors intersecting at an angle $\theta$, as shown in Figure 33.2. Find the angle $\alpha$ in terms of the angle $\theta$.


Figure 33.2 Problem 1.

## Solution

IDENTIFY We will use the law of reflection and geometry to solve this problem
SET UP We begin by sketching the figure and adding angles (Figure 33.3). At the first reflection, the angle of incidence and angle of reflection for the top mirror is $\theta_{a}$. At the second reflection, the angle of incidence and angle of reflection for the bottom mirror is $\theta_{b}$. We determine $\alpha$ by using the law of reflection and triangle identities.


Figure 33.3 Problem 1 Sketch.

EXECUTE Examining the figure, we see that the triangle formed by the three light rays gives a relationship among angles $\theta_{a}, \theta_{b}$, and $\alpha$. At the reflections, the angle between the incident and reflected rays must be twice the incident angle. Summing the angles in the triangle gives $180^{\circ}$ :

$$
\begin{aligned}
& 2 \theta_{a}+2 \theta_{b}+\alpha=180^{\circ} \\
& 2\left(\theta_{a}+\theta_{b}\right)+\alpha=180^{\circ}
\end{aligned}
$$

The triangle to the left of the first triangle (formed by the surfaces of the two mirrors and the ray reflected between them) gives a relationship among angles $\theta_{a}, \theta_{b}$, and $\theta$. The sum of the complements of the two incident angles and $\theta$ is $180^{\circ}$ :

$$
\begin{aligned}
& \theta+\left(90^{\circ}-\theta_{a}\right)+\left(90^{\circ}-\theta_{b}\right)=180^{\circ} \\
& \theta=\theta_{a}+\theta_{b}
\end{aligned}
$$

Combining these results yields

$$
\begin{aligned}
& 2(\theta)+\alpha=180^{\circ} \\
& \alpha=180^{\circ}-2 \theta
\end{aligned}
$$

EVALUATE This problem shows how we must use geometry to solve reflection problems. What happens when $\theta=90^{\circ}$ ? In this case, $\alpha$ becomes $0^{\circ}$ and the two rays are parallel to each other.

## 2: Light through a slab of glass

Light enters a thin sheet of glass, refracts, and exits the glass, as shown in Figure 33.4. The angle of incidence is $60^{\circ}$, the thickness of the glass is 1.2 cm , and the ray is displaced a distance $d=0.80 \mathrm{~cm}$. Find the index of refraction of the glass.


Figure 33.4 Problem 2.

## Solution

IDENTIFY We will use the law of refraction and geometry to solve this problem.
SET UP We are given the angle of incidence and know the index of refraction of air. We need the refraction angle to find the index of refraction of the glass. We find the refraction angle by geometry, using the thickness of the glass and the displacement of the ray. Once we have the refraction angle, we use Snell's law to find the index of refraction of the glass.

EXECUTE We find the angle of refraction by examining the figure. The tangent of the refraction angle is the displacement of the ray divided by the thickness of the glass:

$$
\tan \theta_{g}=\frac{d}{t}, \quad \theta_{g}=33.7^{\circ}
$$

Snell's law is then used to find the index of refraction. For the air-glass interface, we have

$$
n_{a} \sin \theta_{a}=n_{g} \sin \theta_{g}
$$

The index of refraction of air is unity, and the incident angle is $60^{\circ}$. Rearranging terms to find the index of refraction of the glass gives

$$
n_{g}=\frac{n_{a} \sin \theta_{a}}{\sin \theta_{g}}=\frac{(1)\left(\sin 60^{\circ}\right)}{\sin 33.7^{\circ}}=1.56
$$

EVALUATE As a check, we note that the index of refraction is greater than unity as we expected, since the refracted ray was bent towards the normal. The index of refraction for many varieties of glass is around 1.5 , so we conclude that our result is reasonable.

## 3: Light through a prism

A horizontal ray of light is incident on a glass prism as shown in Figure 33.5. The base of the prism is horizontal. The prism $(n=1.35)$ is surrounded by oil $(n=1.48)$. Determine the angle $\theta$ that the exiting light makes with the normal to the right face of the prism.


Figure 33.5 Problem 3.

## Solution

IDENTIFY We will use the law of refraction and geometry to solve this problem.
SET UP We begin by sketching the figure and adding useful angles. Figure 33.6 shows a close-up of the prism. Since both the incoming light ray and the base of the prism are horizontal, the angle between the incoming ray and the side of the prism must be $60.0^{\circ}$. The incident angle $\theta_{a}$ must be $30.0^{\circ}$. The first refracted angle is labeled $\theta_{b}$, the second incident angle is labeled angle $\theta_{c}$, and the second refracted angle is labeled $\theta$. We determine $\theta$ by using the law of refraction and triangle identities.


Figure 33.6 Problem 3.
EXECUTE We apply law of refraction to both interfaces. At the first interface, the ray originates in the oil and refracts into the glass, so

$$
n_{\text {oil }} \sin \theta_{a}=n_{\text {glass }} \sin \theta_{b} \text {. }
$$

We find the $\theta_{b}$, the first angle in the glass:

$$
\theta_{b}=\sin ^{-1}\left(\frac{n_{\text {oil }}}{n_{\text {glass }}} \sin \theta_{a}\right)=\sin ^{-1}\left(\frac{(1.48)}{(1.35)} \sin 30.0^{\circ}\right)=33.24^{\circ} .
$$

To find the angle $\theta_{c}$, we examine the triangle made by the top of the prism and the light ray passing through the prism. The bottom two angles are the complements of $\theta_{b}$ and $\theta_{c}$. The two complement angles and the top angle must add to $180^{\circ}$, so

$$
\begin{aligned}
& \left(90.0^{\circ}-\theta_{b}\right)+\left(60.0^{\circ}\right)+\left(90.0^{\circ}-\theta_{c}\right)=180.0^{\circ}, \\
& \theta_{c}=60.0^{\circ}-\theta_{b}=60.0^{\circ}-\left(33.24^{\circ}\right)=26.76^{\circ}
\end{aligned}
$$

With $\theta_{c}$ known, we can solve for the angle $\theta$ with the formula

$$
n_{\text {glass }} \sin \theta_{c}=n_{\text {oil }} \sin \theta
$$

Solving for $\theta$ yields

$$
\theta=\sin ^{-1}\left(\frac{n_{\text {glass }}}{n_{\text {oil }}} \sin \theta_{c}\right)=\sin ^{-1}\left(\frac{(1.35)}{(1.48)} \sin 26.76^{\circ}\right)=24.2^{\circ} .
$$

The angle of exit is $24.2^{\circ}$.
EVALUATE This solution shows how to combine the law of refraction with geometry to solve a refraction problem. As you can see, you must analyze the problem carefully to find the correct incident and refracted angles. Applying the law of refraction is straightforward.

## 4: Light source in water

A point light source 2.0 m below the surface of a water pool produces a circular pattern of light when viewed from above. Calculate the radius of the circle. Take the index of refraction of water to be 1.33 .

## Solution

IDENTIFY We will solve this problem by using our knowledge of refraction and geometry.

SET UP We begin by sketching the situation as viewed from the side, as shown in Figure 33.7. Light rays from the point source approach the water-air interface and are deflected away from the normal as they pass through the interface. Beyond the critical angle, no light is refracted and all of the light is reflected internally. The edge of the circle corresponds to the critical angle. We will find the critical angle and then find the radius which corresponds to that angle.


Figure 33.7 Problem 4 sketch.
EXECUTE The critical angle for the water-air interface is given by

$$
\sin \theta_{\text {crit }}=\frac{n_{\text {water }}}{n_{\text {air }}}, \quad \theta_{\text {crit }}=48.6^{\circ}
$$

The tangent of this critical angle is the radius of the circle divided by the depth of the light source, or

$$
R=d \tan \theta_{\text {crit }}=(2.0 \mathrm{~m}) \tan 48.6^{\circ}=2.27 \mathrm{~m}
$$

EVALUATE How does a light source located above the air-water interface appear to someone under the water? The light source will appear above the actual location of the light source, but there will be no cutoff angle, as there is no critical angle when the light source is viewed from below.

## 5: Combining several linear polarizers

A number of linear polarizing filters are stacked with the polarizing axis of each successive filter rotated $20^{\circ}$ from the previous filter. If unpolarized light of intensity $5.00 \mathrm{~W} / \mathrm{m}^{2}$ is incident upon the first polarizer and the emerging light has an intensity of $1.52 \mathrm{~W} / \mathrm{m}^{2}$, how many filters are in the stack?

## Solution

IDENTIFY We will solve this problem by using our knowledge of polarization and Malus's law.
SET UP We will use the fact that the first linear polarizer will reduce the unpolarized light intensity by a factor of two and the subsequent filters will reduce the intensity by a factor of $\cos ^{2} \theta$, due to the law of Malus. We will find the intensity emerging from the stack of filters as more filters are added. This exercise will lead to a general expression for the intensity coming from a stack of $n$ filters.

EXECUTE The first linear polarizer will reduce the intensity by a factor of two. The intensity after the first polarizer is thus

$$
I_{1}=\frac{1}{2} I_{0}=2.50 \mathrm{~W} / \mathrm{m}^{2}
$$

The next polarizer will reduce the intensity to

$$
I_{2}=I_{1} \cos ^{2} 20^{\circ}=\left(2.50 \mathrm{~W} / \mathrm{m}^{2}\right) \cos ^{2} 20^{\circ}=2.21 \mathrm{~W} / \mathrm{m}^{2}
$$

This is larger than the final emerging light intensity, so we add another polarizer and find the intensity:

$$
I_{3}=I_{2} \cos ^{2} 20^{\circ}=\left(2.21 \mathrm{~W} / \mathrm{m}^{2}\right) \cos ^{2} 20^{\circ}=1.95 \mathrm{~W} / \mathrm{m}^{2}
$$

Again, the intensity is too large, so we add another polarizer:

$$
I_{4}=I_{3} \cos ^{2} 20^{\circ}=\left(1.95 \mathrm{~W} / \mathrm{m}^{2}\right) \cos ^{2} 20^{\circ}=1.72 \mathrm{~W} / \mathrm{m}^{2}
$$

We are slowly approaching the correct number of polarizers. Let's add one more to find the intensity after five polarizers:

$$
I_{5}=I_{4} \cos ^{2} 20^{\circ}=\left(1.72 \mathrm{~W} / \mathrm{m}^{2}\right) \cos ^{2} 20^{\circ}=1.52 \mathrm{~W} / \mathrm{m}^{2}
$$

We have reached the desired emerging intensity and conclude that there are five polarizers in the stack.
EVALUATE This problem shows how we can systematically add polarizers and find the intensity after each one is added. Those more familiar with mathematics may see a pattern arising. The intensity after $n$ polarizers is

$$
I_{n}=\frac{1}{2} I_{0}\left(\cos ^{2} 20^{\circ}\right)^{n-1} .
$$

The initial and final intensities could be substituted into this expression, and $n$ could be found by taking logarithms of both sides.

## Try It Yourself!

## 1: Coin in a pond

A coin rests on the bottom of a shallow pond of depth 1.0 m . Find the apparent depth of the coin when it is viewed directly from above. Take the index of refraction of water to be 1.33 .

## Solution Checkpoints

IDENTIFY AND SET UP Snell's law and geometry are needed to solve this problem. A sketch of the problem is shown in Figure 33.8.


Figure 33.8 Try It Yourself 1.

The coin appears at a depth $D$ above its actual location. Light leaves the water at a position located a distance $x$ from a point directly above the coin.

EXECUTE The angles of incidence and refraction for a light ray leaving the coin are related by

$$
n_{\text {water }} \sin \theta_{\text {water }}=n_{\text {air }} \sin \theta_{\text {air }} .
$$

The sines are given by

$$
\begin{aligned}
& \sin \theta_{\text {water }}=\frac{x}{\sqrt{x^{2}+d^{2}}}, \\
& \sin \theta_{\text {air }}=\frac{x}{\sqrt{x^{2}+D^{2}}}
\end{aligned}
$$

Substituting gives

$$
\frac{n_{\text {water }} x}{\sqrt{x^{2}+d^{2}}}=\frac{x}{\sqrt{x^{2}+D^{2}}} .
$$

Squaring both sides and letting $x$ go to zero (for viewing from above) produces

$$
D=\frac{d}{n_{\text {air }}}=0.75 \mathrm{~m}
$$

EVALUATE Does the apparent depth increase or decrease as you view the coin from locations at incident angles greater than $0^{\circ}$ ?

## 2: Light through a prism

Light enters a glass prism, is refracted, and exits the prism, as shown in Figure 33.9. Find the exiting angle of the light. Take the index of refraction of the prism to be 1.50 and the initial incident angle to be $30.0^{\circ}$. The prism has equilateral sides.


Figure 33.9 Try It Yourself 2.

## Solution Checkpoints

IDENTIFY AND SET UP Snell's law and geometry are needed to solve this problem. What is the angle between sides of the prism?

EXECUTE The first angle of refraction is

$$
\sin \theta_{1}=n \sin \theta_{2}, \quad \theta_{2}=19.47^{\circ}
$$

Examining the triangles, we see that

$$
\begin{aligned}
& \theta_{2}+\theta_{3}+\left(180^{\circ}-60^{\circ}\right)=180^{\circ} \\
& \theta_{3}=60^{\circ}-\theta_{2} \\
& \theta_{3}=40.53^{\circ}
\end{aligned}
$$

The second refraction gives the exiting angle:

$$
\sin \theta_{4}=n \sin \theta_{3}, \quad \theta_{4}=77.1^{\circ}
$$

EVALUATE Careful analysis of angles and the use of Snell's law led to the solution.

## 3: Light through polarizing filters

Two ideal linear polarizing filters are arranged such that no light passes through them. A third linear polarizer is placed between the two polarizing filters. What is the total transmitted light intensity as a function of the orientation of the new polarizer? Take the intensity of light onto the first filter to be $I_{0}$ and the angle between the first filter and the newly added filter to be $\theta$.

## Solution Checkpoints

IDENTIFY AND SET UP Use Malus's law to solve the problem.
EXECUTE The intensity after the first filter will be

$$
I_{1}=\frac{1}{2} I_{0} .
$$

After the second filter, the intensity drops to

$$
I_{2}=\left(\cos ^{2} \theta\right) I_{1}=\left(\frac{1}{2} \cos ^{2} \theta\right) I_{0}
$$

After the third filter, the intensity drops to

$$
\begin{aligned}
I_{3} & =\cos ^{2}(90-\theta) I_{2} \\
& =\left(\frac{1}{2} \cos ^{2} \theta \sin ^{2} \theta\right) I_{0} \\
& =\left(\frac{1}{2} \cos ^{2} \theta \cos ^{2}(90-\theta)\right) I_{0} \\
& =\left(\frac{1}{8} \sin ^{2} 2 \theta\right) I_{0}
\end{aligned}
$$

EVALUATE When is the transmission maximum?

## Geometric Optics and Optical Instruments

## Summary

In this chapter, we'll use the concepts of reflection and refraction to understand the formation of images in mirrors and lenses. An image is formed by the collection of light rays that converge towards, or appear to diverge from, a point. We'll determine, both graphically and through calculations, the locations of images created by lenses and mirrors. We'll also define virtual and real images, as well as focal points, and we'll learn to quantify the size of an image relative to its object. We'll then go on to explore familiar optical instruments, including cameras, magnifiers, the human eye, microscopes, and telescopes. We'll apply our knowledge of geometric optics to discover how images are formed in these instruments. Optical systems may include one, several, or many lenses and mirrors. We'll focus on systems with one or two components to illustrate their function.

## Objectives

After studying this chapter, you will understand

- How light rays form real and virtual images.
- The magnification of images formed by mirrors and lenses.
- Focal points for mirrors and lenses.
- How to find, both by calculation and graphically, the location and size of an image formed by a spherical mirror.
- How to find, both by calculation and graphically, the location and size of an image formed by a thin lens.
- How to find the locations and sizes of images formed by combinations of multiple lenses.
- Image formation in cameras, magnifiers, the human eye, microscopes, and telescopes.
- Magnification for various optical instruments.
- The fundamental design features of microscopes and telescopes.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Image Formation | Light rays diverging from an object point that reflect or refract from a surface <br> may form an image if the outgoing rays appear to have diverged from an <br> image point. The image formed is a virtual image if the rays don't converge <br> at a point in space and is a real image if the rays converge at a point in space. <br> The lateral magnification $m$ is the ratio of the image height $y^{\prime}$ to the object <br> height $y:$ |
| $\qquad m=\frac{y^{\prime}}{y}$. |  |

For positive $m$, the object is erect, and for negative $m$, the image is inverted.

## Reflection at a Spherical Surface

Graphical Method for Image Formation in Mirrors

Images are formed in spherical mirrors through the law of reflection. Parallel rays will reflect and converge at a focal point of a concave mirror and will appear to diverge at the focal point of a convex mirror. The focal length $f$ of a spherical mirror is the distance from the vertex to the focal point and is equal to half the radius of curvature $(f=R / 2)$. The object distance $s$, image distance $s^{\prime}$, and focal length $f$ are related by

$$
\frac{1}{s}+\frac{1}{s^{\prime}}=\frac{1}{f}
$$

Four principal rays can be drawn to find the size and location of the image formed by a mirror:

1. A ray reflected parallel to the optic axis passes through the focal point of a concave mirror or appears to originate at the virtual focal point of a convex mirror.
2. A ray passing through (or towards) the focal point is reflected parallel to the optic axis.
3. A ray passing along the radius through or away from the center of curvature intersects the surface normally and is reflected back along its original path.
4. A ray going to the vertex of the mirror is reflected, forming equal angles with the optic axis.
The location of the image is the point where the rays intersect or appear to originate.

## Refraction at a Spherical Surface Thin Lenses

Images are formed in spherical lenses through the law of refraction. A thin lens has two spherical surfaces close enough together that the distance between the surfaces can be ignored. A thin lens has two focal points located equidistant from the lens on either side. The thin-lens equation describes the behavior of the lens:

$$
\frac{1}{\mathrm{~s}}+\frac{1}{s^{\prime}}=\frac{1}{f}=(n-1)\left(\frac{1}{R_{1}}-\frac{1}{R_{2}}\right)
$$

Here, $n$ is the index of refraction of the lens material, and $R_{1}$ and $R_{2}$ are the radii of curvature of the first and second surfaces, respectively. The image magnification for a thin lens is given by

$$
m=-\frac{s^{\prime}}{s}
$$

| Graphical Method for Image |
| :--- | :--- |
| Formation in Thin Lenses | | Three principal rays can be drawn to find the size and location of the image |
| :--- |
| formed by a thin lens: |
| 1. A ray refracted parallel to the optic axis passes through the second focal |
| point of a converging lens or appears to originate from the second focal point |
| of a diverging lens. |
| 2. A ray passing through the center of the lens does not deviate appreciably |
| from its path. |
| 3. A ray passing through (or towards) the first focal point is refracted and |
| emerges parallel to the optic axis. |
| The location of the image is the point where the rays intersect or appear to |
| originate. |$\quad$| A camera forms a real, inverted, often reduced image of an object on a light- |
| :--- |
| sensitive surface. The intensity of the light striking the surface is controlled |
| by the shutter speed and aperture. The intensity is inversely proportional to |
| the square of the $f$-number of the lens: |

For the average viewer with a $25-\mathrm{cm}$ near point, the angular magnification becomes

$$
M=\frac{25 \mathrm{~cm}}{f}
$$

## Microscopes and Telescopes

In a compound microscope, the first lens (the objective lens) forms an image that becomes the object of the second lens (the eyepiece). The eyepiece forms a virtual image, often at infinity, of the first image. The overall magnification is the product of the magnifications of the two lenses; that is,

$$
M=m_{1} M_{2}=\frac{(25 \mathrm{~cm}) s_{1}^{\prime}}{f_{1} f_{2}}
$$

where $s^{\prime}{ }_{1}$ is the image distance of the objective.
A telescope operates similarly, but on objects that are far away. Since the object is far away, the image is formed at the focal plane of the objective lens, and the focal point of the eyepiece is set to coincide with the objective's focal point. The angular magnification of a telescope is

$$
M=-\frac{f_{1}}{f_{2}}
$$

## Conceptual Questions

## 1: Image in a plane mirror

You view the image of an object in the plane mirror shown in Figure 34.1. Can you see an image? If so, at what position is the image formed? If not, why not?


Figure 34.1 Question 1.

## Solution

IDENTIFY, SET UP, AND EXECUTE Images in plane mirrors are formed equidistant from the mirror; the image and object distances are the same in magnitude. The image of the object is located behind the mirror, at point $E$. Note that the mirror doesn't need to extend to the region where the object and image are located; the collection of light rays emanating from the object that reflects from the mirror into your eyes forms the image at $E$.

EVALUATE We see that images can be formed by a mirror that only partially covers a region. The same holds true for lenses: Partial lenses will also form images, but with reduced intensity compared with that of a full lens.

## 2: What does a fish see?

A fish looks at a spectator from the inside of an aquarium. The spectator is 1.0 m from the side of the aquarium. How far from the side of the aquarium does the spectator appear to the fish, at 1.0 m , or closer or farther than 1.0 m ?

## Solution

IDENTIFY, SET UP, AND EXECUTE Since the spectator is in (air $n_{\text {air }}<n_{\text {water }}$ ), the light rays coming from the person will refract towards the normal as they pass into the aquarium. The person will appear farther from the fish than the $1.0-\mathrm{m}$ distance. We can confirm this fact with the sketch shown in Figure 34.2. We see that the rays emanating from the person are bent towards the normal as they enter the aquarium. Tracing those rays backwards yields an image farther from the side of the aquarium.


Figure 34.2 Question 2.
EVALUATE This situation is opposite to the one illustrated in Figure 34.3 in the textbook: When the person looks into a medium with a higher index of refraction, the image is formed closer to the interface than the object is located.

## 3: Finding the focal points

An object is placed in front of a thin lens. Two light rays from the object pass through the lens and refract as shown in Figure 34.3, crossing at point $A$ to the right of the lens. Find the two focal points, located on either side of the lens.


Figure 34.3 Question 3.

## Solution

IDENTIFY, SET UP, AND EXECUTE The image formed by the lens must be at point $A$, where the two rays cross. We find the focal points by constructing two principal rays. Recall that a principal ray passing parallel to the optic axis will refract through the second focal point. We construct this ray by taking a ray parallel to the axis and refracting it such that it passes through point $A$, where the second focal point is when it crosses the optic axis. Recall also that a ray passing through the first focal point will refract parallel to the optic axis. We construct this ray by tracing a ray parallel to the axis backwards from point $A$ to the lens. The ray must originate at the object, where the first focal point is when it crosses the optic axis. These rays are illustrated in Figure 34.4.


Figure 34.4 Question 3 sketch.

EVALUATE We see that both focal points are equidistant from the lens and on opposite sides, consistent with our expectation. You might expect that the focal points are where the two original light rays crossed the optic axis. However, as we've seen, we must use our principal rays to locate the focal points. We can draw many light rays that cross the optic axis at various points along the axis, but only two pass through the focal points.

## 4: Diving masks

When snorkeling or scuba diving, you wear a diving mask that leaves a region of air between the mask and your eyes. If you take the mask off, you cannot see underwater objects clearly. Why?

## Solution

IDENTIFY, SET UP, AND EXECUTE The indexes of refraction for the components of the eye are close to the index of refraction for water. Therefore, much of the refraction of light entering the eye occurs at the cornea. When your eye is in air, the shape of the cornea allows images to be formed in your eye. When your eye is submerged in water, much less focusing occurs at the cornea, and you cannot focus on underwater objects.

EVALUATE The eyes of fish must have a larger curvature in order to image underwater objects. Fish cannot use the difference in the air-eye indexes of refraction to image objects.

## 5: Finding a strange lens

You find a lens and discover that it is thicker at the center than at the edges. You also find that the thickness at the edges changes: The lens is thicker at the top and bottom than at the sides. What kind of vision defect is this lens designed to correct?

## Solution

IDENTIFY, SET UP, AND EXECUTE Since the lens is thicker at the center than at the edges, it is a converging lens with a positive focal length. Converging lenses are used to correct for farsightedness, or hyperopia. The changing thickness around the edge indicates that the lens has different refractive powers along different axes. Therefore, this lens must also correct for astigmatism, or differences in the focus along perpendicular planes due to a nonspherical cornea. The lens corrects for both farsightedness and astigmatism.

EVALUATE Cylindrical lenses correct for both distance and astigmatism.

## Problems

## 1: Image in a convex mirror

A convex mirror has a radius of curvature of absolute value 25 cm . Determine graphically the image position of a real object placed 45 cm from the vertex of the mirror. Confirm the graphical analysis by computing the image distance and magnification.

## Solution

IDENTIFY The problem requires us to find the solution both graphically and by calculations. This will go a long way toward making sure that our result is correct.

SET UP AND EXECUTE The graphical construction is shown in Figure 34.5. The mirror is convex, so its radius of curvature and focal point are located behind it and to the right. Four light rays are drawn. The first ray passes parallel to the optic axis and is reflected so as to appear to originate from the focal point of the mirror. The second ray is directed at the focal point and reflects from the mirror's surface and parallel to the optic axis. The third ray is directed at the center of curvature and reflects back along its original path. The fourth ray is directed towards the vertex and reflects at the vertex such that the incoming and outgoing rays form equal angles with the optic axis.


Figure 34.5 Problem 1 sketch.
We find the location of the image by tracing the four reflected rays backwards and noting where they meet. As we see, the four rays intersect behind the mirror, forming a virtual image just to the left of the focal point. The image is upright and smaller than the original.

We check the graphical analysis by computing the image distance and magnification. The image distance is given by

$$
\frac{1}{s}+\frac{1}{s^{\prime}}=\frac{1}{f}
$$

The focal length of a convex mirror is half the radius of curvature of the mirror. The radius of curvature is negative for a convex mirror, so the focal length of this mirror is $f=-12.5 \mathrm{~cm}$. We find the image distance by rearranging terms and substituting:

$$
\frac{1}{s^{\prime}}=\frac{1}{f}-\frac{1}{s}=\frac{1}{(-12.5 \mathrm{~cm})}-\frac{1}{(45 \mathrm{~cm})}, \quad s^{\prime}=-9.78 \mathrm{~cm} .
$$

The magnification is then

$$
m=-\frac{s^{\prime}}{s}=-\frac{-9.78 \mathrm{~cm}}{45 \mathrm{~cm}}=+0.22
$$

The image is located 9.78 cm behind the mirror, which is a little to the left of the focal point. The image is virtual, is erect (the magnification is positive), and is $22 \%$ of the original size. This agrees with our graphical analysis.

EVALUATE We have seen throughout this course that sketches aid in our analysis. Solving image formation problems by both graphical methods and calculations helps ensure accurate results. When the two analyses don't agree, we must check our methods and correct any mistakes.

## 2: A thin-lens problem

You are given a converging lens with $f=+15.0 \mathrm{~cm}$ and are asked to position it so that it forms a virtual, erect image of an object that is three times taller than the original object. Where do you place the lens? Confirm your calculations by sketching a principal-ray diagram.

## Solution

IDENTIFY AND SET UP We can solve the problem either by graphical methods or by computing values. Since we are unsure of the location of the object, we begin by computing values. The magnification and thin-lens equations will be used to solve the problem. A graphical check will confirm our solution.

EXECUTE The magnification equation will give us a relation between the locations of the object and image. For a thin lens,

$$
m=-\frac{s^{\prime}}{s}
$$

The image is to be erect $(m>0)$ and three times higher $(|m|=3)$ than the object. Accordingly, we have

$$
s^{\prime}=-m s=-3 s
$$

The thin-lens equation relates the object distance, image distance, and focal length:

$$
\frac{1}{s}+\frac{1}{s^{\prime}}=\frac{1}{f}
$$

We substitute for the image distance:

$$
\frac{1}{s}+\frac{1}{-3 s}=\frac{1}{f}
$$

Multiplying all terms by $-3 s f$ gives

$$
\begin{gathered}
-3 f+f=-2 f=-3 s \\
s=\frac{2}{3} f=\frac{2}{3}(+15.0 \mathrm{~cm})=+10.0 \mathrm{~cm}
\end{gathered}
$$

The image distance is therefore

$$
s^{\prime}=-3 s=-3(10.0 \mathrm{~cm})=-30.0 \mathrm{~cm}
$$

The principal-ray diagram is sketched in Figure 34.6. The lens is converging, and the object is located to the left of the lens, between the first focal point and the lens. The first ray from the object passes parallel to the optic axis and refracts through the second focal point of the lens. The second ray passes through the center of the lens undeviated. The third ray appears to originate from the first focal point and is refracted, emerging parallel to the optic axis. These three rays do not intersect in real space, so we trace them back to where they appear to intersect. We see that they appear to intersect to the left of the first focal point, thus creating a virtual image. The image is erect and larger than the object. We conclude that our graphical analysis leads to results that are consistent with our earlier analysis.


Figure 34.6 Problem 2 sketch.

We need to place the lens 10.0 cm from the object in order to create an erect, virtual image at an image distance of -30.0 cm .

EVALUATE We could also have begun by sketching various scenarios to identify the solution. It may have taken several attempts, but we would have found the solution. Being able to start with either a principal-ray diagram or the thin-lens equation is a valuable technique.

## 3: Focal length of a spherical mirror

A real image is formed when an object is placed 10.0 cm from a concave spherical mirror. When the object is moved 2.0 cm farther from the mirror, the image moves 16.0 cm closer to the mirror. Find the focal length of the mirror.

## Solution

IDENTIFY We will apply the mirror equation to each situation to solve for the focal length of the mirror, the target variable.
SET UP We first sketch the two situations shown in Figure 34.7. On top, we see the original situation; on the bottom, we see the situation after the object is moved 2.0 cm farther from the mirror. We'll proceed to calculate the solution from the mirror equation for the two situations.


Figure 34.7 Problem 3 sketch.
EXECUTE The mirror equation relates the object distance, image distance, and focal length for spherical mirrors:

$$
\frac{1}{s}+\frac{1}{s^{\prime}}=\frac{1}{f}
$$

When the object is located at 10.0 cm , we have

$$
\frac{1}{(10.0 \mathrm{~cm})}+\frac{1}{s^{\prime}}=\frac{1}{f}
$$

There are two unknowns in this equation. We'll find a second relationship that describes the situation after the object has been moved and then combine the two equations to solve for the focal length. When the object is moved to 12.0 cm , we have

$$
\frac{1}{(12.0 \mathrm{~cm})}+\frac{1}{s^{\prime}-16.0 \mathrm{~cm}}=\frac{1}{f}
$$

The right-hand sides of the last two equations are equal, so we set them equal to each other:

$$
\frac{1}{(10.0 \mathrm{~cm})}+\frac{1}{s^{\prime}}=\frac{1}{(12.0 \mathrm{~cm})}+\frac{1}{s^{\prime}-16.0 \mathrm{~cm}}
$$

Simplifying yields

$$
\frac{1}{(60.0 \mathrm{~cm})}+\frac{1}{s^{\prime}}=\frac{1}{s^{\prime}-16.0 \mathrm{~cm}}
$$

We multiply all terms by $(60.0 \mathrm{~cm}) s^{\prime}\left(s^{\prime}-16.0 \mathrm{~cm}\right)$, which gives

$$
s^{\prime}\left(s^{\prime}-16.0 \mathrm{~cm}\right)+(60.0 \mathrm{~cm})\left(s^{\prime}-16.0 \mathrm{~cm}\right)=(60.0 \mathrm{~cm}) s^{\prime}
$$

or

$$
s^{\prime 2}-(16.0 \mathrm{~cm}) s^{\prime}-(960.0 \mathrm{~cm})=0
$$

We can use the quadratic formula or factor the preceding equation to solve for the image distance. Choosing factoring, we see that the equation becomes

$$
\left(s^{\prime}-40.0 \mathrm{~cm}\right)\left(s^{\prime}+24.0 \mathrm{~cm}\right)=0, \quad s^{\prime}=+40.0 \mathrm{~cm} \quad \text { or } \quad-24.0 \mathrm{~cm}
$$

The image is real, so it must be located at +40.0 cm . We can now substitute to find the focal length:

$$
\frac{1}{(10.0 \mathrm{~cm})}+\frac{1}{(40.0 \mathrm{~cm})}=\frac{1}{f}, \quad f=+8.0 \mathrm{~cm}
$$

The spherical mirror has a focal length of +8.0 cm .
EVALUATE This problem illustrates how we can apply the mirror equation to several related situations to find our solution. The mirror equation is straightforward, but in this case our solution required solving a quadratic formula.

Practice Problem: Check the results with a principal-ray diagram.

## 4: Combination of lenses

Most optical instruments contain multiple lenses. The image formed by the first lens becomes the object in the second lens and so on. To illustrate how these systems work, examine the following combination of lenses:

A $2.0-\mathrm{cm}$-tall object is placed 30.0 cm to the left of a converging lens with a focal length of +20.0 cm . A second converging lens with a focal length of +15.0 cm is placed 85.0 cm to the right of the first lens along the same optic axis. Find the location and height of the final image formed by the combination of lenses.

## Solution

IDENTIFY AND SET UP We'll first find the location of the image created by the first lens and then use that as our object in the second lens. The sketch shown in Figure 34.8 will aid our analysis. The top principal-ray diagram shows that the image from the first lens will be between the two lenses. We'll need to find the new object distance $\left(s_{2}\right)$, referenced by its relation to the second lens. We expect the final image to be real and erect. Each lens will magnify the image, and the total magnification will be the product of the two individual magnifications.

```
Lens 1
```



Figure 34.8 Problem 4 sketch.
EXECUTE The location of the first image is found with the thin-lens equation:

$$
\frac{1}{s}+\frac{1}{s^{\prime}}=\frac{1}{f}
$$

For the first lens, we use the subscript 1 to identify the quantities. The image distance is

$$
\frac{1}{s_{1}^{\prime}}=\frac{1}{f_{1}}-\frac{1}{s_{1}}=\frac{1}{(+20.0 \mathrm{~cm})}-\frac{1}{(+30.0 \mathrm{~cm})}, \quad s_{1}^{\prime}=+60.0 \mathrm{~cm}
$$

The first image becomes the object in the second lens. Since the second lens is 85.0 cm to the right of the first lens, the image of the first lens is formed 25.0 cm to the left of the second lens, or $s_{2}=+25.0 \mathrm{~cm}$. The second image distance is

$$
\frac{1}{s_{2}^{\prime}}=\frac{1}{f_{2}}-\frac{1}{s_{2}}=\frac{1}{(+15.0 \mathrm{~cm})}-\frac{1}{(+25.0 \mathrm{~cm})}, \quad s_{2}^{\prime}=+37.5 \mathrm{~cm} .
$$

The final image is formed 37.5 cm to the right of the second lens.

The size of the image is found by combining the magnifications. The magnification for the first lens is

$$
m_{1}=-\frac{s_{1}^{\prime}}{s_{1}}=-\frac{60.0 \mathrm{~cm}}{30.0 \mathrm{~cm}}=-2.0 .
$$

The magnification for the second lens is

$$
m_{2}=-\frac{s_{2}^{\prime}}{s_{2}}=-\frac{37.5 \mathrm{~cm}}{25.0 \mathrm{~cm}}=-1.5 .
$$

The total magnification is the product of the two magnifications:

$$
m_{\text {total }}=m_{1} m_{2}=(-2.0)(-1.5)=+3.0
$$

The image is erect and three times larger than the object. The image is 6.0 cm tall.
EVALUATE Designing optical instruments involves analyzing combinations of multiple lenses. This problem shows how to combine two lenses. In this case, the first image was formed between the two lenses. In some cases, the first image can be formed to the right of the second lens, creating a virtual object for the second lens. However, even in these more complicated lens systems, our analysis tools will lead to the correct solution.

Practice Problem: Confirm the location of the final image by drawing a principal-ray diagram for the second lens.

## 5: Travel of a camera lens

A camera must be able to image objects from as close as 25 cm away to objects infinitely far away. If a camera has a lens with a focal length of 50 mm , find the travel, or distance the lens must move, in order to create images of objects at any distance.

## Solution

IDENTIFY AND SET UP We will use the thin-lens equation to find the image distances for the two different extreme object distances. By comparing the two, we'll find the travel of the lens.

EXECUTE The thin-lens equation relates the object distance, image distance, and focal length:

$$
\frac{1}{s}+\frac{1}{s^{\prime}}=\frac{1}{f}
$$

For the close object distance, the image distance is

$$
\frac{1}{s^{\prime}}=\frac{1}{f}-\frac{1}{s}=\frac{1}{(5.0 \mathrm{~cm})}-\frac{1}{(25 \mathrm{~cm})}, \quad s^{\prime}=6.25 \mathrm{~cm}
$$

For the infinitely distant object, the image is formed at the focal point, so the image distance is 5.0 cm . The difference in the two distances is the travel:

$$
\Delta s=s_{\text {close }}^{\prime}-s_{\text {far }}^{\prime}=6.25 \mathrm{~cm}-5.0 \mathrm{~cm}=1.25 \mathrm{~cm} .
$$

The lens must be able to travel 1.25 cm , or 12.5 mm , to accommodate all object distances.
EVALUATE This problem illustrates why camera lenses must be adjusted to form clear images.
Practice Problem: Would a lens with a longer focal length require more or less travel? Answer: A longer focal length lens would require a longer travel to accommodate all object distances.

## 6: Designing corrective lenses

Determine the power of the corrective lenses required for (a) a myopic eye with far point at 40.0 cm and (b) a hyperopic eye with near point at 40.0 cm .

## Solution

IDENTIFY We will find the corrective lenses that create images at the near or far points of the eyes with the given properties.

SET UP For a myopic eye, the lens must form a virtual image at the far point of the eye for objects that are far away. Thus, in this case, when the object distance is infinite, the image distance must be -40.0 cm .

For a hyperopic eye, the lens must form a virtual image at the near point of the eye for objects that are close. Thus, in this case, when the object distance is 25.0 cm , the image distance must be -40.0 cm .

We'll use the thin-lens equation to determine the focal lengths for each case. The corrective power is the inverse of the focal length, in meters.

EXECUTE The thin-lens equation relates the object distance, image distance, and focal length:

$$
\frac{1}{s}+\frac{1}{s^{\prime}}=\frac{1}{f}
$$

For the myopic eye, the object distance is infinite and the image distance is -40.0 cm :

$$
\frac{1}{f}=\frac{1}{s}+\frac{1}{s^{\prime}}=\frac{1}{\infty}+\frac{1}{(-0.400 \mathrm{~m})}=-2.5 \text { diopters. }
$$

The myopic eye needs a diverging lens of -2.5 diopters.
For the hyperopic eye, when the object distance is 25.0 cm , the image distance is -40.0 cm :

$$
\frac{1}{f}=\frac{1}{s}+\frac{1}{s^{\prime}}=\frac{1}{(0.25 \mathrm{~m})}+\frac{1}{(-0.400 \mathrm{~m})}=+1.5 \text { diopters. }
$$

The hyperopic eye needs a converging lens of +1.5 diopters.
EVALUATE This problem contrasts the two types of distance vision defects. For each, we use the appropriate corrective lens to create an image at the farthest or closest point the defective eye can image.

## 7: Designing a microscope

You are to design a microscope that consists of a $10.0-\mathrm{mm}$ objective lens and a tube that will hold the lenses 15.0 cm apart. If the microscope is designed to view an object that is placed 3.0 mm beyond the objective lens's focal point and create its final image at infinity, what is the focal length of the eyepiece? What is the magnification of the microscope?

## Solution

IDENTIFY We'll use the design parameters of the microscope to solve the problem.
SET UP For the final image to be at infinity, the first image must coincide with the focal point of the objective. We'll find the location of the first image to determine the focal length of the eyepiece. We can then find the overall magnification by using the microscope formula.

EXECUTE The location of the image created by the objective lens is found from the thin-lens equation. The object distance is the focal length plus 3.0 mm , or 13.0 mm . The image distance is then

$$
\frac{1}{s_{1}^{\prime}}=\frac{1}{f_{1}}-\frac{1}{s_{1}}=\frac{1}{10.0 \mathrm{~mm}}-\frac{1}{(13.0 \mathrm{~mm})}, \quad s_{1}^{\prime}=43.3 \mathrm{~mm} .
$$

The distance between the lenses is 15.0 cm , so the image is formed

$$
s_{2}=15.0 \mathrm{~cm}-4.33 \mathrm{~cm}=10.7 \mathrm{~cm}
$$

from the eyepiece. The eyepiece must have a focal length of 10.7 cm . The magnification of the microscope is

$$
M=\frac{(25 \mathrm{~cm}) s_{1}^{\prime}}{f_{1} f_{2}}=\frac{(25 \mathrm{~cm})(4.3 \mathrm{~cm})}{(1.00 \mathrm{~cm})(10.7 \mathrm{~cm})}=10.0 \times .
$$

The eyepiece must have a focal length of 10.7 cm , and the overall magnification is a factor of 10 .
EVALUATE This problem illustrates the design of a microscope. The objective creates an image inside the microscope tube, and the eyepiece in turn magnifies the objective's image. The eyepiece is positioned so that its image is located at infinity. That is, the focal point of the eyepiece is positioned near the objective's image location.

## 8: Exploring a telescope

An astronomical telescope is made of two lenses separated by 2.00 m . The focal length of the eyepiece is 12.0 cm . What is the angular magnification of the telescope? The moon subtends an angle of approximately $1 / 2^{\circ}$ at the earth. What is the diameter of the moon's image produced by the objective lens of the telescope?

## Solution

IDENTIFY We will use the design features of a telescope to solve the problem.
SET UP Astronomical telescopes are designed to view distant objects. The image from the objective lens is formed at the objective's focal point. The eyepiece then magnifies that image, creating a final image at infinity. This is accomplished by setting the eyepiece's focal point at the focal point of the objective. The distance between the two lenses must be the sum of the focal lengths. We find the focal length of the objective by subtracting the focal length of the eyepiece from the overall length of the telescope. With the focal length of the objective known, we can then find the magnification and diameter of the moon's image.

EXECUTE The focal length of the objective is

$$
f_{1}=200 \mathrm{~cm}-f_{2}=200 \mathrm{~cm}-12 \mathrm{~cm}=188 \mathrm{~cm} .
$$

The angular magnification is then

$$
M=-\frac{f_{1}}{f_{2}}=-\frac{(188 \mathrm{~cm})}{(12 \mathrm{~cm})}=-15.7 \times
$$

The diameter of the moon's image is the magnification times the original diameter:

$$
\left|y^{\prime}\right|=M y=\frac{s^{\prime}}{s} y .
$$

The image is formed at the focal point, so the image distance is the focal length of the objective. The angle the moon subtends without the telescope is the tangent of the moon's diameter divided by its distance, which, for small angles, reduces to

$$
\theta=\tan ^{-1} \frac{y}{s} \approx \frac{y}{s}
$$

Substituting gives

$$
\left|y^{\prime}\right|=\frac{s^{\prime}}{s} y=\theta f=\left(\frac{1}{2}^{\circ}\right)\left(\frac{\pi}{180^{\circ}}\right)(188 \mathrm{~cm})=1.64 \mathrm{~cm}
$$

The diameter of the moon's image is 1.64 cm .
EVALUATE In this problem, we see how a telescope is designed and forms images. A telescope is similar to a microscope, but with the objective's image being formed at the objective's focal point.

## Try It Yourself!

## 1: A concave mirror

A concave spherical refracting surface forms an image of an object placed 1.6 m from the surface in air. The index of refraction of the material is 1.50 and the surface has a radius with an absolute value of 0.80 m . Find the location of the image and its magnification.

## Solution Checkpoints

IDENTIFY AND SET UP Use the relations for refraction at a spherical surface to solve the problem. Carefully assign and evaluate the signs of variables. Use a principal-ray diagram to check your results.

EXECUTE The general formula for finding the location of the image is

$$
\frac{n_{a}}{s}+\frac{n_{b}}{s^{\prime}}=\frac{n_{b}-n_{a}}{R} .
$$

In this case, the object distance is 1.6 m and the object is in air, so $n_{a}=1.0$, and the radius of curvature is negative. (Why?) Substituting gives $s^{\prime}=1.2 \mathrm{~m}$.

The magnification is given by

$$
m=-\frac{n_{a} s^{\prime}}{n_{b} s}
$$

Substituting gives a magnification of +0.5 .
EVALUATE Is your ray diagram consistent with your numeric result? You should check both position and magnification with the diagram.

## 2: Two thin lenses

Two thin lenses with focal lengths $f_{1}=+3.0 \mathrm{~cm}$ and $f_{2}=-5.0 \mathrm{~cm}$ are placed in contact with each other. An object situated 4.0 cm from the lens combination, closer to the lens with the $+3.0=\mathrm{cm}$ focal length. Find the location of the final image.

## Solution Checkpoints

IDENTIFY AND SET UP Use the thin-lens relations to solve the problem. Carefully assign and evaluate the signs of variables. Use a principal-ray diagram to check your results.

EXECUTE The thin lens formula to find the location of the first image is

$$
\frac{1}{s_{1}}+\frac{1}{s_{1}^{\prime}}=\frac{1}{f_{1}}
$$

Substituting the given values, we find that the image due to the first lens is located at 12.0 cm . This image becomes the object for the second lens. The object distance is -12.0 cm , since it is located on the real side of the lens. Use this result in the thin-lens formula for the second lens:

$$
\frac{1}{s_{2}}+\frac{1}{s_{2}^{\prime}}=\frac{1}{f_{2}}
$$

Substituting the given values results in a final image location of -8.57 cm . This is a virtual image located on the side of the first lens.

EVALUATE Does your ray diagram agree with your result? You should have used two principal-ray diagrams to check.

## 3: Designing a microscope

A microscope has an objective lens with focal length $f_{o}=1.6 \mathrm{~cm}$ and an eyepiece lens with focal length $f_{e}=2.5 \mathrm{~cm}$. What is the magnification of the microscope if an object is placed 0.10 cm from the focal point of the objective lens?

## Solution Checkpoints

IDENTIFY AND SET UP Use the magnification equation for a microscope to solve the problem.
EXECUTE The magnification for a microscope is given by

$$
M=\frac{(25 \mathrm{~cm}) s_{1}^{\prime}}{f_{1} f_{2}}
$$

From the thin-lens equation, the image distance $s_{1}^{\prime}$ is found to be 27.2 cm . This yields a magnification of 170 .

EVALUATE Did you remember to include both the objective focal length and the 0.1 cm in calculating the object distance?

## 35 <br> Interference

## Summary

In this chapter, we explore the wave nature of light. Earlier we learned that light is an electromagnetic wave, but we have not yet encountered situations that demonstrate the wave nature of light. In this chapter, we will see how light can interfere, one property that confirms the wave nature of light. The study of phenomena associated with the wave nature of light is called physical optics.

## Objectives

After studying this chapter, you will understand

- How waves combine or interfere in space.
- The definition of coherent light sources, interference, and phase.
- How constructive and destructive interference leads to interference patterns.
- How to analyze two-source interference and thin-film interference.
- The intensity of light at various points in an interference pattern.
- How to use interference to measure tiny distances.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Interference | Light emanating from two coherent sources and overlapping in a region <br> forms an interference pattern. Constructive interference occurs when two <br> waves arrive at a point in phase. Destructive interference occurs when two <br> waves arrive at a point exactly half a cycle out of phase. |
| Two-Source Interference | For two light sources in phase, constructive interference occurs at points <br> where the path-length difference is zero or an integral number of wave- <br> lengths; destructive interference occurs when the path-length difference is a <br> half-integral number of wavelengths. For two light sources located a distance <br> $d$ apart, the condition for constructive interference is |
| $d \sin \theta=m \lambda, \quad m=0, \pm 1, \pm 2, \ldots$, |  |

and the condition for destructive interference is

$$
d \sin \theta=\left(m+\frac{1}{2}\right) \lambda, \quad m=0, \pm 1, \pm 2, \ldots
$$

## Intensity in Interference Patterns

## Thin-Film Interference

When light is reflected from both sides of a thin film of thickness $t$, interference occurs. Constructive interference occurs when

$$
2 t=m \lambda, \quad m=0,1,2, \ldots,
$$

in cases where both waves or neither wave is phase shifted by a half-cycle and where $\lambda$ is the wavelength in the thin film. If a half-cycle phase shift occurs at only one surface, the given equation is the condition for destructive interference. A half-cycle phase shift occurs during reflection whenever the index of refraction of the reflecting material is greater than that of the medium the wave is traveling through.

## Conceptual Questions

## 1: Two-source interference

Coherent light passes through two thin slits and produces an interference pattern. If one slit is covered with a glass plate that introduces a half-cycle phase shift between the two emerging waves, how is the interference pattern altered?

## Solution

IDENTIFY, SET UP, AND EXECUTE At the center of the interference pattern, both waves travel the same distance, but are out of phase, creating a dark spot due to destructive interference. At locations where there was constructive interference, there is now destructive interference due to the addition of the glass plate. At locations where there was destructive interference, constructive interference occurs
due to the glass plate. We conclude that the pattern of alternating bright and dark locations reverses after the glass plate is added.

EVALUATE The addition of the glass plate does not change the path differences that create the interference pattern, but rather only exchanges the pattern of bright and dark locations.

## 2: Newton's rings

When a planoconvex lens is placed on top of a plane glass surface, circular interference fringes known as Newton's rings appear. Why is the center of this pattern black?

## Solution

IDENTIFY, SET UP, AND EXECUTE Figure 35.1 shows an enlarged view of the situation. The interference pattern is created by light that reflects off of the bottom surface of the lens interfering with light that passes through the lens and reflects off of the plane glass surface. Light that reflects off of the bottom lens surface does not undergo a phase shift, since the index of refraction of the air is less than that of glass. Light reflecting off of the plane glass surface does undergo a half-cycle phase shift, since the index of refraction of the glass is greater than that of air. The relative phase of the two interfering light waves is different by one half cycle. At the center of the lens, there is essentially no path difference, so the two waves interfere destructively, creating a dark spot.


Figure 35.1 Question 2.
EVALUATE The analysis of phase differences is the first step in the investigation of thin-film interference. We could continue our investigation of this problem and identify regions of bright and dark fringes by determining the conditions for constructive and destructive interference.

## 3: Interference pattern on a slide

A thin film of oil is sprayed on a vertically oriented, smooth, flat glass slide. After the spraying, some of the oil moves to the bottom of the slide, leaving a wedge-shaped coating of oil on the slide (Figure 35.2). Sketch the interference pattern that occurs when the slide is viewed from the oil side with monochromatic light. The index of refraction of the oil is 1.3 and of the glass is 1.4.


Figure 35.2 Question 3.

## Solution

IDENTIFY, SET UP, AND EXECUTE A thin-film interference pattern will occur because the light that reflects off of the oil surface interferes with the light that reflects off of the glass surface. The pattern will consist of bright and dark regions corresponding to constructive and destructive interference, respectively.

The wedge-shaped oil coating will exhibit alternating bright and dark horizontal bands. The bands are horizontal because, by assumption, the thickness of the oil is uniform across the slide.

We also learn that the top band across the slide is a bright band. This is because both the light reflected off of the oil surface and the light reflected off of the glass slide have a half-cycle phase shift due to their reflecting off materials with higher indexes of refraction. There is no net phase difference between the two reflected waves and there is no path difference at the top of the slide. The two waves interfere constructively at the top, creating a bright fringe. The pattern is shown in Figure 35.3.


Figure 35.3 Question 3.
EVALUATE This problem contrasts with the previous one, in which there was a dark fringe at the center. Analyzing the reflected light waves for phase changes is crucial in thin-film problems.

## Problems

## 1:Separation of interference lines

Monochromatic light of wavelength 550 nm is incident upon two slits separated by 0.15 mm . The resulting interference pattern is observed on a screen 2.0 m away. What is the linear separation between the third bright line and the fifth dark line on the screen?

## Solution

IDENTIFY We will use the two-slit interference relations to find the distance between the two lines on the screen.

SET UP The linear separation on the screen is found by first finding the angular separation. We'll use the constructive and destructive two-source interference equations to find the angular locations of the two lines and then apply trigonometry to find their positions on the screen. Bright lines correspond to constructive interference, dark lines to destructive interference. The difference between the two positions is the linear separation.

EXECUTE The condition for constructive interference is

$$
d \sin \theta=m \lambda, \quad m=0, \pm 1, \pm 2, \ldots
$$

We take the third bright line to be the third bright line past the central maximum. The third bright line will be at an angle given by

$$
\sin \theta \approx \theta=\frac{3 \lambda}{d}=\frac{3\left(550 \times 10^{-9} \mathrm{~m}\right)}{\left(1.5 \times 10^{-4} \mathrm{~m}\right)}=0.0110 \mathrm{rad}
$$

where we replaced $\sin \theta$ with $\theta$ in radians for small $\theta$. The third bright line will be located at a distance from the central maximum given by

$$
y_{3}=R \tan \theta \approx R \theta=(2.0 \mathrm{~m})(0.0110 \mathrm{rad})=2.2 \mathrm{~cm}
$$

where $R$ is the distance to the screen and we again used the small-angle approximation.
The condition for destructive interference is

$$
d \sin \theta=\left(m+\frac{1}{2}\right) \lambda, \quad m=0, \pm 1, \pm 2, \ldots
$$

The fifth dark line corresponds to $m=4$, which will be at an angle given by

$$
\sin \theta \approx \theta=\frac{\left(4+\frac{1}{2}\right) \lambda}{d}=\frac{4.5\left(550 \times 10^{-9} \mathrm{~m}\right)}{\left(1.5 \times 10^{-4} \mathrm{~m}\right)}=0.0165 \mathrm{rad}
$$

where we again replaced $\sin \theta$ with $\theta$ in radians for small $\theta$. The fifth dark line will be located at a distance from the central maximum given by

$$
y_{5}=R \tan \theta \approx R \theta=(2.0 \mathrm{~m})(0.0165 \mathrm{rad})=3.3 \mathrm{~cm}
$$

The difference in position on the screen between the third bright and fifth dark lines is the difference

$$
\Delta y=y_{5}-y_{3}=(3.3 \mathrm{~cm})-(2.2 \mathrm{~cm})=1.1 \mathrm{~cm}
$$

The two lines are separated by 1.1 cm .

EVALUATE This problem illustrates how to interpret and solve two-source interference problems. We see that we must carefully interpret the statement of the problem and carefully count the interference lines to find the solution.

## 2: Thin-film interference pattern

The flat surface of a glass block with a refractive index of 1.53 is coated with a thin, ransparent film of refractive index 1.63 and thickness 630.0 nm . Assuming that the visible spectrum extends from 400.0 to 700.0 nm , what visible wavelength(s) of light will appear intensified in the reflected beam?

## Solution

IDENTIFY Light will reflect off the top of the film and from the film-glass interface, creating a thinfilm interference pattern. The target variables are the visible wavelengths that are seen from above.

SET UP The wavelengths that will be intensified will be those undergoing constructive interference. We will solve the problem by varying $n$ and identifying wavelengths in the visible range that interfere constructively in the layer of film.

EXECUTE To determine the proper thin-film interference equation to use, we first check the indexes of refraction of the reflecting materials. At the air-film interface, there is a half-cycle phase change because the film has a higher index of refraction than that of air. At the film-glass interface, there is no phase change, because the glass has a lower index of refraction than that of the film. Therefore, there is a net phase change of one half cycle, and the proper constructive interference equation is

$$
2 t=\left(m+\frac{1}{2}\right) \lambda .
$$

Recall that the wavelength in the equation is the light's wavelength in the medium. The wavelength of light in air is the wavelength of light in the medium times the index of refraction of the medium:

$$
\lambda_{0}=\lambda n=\frac{2 t n}{\left(m+\frac{1}{2}\right)}
$$

We now vary $m$ and find all wavelengths in the visible part of the spectrum:

$$
\begin{array}{ll}
\lambda_{0}=\lambda n=\frac{2 t n}{\left(\frac{1}{2}\right)}=\frac{2(630.0 \mathrm{~nm})(1.63)}{\left(\frac{1}{2}\right)}=4108 \mathrm{~nm} & (m=0), \\
\lambda_{0}=\lambda n=\frac{2 \mathrm{tn}}{\left(1+\frac{1}{2}\right)}=\frac{2(630.0 \mathrm{~nm})(1.63)}{\left(\frac{3}{2}\right)}=1369 \mathrm{~nm} & (m=1), \\
\lambda_{0}=\lambda n=\frac{2 t n}{\left(2+\frac{1}{2}\right)}=\frac{2(630.0 \mathrm{~nm})(1.63)}{\left(\frac{5}{2}\right)}=822 \mathrm{~nm} & (m=2), \\
\lambda_{0}=\lambda n=\frac{2 t n}{\left(3+\frac{1}{2}\right)}=\frac{2(630.0 \mathrm{~nm})(1.63)}{\left(\frac{7}{2}\right)}=587 \mathrm{~nm} & (m=3), \\
\lambda_{0}=\lambda n=\frac{2 t n}{\left(4+\frac{1}{2}\right)}=\frac{2(630.0 \mathrm{~nm})(1.63)}{\left(\frac{9}{2}\right)}=456 \mathrm{~nm} & (m=4), \\
\lambda_{0}=\lambda n=\frac{2 t n}{\left(5+\frac{1}{2}\right)}=\frac{2(630.0 \mathrm{~nm})(1.63)}{\left(\frac{11}{2}\right)}=373 \mathrm{~nm} \quad(m=5)
\end{array}
$$

We see that only for $m=3$ and $m=4$ do we get wavelengths in the visible spectrum. We conclude that light of wavelengths 587 nm and 456 nm is intensified in the reflected beam.

EVALUATE When working with thin-film interference, we must determine the type of interference (constructive or destructive), as well as the phase changes at the reflecting interfaces. Once these are established, we proceed directly to the solution.

Practice Problem: What wavelengths would appear intensified if a film with an index of refraction of 1.44 is substituted in this problem? Answer: $605 \mathrm{~nm}, 454 \mathrm{~nm}$.

## 3: Nonglare coating

A thin layer of water ( $n=1.33$ ) on top of a layer of glass ( $n=1.5$ ) produces a nonglare optical coating viewed normally by light of 600 nm wavelength. What is the minimum thickness of the water layer?

## Solution

IDENTIFY We will use the condition for thin-film interference to solve this problem. Light reflecting off of the top of the water will interfere with light that refracts through the water and reflects off of the water-glass interface.

SET UP To form a nonglare coating, the light must interfere destructively upon reflection from the surface. Light that reflects from the top of the water layer will have a half-cycle phase shift, since the index of refraction of water is greater than that of air. Light that refracts through the water and reflects off of the top of the glass also has a half-cycle phase shift, because the index of refraction of glass is greater than that of water. The combination of the two kinds of light results in no relative phase change between the interfering waves.

EXECUTE The condition for destructive interference is then

$$
2 t=\left(m+\frac{1}{2}\right) \lambda
$$

Recall that the wavelength in the equation is the wavelength of light in the medium. The wavelength of the light in air is the wavelength of the light in the medium times the index of refraction of the medium. Thus,

$$
t=\frac{\lambda_{0}\left(m+\frac{1}{2}\right)}{2 n}
$$

The minimum thickness of water corresponds to the condition where $m=0$, giving

$$
t=\frac{\lambda_{0}\left(\frac{1}{2}\right)}{2 n}=\frac{(600 \mathrm{~nm})}{4(1.33)}=113 \mathrm{~nm} .
$$

A water layer 113 nm thick will lead to destructive interference for $600-\mathrm{nm}$ light.
EVALUATE To make a reflective coating, you could double the thickness of the water layer. Of course, using water as a nonreflective coating isn't the best solution, because water evaporates, so the nonreflective properties would last only a limited time.

## Try It Yourself!

## 1: Interference fringes

Calculate the separation between interference fringes obtained in Young's experiment when 550-nm light is shined upon two slits separated by 0.22 mm and viewed on a screen that is 3.0 m away from the slits.

## Solution Checkpoints

IDENTIFY AND SET UP Sketch the problem. What is the condition for constructive interference?
EXECUTE The condition for constructive interference is given by

$$
d \sin \theta=m \lambda, \text { for } m=0,1,2, \ldots
$$

The distance to the screen is large, so

$$
\sin \theta \cong \tan \theta=\frac{y}{R},
$$

where $y$ is the distance from the axis to the constructive fringe and $R$ is the distance between the slits and the screen. Bright fringes appear for values of $y_{n}$ equal to

$$
y_{n}=n \frac{\lambda R}{D} .
$$

Evaluating the expression gives a distance of 7.5 mm between fringes.
EVALUATE What is the distance between dark fringes?

## 2: Soap film

A soap film with index of refraction equal to 1.35 is viewed normally. Calculate the minimum thickness of the film that will give constructive interference for yellow light of wavelength 575 nm .

## Solution Checkpoints

IDENTIFY AND SET UP Sketch the problem. What two rays interfere? Is a half-cycle phase shift introduced for one or both reflected light rays?

EXECUTE Light reflected from the top surface of the soap filmhas a half-cycle phase shift introduced. Constructive interference occurs when

$$
2 n t=\left(m+\frac{1}{2}\right) \lambda
$$

The minimum thickness corresponds to $m=0$. The minimum thickness for constructive interference is 107 nm .

EVALUATE Why was the factor $n$ introduced in the path difference?

## 36 <br> Diffraction

## Summary

In this chapter, we continue our exploration of the wave nature of light. Here, we'll examine how light bends when it passes an edge or corner. This bending, or diffraction, results from the wave nature of light. We will investigate phenomena such as single-slit diffraction, diffraction gratings, X-ray diffraction, resolving power, and holography.

## Objectives

After studying this chapter, you will understand

- How a coherent light source shining on an edge creates a diffraction pattern.
- The diffraction pattern formed by light passing through a narrow slit.
- How to calculate intensity as a function of angle in single-slit diffraction.
- The diffraction pattern created by a diffraction grating.
- How to probe atomic arrangements of crystals through X-ray diffraction.
- How to apply Rayleigh's criterion to find the limit of resolution of optical systems.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Diffraction | Diffraction occurs when light passes through an aperture or around an edge. <br> Fraunhofer diffraction occurs when the source and observer are far from the <br> obstructing surface and the outgoing rays can be considered to be parallel. <br> Fresnel diffraction occurs when the observer or source is close to the <br> obstructing surface. |
| Single-Slit Diffraction | For a single narrow slit of width $a$ illuminated with monochromatic light, <br> destructive interference occurs at angles satisfying the relationship |
| $\qquad \sin \theta=\frac{m \lambda}{a}, \quad m= \pm 1, \pm 2, \pm 3, \ldots$ |  |

The intensity of the diffraction patter is given by

$$
I=I_{0}\left\{\frac{\sin [\pi a(\sin \theta) / \lambda]}{\pi a(\sin \theta) / \lambda}\right\}^{2}
$$

| Diffraction Grating | A diffraction grating consists of many thin parallel slits spaced a distance $d$ <br> apart. Maximum intensity occurs when |
| :--- | :--- |
| $\qquad$$d \sin \theta=m \lambda, \quad m=0, \pm 1, \pm 2, \ldots$ |  |
| X-ray Diffraction | The pattern is similar to that seen with two-source interference, but with the <br> maxima very sharp and narrow. |
| A crystal serves as a three-dimensional diffraction grating for X-rays with <br> wavelengths of the same order of magnitude as the spacing of the atoms in <br> the crystal. For a set of crystal planes spaced a distance $d$ apart, maximum <br> intensity occurs when |  |
| $\quad 2 d \sin \theta=m \lambda, \quad m=0,1,2, \ldots$. |  |

## Circular Apertures and <br> Resolving Power

$$
2 d \sin \theta=m \lambda, \quad m=0,1,2, \ldots
$$

This condition is referred to as the Bragg condition.
A circular aperture of diameter $D$ creates a diffraction pattern consisting of a central bright spot, called the Airy disk, surrounded by a series of concentric dark and bright rings. The limit of resolution is defined by the angular size of the first dark ring, given by

$$
\sin \theta_{1}=\frac{1.22 \lambda}{D}
$$

Rayleigh's criterion states that two point objects are just barely resolved when their angular separation is at the limit of resolution.

## Conceptual Questions

## 1: Light as a particle

If light were composed of tiny particles, like grains of sand, would a diffraction pattern be formed when the light went through a thin slit?

## Solution

IDENTIFY, SET UP, AND EXECUTE A diffraction pattern is created by waves interfering at the diffractive edge or aperture. Imagine many adjacent waves at the thin slit interfering with each other. These adjacent waves will produce a diffraction pattern.

Particles neither interfere nor create diffraction patterns. If light were made of tiny particles, then light would not exhibit wave behavior and create diffraction patterns. Since light does make diffraction patterns, it must have a wave nature.

EVALUATE This problem helps illustrate the difference between wave and particle phenomena. Waves and particles have distinct behaviors; by finding evidence of these behaviors, we identify objects as waves or particles.

In Chapter 38, we shall see that light can also exhibit particle behavior. Wave-particle duality is a fundamental component of modern physics.

## 2: Can sound interfere and diffract?

Can sound waves exhibit interference and diffraction?

## Solution

IDENTIFY, SET UP, AND EXECUTE Sound waves interfere, although it can be difficult to observe their interference. You may have noticed that the sound level at a concert changes as you move around in front of the speakers. The changing level is due to interference.

Sound waves can certainly diffract when they go around a corner. You experience sound-wave diffraction whenever you listen to a conversation occurring around a corner or hear music from down the hall. If sound waves didn't diffract, then you would be able to hear only conversations that occur within view of your location.

EVALUATE Wave phenomena, including interference and diffraction, occur with all forms of waves, not just light waves. Familiarity with the various forms of wave interaction helps us build an understanding of electromagnetic wave interactions.

## Problems

## 1: Wavelength of light

Monochromatic light from a laser is incident on a slit 0.550 mm wide. On a screen 1.50 m away from the slit, the distance between the first minima on either side of the central maximum is 2.35 mm . Determine the wavelength of light emitted by the laser.

## Solution

IDENTIFY We will use our knowledge of single-slit diffraction to solve this problem.
SET UP Diffractive patterns are symmetric, so the distance from the central maximum to either of the first minima is half the distance between the first minima on either side of the maximum. The distance to the screen is much larger than the separation of the minima, so the angle is very small and we replace $\sin \theta$ with $y / 2 R$.

EXECUTE The condition for dark fringes is

$$
\sin \theta=\frac{m \lambda}{a} .
$$

Substituting for the sine gives

$$
\frac{y}{2 R}=\frac{m \lambda}{a}
$$

where $y$ is the distance between the two minima on either side of the central maximum and $R$ is the distance to the screen. Note that we have divided the distance to the screen by $1 / 2$, as that is the distance between the central maximum and the first minimum. The first minimum corresponds to $m=1$. Rearranging terms and solving for the wavelength gives

$$
\lambda=\frac{a y}{m 2 R}=\frac{(0.550 \mathrm{~mm})(2.35 \mathrm{~mm})}{2(1)(1.50 \mathrm{~m})}=431 \mathrm{~nm} .
$$

The laser emits 431-nm light.
EVALUATE This problem illustrates how we can use macroscopic objects to measure very small wavelength of light. Both interference and diffraction are used to measure small quantities throughout the field of physics.

## 2: Intensity from a single slit

For the single-slit diffraction pattern, find the ratio of the intensity at the first maximum to that at the center.

## Solution

IDENTIFY We will use the diffraction intensity equation to solve the problem.
SET UP The position of the first maximum is not easy to determine in diffraction, because the intensity function is not linear. The location is not where the sine function in the numerator is zero, since there is a second sine function in the denominator. To solve the problem, we will use an approximate value of the maximum to estimate the ratio of the intensities.

EXECUTE The intensity due to a single slit is given by

$$
I=I_{0}\left\{\frac{\sin [\pi a(\sin \theta) / \lambda]}{\pi a(\sin \theta) / \lambda}\right\}^{2}
$$

We set the argument of the numerator to $3 \pi / 2$ to find the first maximum; that is, we set

$$
\frac{\pi a(\sin \theta)}{\lambda}=\frac{3 \pi}{2}
$$

The resulting intensity is then

$$
I=I_{0}\left\{\frac{\sin [3 \pi / 2]}{3 \pi / 2}\right\}^{2}=I_{0}\left\{\frac{4}{9 \pi^{2}}\right\}=0.045 I_{0}
$$

We see that the first maximum has only about $4.5 \%$ of the intensity of the central maximum.
EVALUATE Careful analysis shows that the true intensity at the first maximum is $4.72 \%$ of the intensity of the central maximum. We see that our approximation was reasonable in this problem. We also see how rapidly the maximum intensity decreases as one moves away from the central maximum.

## 3: The diffraction grating

Light of 665 nm produces a third-order bright fringe band at an angle of $75.0^{\circ}$ after passing through a diffraction grating. Find the locations of the first-, second-, and fourth-order bright bands.

## Solution

IDENTIFY We will use the diffraction grating equation to solve the problem.
SET UP To find the locations of the bright bands, we must first find the slit separation of the grating. We will use the information provided for the third-order bright band to find the slit separation. We then proceed to find the angular positions of the bands.

EXECUTE The condition for bright bands for a diffraction grating is

$$
d \sin \theta=m \lambda, \quad m=0, \pm 1, \pm 2, \ldots
$$

Knowing the angle allows us to find the slit separation:

$$
d=\frac{m \lambda}{\sin \theta}=\frac{(3)(665 \mathrm{~nm})}{\sin 75.0^{\circ}}=2065 \mathrm{~nm}
$$

The slits are separated by 2065 nm , or equivalently, there are 4842 lines per millimeter. We can now find the angles for the other orders. The first-order bright band is located at

$$
\theta=\sin ^{-1} \frac{m \lambda}{d}=\sin ^{-1} \frac{(1) 665 \mathrm{~nm}}{2065 \mathrm{~nm}}=18.8^{\circ}
$$

The second-order bright band is located at

$$
\theta=\sin ^{-1} \frac{m \lambda}{d}=\sin ^{-1} \frac{(2) 665 \mathrm{~nm}}{2065 \mathrm{~nm}}=40.0^{\circ}
$$

The fourth-order bright band is located where

$$
\sin \theta=\frac{m \lambda}{d}=\frac{(4) 665 \mathrm{~nm}}{2065 \mathrm{~nm}}=1.29
$$

There is no angle whose sine is 1.29 , so the fourth-order bright band is not visible. The first- and secondorder bright bands are located at $18.8^{\circ}$ and $40.0^{\circ}$, respectively.

EVALUATE Diffraction gratings allow us to determine the wavelength of light accurately. One can start with a known wavelength to calibrate the grating and then use the grating to find the wavelength of an unknown light.

## 4: Bragg scattering

An X-ray photon with a wavelength of 0.124 nm is incident on a single crystal of sodium chloride $(\mathrm{NaCl})$. If constructive interference is observed at $12.7^{\circ}$, what is the minimum lattice spacing of the crystal?

## Solution

IDENTIFY We will use the Bragg condition to find the lattice spacing of the crystal.

SET UP The minimum lattice spacing correlates to the first-order constructive interference condition. This means that we will need to use $m=1$ in the Bragg relation.

EXECUTE The Bragg condition for constructive interference is given by

$$
2 d \sin \theta=m \lambda
$$

Rearranging terms to solve for $d$ gives

$$
d=\frac{m \lambda}{2 \sin \theta}
$$

Substituting $m=1$ and solving yields

$$
d=\frac{m \lambda}{2 \sin \theta}=\frac{(1)(0.124 \mathrm{~nm})}{2 \sin 12.7^{\circ}}=0.282 \mathrm{~nm} .
$$

The minimum lattice spacing for the NaCl crystal is 0.282 nm .
EVALUATE This problem illustrates how we can use X-ray diffraction to measure the structure of crystals. Because the lattice spacing is much smaller than the wavelength of visible light, X-rays must be used to probe the atomic spacing.

## 5: Designing a telescope

Pluto's moon Charon was discovered in 1978. What minimum-diameter reflecting telescope was needed to discover Charon? Take the distance to Charon to be $5.9 \times 10^{12} \mathrm{~m}$ and the diameter of Charon to be $1.2 \times 10^{6} \mathrm{~m}$.

## Solution

IDENTIFY We will use Rayleigh's criterion to find the resolving power of a telescope.
SET UP To discover Charon, one must be able to resolve it in the telescope. We will use Rayleigh's criterion to find the minimum diameter of the telescope, and we will use the average wavelength of light from the sun ( $\lambda=550 \mathrm{~nm}$ ) in our calculations.

EXECUTE Rayleigh's criterion gives the limit of resolution for an optical instrument:

$$
\theta_{\mathrm{res}}=1.22 \frac{\lambda}{D}
$$

The angular size of Charon is

$$
\theta=\frac{s}{r}=\frac{1.2 \times 10^{6} \mathrm{~m}}{5.9 \times 10^{12} \mathrm{~m}}=2.03 \times 10^{-7} \mathrm{rad}
$$

where we used the diameter and the distance to the planet. The minimum diameter of the telescope lens is

$$
D=1.22 \frac{\lambda}{\theta}=1.22 \frac{550 \times 10^{-9} \mathrm{~m}}{2.03 \times 10^{-7} \mathrm{rad}}=3.30 \mathrm{~m}
$$

The minimum diameter of the telescope lens needed to resolve Charon is 3.30 m .

EVALUATE We see that Charon's small angular size requires a large-diameter telescope to observe it. Also, the amount of light reflecting off the satellite is very small. Both of these challenges led to the discovery of Charon only rather recently.

## Try It Yourself!

## 1: A diffraction grating

Light of wavelength 550 nm strikes a diffraction grating. The first maximum away from the central bright spot occurs at $15.96^{\circ}$. (a) Find the value of the spacing in the grating. (b) At what angle should one look to find the second maximum for light of wavelength 500 nm ?

## Solution Checkpoints

IDENTIFY AND SET UP Which equation should be used? What value of $m$ should the first and second maxima correspond to?

EXECUTE (a) The spacing in the grating is found from

$$
d=\frac{m \lambda}{\sin \theta} .
$$

Using $m=1$, you should find that the spacing in the grating is $2.00 \times 10^{-6} \mathrm{~m}$.
(b) The angle for the second maximum for $500-\mathrm{nm}$ light is

$$
\theta=\sin ^{-1}\left(\frac{2 \lambda}{d}\right)
$$

The angle is $30.0^{\circ}$.
EVALUATE Can you find the angle for the third maximum with light of wavelength 400 nm ?

## 2: Viewing a binary star system

Blue light ( 486 nm ) from a binary star system located 100 light-years from earth is observed through a telescope. What must the minimum diameter of the primary lens be in order to determine that there are two stars in the binary system if the two stars are located $5.0 \times 10^{-5}$ light-year apart?

## Solution Checkpoints

IDENTIFY AND SET UP Use Rayleigh's criterion to solve the equation for the minimum diameter.
EXECUTE The angular size of separation of the two stars is

$$
\theta=\frac{s}{r}=\frac{5.0 \times 10^{-5} \mathrm{LY}}{100 \mathrm{LY}}=5 \times 10^{-7} \mathrm{rad}
$$

The minimum diameter of the telescope lens is

$$
D=1.22 \frac{\lambda}{\theta}=1.22 \frac{486 \times 10^{-9} \mathrm{~m}}{5.0 \times 10^{-7} \mathrm{rad}}=1.19 \mathrm{~m}
$$

EVALUATE We see that we are able to resolve very distant binary star systems if we have a telescope lens of sufficiently large diameter.

## Relativity

## Summary

In this chapter, we study the special theory of relativity introduced by Einstein in 1905. The theory is based on two postulates: The laws of physics are the same in every inertial reference frame, and the speed of light in vacuum is the same in all reference frames. These simple postulates have far-reaching implications. We will explore how two observers moving relative to each other might not measure the same time or length, how two events might not be simultaneous in all reference frames, and how we must modify the principles of momentum and energy to suit special relativity. You will find that your intuition is often unreliable in these types of situations. We will learn to develop new tools to analyze relativity problems.

## Objectives

After studying this chapter, you will understand

- The two postulates of Einstein's special theory of relativity.
- How to identify inertial frames of reference.
- When to apply the special theory of relativity.
- The definition of proper time and proper length.
- How to apply time dilation and length contraction to various problems.
- How to use Lorentz transformations to find quantities in various reference frames.
- The concepts of, and how to apply, relativistic momentum and relativistic energy.
- The concept of rest energy.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Relativity and Simultaneity | All fundamental laws of physics have the same form in inertial reference <br> frames. The speed of light in vacuum is the same in all inertial frames of ref- <br> erence and is independent of the motion of the source. Simultaneity is not <br> absolute: Two events occurring simultaneously in one frame might not appear <br> simultaneous in a second frame. |
| Time Dilation | The proper time $\Delta t_{0}$ is the time interval between two events that occur at the <br> same spatial point in a frame of reference. If this frame moves with a constant <br> velocity $u$ relative to a second frame, the time interval $\Delta t$ between events <br> observed in the second frame is longer: |

$$
\begin{aligned}
& \Delta t=\frac{\Delta t_{0}}{\sqrt{1-u^{2} / c^{2}}}=\gamma \Delta t_{0} \\
& \gamma=\frac{1}{\sqrt{1-u^{2} / c^{2}}}
\end{aligned}
$$

This effect is known as time dilation.

## Length Contraction

\author{

## Lorentz Transformations

}

## Doppler Effect for

 Electromagnetic WavesThe proper length $l_{0}$ is the distance between two points at rest in a frame of reference. If this frame moves with a constant velocity $u$ relative to a second frame, the distance $l$ measured parallel to the frame's velocity in the second frame is shorter:

$$
l=l_{0} \sqrt{1-u^{2} / c^{2}}=\frac{l_{0}}{\gamma} .
$$

This effect is known as length contraction.
Lorentz transformations relate the coordinates and time of an event in an inertial coordinate system $S$ to the coordinates and time of a second inertial frame $S^{\prime}$ moving with constant velocity $u$ relative to the first frame. The Lorentz transformations are

$$
\begin{aligned}
x^{\prime} & =\frac{x-u t}{\sqrt{1-u^{2} / c^{2}}}=\gamma(x-u t), \quad y^{\prime}=y, \quad z^{\prime}=z \\
t^{\prime} & =\frac{t-u x / c^{2}}{\sqrt{1-u^{2} / c^{2}}}=\gamma\left(t-u x / c^{2}\right) .
\end{aligned}
$$

For one-dimensional motion, the velocities in the two systems are related by

$$
\begin{aligned}
& v_{x}^{\prime}=\frac{v_{x}-u}{1-u v_{x} / c^{2}} \\
& v_{x}=\frac{v_{x}^{\prime}+u}{1+u v_{x}^{\prime} / c^{2}}
\end{aligned}
$$

The Doppler effect is a shift in the frequency of light from a source that is moving relative to an observer. For a source moving towards an observer at speed $u$, the frequency shift is given by

$$
f=\sqrt{\frac{c+u}{c-u}} f_{0} .
$$

Relativistic Momentum and Energy

The relativistic momentum and energy are respectively given by

$$
\vec{p}=\frac{m \vec{v}}{\sqrt{1-v^{2} / c^{2}}}=\gamma m \vec{v}, \quad K=\frac{m c^{2}}{\sqrt{1-v^{2} / c^{2}}}-m c^{2}=(\gamma-1) m c^{2}
$$

The total energy of a particle is the sum of its kinetic energy and its rest mass energy $E_{\text {rest }}=m c^{2}$, given by

$$
E=K+m c^{2}=\frac{m c^{2}}{\sqrt{1-v^{2} / c^{2}}}=\gamma m c^{2}
$$

The total energy can also be expressed as a function of momentum and rest mass:

$$
E^{2}=\left(m c^{2}\right)^{2}+(p c)^{2}
$$

## Conceptual Questions

## 1: When does relativity become important?

A train is speeding past a platform. At what speed must the train be traveling for the proper time measured on the train to differ from the time measured on the platform by $0.1 \%$ ?

## Solution

IDENTIFY, SET UP, AND EXECUTE The time dilation relation will solve this problem. The proper time $\left(\Delta t_{0}\right)$ on the train will transform to the time measured on the platform $(\Delta t)$ by

$$
\Delta t=\frac{\Delta t_{0}}{\sqrt{1-v^{2} / c^{2}}}
$$

We want to find the velocity corresponding to a $0.1 \%$ difference in time, so we set $\Delta t$ to $1.001 \Delta t_{0}$ :

$$
\Delta t=1.001 \Delta t_{0}=\frac{\Delta t_{0}}{\sqrt{1-v^{2} / c^{2}}}
$$

Solving gives

$$
v=\sqrt{1-\left(\frac{1}{1.001}\right)^{2}} c=0.045 c
$$

A $0.1 \%$ time difference will occur when the train is moving at $0.045 c$, or $1.3 \times 10^{7} \mathrm{~m} / \mathrm{s}$.
EVALUATE For a very small relativistic change, the train must be moving at over 13 million meters per second, or 48 million kph . It is highly unlikely that we will see such effects on any train we will ever travel on.

## 2: Proper time and length

You are traveling on a spacecraft moving at $0.1 c$ past a space station. If you are holding your physics textbook and flashing LED, who measures the proper length of the book and the proper time interval for the flashing LED, you or an observer on the space station?

## Solution

IDENTIFY, SET UP, AND EXECUTE The proper length is measured in a frame in which the book is at rest. Therefore, you measure the proper length of your physics textbook. The proper time is also measured with respect to a frame at rest. Again, you measure the proper time for the flashing LED.

EVALUATE Identifying the correct frame is of ten most challenging step in a relativity problem. Here, the observer on the station would measure a length-contracted physics textbook and a timedilated flashing LED.

## 3: Increasing energy

An electron is traveling at a speed of $0.95 c$. Can its energy be increased by more than $5 \%$ ? by more than $25 \%$ ? by more than $500 \%$ ?

## Solution

IDENTIFY, SET UP, AND EXECUTE Relativistic energy for the electron is given by

$$
E=\frac{m c^{2}}{\sqrt{1-v^{2} / c^{2}}} .
$$

As $v$ gets closer and closer to $c$, the denominator gets closer to zero and the energy increases to infinity. For this reason, there is no limit on the maximum energy of the electron.

EVALUATE For the electron to achieve an increase in energy of $500 \%$, it will need to travel at 0.998 c.

## Problems

## 1 : Marley and Cassie

Marley, a twin, takes a round-trip journey to a distant star 12 light years from earth at a speed of $0.93 c$. Cassie, the second twin, remains on earth. What is the age difference between Marley and Cassie when Marley returns?

## Solution

IDENTIFY We need to determine the proper time and use time dilation to solve the problem.
SET UP We need to find the time taken for the journey, as observed by both twins. The elapsed time for Cassie, the twin on earth, will be the distance of the journey divided by the speed. The elapsed time for Marley, the traveling twin, will be the elapsed time on earth transformed to the proper time of the traveler.

EXECUTE The elapsed time for the twin on earth (Cassie) is the distance traveled divided by the speed:

$$
\Delta t=\frac{c(24 \text { years })}{0.93 c}=25.8 \text { years }
$$

Note that the round-trip journey takes 25.8 years. The time in the traveling twin's (Marley's) reference frame is the proper time, since she is the observer at rest. We find this time interval from the time dilation equation:

$$
\begin{aligned}
\Delta t & =\frac{\Delta t_{0}}{\sqrt{1-u^{2} / c^{2}}} \\
\Delta t_{0} & =\Delta t \sqrt{1-u^{2} / c^{2}}=(25.8 \text { years }) \sqrt{1-(0.93 c)^{2} / c^{2}}=9.48 \text { years }
\end{aligned}
$$

Marley ages 9.48 years, while Cassie ages 25.8 years. The difference between Marley's and Cassie's ages will be 16.3 years after Marley returns.

EVALUATE This is the classic twin paradox. We see that Marley, the traveling twin, ages much more slowly than Cassie, the stationary twin. Critical to this problem was identifying which frame was associated with the proper time.

CAUTION Find the Proper Time Carefully: Determining the proper time is critical in time dilation problems. The proper time is the time measured in a frame that is moving with the clock.

## 2: Moving spacecraft

A spacecraft moves past the earth at a speed of 0.850 c. A student on earth measures the length of the moving spacecraft to be 97.0 m . How long does the spacecraft appear to the crew on the ship?

## Solution

IDENTIFY We'll use the length contraction equation to solve the problem.
SET UP The proper length is the length measured in a frame in which the spacecraft is at rest. The crew will measure the spacecraft when it is at rest, which is the proper length. The student measures the relativistically contracted length.

EXECUTE The length contraction equation is

$$
l=l_{0} \sqrt{1-u^{2} / c^{2}}
$$

The student measures the contracted length and we need to find the proper length. Solving for $l_{0}$ yields

$$
l_{0}=\frac{l}{\sqrt{1-u^{2} / c^{2}}}=\frac{(97.0 \mathrm{~m})}{\sqrt{1-(0.85 c)^{2} / c^{2}}}=184 \mathrm{~m}
$$

The spacecraft's crew would measure the ship to be 184 m long.
EVALUATE This is a classic length contraction problem. The length we found was longer than the measured length because the measured length was the contracted length. You must carefully determine which length is the proper length.

CAUTION Find the Proper Length Carefully: Determining the proper length is critical in length contraction problems. The proper length is the length measured in a frame in which the body is at rest.

## 3: Spaceship relativity

A spaceship moving away from the earth at a speed of 0.60 c fires a $5.0-\mathrm{m}$-long missile in its direction of motion with a speed of 0.20 c relative to the spaceship. A crew member on the ship observes that the firing takes 10.0 s . (a) What is the missile's speed relative to earth? (b) What is the length of the missile prior to firing, as observed on earth? (c) What is the time interval of the firing event, as measured on earth?

## Solution

IDENTIFY AND SET UP We'll use a Lorentz transformation, time dilation, and length contraction to determine the values in the earth's frame.

EXECUTE From the Lorentz transformation, the speed of the missile relative to earth is found to be

$$
v=\frac{v^{\prime}+u}{1+u v^{\prime} / c^{2}}=\frac{(0.20 c)+(0.60 c)}{1+(0.60 c)(0.20 c) / c^{2}}=0.71 c
$$

where we have set the moving frame $\left(S^{\prime}\right)$ to be the frame of the spaceship moving at speed $u=0.60 c$ and the speed of the missile to be $v^{\prime}=0.20 c$.

The length of the missile as measured on the earth is found from the length contraction relation. The known length was measured on the spaceship and is therefore the proper length. The length on earth is

$$
l=l_{0} \sqrt{1-u^{2} / c^{2}}=(5.0 \mathrm{~m}) \sqrt{1-(0.60 c)^{2} / c^{2}}=4.0 \mathrm{~m} .
$$

The time elapsed for the firing of the missile, as measured on the earth, is found from the time dilation relation. The known firing time was measured on the spaceship and is therefore the proper time. The firing time on earth is

$$
\Delta t=\frac{\Delta t_{0}}{\sqrt{1-u^{2} / c^{2}}}=\frac{(10.0 \mathrm{~s})}{\sqrt{1-(0.60 c)^{2} / c^{2}}}=12.5 \mathrm{~s}
$$

As measured on earth, the speed of the missile is $0.71 c$, the length of the missile is 4.0 m , and the time taken to fire the missile is 12.5 s .

EVALUATE This problem combined several aspects of relativity. In contrast to the previous problems, the given time and length were in the proper frame.

## 4: Accelerating an electron

An electron is accelerated from rest to a velocity of $0.9 c$ by a potential difference. Calculate the potential difference. The rest energy of the electron is 0.511 MeV .

## Solution

IDENTIFY We will use the relativistic energy to find the solution.
SET UP The change in kinetic energy is the change in the electric potential energy. We'll find the potential needed to create the corresponding change in kinetic energy. Recall that 1 MeV is 1 million electron volts, a unit of energy.

EXECUTE The initial energy of the electron at rest is 0.511 MeV . The energy of the electron at $0.9 c$ is

$$
E=\left(\frac{m_{\text {electron }} c^{2}}{\sqrt{1-v^{2} / c^{2}}}\right)=\left(\frac{0.511 \mathrm{MeV}}{\sqrt{1-(0.9)^{2}}}\right)=1.172 \mathrm{MeV}
$$

The change in energy is due to the change in electric potential energy. Algebraically, this is

$$
\begin{aligned}
e \Delta V & =\Delta E \\
& =E-m c^{2} \\
& =1.172 \mathrm{MeV}-0.511 \mathrm{MeV}=0.661 \mathrm{MeV}
\end{aligned}
$$

A potential difference of $661,000 \mathrm{~V}$ is needed to accelerate the electron to $0.9 c$.
EVALUATE We see how we must combine our knowledge of energy in general with our knowledge of relativistic energy in this problem.

## 5: Creating a particle

Two protons moving towards each other with equal speeds collide and produce an $\eta_{0}$ particle. (a) If the two protons and the $\eta_{0}$ are at rest after the collision, find the initial speed of the protons. (b) What is the kinetic energy of each proton? (c) What is the rest energy of the $\eta_{0}$ particle? The rest mass of each proton is $1.67 \times 10^{-27} \mathrm{~kg}$, and the rest mass of the $\eta_{0}$ is $9.75 \times 10^{-28} \mathrm{~kg}$.

## Solution

IDENTIFY AND SET UP We will use the relativistic energy relations to find the solution to the problem. We must include the rest energy of the particles.

EXECUTE Conservation of mass and energy requires that the energy before the interaction be the same as the energy after the interaction. After the interaction, there is only rest energy; before, there is only the total energy of the two protons. Energy-mass conservation gives

$$
2\left(\frac{m_{\text {proton }} c^{2}}{\sqrt{1-v^{2} / c^{2}}}\right)=2 m_{\text {protoon }} c^{2}+m_{\eta_{0}} c^{2} .
$$

We need to solve this equation for the velocity of the protons. Rearranging terms yields

$$
\begin{aligned}
& \sqrt{1-v^{2} / c^{2}}=\frac{2 m_{\text {proton }}}{2 m_{\text {proton }}+m_{\eta_{0}}}=\frac{2\left(1.67 \times 10^{-27} \mathrm{~kg}\right)}{2\left(1.67 \times 10^{-27} \mathrm{~kg}\right)+\left(9.75 \times 10^{-28} \mathrm{~kg}\right)}=0.774 \\
& 1-v^{2} / c^{2}=(0.774)^{2}=0.5991 \\
& v=\sqrt{(1-0.5991)} c=0.633 c
\end{aligned}
$$

The kinetic energy of each proton is

$$
\begin{aligned}
K & =\frac{m_{\text {proton }} c^{2}}{\sqrt{1-v^{2} / c^{2}}}-m_{\text {proton }} c^{2}=\frac{\left(1.67 \times 10^{-27} \mathrm{~kg}\right) c^{2}}{\sqrt{1-(0.633 c)^{2} / c^{2}}}-\left(1.67 \times 10^{-27} \mathrm{~kg}\right) c^{2} \\
& =4.38 \times 10^{-11} \mathrm{~J}=274 \mathrm{MeV} .
\end{aligned}
$$

The rest mass of the $\eta_{0}$ particle is

$$
E=m_{\eta_{0}} c^{2}=\left(9.75 \times 10^{-28} \mathrm{~kg}\right) c^{2}=8.78 \times 10^{-11} \mathrm{~J}=548 \mathrm{MeV}
$$

The initial speed of the protons is $0.663 c$. The initial kinetic energy of each proton is 274 MeV , and the rest energy of the $\eta_{0}$ particle is 548 MeV .

EVALUATE In this problem, we see how the kinetic energy of the two protons converts to the mass of the $\eta_{0}$ particle. Each proton loses 274 MeV , for a total of 548 MeV . This energy is converted into the rest energy of the $\eta_{0}$ particle.

## Try It Yourself!

## 1: Traveling twin

One of a pair of twins makes a round-trip journey to a distant star, traveling at $0.95 c$. If this twin ages 6.57 years during the journey, how much does his twin who remains on earth age during the trip? How far from earth did the twin travel?

## Solution Checkpoints

IDENTIFY AND SET UP Use time dilation to solve the problem. Which twin measures proper time?
EXECUTE The time in the traveling twin's reference frame is the proper time, since he is the observer at rest. The twin on earth ages

$$
\Delta t=\frac{\Delta t_{0}}{\sqrt{1-u^{2} / c^{2}}}=21.05 \text { years }
$$

The distance traveled is the velocity times the elapsed time, given by

$$
\Delta x=\frac{(0.95 c)(21.05 \text { years })}{c}=20.0 \text { light years. }
$$

The total distance traveled is twice the maximum distance from earth. The traveling twin traveled 10 light years away from earth and back.

EVALUATE How can you check these results?

## 2: Traveling electrons

Two electrons traveling in the same direction have respective energies of 1.0 MeV and 2.0 MeV in reference frame $S$. Find the velocity of each electron in $S$. Find the velocity of the $2.0-\mathrm{MeV}$ electron relative to the $1.0-\mathrm{MeV}$ electron. Use 0.511 MeV for the rest energy of the electron.

## Solution Checkpoints

IDENTIFY AND SET UP Use relativistic energy and velocity to solve the problem.
EXECUTE The velocity of each electron is found from

$$
E=\left(\frac{m_{\text {electron }} c^{2}}{\sqrt{1-u^{2} / c^{2}}}\right)
$$

The $1.0-\mathrm{MeV}$ electron has velocity $u_{1}=0.8602 c$, and the $2.0 \mathrm{Me}-\mathrm{V}$ electron has velocity $u_{2}=$ 0.9669 c. The relative velocities are found from

$$
v_{x}^{\prime}=\frac{v_{x}-u}{1-u v_{x} / c^{2}} .
$$

This equation is applied to a frame $S^{\prime}$ in which the $1.0-\mathrm{MeV}$ electron is at rest. With this reference frame, $v_{x}$ is the velocity of the $2.0-\mathrm{MeV}$ electron in $S(0.9669 c), u$ is the velocity of $S^{\prime}$ with respect to $S$ $(0.8602 c)$, and $v_{x}^{\prime}$ is the velocity of the $2.0-\mathrm{MeV}$ electron in $S^{\prime}$ (the target variable). Substituting and solving leads to $v_{x}{ }^{\prime}=0.636 c$.

EVALUATE The challenge in this problem was to interpret the relative velocity variables carefully. With relativity, relative velocities are no longer simply vector sums of velocities.

## 3: Neutron decay

When a neutron spontaneously decays into a proton, an electron, and a neutrino, the decay products are found to have a total kinetic energy of 7.81 MeV . The proton has a mass of $1.673 \times 10^{-27} \mathrm{~kg}$, the electron has a mass of $9.110 \times 10^{-31} \mathrm{~kg}$, and the neutrino has no mass. What is the mass of the neutron?

## Solution Checkpoints

IDENTIFY AND SET UP Equate the energy of the reaction products to the rest energy of the neutron to solve the problem.

EXECUTE Energy conservation leads to

$$
m_{\mathrm{n}} c^{2}=m_{\mathrm{p}} c^{2}+m_{\mathrm{e}} c^{2}+E_{k} .
$$

Converting the kinetic energy to joules gives a kinetic energy of $1.25 \times 10^{-13} \mathrm{~J}$. Substituting values into the energy relation gives a neutron mass of $1.675 \times 10^{-27} \mathrm{~kg}$.
EVALUATE Can a proton spontaneously decay into a neutron?

## Photons, Electrons, and Atoms

## Summary

We will explore the quantum mechanics of photons, electrons, and atoms in this chapter. We will see that light can behave as a stream of individual particles, or photons, that have energy and momentum. These photons will lead us to the discovery that atoms are quantized, or have distinct, discrete energy levels. This finding in turn will lead us to the discovery that electrons can behave as waves. In essence, we will be learning about the very unusual and nonintuitive atomic world. We will examine the discoveries that led to this new interpretation, including the photoelectric effect, atomic spectra, the Bohr model of the atom, Compton scattering, and blackbody radiation. The basic quantum mechanics learned in this chapter will be applied in the remaining chapters of the text.

## Objectives

After studying this chapter, you will understand

- How the photoelectric effect confirmed the photon nature of light.
- How atomic line spectra reveal the energy levels of atoms.
- The Bohr model of the hydrogen atom and how energy levels are quantized.
- How the laser operates.
- How to use the photon nature of light to interpret the Compton scattering of X-rays.
- How blackbody radiation shows that electromagnetic radiation is quantized.
- Wave-particle duality.


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Photons | Electromagnetic radiation exhibits both wave and particle behavior. Photons <br> carry units of electromagnetic radiation. The energy of a photon is given by |

$$
E=h f=\frac{h c}{\lambda},
$$

where $h=6.626 \times 10^{-34} \mathrm{~J}$ s is Planck's constant, $f$ is the frequency of the photon, and $\lambda$ is the wavelength of the photon. The momentum of a photon is given by

$$
p=\frac{E}{c}=\frac{h f}{c}=\frac{h}{\lambda} .
$$

## Photoelectric Effect

The photoelectric effect describes how a photon striking a surface can eject an electron from that surface. The photon must have sufficient energy (greater than the work function $\phi$ of the material of the surface) for the electron to escape. Mathematically,

$$
e V_{0}=h f-\phi .
$$

## Energy Levels and Atomic Line Spectra

An atom making a transition from a higher energy $E_{i}$ to a lower energy $E_{f}$ will emit the energy difference through a photon of energy:

$$
h f=\frac{h c}{\lambda}=E_{i}-E_{f} .
$$

Energy differences can be detected through an atom's spectral lines. For hydrogen, the energy levels are given by

$$
E_{n}=-\frac{h c R}{n^{2}}=-\frac{13.6 \mathrm{eV}}{n^{2}}, \quad n=1,2,3, \ldots,
$$

where $R=1.097 \times 10^{7} / \mathrm{m}$ is the Rydberg constant.

The Nuclear Atom and the Bohr Model

Rutherford discovered that an atom has a very small positively charged nucleus at its center. Bohr successfully modeled the hydrogen atom as having a lone proton as its nucleus, surrounded by an electron that revolved in certain allowed (quantized) orbits. The electron's allowed angular momenta are given by

$$
L=m v_{n} r_{n}=n \frac{h}{2 \pi}, \quad n=1,2,3, \ldots,
$$

where $n$ is the principal quantum number. The electron's radius and orbital speed are also quantized:

$$
\begin{aligned}
& r_{n}=\mathcal{E}_{0} \frac{n^{2} h^{2}}{\pi m e^{2}}=n^{2} a_{0}=n^{2}\left(5.29 \times 10^{-11} \mathrm{~m}\right), \\
& v_{n}=\frac{e^{2}}{\mathcal{E}_{0} 2 n h}=\frac{21.9 \times 10^{6} \mathrm{~m} / \mathrm{s}}{n}
\end{aligned}
$$

## The Laser

The laser operates on the principle of stimulated emission, in which many photons with identical wavelengths and phases are emitted. For the laser to operate, a nonequilibrium population inversion must exist in which more atoms are in a higher energy state than a lower energy state. The stimulated emission occurs as the higher energy state decays to the lower energy state, triggered by photons passing through the material.

| X-Rays and Compton Scattering | X-rays are high-energy, small-wavelength photons that can be produced when <br> electrons strike a target. X-rays may scatter from electrons bound to a nucleus <br> in a process called Compton scattering. After scattering, the X-rays have less <br> energy and a longer wavelength. The change in an X-ray's wavelength is <br> given by |
| :--- | :--- |
| $\qquad \Delta \lambda=\lambda^{\prime}-\lambda=\frac{h}{m c}(1-\cos \phi)$, |  |
| Blackbody Radiation | where $m$ is the electron's mass and $\phi$ is the scattering angle. This discovery <br> helped prove that light, X-rays, and all electromagnetic radiation are made of <br> discrete energy packets called photons. |
|  | The total radiated intensity from a blackbody surface is described by the <br> Stefan-Boltzmann law $\quad I=\sigma T^{4}$, |
| where $\sigma=5.67 \times 10^{-8} \mathrm{~W} / \mathrm{m}^{2}$ is the Stefan-Boltzmann constant. |  |

## Conceptual Questions

## 1: Finding the stopping potential

Figure 38.1 shows a graph of the stopping potential as a function of the frequency of incident light illuminating a metal surface. Find the photoelectric work function for this surface.


Figure 38.1 Question 1.

## Solution

IDENTIFY, SET UP, AND EXECUTE The graph of the stopping potential as a function of the frequency of light is a straight line given by

$$
V_{0}=\frac{h}{e} f-\frac{\phi}{e} .
$$

The vertical intercept of this graph is the negative of the work function, divided by $e$. Here, the vertical intercept is -2.0 V , so the work function is 2.0 eV .

EVALUATE We'll often graph the stopping potential as a function of the frequency of light in order to find the work function. We see how extracting the work function is simple.

## 2: Interpreting an energy-level diagram

The energy-level diagram of the hypothetical one-electron element nerdium is shown in Figure 38.2. The potential energy is taken to be zero for an electron an infinite distance from the atom. If a $10.5-\mathrm{eV}$ photon interacts with the nerdium atom in its ground state, what will happen? If a $6.5-\mathrm{eV}$ photon interacts with the nerdium atom in its ground state, what will happen?


Figure 38.2 Question 2.

## Solution

IDENTIFY, SET UP, AND EXECUTE The $10.5-\mathrm{eV}$ photon has energy equal to the energy difference between the $n=1$ and the $n=3$ levels of nerdium. This photon will be absorbed, leaving the nerdium in the $n=3$ excited state.

The $6.5-\mathrm{eV}$ photon has energy less than the energy difference between the $n=1$ and the $n=2$ levels of nerdium. This photon will not be absorbed by the nerdium atom.

EVALUATE We see that atoms do not behave classically. In a classical system, both photons would be absorbed. In the quantized system, only photons with the proper energy difference are absorbed.

## 3: Compton effect

Why is there no wavelength shift for forward scattering $\left(\theta=0^{\circ}\right)$ in the Compton effect?

## Solution

IDENTIFY, SET UP AND EXECUTE When $\theta=0^{\circ}$ is substituted into the Compton scattering equation, the wavelength shift is found to be zero. A photon scattered at $0^{\circ}$ does not interact with an electron; its momentum and energy remain unchanged. Since the photon does not interact, it continues undeflected with the same wavelength.

EVALUATE The photon's energy and momentum must remain constant due to conservation of energy and momentum.

## Problems

## 1: Photoelectric effect for tungsten

The photoelectric work function for sodium is 2.7 eV . If light of frequency $9.0 \times 10^{14} \mathrm{~Hz}$ falls on sodium, find (a) the stopping potential, (b) the kinetic energy of the most energetic electrons ejected, and (c) the speed of those electrons.

## Solution

IDENTIFY AND SET UP We will use energy conservation to find the stopping potential. The most energetic electrons will have energy equal to the stopping potential energy.

EXECUTE The stopping potential energy is equal to the photon's energy minus the work function:

$$
e V_{0}=h f-\phi .
$$

The stopping potential is then

$$
V_{0}=\frac{h f-\phi}{e}=\frac{\left(9.0 \times 10^{14} \mathrm{~Hz}\right)\left(6.626 \times 10^{-34} \mathrm{~J} \cdot \mathrm{~s}\right)-(2.7 \mathrm{eV})\left(1.6 \times 10^{-19} \mathrm{~J} / \mathrm{eV}\right)}{1.6 \times 10^{-19} \mathrm{C}}=1.03 \mathrm{~V}
$$

The maximum kinetic energy of an electron is then

$$
K=e V_{0}=e(1.03 \mathrm{~V})=1.03 \mathrm{eV}
$$

and the speed of that electron is

$$
v=\sqrt{\frac{2 K}{m}}=\sqrt{\frac{2(1.03 \mathrm{eV})\left(1.6 \times 10^{-19} \mathrm{~J} / \mathrm{eV}\right)}{\left(9.11 \times 10^{-31} \mathrm{~kg}\right)}}=6.01 \times 10^{5} \mathrm{~m} / \mathrm{s}
$$

The stopping potential is 1.03 V , the maximum kinetic energy of the electrons that are ejected is 1.03 eV , and the speed of the electrons ejected with maximum kinetic energy is $6.01 \times 10^{5} \mathrm{~m} / \mathrm{s}$.

EVALUATE We rely on basic energy conservation laws to solve problems involving the photoelectric effect. Note how we've used energy relations which imply that light is a particle.

## 2: Stopping potential

The stopping potential for photoelectrons ejected from a surface by $375-\mathrm{nm}$ photons is 1.870 V . Calculate the stopping potential for $600-\mathrm{nm}$ photons.

## Solution

IDENTIFY AND SET UP We will use energy conservation to find the work function and then use that to find the stopping potential for the $600-\mathrm{nm}$ photons, the target variable.

EXECUTE The stopping potential energy is equal to the photon's energy minus the work function:

$$
e V_{0}=h f-\phi
$$

The work function is found from the $375-\mathrm{nm}$ photon response. Substituting and solving gives

$$
\begin{aligned}
\phi & =\frac{h c}{\lambda}-e V_{0} \\
& =\frac{\left(6.626 \times 10^{-34} \mathrm{~J} \cdot \mathrm{~s}\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}{\left(375 \times 10^{-9} \mathrm{~m}\right)}-\left(1.6 \times 10^{-19} \mathrm{C}\right)(1.870 \mathrm{~V}) \\
& =2.31 \times 10^{-19} \mathrm{~J}=1.44 \mathrm{eV}
\end{aligned}
$$

For the $600-\mathrm{nm}$ photons, we find the stopping potential from

$$
\begin{aligned}
e V_{0}^{\prime} & =\frac{h c}{\lambda^{\prime}}-\phi \\
& =\frac{\left(6.626 \times 10^{-34} \mathrm{~J} \cdot \mathrm{~s}\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}{\left(600 \times 10^{-9} \mathrm{~m}\right)}-\left(2.31 \times 10^{-19} \mathrm{~J}\right) \\
& =1.00 \times 10^{-19} \mathrm{~J} .
\end{aligned}
$$

The stopping potential for the $600-\mathrm{nm}$ photons is then

$$
V_{0}^{\prime}=\frac{1.00 \times 10^{-19} \mathrm{~J}}{1.6 \times 10^{-19} \mathrm{C}}=0.627 \mathrm{~V}
$$

EVALUATE The stopping potential depends on the wavelength of light used, but the work function is constant for a material.

## 3: Transitions in the hydrogen atom

A hydrogen atom initially in the ground state absorbs a photon, exciting it to the $n=5$ state. Later, the atom makes a transition to the $n=2$ state. Find the wavelength of the photon that is absorbed when the atom goes from the ground state to the $n=5$ state. Then find the wavelength of the photon that is emitted when the atom goes from the $n=5$ state to the $n=2$ state.

## Solution

IDENTIFY We'll use the energy difference to find the wavelength of the photon.
SET UP The photon that is absorbed must have energy equal to the energy difference between the ground state $(n=1)$ and the $n=5$ state. The transition back to the $n=2$ state will decrease the energy, so a photon is emitted. This photon's energy will be equal to the energy difference between the $n=5$ and $n=2$ states. We'll use that energy to find the wavelength of the emitted photon.

EXECUTE The energy difference between the ground state $(n=1)$ and the $n=5$ state is

$$
E_{i}-E_{f}=-\frac{13.6 \mathrm{eV}}{n^{2}}-\left(-\frac{13.6 \mathrm{eV}}{n^{2}}\right)=-\frac{13.6 \mathrm{eV}}{1^{2}}+\frac{13.6 \mathrm{eV}}{5^{2}}=-13.1 \mathrm{eV}=-2.09 \times 10^{-18} \mathrm{~J}
$$

The negative sign indicates that energy is absorbed. The photon's wavelength is

$$
\lambda=\frac{h c}{E_{i}-E_{f}}=\frac{\left(6.626 \times 10^{-34} \mathrm{~J} \cdot \mathrm{~s}\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}{2.09 \times 10^{-18} \mathrm{~J}}=95 \mathrm{~nm}
$$

The energy difference in the transition from the $n=5$ state to the $n=2$ state is

$$
E_{i}-E_{f}=-\frac{13.6 \mathrm{eV}}{n^{2}}-\left(-\frac{13.6 \mathrm{eV}}{n^{2}}\right)=-\frac{13.6 \mathrm{eV}}{5^{2}}+\frac{13.6 \mathrm{eV}}{2^{2}}=2.86 \mathrm{eV}=4.57 \times 10^{-19} \mathrm{~J}
$$

The positive sign indicates that energy is emitted. This photon's wavelength is

$$
\lambda=\frac{h c}{E_{i}-E_{f}}=\frac{\left(6.626 \times 10^{-34} \mathrm{~J} \cdot \mathrm{~s}\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}{4.57 \times 10^{-19} \mathrm{~J}}=435 \mathrm{~nm}
$$

A 95-nm-wavelength photon is absorbed to excite the ground-state hydrogen atom to the $n=5$ state. A $435-n m$-wavelength photon is emitted to "de-excite" the excited-state hydrogen atom to the $n=2$ state.

EVALUATE Imagine that instead of predicting the wavelengths of the photons, you were given the wavelengths and had to deduce the energies. This is what physicists did in the early 1900s to unravel the structure of the atom.

## 4: X-ray scattering

X-rays of frequency $9.0 \times 10^{18} \mathrm{~Hz}$ Compton-scatter off electrons. What are the wavelength and energy of X-rays scattered to an angle of $135^{\circ}$ ?

## Solution

IDENTIFY AND SET UP The Compton-scattering formula gives the change in wavelength for the X-rays under consideration. We'll convert the X-ray frequency to a wavelength before finding the change in wavelength.

EXECUTE The X-rays have wavelength

$$
\lambda=\frac{c}{f}=\frac{3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}}{9.0 \times 10^{18} \mathrm{~Hz}}=3.33 \times 10^{-11} \mathrm{~m}
$$

The change in wavelength due to Compton scattering is

$$
\begin{aligned}
\Delta \lambda & =\lambda^{\prime}-\lambda=\frac{h}{m c}(1-\cos \phi)=\frac{\left(6.626 \times 10^{-34} \mathrm{~J} \cdot \mathrm{~s}\right)}{\left(9.11 \times 10^{-31} \mathrm{~kg}\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}\left(1-\cos 135^{\circ}\right) \\
& =4.15 \times 10^{-12} \mathrm{~m}
\end{aligned}
$$

The scattered X-ray's wavelength is

$$
\lambda^{\prime}=\lambda+\Delta \lambda=\left(3.33 \times 10^{-11} \mathrm{~m}\right)+\left(4.15 \times 10^{-12} \mathrm{~m}\right)=3.75 \times 10^{-11} \mathrm{~m}
$$

The energy of the scattered X-ray is

$$
E=\frac{h c}{\lambda}=\frac{\left(6.626 \times 10^{-34} \mathrm{~J} \cdot \mathrm{~s}\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}{\left(3.75 \times 10^{-11} \mathrm{~m}\right)}=5.31 \times 10^{-15} \mathrm{~J}=33.2 \mathrm{keV}
$$

The scattered X-ray has a wavelength of 37.5 pm and an energy of 33.2 keV .
EVALUATE Compton scattering changes the wavelength and energy by a relatively small amount. Here, the energy decreased by about $11 \%$. For longer wavelength photons, the effect is much smaller.

## 5: Compton scattering

Compare the maximum relative frequency change for $550-\mathrm{nm}$ photons to $0.025-\mathrm{nm}$ X-rays. For which type of photon will Compton scattering be more easily observed?

## Solution

IDENTIFY AND SET UP We will calculate the maximum frequency shift due to Compton scattering for the two photons. The maximum frequency shift corresponds to a scattering angle of $180^{\circ}$.

EXECUTE The change in wavelength due to Compton scattering is

$$
\Delta \lambda=\lambda^{\prime}-\lambda=\frac{h}{m c}(1-\cos \phi)
$$

Converting the formula to frequencies gives

$$
\frac{1}{f^{\prime}}-\frac{1}{f}=\frac{h}{m c^{2}}(1-\cos \phi)
$$

Rearranging terms to find the relative frequency change, we obtain

$$
\frac{f^{\prime}-f}{f^{\prime}}=\frac{\Delta f}{f^{\prime}}=\frac{h f}{m c^{2}}(1-\cos \phi)
$$

The maximum frequency shift occurs when the photon is backscattered, or $\phi=180^{\circ}$. We find the relative frequency change for the two photons. For the $550-\mathrm{nm}$ photons,

$$
\frac{\Delta f}{f^{\prime}}=\frac{h f}{m c^{2}}(1-\cos \phi)=\frac{2 h}{m c \lambda}=\frac{2\left(6.626 \times 10^{-34} \mathrm{~J} \cdot \mathrm{~s}\right)}{\left(9.11 \times 10^{-31} \mathrm{~kg}\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)(550 \mathrm{~nm})}=8.8 \times 10^{-6}
$$

For the 0.025 -nm X-rays,

$$
\frac{\Delta f}{f^{\prime}}=\frac{2 h}{m c \lambda}=\frac{2\left(6.626 \times 10^{-34} \mathrm{~J} \cdot \mathrm{~s}\right)}{\left(9.11 \times 10^{-31} \mathrm{~kg}\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)(0.025 \mathrm{~nm})}=0.18
$$

X-ray Compton scattering is much easier to observe than Compton scattering with visible photons.
EVALUATE We noted in the previous problem that the Compton effect produces small changes in the scattered photon's energies. We see that effect clearly in the results of this problem, noting that the X-rays have small, but measurable frequency (and energy) changes while the visible-light photons have even smaller, not readily measurable changes in their frequency (and energy).

## Try It Yourself!

## 1: Threshold wavelength

Find the threshold wavelength of light that would produce photoelectrons for a silver surface. The work function for silver is 4.8 eV .

## Solution Checkpoints

IDENTIFY AND SET UP Use energy conservation to solve the problem.
EXECUTE Energy conservation for the photoelectric effect gives

$$
e V_{0}=h f-\phi
$$

At threshold, the photoelectrons have zero kinetic energy. Solving for the wavelength yields

$$
\lambda=\frac{h c}{\phi}=259 \mathrm{~nm}
$$

EVALUATE What is the threshold wavelength for a cesium surface, for which the work function is 1.8 eV ?

## 2: Atomic spectra

An atom has, in addition to the ground-state energy $E_{0}$ at zero energy, energy levels at $E_{1}=$ $10.20 \mathrm{eV}, \mathrm{E}_{2}=12.09 \mathrm{eV}$, and $E_{3}=12.75 \mathrm{eV}$. If the atom is excited from the ground state to the state with an energy of 12.75 eV , find all possible wavelengths in the atom's spectrum.

## Solution Checkpoints

IDENTIFY AND SET UP Find all possible transitions between the various levels, and calculate the corresponding wavelengths.

EXECUTE For any level, the energy is given by

$$
\Delta E=\frac{h c}{\lambda}
$$

The excited atom in the $E_{3}$ level decays to the $E_{2}, E_{1}$, and $E_{0}$ levels, by emitting photons with wavelengths of $1880 \mathrm{~nm}, 487 \mathrm{~nm}$, and 97.3 nm , respectively. The $E_{2}$ level decays to the $E_{1}$ and $E_{0}$ levels, emitting photons with wavelengths of 657 nm and 102 nm , respectively. Finally, the $E_{1}$ level decays to the $E_{0}$ level, emitting a photon with wavelength of 122 nm .

EVALUATE To calculate the complete spectrum, you must consider all possible decays and their corresponding wavelengths. You may also consider the opposite situation, in which you are given the spectrum and need to calculate the possible levels, much as the physicists who first disentangled atomic spectra had to do.

## 3: Compton scattering

A 50-keV X-ray strikes an electron at rest and scatters. If the X-ray is scattered at an angle of $90^{\circ}$, find (a) the change in wavelength of the X-ray, (b) the energy of the X-ray after scattering, and (c) the velocity of the electron after scattering.

## Solution Checkpoints

IDENTIFY AND SET UP Use the Compton-scattering equation and energy conservation to solve the problem.

EXECUTE (a) Compton scattering gives the change in wavelength for the X-ray:

$$
\Delta \lambda=\lambda^{\prime}-\lambda=\frac{h}{m c}(1-\cos \phi)
$$

For this problem, the scattering angle is $90^{\circ}$, the wavelength shift is $2.42 \times 10^{-12} \mathrm{~m}$, and the scattered X-ray has a wavelength of $2.73 \times 10^{-11} \mathrm{~m}$.
(b) The new frequency is found from

$$
f^{\prime}=c / \lambda^{\prime}
$$

The new energy is 45.6 keV .
(c) The energy lost by the X-ray becomes kinetic energy of the electron. The kinetic energy is much smaller than the electron rest energy, so we can use classical energy expressions:

$$
\frac{1}{2} m v^{2}=\Delta E=4.44 \mathrm{keV}
$$

The electron acquires a velocity of $3.96 \times 10^{7} \mathrm{~m} / \mathrm{s}$.
EVALUATE Is the electron's velocity small enough to justify using classical energy expressions?

## The Wave Nature of Particles

## Summary

We will explore the wave nature of particles in this chapter. Wave-particle duality has shown us that light behaves both like a wave and like a particle. We will investigate how subatomic particles behave like waves in this chapter, opening our exploration of quantum mechanics. Our analysis of particles will transform to an analysis of wave functions-an analysis that will be used to predict the probability of a particle being in a specific location at a specific time. We will see how Schrödinger's equation is used to solve for wave functions and becomes the equivalent of Newton's law for quantum mechanics.

## Objectives

After studying this chapter, you will understand

- How electrons and other subatomic particles behave like waves.
- How to use the Heisenberg uncertainty principle to interpret atomic phenomena.
- How particles are described in terms of wave functions.
- The use of the Schrödinger equation to determine the behavior of particles.

Concepts and Equations

| Description |  |
| :--- | :--- |
| De Broglie Waves | Electrons and other particles have wave properties. Particles are described as <br> waves having a de Broglie wavelength |
| $\lambda=\frac{h}{p}=\frac{h}{m v}$. |  |

As waves, particles are inherently spread-out entities described by their wave functions.

The diffraction of an electron beam from the surface of a metallic crystal confirms the wave nature of particles. The wavelength of a nonrelativistic electron accelerated by a potential difference $V$ is given by

$$
\lambda=\frac{h}{p}=\frac{h}{\sqrt{2 m e V}} .
$$

Images from electron microscopes have much higher resolution than images from visible light, due to the small wavelength of the electrons.
Heisenberg Uncertainty Principle

## Wave Functions

## The Schrödinger Equation

Heisenberg's uncertainty principle states that one cannot determine both the precise position and the precise momentum (or the precise energy and the precise time) of a particle. The uncertainties in each of the two quantities are related as

$$
\Delta x \Delta p_{x} \geq \hbar, \quad \Delta E \Delta t \geq \hbar, \quad \hbar=\frac{h}{2 \pi}
$$

Similar expressions hold for the $y$ and $z$ components.
The wave function $\Psi(x, y, z, t)$ for a particle contains all of the information about the particle. The quantity $|\Psi(x, y, z, t)|^{2}$ is the probability distribution function that determines the relative probability of finding a particle near a given position at a given time. For a particle in a definite energy state, called a stationary state, the wave function can be separated into a spatial component and a temporal component:

$$
\Psi(x, y, z, t)=\psi(x, y, z) e^{-i E t / \hbar}
$$

The Schrödinger equation can be used to determine the wave function for a particle moving in one dimension in the presence of a potential-energy function $U(x)$. The Schrödinger equation is

$$
-\frac{\hbar^{2}}{2 m} \frac{d^{2} \psi(x)}{d x^{2}}+U(x) \psi(x)=e \psi(x)
$$

## Conceptual Questions

## 1: Tennis-ball wave function

Tennis balls are made of electrons and other subatomic particles, so a tennis ball can be described in terms of a combination of wave functions. If tennis balls are wave functions, couldn't they exhibit destructive interference and disappear before reaching your racket?

## Solution

IDENTIFY, SET UP, AND EXECUTE It is true that the particles which make up the tennis ball are described by wave functions. The wave functions describe the particles as spread out in a region of space, and the waves can interfere destructively. However, the wavelength of the tennis ball is so small that any interference effects that might arise are below the threshold of visibility.

EVALUATE As we look at the strange subatomic world, we need to make sure that our interpretations and predictions relating to the macroscopic world remain valid. Quantum mechanics, as we see in this problem, apply to the macroscopic world, but the effects are so small that they do not change our existing interpretation.

## 2: Particles from nowhere

Does the Heisenberg uncertainty principle imply that particles could be created for very short amounts of time? How long could an electron-positron pair exist without violating the uncertainty principle?

## Solution

IDENTIFY, SET UP, AND EXECUTE The uncertainty principle for energy and time states that the product of the uncertainty in energy and the uncertainty in time is greater than $h / 2 \pi$. The uncertainty in energy can be very large, as long as the uncertainty in time is very small. The energy needed to create an electron-positron pair is roughly 1 MeV , so the uncertainty in time would be

$$
\Delta t \geq \frac{\hbar}{\Delta E}=\frac{\left(6.626 \times 10^{-34} \mathrm{~J} \cdot \mathrm{~s}\right)}{2 \pi\left(10^{6} \mathrm{eV}\right)\left(1.6 \times 10^{-19} \mathrm{~J} / \mathrm{eV}\right)}=6.6 \times 10^{-22} \mathrm{~s}
$$

If the electron and positron are created and annihilated within $6.6 \times 10^{-22} \mathrm{~s}$, then the uncertainty principle is not violated.

EVALUATE Creating particles from vacuum is not just a theoretical exercise. In what is called the quantum vacuum, particles are created and annihilated on a continual basis. Using particle accelerators, physicists have proven the existence of these particles by scattering other particles off of them during their short lifetimes. These so-called "sea" particles (since they are created in the vacuum "sea") are responsible for important phenomena at the subatomic scale.

## Problems

## 1: Wavelength of a golf ball

An uncharged golf ball of mass 0.1 kg is put into orbit at the earth's surface. Find the de Broglie wavelength of the golf ball.

## Solution

IDENTIFY AND SET UP The de Broglie wavelength is related to the momentum. We will find the momentum and then the wavelength.

EXECUTE The force acting on the golf ball in orbit is the gravitational force between the ball and the earth. The ball undergoes centripetal acceleration, so

$$
\frac{G M m}{r^{2}}=m g=\frac{m v^{2}}{r}
$$

We use this expression to find the velocity of the golf ball. Rearranging terms to solve for the velocity gives

$$
v=\sqrt{g R_{E}}
$$

where $R_{E}$ is the radius of the earth. The momentum is the mass of the golf ball times the velocity. The wavelength is then

$$
\begin{aligned}
\lambda & =\frac{h}{p}=\frac{h}{m v}=\frac{h}{m \sqrt{g R_{E}}} \\
& =\frac{6.626 \times 10^{-34} \mathrm{Js}}{(0.1 \mathrm{~kg}) \sqrt{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)\left(6.4 \times 10^{6} \mathrm{~m}\right)}}=8.37 \times 10^{-37} \mathrm{~m}
\end{aligned}
$$

EVALUATE We see that the wavelength is much smaller than the size of an atomic nucleus (around $10^{-15} \mathrm{~m}$ ), so you don't have to worry about the wave nature of a golf ball the next time you play a round of miniature golf!

## 2: Quantum number of a golf ball

Suppose the uncharged golf ball (of mass 0.1 kg ) of Problem 1 is viewed as a quantum object in orbit around a nucleus. If its angular momentum is quantized as in the Bohr atom, what is the associated quantum number?

## Solution

IDENTIFY AND SET UP We will apply the Bohr model to the golf ball orbiting the earth to solve the problem.

EXECUTE Quantizing the angular momentum gives

$$
m v r=n \hbar .
$$

In the previous problem, we found that

$$
v=\sqrt{g R_{E}}
$$

Combining these results, we find the following expression for the quantized radius:

$$
r_{n}=\frac{n^{2} \hbar^{2}}{m^{2} R_{E}^{2} g}
$$

For $r_{n}=R_{E}$, we find the value of $n$ :

$$
\begin{aligned}
n & =\sqrt{\frac{m^{2} R_{E}^{3} g}{\hbar^{2}}} \\
& =\sqrt{\frac{(0.1 \mathrm{~kg})^{2}\left(6.4 \times 10^{6} \mathrm{~m}\right)^{3}\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{\left(1.05 \times 10^{-34} \mathrm{Js}\right)^{2}}}=4.8 \times 10^{43}
\end{aligned}
$$

EVALUATE With this large value of $n$, we see how the allowed energy levels and values of $r$ are essentially continuous. It should be clear at this point why the quantum nature of the universe was not discovered in the macroscopic world.

## 3: A thermal neutron

A thermal neutron is a neutron with mean kinetic energy of $3 / 2 k_{B} T$, where $k_{B}$ is Boltzmann's constant and $T$ is room temperature ( 300 K ). What is the wavelength of a thermal neutron?

## Solution

IDENTIFY AND SET UP We will find the thermal neutron's momentum and then determine the de Broglie wavelength of the neutron.

EXECUTE The kinetic energy of the neutron is given by

$$
K=\frac{p^{2}}{2 m}=\frac{3}{2} k_{B} T .
$$

Substituting for the momentum, we have

$$
\frac{1}{2 m}\left(\frac{h}{\lambda}\right)^{2}=\frac{3}{2} k_{B} T .
$$

Solving for $\lambda$ gives

$$
\begin{aligned}
\lambda & =\frac{h}{\sqrt{3 m k_{B} T}} \\
& =\frac{\left(6.626 \times 10^{-34} \mathrm{Js}\right)}{\sqrt{3\left(1.67 \times 10^{-27} \mathrm{~kg}\right)\left(1.38 \times 10^{-23} \mathrm{~J} / \mathrm{K}\right)(300 \mathrm{~K})}}=0.145 \mathrm{~nm}
\end{aligned}
$$

EVALUATE Thermal neutrons have very small wavelengths. Consequently, they are used in research laboratories to explore the atomic structure of materials through diffraction.

## 4: Lifetime of a molecule

The full-width, half-maximum intensity for a spectral line characteristic of a $\mathrm{pH}^{2}$ molecule in an excited rotational energy level is $6 \times 10^{9} \mathrm{~Hz}$. Estimate the lifetime of the molecule in this state.

## Solution

IDENTIFY AND SET UP We will use the Heisenberg uncertainty principle to solve the problem.
EXECUTE The uncertainty relationship is given by

$$
(\Delta E)(\Delta t) \geq \hbar
$$

With this expression, we can estimate the lifetime of the state if we know the energy associated with the state. The energy is found from the formula

$$
\Delta E=h \Delta f
$$

Rearranging terms and solving for the lifetime gives

$$
\Delta t \geq \frac{\hbar}{h \Delta f}=\frac{1}{2 \pi \Delta f}=\frac{1}{2 \pi\left(6 \times 10^{9} / \mathrm{s}\right)}=2.65 \times 10^{-11} \mathrm{~s}
$$

EVALUATE This problem illustrates how we can use the uncertainty principle to estimate the lifetimes of atomic states.

## Try It Yourself!

## 1: Electron diffraction

An electron is accelerated through a potential difference of 1000 V and passes through a thin slit before striking a photographic film 0.5 m away. What should the size of the slit be in order for the first minimum in the electron diffraction pattern to be 0.1 mm from the center of the pattern?

## Solution Checkpoints

IDENTIFY AND SET UP Use diffraction to solve the problem. How do you calculate the wavelength of the electron?

EXECUTE The wavelength of the electron is

$$
\lambda=\frac{h}{\sqrt{2 m e V}}=3.88 \times 10^{-11} \mathrm{~m}
$$

The first minimum is found when

$$
\sin \theta=\frac{\lambda}{d}
$$

Using the small-angle approximation, we find that the width is 194 nm .
EVALUATE How could a slit of width 194 nm be created?

## 2: Electron orbiting the nucleus

Calculate (a) the speed of an electron whose wavelength is $2 \pi r_{0}$ (the circumference of its orbit, where $\left.r_{0}=0.053 \mathrm{~nm}\right)$ and (b) the speed of an electron in the first Bohr orbit of a hydrogen atom.

## Solution Checkpoints

IDENTIFY AND SET UP Use the de Broglie wavelength to find the speed for part (a). Use the velocity from the Bohr model to solve part (b).

EXECUTE (a) The de Broglie wavelength leads to

$$
v=\frac{h}{m \lambda}=\frac{h}{m 2 \pi r_{0}}
$$

Substituting values, we find that the velocity is $2.19 \times 10^{6} \mathrm{~m} / \mathrm{s}$.
(b) According to the Bohr model, the velocity is

$$
v_{n}=\frac{e^{2}}{2 \epsilon_{0} h n}
$$

Substituting values for $n=1$, we again find that the velocity is $2.19 \times 10^{6} \mathrm{~m} / \mathrm{s}$.
EVALUATE Why are the two values the same?

## 40 <br> Quantum Mechanics

## Summary

In this chapter, we will explore the quantum mechanics of particles trapped in bound states, such as electrons orbiting atoms. We will apply the Schrödinger equation to find wave function solutions of it in a variety of cases, including a particle confined to a box, a particle in a square well, and a particle in a harmonic oscillator potential. We will also investigate phenomena forbidden by Newtonian mechanics, such as quantum mechanical tunneling.

## Objectives

After studying this chapter, you will understand

- How to calculate the wave functions for a particle in a box.
- How to determine the wave function for a particle in a potential well.
- The definition of tunneling and barrier penetration.
- How to solve the problem of a particle in a harmonic oscillator potential.
- How to solve quantum mechanics problems in three dimensions.

Concepts and Equations

| Term | Description |
| :--- | :--- |
| Particle in a Box | The energy levels for a particle of mass $m$ in an infinitely deep square well <br> potential of width $L$ are given by |
| $\qquad$$E_{n}=\frac{p_{n}^{2}}{2 m}=\frac{n^{2} h^{2}}{8 m L^{2}}, \quad n=1,2,3, \ldots$ <br> The corresponding normalized particle wave functions are given by <br> $\psi_{n}(x)=\sqrt{\frac{2}{L}} \sin \frac{n \pi x}{L}, \quad n=1,2,3, \ldots$. |  |
| Wave Functions and Normalization | To be a solution of the Schrödinger equation, the wave function and its deriv- <br> ative must be continuous everywhere. Wave functions are usually normalized <br> such that the probability of finding the particle somewhere is unity: |
| Finite Potential Well | The energy levels in a potential well of finite depth are lower than those in an |
| infinite well. The levels are obtained by matching wave functions at the well |  |
| walls and satisfying continuity of the wave function and its derivative. |  |, $1 . \quad$| Because of a process called tunneling, there is a finite, nonzero probability |
| :--- |
| that a particle will penetrate a potential-energy barrier even if its initial |
| kinetic energy is less than the height of the barrier. |

## Conceptual Questions

## 1: Probability in a box

Consider the allowed energy states of a particle in a box. How does the probability of finding the particle in the left half of the box compare with the probability of finding the particle in the right half of the box, for any energy state?

## Solution

IDENTIFY, SET UP, AND EXECUTE The box is symmetric; therefore, the probability of finding a particle in the left half of the box is equal to the probability of finding the particle in the right half of the box, regardless of the energy state.

EVALUATE Building intuition about quantum mechanics and particle probabilities will help you prepare for the later chapters of the text.

## 2: Interpreting wave functions

Consider the wave function for a particle in a box, shown in Figure 40.1. Where are the locations at which the particle is most likely to be found? Where are the locations at which the particle is least likely to be found?


Figure 40.1 Question 2.

## Solution

IDENTIFY, SET UP, AND EXECUTE The figure shows the wave function for the $n=3$ state. The wave function has three nodes and crosses zero twice in the box. The probability function is the square of the wave function. When the wave function is squared, there will be three maxima, located at $L / 6$, $L / 2$, and $5 L / 6$. The particle is most likely to be found at these three locations.

There are two minima in the probability function, located at $L / 3$ and $2 L / 3$, corresponding to locations where the particle is least likely to be found. In addition, the particle will not likely be found at the edges of the box: 0 and $L$.

EVALUATE It is important to remember that the probability function is the square of the wave function. Negative values of a wave function do not correspond to locations where the particle is least likely to be found. Here, the minimum of the wave function is one of the most likely places for the particle to be found.

## Problems

## 1: Electron in a well

Suppose an electron is confined to an infinitely deep, one-dimensional potential well of length $L$. Calculate the value of $L$ required for the frequency of a photon emitted in a transition from the $n=2$ state to the $n=1$ state to be equal to the frequency of a photon emitted in a transition from the $n=2$ state to the $n=1$ state in the Bohr hydrogen model.

## Solution

IDENTIFY AND SET UP We will use energy levels for a particle in a box and the Bohr model to solve the problem.

EXECUTE The energy levels for a particle in an infinitely deep well of length $L$ are given by

$$
E_{n}=\frac{n^{2} h^{2}}{8 m L^{2}}
$$

The energy needed for the transition between the $n=2$ and $n=1$ states is then

$$
\Delta E_{\mathrm{box}}=\frac{h^{2}}{8 m L^{2}}\left(2^{2}-1^{2}\right) .
$$

Solving for $L$, we have

$$
L=\sqrt{\frac{3 h^{2}}{8 m \Delta E}}
$$

We need the energy difference, which we find from the Bohr model. The energy for the $n$ level in the Bohr model is given by

$$
E_{n}=-\frac{h c R}{n^{2}}=-\frac{13.60 \mathrm{eV}}{n^{2}}
$$

Solving for the energy difference gives

$$
\Delta E_{\mathrm{H} \text { atom }}=E_{2}-E_{1}=\frac{-13.60 \mathrm{eV}}{2^{2}}-\frac{-13.60 \mathrm{eV}}{1^{2}}=10.2 \mathrm{eV}
$$

Substituting this result into the earlier result allows us to solve for $L$ :

$$
L=\sqrt{\frac{3\left(6.626 \times 10^{-34} \mathrm{Js}\right)^{2}}{8\left(9.1 \times 10^{-31} \mathrm{~kg}\right)(10.2 \mathrm{eV})\left(1.6 \times 10^{-19} \mathrm{~J} / \mathrm{eV}\right)}}=3.3 \times 10^{-10} \mathrm{~m}
$$

EVALUATE How does this result compare with the circumference of the first Bohr orbit? It approximately matches it, indicating that a square well potential serves as a reasonable model of the hydrogen atom.

## 2: Finding the wave function

Find the wave function of a particle in a box centered at $x=0$ with walls at $x= \pm L / 2$.

## Solution

IDENTIFY AND SET UP We will solve the Schrödinger equation to solve the problem. We will need to check the boundary conditions to ensure that our solution is correct.

EXECUTE The Schrödinger equation for a potential $U=0$ is given by

$$
-\frac{\hbar^{2}}{2 m} \frac{d^{2} \psi}{d x^{2}}=E \psi
$$

The solutions are given by

$$
\begin{aligned}
& \psi_{1}=\sin k x \\
& \psi_{2}=\cos k x .
\end{aligned}
$$

We substitute these solutions into the Schrödinger equation and find that

$$
E=\frac{\hbar^{2} k^{2}}{2 m}
$$

We now check the boundary conditions (i.e., the wave functions go to zero at the walls). For the first wave function, we have

$$
\psi_{1}(-L / 2)=\psi_{1}(L / 2)=0 .
$$

From this, we see that argument of the sin function must be integer multiples of $\pi$, or

$$
\begin{aligned}
& \frac{k L}{2}=m \pi, \quad m=1,2, \ldots \\
& k=\frac{2 m \pi}{L}=\frac{n \pi}{L} \quad \text { for even } n
\end{aligned}
$$

Applying the boundary condition to the second wave function gives

$$
\psi_{2}(-L / 2)=\psi_{2}(L / 2)=0
$$

This limits the values of $k$ to

$$
\begin{aligned}
& \frac{k L}{2}=\frac{n \pi}{2}, \\
& k=\frac{n \pi}{L} \quad \text { for odd } n .
\end{aligned}
$$

We use these results to find the energy levels. With the given $k$ 's, the energy levels are

$$
E_{n}=\frac{\hbar^{2} k^{2}}{2 m}=\frac{h^{2} n^{2}}{8 m L^{2}}
$$

The corresponding wave functions are

$$
\begin{array}{ll}
\psi_{n}=\sin \frac{n \pi}{L} x & \text { for even } n \\
\psi_{n}=\cos \frac{n \pi}{L} x & \text { for odd } n
\end{array}
$$

EVALUATE We see that the result is similar to the box wave functions between 0 and $L$, but with the wave functions displaced by $L / 2$.

## 3: Tunneling of a car

A car of mass 3000 kg rolls without friction on a level track and approaches a hill of height 1.0 m and width 1.0 m . It has enough kinetic energy so that it will rise to a height of 0.5 m and then return to the track. What is the probability that the car will tunnel through the hill?

## Solution

IDENTIFY AND SET UP We'll use the probability-of-tunneling function, equation 40.21 , in the text, to solve the problem.

EXECUTE Given kinetic energy $E$, the probability of tunneling through a barrier of height $U_{0}$ and width $L$ is given by the transmission coefficient, approximated by

$$
T=G e^{-2 \kappa L}
$$

where

$$
G=16 \frac{E}{U_{0}}\left(1-\frac{E}{U_{0}}\right), \quad \kappa=\frac{\sqrt{2 m\left(U_{0}-E\right)}}{\hbar} .
$$

For this problem,

$$
U_{0}=m g Y, \quad E=\frac{U_{0}}{2} .
$$

Substituting, we find that

$$
G=16 \frac{1}{2}\left(1-\frac{1}{2}\right)=4
$$

and

$$
\kappa=\frac{m \sqrt{g Y}}{\hbar} .
$$

The exponential factor becomes

$$
2 \kappa L=\frac{2 m \sqrt{g Y} L}{\hbar}=\frac{2(3000 \mathrm{~kg}) \sqrt{\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)(1.0 \mathrm{~m})}(1.0 \mathrm{~m})}{1.05 \times 10^{-34} \mathrm{~J} \mathrm{~s}}=1.8 \times 10^{38} .
$$

The transmission factor is therefore

$$
T=4 \frac{1}{e^{1.8 \times 10^{38}}} \ll 1 .
$$

evaluate Clearly, this probability is extremely small. It is so small that you would never expect to see the car tunneling through the hill, much as common sense tells us.

## Try It Yourself!

## 1: Particle in a well

A particle of mass $m$ is confined by a potential well of depth $V_{0}$ and width $d$. For a given value of $d$, what is the minimum value of $V_{0}$ necessary to confine the particle?

## Solution Checkpoints

IDENTIFY AND SET UP Use the uncertainty principle to solve the problem. Draw a sketch of the potential-energy function.

EXECUTE For the particle to be confined to the well, we must set the uncertainty in position to the width. The uncertainty in momentum is then

$$
\Delta p_{x}=\frac{\hbar}{\Delta x}=\frac{\hbar}{d}
$$

This gives a corresponding uncertainty in kinetic energy:

$$
\Delta E_{k}=\frac{\left(\Delta p_{x}\right)^{2}}{2 m}=\frac{\hbar^{2}}{2 m d^{2}}
$$

How much energy must the particle have to be confined? The total energy must be less than zero, or

$$
E=\Delta E_{k}-V_{0}<0
$$

Substituting into this expression gives the values of $V_{0}$ necessary for confinement:

$$
V_{0}>\frac{\hbar^{2}}{2 m d^{2}}
$$

The minimum value of $V_{0}$ occurs when it is equal to the expression.
EVALUATE What is the minimum depth required to produce a bound state for a proton in a nucleus $\left(d=3.0 \times 10^{-15} \mathrm{~m}\right) ? 2.3 \mathrm{MeV}$.

## 2: Electron in a box

An electron trapped in a box has a ground-state energy of 10 eV . How big is the box?

## Solution Checkpoints

IDENTIFY AND SET UP Use energy levels for a particle in a box.
EXECUTE The energy levels for a particle in an infinitely deep well of length $L$ are given by

$$
E_{n}=\frac{n^{2} h^{2}}{8 m L^{2}}
$$

The ground state corresponds to $n=1$. Substituting into the equation gives $L=1.94 \times 10^{-10} \mathrm{~m}$.
EVALUATE Does this result seem reasonable? Is it comparable to the size of the hydrogen atom?

## 41

## Atomic Structure

## Summary

In this chapter, we will apply our knowledge of quantum mechanics and the Schrödinger equation to atoms in order to understand their structure. We will leam how the quantization of angular momentum is a natural result of our investigation. In addition, we will see how atoms are described in terms of their quantum numbers, and we will learn that electrons have an intrinsic spin quantum number. We'll also learn how the Pauli exclusionary principle prevents two particles from occupying the same quantum mechanical state.

## Objectives

After studying this chapter, you will understand

- How to describe the states of the hydrogen atom in terms of quantum numbers.
- How the Zeeman effect describes the orbital motion of atomic electrons in a magnetic field.
- That electrons have intrinsic spin angular momentum.
- How to analyze the structure of many-electron atoms.
- How X-rays emitted by atoms unveil the inner structure of atoms.


## Concepts and Equations

| Term | Description |
| :---: | :---: |
| The Hydrogen Atom | The Schrödinger equation predicts the same energy levels as the Bohr model: $E_{n}=-\frac{1}{\left(4 \pi \mathcal{E}_{0}\right)^{2}} \frac{m_{r} e^{4}}{2 n^{2} \hbar^{2}}=-\frac{13.60 \mathrm{eV}}{n^{2}}$ <br> It also gives the possible magnitudes of orbital angular momentum as $L=\sqrt{l(l+1)} \hbar, \quad l=0,1,2, \ldots, n-1$ <br> and the $z$ component of the orbital angular momentum as $L_{z}=m_{l} \hbar, \quad m_{l}=0, \pm 1, \pm 2, \ldots, \pm l .$ |
| The Zeeman Effect | The interaction energy of an electron with magnetic quantum number $m_{l}$ in a magnetic field along the $+z$ axis is given by $U=-\mu_{z} B=m_{l} \frac{e \hbar}{2 m} B=m_{l} \mu_{B} B, \quad m_{l}=0, \pm 1, \pm 2, \ldots, \pm l,$ <br> where $v_{B}$ is the Bohr magneton. |
| Electron Spin | Electrons have intrinsic spin angular momentum of magnitude $S$, given by $S=\sqrt{\frac{1}{2}\left(\frac{1}{2}+1\right)} \hbar=\sqrt{\frac{3}{4}} \hbar$ <br> The $z$ component of the spin angular momentum has values $S_{z}=m_{s} \hbar, \quad m_{s}= \pm \frac{1}{2} .$ |
| Electrons in Atoms | In a hydrogen atom, the quantum numbers $\left(n, l, m_{l}, m_{s}\right)$ specify the quan-tum-mechanical state of the atom and have allowed values given by $\begin{gathered} n \geq 1, \quad 0 \leq l \leq n-1, \\ \left\|m_{l}\right\| \leq l, \quad m_{s}= \pm \frac{1}{2} . \end{gathered}$ |
| X-ray Spectra | Moseley's law states that the frequency of the $K_{\alpha} \mathrm{X}$-ray from a target with atomic number $Z$ is given by $f=\left(2.48 \times 10^{15} \mathrm{~Hz}\right)(Z-1)^{2} .$ |

## Conceptual Questions

## 1: Atoms without the Pauli exclusion principle

What would the electron configuration of the ground state of calcium be if the Pauli exclusion principle did not hold?

## Solution

IDENTIFY, SET UP, AND EXECUTE Without the exclusion principle, all electrons could occupy the same lowest energy state. The ground state of calcium would be $1 s^{20}$.

EVALUATE Without the Pauli exclusion principle, all ground-state atoms would have all their electrons in the $1 s$ state. The world, and the universe, would be a rather boring place, since all atoms would have similar chemical properties.

## 2: Identify the atom

Determine the element corresponding to the following ground-state electron configurations:
(a) $\quad 1 s^{2} 2 s^{2} 2 p^{6} 3 s^{2} 3 p^{3}$
(b) $\quad 1 s^{2} 2 s^{2} 2 p^{6} 3 s^{2} 3 p^{6} 4 s$
(c) $1 s^{2} 2 s^{2} 2 p^{6} 3 s^{2} 3 p^{6} 3 d^{3} 4 s^{2}$

## Solution

IDENTIFY, SET UP, AND EXECUTE We could determine the elements in all of the preceding cases by comparing the configurations against Table 41.3 in the textbook. A simpler solution is found by counting the electrons in each configuration and finding the element with the correct $Z$.

Element (a) has 15 electrons and is phosphorus. Element (b) has 19 electrons and is potassium. Element (c) has 23 electrons and is vanadium.

EVALUATE As we become familiar with electron configurations, the problems become easier. Our next step would be to examine excited states of elements. How could we identify the element in those situations? We would follow the same procedure: Count the electrons and correlate with the $Z$. The excited states contain the same number of electrons, with some of the electrons occupying higher substates.

## Problems

## 1: Possible states of hydrogen

An electron is in the hydrogen atom with $n=6$. Find the possible values of $L$ and $L_{z}$ for this electron.

## Solution

IDENTIFY AND SET UP For $n=6$, the largest possible value of $l$ is 5 and the largest positive value of $m_{z}$ is 5 . The possible values of $L$ and $L_{z}$ are found from the angular momentum relations.

EXECUTE The values of $L$ and $L_{z}$ are given by

$$
L=\sqrt{l(l+1)} \hbar, \quad L_{z}=m_{l} \hbar
$$

In this problem, $l$ ranges from 0 to 5 and $m_{l}$ ranges from 0 to $\pm 5$. The values are then

$$
\begin{array}{lll}
L=\sqrt{l(l+1)} \hbar=0, & L_{z}=m_{l} \hbar=0 & (n=0), \\
L=\sqrt{l(l+1)} \hbar=\sqrt{1(1+1)} \hbar=\sqrt{2} \hbar, & L_{z}=0, \pm \hbar & (n=1), \\
L=\sqrt{l(l+1)} \hbar=\sqrt{2(2+1)} \hbar=\sqrt{6} \hbar, & L_{z}=0, \pm \hbar, \pm 2 \hbar & (n=2), \\
L=\sqrt{l(l+1)} \hbar=\sqrt{3(3+1)} \hbar=\sqrt{12} \hbar, & L_{z}=0, \pm \hbar, \pm 2 \hbar, \pm 3 \hbar & (n=3), \\
L=\sqrt{l(l+1)} \hbar=\sqrt{4(4+1)} \hbar=\sqrt{20} \hbar, & L_{z}=0, \pm \hbar, \pm 2 \hbar, \pm 3 \hbar, \pm 4 \hbar, & (n=4), \\
L=\sqrt{l(l+1)} \hbar=\sqrt{5(5+1)} \hbar=\sqrt{30} \hbar, & L_{z}=0, \pm \hbar, \pm 2 \hbar, \pm 3 \hbar, \pm 4 \hbar, \pm 5 \hbar & (n=5) .
\end{array}
$$

EVALUATE As the principal quantum number increases, the number of possible states increases rapidly.

## 2: Electron configuration of gallium

Write the ground-state electron configuration for gallium. What next-smaller and next-larger Z's have chemical properties similar to those of gallium?

## Solution

IDENTIFY AND SET UP Gallium has an atomic number of 31, so we must fill the lowest 31 electron states. Each $s$ subshell can accommodate 2 electrons, each $p$ substate can accommodate 6 electrons, and the $d$ subshells can accommodate 10 electrons.

EXECUTE Gallium's $1 s, 2 s$, and $2 p$ subshells hold $2+2+6=10$ electrons. The $n=3$ subshells$3 s, 3 p$, and $3 d$-hold $2+6+10=18$ electrons, for a total of 28 electrons in the first 6 subshells. This leaves three electrons. These electrons go into the $4 s$ and $4 p$ subshells: two in the $4 s$ subshell and one in the $4 p$ subshell.

The next-lower $Z$ with chemical properties similar to those of gallium is aluminum since its outer shells are filled with two electrons in the $3 s$ subshell and one in the $3 p$ subshell. The next-larger $Z$ with chemical properties similar to those of gallium is indium, since its outer shells are filled with two electrons in the $5 s$ subshell and one in the $5 p$ subshell.

EVALUATE We can see that gallium, aluminum, and indium are chemically similar, since they occupy different rows of the same column in the periodic table.

## 3: Zeeman effect

The difference in energies of a hypothetical atom between its $2 p$ and $3 s$ levels is 1.2 eV . How large a magnetic field would be required to raise the energy of the highest possible state of the $2 p$ level to that of the lowest possible $3 s$ state due to the electron spin energies?

## Solution

IDENTIFY AND SET UP The change in energy for a level due to electron spin energies is given by the Zeeman effect. The highest-energy $2 p$ state is the energy of the $2 p$ state plus the energy difference due to the spin energy. The lowest-energy $3 s$ state is the energy of the $3 s$ state minus the energy difference due to the spin energy. We will set these equal to each other to solve the problem.

EXECUTE The energy difference due to a magnetic moment in a magnetic field is given by

$$
U=-\vec{\mu} \cdot \vec{B}
$$

The highest-possible-energy $2 p$ state in a magnetic field is

$$
E_{2 p}=E_{2 p}(0)+\mu_{B} B
$$

where $E_{2 \mathrm{p}}(0)$ is the energy of the $2 p$ state without the electron spin interaction. The lowest-possibleenergy $3 s$ state in a magnetic field is

$$
E_{3 s}=E_{3 s}(0)-\mu_{B} B
$$

We set these energies equal to each other, giving

$$
\begin{aligned}
& E_{2 p}=E_{3 s}, \\
& E_{2 p}(0)+\mu_{B} B=E_{3 s}(0)-\mu_{B} B .
\end{aligned}
$$

The energy difference between the $2 p$ and $3 s$ states is 1.2 eV , so we solve for that energy difference:

$$
E_{3 s}(0)-E_{2 p}(0)=2 \mu_{B} B=1.2 \mathrm{eV}
$$

Solving for the magnetic field, we have

$$
B=\frac{1.2 \mathrm{eV}}{2 \mu_{B}}=\frac{(1.2 \mathrm{eV})\left(1.6 \times 10^{-19} \mathrm{~J} / \mathrm{eV}\right)}{2\left(9.27 \times 10^{-24} \mathrm{~J} / \mathrm{T}\right)}=1.04 \times 10^{4} \mathrm{~T}
$$

EVALUATE The magnetic field required is enormous, larger than the largest steady-state field produced by a superconducting magnet. This problem illustrates how electron spin energy is relatively small, although it can be measured (albeit indirectly).

## 4: Calculating energy differences from X-ray transitions

The X-ray transitions $K_{\alpha}$ and $L_{\alpha}$ are shown in Figure 41.1. The energies of the X-ray photons emitted in those transitions are shown for five elements in Table 41.1. Calculate the energy differences between the $n=2$ and $n=3$ levels and the $n=3$ and $n=2$ levels, using the data provided.


Figure 41.1 Problem 4.

| Element | $\boldsymbol{Z}$ | $\boldsymbol{K}_{\boldsymbol{\alpha}} \mathbf{( k e V )}$ | $\boldsymbol{L}_{\boldsymbol{\alpha}} \mathbf{( k e V )}$ |
| :---: | :---: | :---: | :---: |
| Mn | 25 | 6.51 | 0.721 |
| Zn | 30 | 9.57 | 1.11 |
| Br | 35 | 13.3 | 1.60 |
| Zr | 40 | 17.7 | 2.06 |
| Rh | 45 | 22.8 | 2.89 |

Table 41.1 Problem 4.

## Solution

IDENTIFY AND SET UP To solve the problem, we will use the definition of X-ray transitions. The $K_{\alpha}$ line arises from the energy difference between the $n=3$ and $n=1$ states, so it is equal to $E_{3}-E_{1}$. The $L_{\alpha}$ line arises from the energy difference between the $n=3$ and $n=2$ states, so it is equal to $E_{3}-E_{2}$. We will combine these results to find the Solution.

EXECUTE The energy difference between the $n=2$ and $n=1$ states is found by combining information from both lines. Specifically,

$$
E_{2}-E_{1}=\left(E_{3}-E_{1}\right)-\left(E_{3}-E_{2}\right)=E_{K_{\alpha}}-E_{L_{\alpha}} .
$$

The energy difference between the $n=3$ and $n=2$ states is found directly from the $L_{\alpha}$ line. The results for the energy differences are therefore as follows:

| Element | $\boldsymbol{E}_{\mathbf{3}}-\boldsymbol{E}_{\mathbf{2}}(\mathbf{k e V})$ | $\boldsymbol{E}_{\mathbf{3}}-\boldsymbol{E}_{\mathbf{2}}(\mathbf{k e V})$ |
| :---: | :---: | :---: |
| Mn | 5.79 | 0.721 |
| Zn | 8.46 | 1.11 |
| Br | 11.7 | 1.60 |
| Zr | 15.6 | 2.06 |
| Rh | 19.9 | 2.89 |

Table 41.1 Problem 4.1
EVALUATE Examining the trends in these elements, we see that the energy differences increase with increasing atomic number. What does that tell us about atomic structure?

## Try It Yourself!

## 1: Filling shells

Write down the expected electron configuration for hydrogen atom states (a) for an atom with 18 electrons and (b) for an atom with 22 electrons. (c) If an atom has 22 electrons, how many additional electrons would be required to complete a closed shell?

## Solution Checkpoints

IDENTIFY AND SET UP Fill the electron shells from lowest to highest. How many electrons fit in each shell?

EXECUTE (a) Two electrons will fill the $1 s$ shell, and then two electrons will fill the $2 s$ shell. Continue filling until you have

$$
1 s^{2} 2 s^{2} 2 p^{6} 3 s^{2} 3 p^{6}
$$

for the 18 -electron atom.
(b) Starting with the 18 electrons of part (a), add 4 more electrons to the next shell. This will give

$$
1 s^{2} 2 s^{2} 2 p^{6} 3 s^{2} 3 p^{6} 3 d^{4}
$$

for the 22-electron atom.
(c) The $3 d$ shell holds 10 electrons and the 22 -electron atom has 4 atoms in the $3 d$ shell, so 6 additional electrons are required to fill the $3 d$ shell completely.

EVALUATE Remember that each shell hold $2(2 l+1)$ electrons.

## 2: Classical versus quantum physics

(a) Calculate the classical angular precession frequency of an electron in a constant magnetic field. (b) If the electron's angular momentum is quantized, find the allowed values of the possible orbits. (c) For an electron spin in a constant magnetic field, find the difference in energy of its two states and the corresponding angular frequency.

## Solution Checkpoints

IDENTIFY AND SET UP The classical precession frequency is found by looking at the forces acting on the electron. Quantized angular momentum requires that the angular momentum have only discrete values.

EXECUTE (a) Classically, an electron is acted upon by a magnetic force in a constant magnetic field and follows a circular path. This gives

$$
e v B=\frac{m v^{2}}{R} .
$$

The angular frequency is then

$$
\omega_{\mathrm{c}}=\frac{e B}{m} .
$$

(b) For the angular momentum to be quantized, we must have

$$
m v R=n \hbar .
$$

Using the results from part (a), we find that the allowed values of the radius are

$$
R_{n}=\sqrt{\frac{n \hbar}{e B}}
$$

for integer values of $n$.
(c) The energy difference between the two states is twice the value of the Zeeman effect:

$$
\Delta E=2 \mu_{B} B .
$$

This corresponds to a frequency difference of

$$
f=\frac{\Delta E}{h}=\frac{2 \mu_{B} B}{h} .
$$

EVALUATE How do the two frequencies compare? Check that they are the same by converting the classical angular frequency to ordinary frequency and replacing the Bohr magneton with its definition.

## Molecules and Condensed Matter

## Summary

In this chapter, we will extend our application of quantum mechanics from atoms to molecules and larger structures of atoms. We will investigate chemical behavior, molecular bonds, and the large-scale assembly of atoms into crystalline solids. Our quantum-mechanical foundation will allow us to examine semiconductors and superconductors, two types of materials having profound effects on science and society today.

## Objectives

After studying this chapter, you will understand

- Forms of bonds that hold atoms together.
- The rotational and vibrational dynamics of molecules.
- How atoms form crystalline structures.
- How to use the energy-band structure to explain electrical properties of solids.
- The basis of semiconducting materials and how semiconductor devices operate.
- The basis of superconducting materials


## Concepts and Equations

| Term | Description |
| :--- | :--- |
| Molecular Bonds and <br> Molecular Spectra | Molecules bind through ionic, covalent, van der Waals, and hydrogen bonds. <br> For a diatomic molecule, the rotational energy levels are given by |
| $\qquad$$E_{l}=l(l+1) \frac{\hbar^{2}}{2 I} \quad l=0,1,2, \ldots$ <br> $I=m_{r} r_{0}^{2}, \quad m_{r}=\frac{m_{1} m_{2}}{m_{1}+m_{2}}$, |  |

where $I$ is the moment of inertia of the molecule, $m_{r}$ is the reduced mass of the molecules, and $r_{0}$ is the distance between the atoms. The vibrational levels are given by

$$
E_{n}=\left(n+\frac{1}{2}\right) \hbar \omega=\left(n+\frac{1}{2}\right) \hbar \sqrt{\frac{k^{\prime}}{m_{r}}} \quad n=0,1,2, \ldots
$$

## Free-electron Model of Metals

In the free-electron model of metals, electrons are treated as free particles in a conductor. The density of states is given by

$$
g(E)=\frac{(2 m)^{3 / 2} V}{2 \pi^{2} \hbar^{3}} E^{1 / 2}
$$

The probability that an energy state $E$ is occupied is given by the Fermi-Dirac distribution,

$$
f(E)=\frac{1}{e^{\left(E-E_{F}\right) k T}+1}
$$

where $E_{F}$ is the Fermi energy.

## Semiconductors

## Semiconducting Devices

A semiconductor is a material with electrical resistivity intermediate between that of a good conductor and a good insulator. In $n$-type semiconductors, the conductivity is due to the motion of electrons. In $p$-type semiconductors, the conductivity is due to the motion of holes, or vacancies of electrons. Semiconductors have energy gaps of about 1 eV between their valence and conduction bands.

A diode is made of two semiconductors, one $p$-type and one $n$-type, that can behave much like a switch, conducting above a threshold voltage and insulating below. A transistor is made of two $p-n$ junctions. The current-voltage relationship for an ideal $p-n$ junction is given by

$$
I=I_{S}\left(e^{e V k T}-1\right)
$$

## Conceptual Questions

## 1: Types of bonds

What kind of chemical bonds hold the following objects together: (a) NaCl molecules, (b) $\mathrm{N}_{2}$ molecules, and (c) copper atoms in a wire.

## Solution

IDENTIFY, SET UP, AND EXECUTE (a) NaCl molecules are made of oppositely charged atoms, so ionic bonding holds the molecules together. The sodium atom gives its one $3 s$ electron to the chlorine atom, filling a vacancy in the chlorine's $3 p$ shell.
(b) Since two nitrogen atoms form a nitrogen molecule, the molecule is made of two similarly charged atoms and cannot be bound by an ionic bond. The nitrogen molecule is bound by a covalent bond.
(c) The copper atoms in a wire are bound by metallic bonds, since the copper is arranged in a crystalline structure in the wire.

EVALUATE This small sample of problems involving bonds illustrates the variations on chemical bonding in molecules and solids.

## 2: Covalent bonds without electron spin

How many electrons would participate in a covalent bond if electrons had no spin?

## Solution

IDENTIFY, SET UP, AND EXECUTE If electrons had no spin, only one electron could occupy a particular space, due to the exclusion principle. The usual covalent bond has two electrons sharing the same spatial state, but not the same spin state, in order to satisfy the exclusion principle. Without spin, only one electron could participate in a covalent bond.

EVALUATE Spin plays an important role in atomic and molecular physics, as we see here in its role in covalent bonds. This exercise should elucidate the role spin plays in molecular binding. True spinless particles exist, but they do not obey the exclusion principle.

## Problems

## 1: Predicting states

Light of wavelength $5.0 \mu \mathrm{~m}$ strikes and is absorbed by a molecule. Is this process most likely to alter the rotational, vibrational, or atomic energy levels of the molecule? If the light had a wavelength of 3.7 mm , which energy levels would most likely be affected?

## Solution

IDENTIFY AND SET UP The energy difference for atomic energy levels is generally several eVs, for vibrational levels is generally several 0.1 eVs , and for rotational levels is generally several 0.001 eVs . We will calculate the energy of the light to determine which type of transition it may affect.

EXECUTE The energy of a photon is

$$
E=\frac{h c}{\lambda} .
$$

Since we are calculating electron-volts, it is convenient to use $h=4.136 \times 10^{-15} \mathrm{eVs}$. The energy of the $5.0-\mu \mathrm{m}$ light is then

$$
E=\frac{h c}{\lambda}=\frac{\left(4.136 \times 10^{-15} \mathrm{~J} \cdot \mathrm{~s}\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}{5.0 \times 10^{-6} \mathrm{~m}}=0.25 \mathrm{eV}
$$

0.25 eV corresponds to the energy difference of a vibrational transition. For the $3.7-\mathrm{mm}$ light, the energy is

$$
E=\frac{h c}{\lambda}=\frac{\left(4.136 \times 10^{-15} \mathrm{~J} \cdot \mathrm{~s}\right)\left(3.00 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}{3.7 \times 10^{-3} \mathrm{~m}}=0.0003 \mathrm{eV}
$$

0.0003 eV corresponds to the energy difference of a rotational transition.

The $5.0 \mu \mathrm{~m}$ light should affect the vibrational level, and the $3.7-\mathrm{mm}$ light should affect the rotational level.

EVALUATE These different ranges of transition energies indicate the likelihood of causing a molecule to rotate, vibrate, or change its atomic state. The situation corresponds to a similar one in the classical world, where it is easy to rotate an object, harder to induce a vibration, and much harder to have it change state.

Practice Problem: If the light had a wavelength of 140 nm , which energy levels would most likely be affected? Answer: Light of this wavelength has an energy of 8.9 eV , corresponding to an atomic transition.

## 2: A diatomic molecule

The $\mathrm{D}_{2}$ molecule is made up of two deuterons (a proton plus a neutron) and two electrons. The spacing between the nuclei in the molecule is approximately $7.5 \times 10^{-11} \mathrm{~m}$. Calculate the moment of inertia of the $\mathrm{D}_{2}$ molecule about its center of mass and the rotational energy of its ground state and first two excited states.

## Solution

IDENTIFY AND SET UP We will use the formulae for the moment of inertia and rotational energy to solve the problem. We will ignore the masses of the electrons, as they are negligible compared with the mass of the nuclei. We also will take the mass of the neutron to be that of the proton, as they have similar masses.

EXECUTE The moment of inertia is

$$
I=\sum m_{i} r_{i}^{2}
$$

There are four masses, each located at a point half the interatomic distance from the center of mass. This gives

$$
\begin{aligned}
I & =2 m_{\mathrm{H}}\left(\frac{d}{2}\right)^{2}+2 m_{\mathrm{H}}\left(\frac{d}{2}\right)^{2}=m_{\mathrm{H}} d^{2} \\
& =\left(1.67 \times 10^{-27} \mathrm{~kg}\right)\left(7.5 \times 10^{-11} \mathrm{~m}\right)^{2}=9.39 \times 10^{-48} \mathrm{~kg} \mathrm{~m}^{2}
\end{aligned}
$$

The energy levels for rotational states are given by

$$
\begin{aligned}
E & =l(l+1) \frac{(\hbar)^{2}}{2 I} \\
& =l(l+1)\left(3.70 \times 10^{-3} \mathrm{eV}\right)
\end{aligned}
$$

The ground-state energy has $l=0$, so $E_{0}=0$; the first excited-state energy has $l=1$, so $E_{1}=7.40 \times 10^{-3} \mathrm{eV}$; and the second excited-state energy has $l=2$, so $E_{2}=2.22 \times 10^{-3} \mathrm{eV}$.

EVALUATE What frequency of light will be emitted when the first excited state decays to the ground state?

## 3: Fermi energy of a particle in a box

Suppose the energy values for a particle in a box are given by

$$
E_{n}=E_{0} n^{2},
$$

where $E_{0}$ is a constant and $n$ is the quantum number of the state. (a) If there are 50 electrons in such states, find the Fermi energy at a temperature of 0 K . (b) What is the ratio of the average energy of the electrons to the Fermi energy?

## Solution

IDENTIFY AND SET UP The Fermi energy is the energy of the last filled state, so we will find the energy of the highest filled state to solve part (a). To solve part (b), we'll add up the energies of all the states and take the ratio.

EXECUTE (a) The 50 electrons will occupy 25 states, since 2 electrons fill each state. The energy of the 25th state is

$$
E_{F 0}=E_{0}(25)^{2}=625 E_{0} .
$$

(b) The total energy is found by summing the contributions from each state, given by

$$
E_{T}=2 E_{0}[1+4+9+\cdots]=2 E_{0} \sum_{n=1}^{25} n^{2} .
$$

We can solve this equation by summing up the 25 terms, or we can use a summation rule. The sum over $N$ squares is given by

$$
\sum_{n=1}^{N} n^{2}=\frac{N(N+1)(2 N+1)}{6} .
$$

Evaluating this sum for $N=25$, we have

$$
E_{T}=2 E_{0} \sum_{n=1}^{25} n^{2}=2 E_{0} \frac{25(25+1)(2(25)+1)}{6}=11,050 E_{0} .
$$

The average energy is the total energy divided by 50 :

$$
E_{\mathrm{avg}}=\frac{E_{T}}{50}=\frac{11,050}{50} E_{0}=221 E_{0}=\left(\frac{221}{625}\right) E_{\mathrm{F} 0} .
$$

The ratio of the average energy per electron to the Fermi energy is 0.354 .
evaluate What does the ratio of the average energy per electron to the Fermi energy tell us about the system? Since the ratio is about $1 / 3$, it indicates that many of the electrons are in energy states having less than half the Fermi energy.

## 4: Fermi-Dirac statistics

Given the Fermi-Dirac probability distribution and the fact that sodium has a Fermi energy of 3.15 eV , find the ratio of the width $\Delta E$ to $E_{F}$ at 273 K , where $\Delta E=E(0.2)-E(0.8)$. The quantity $E(0.2)$ is the energy for which the occupation probability is 0.2 .

## Solution

IDENTIFY AND SET UP Use the Fermi-Dirac probability function to solve the problem.
EXECUTE The Fermi-Dirac probability function $f(E)$ is given by

$$
f(E)=\frac{1}{e^{\left(E-E_{F}\right) / k T}+1}
$$

Rearranging terms to solve for $E$ gives

$$
E=E_{F}+k T \ln \left(\frac{1-f(E)}{f(E)}\right)
$$

Solving for $E(0.2)$ and $E(0.8)$ yields

$$
\begin{aligned}
& E(0.2)=E_{F}+k T \ln \left(\frac{1-0.2}{0.2}\right)=E_{F}+k T \ln 4 \\
& E(0.8)=E_{F}+k T \ln \left(\frac{1-0.8}{0.8}\right)=E_{F}-k T \ln 4
\end{aligned}
$$

Forming the ratio requested in the problem gives

$$
\frac{\Delta E}{E_{F}}=\frac{2 k T \ln 4}{E_{F}}=\frac{2\left(8.617 \times 10^{-5} \mathrm{eV} / \mathrm{K}\right)(273 \mathrm{~K}) \ln 4}{(3.15 \mathrm{eV})}=0.02
$$

EVALUATE What does this ratio tell us about the system? Near room temperature, changing the energy by $2 \%$ of the Fermi energy changes the occupation probability for sodium from $80 \%$ to $20 \%$. This is a very dramatic change over a small distance (i.e. the Fermi distribution function is very sharp).

## Try It Yourself!

## 1: Rotational transitions

(a) If a molecule with moment of inertia $I$ is induced to make a pure rotational transition from a state $L$ to a state $L+1$, what frequency of radiation is needed? (b) If the same molecule makes a transition from the state $L$ to the state $L-1$, what frequency of radiation is emitted?

## Solution Checkpoints

IDENTIFY AND SET UP Use rotational energy levels to solve the problem.
EXECUTE (a) The rotational energy levels are given by

$$
E_{l}=l(l+1) \frac{\hbar^{2}}{2 I} .
$$

To go from state $L$ to state $L+1$, the energy $h f$ supplied must be equal to

$$
\begin{aligned}
h f & =E_{L+1}-E_{L} \\
& =2(L+1) \frac{\hbar^{2}}{2 I} .
\end{aligned}
$$

(b) To go from state $L$ to state $L-1$, the energy $h f$ released is equal to

$$
\begin{aligned}
h f^{\prime} & =E_{L}-E_{L-1} \\
& =2 L \frac{\hbar^{2}}{2 I} .
\end{aligned}
$$

EVALUATE How do the spacing between different possible values of $f$ and $f^{\prime}$ compare?

## 2: Promoting to the conduction band

Electromagnetic radiation can promote an electron from the top of a nearly complete valence band to the bottom of an unfilled conduction band. The lowest frequency for which this is possible is $f_{m}=2.75 \times 10^{14} \mathrm{~Hz}$ for silicon and $f_{m}=1.79 \times 10^{14} \mathrm{~Hz}$ for germanium, both at room temperature. Calculate the energy gap between the valence and conduction bands for the two materials.

## Solution Checkpoints

IDENTIFY AND SET UP What equation relates energy and frequency for electromagnetic radiation?
EXECUTE The energy gap is found from the frequency by the relationship

$$
\Delta E=h f_{m}=E_{\text {gap }} .
$$

Evaluating this equation gives an energy gap of 1.13 eV for silicon and an energy gap of 0.736 eV for germanium.

EVALUATE What practical purpose can this material carry out? Could it be used to detect light? How?

## Nuclear Physics

## Summary

In this chapter, we look deep inside the atom into the nucleus and the particles that make up the nucleus. We will investigate how protons and neutrons combine to influence the size, mass, and stability of a nucleus. We will then turn to radioactive decay and nuclear reactions to examine the processes that shape our understanding of nuclei. Beneficial uses of radioactivity, fission, and fusion will also be examined.

## Objectives

After studying this chapter, you will understand

- How protons and neutrons form nuclei.
- Some fundamental properties of nuclei, including radii, densities, spins, and magnetic moments.
- How binding energies depend on the number of protons and neutrons in the nucleus.
- How nuclei undergo radioactive decay and how they decay at different rates.
- The basics of nuclear reactions and how to predict reaction energies.

Concepts and Equations

| Term | Description |
| :--- | :--- |
| Properties of the Nucleus | A nucleus is composed of $A$ nucleons $(Z$ protons and $N$ neutrons $)$, is roughly <br> spherical in shape, and has a radius that depends on $A$ according to the formula |
|  | $R=R_{0} A^{1 / 3}$, |
| where $R_{0}=1.2 \mathrm{fm}$. |  |

where $z$ is the number of protons in the nucleus, $M_{\mathrm{H}}$ is atomic mass of neutral hydrogen, $N$ is the number of neutrons in the nucleus, $m_{n}$ is the mass of a neutron, ${ }_{2}^{A} M$ is the mass of the neutral atom containing the nucleus, and $c$ is the speed of light in vacuum. Nuclei are held together by the attractive nuclear force overcoming the repulsive electrical force. The nuclear force favors paired nucleons with opposite spin and pairs of pairs. Most stable nuclei have more neutrons than protons. Both the shell model and the liquid drop model are used to describe the properties of the nucleus.

| Radioactivity and Radiation | Unstable nuclei undergo radioactive decay, most commonly through alpha <br> $(\alpha)$ and beta $(\beta)$ particle emission and sometimes through gamma-ray emis- |
| :--- | :--- |
| sion. Alpha particles are two protons and two neutrons bound together. The |  |
| most common beta particles (beta-minus particles) are electrons. Radioactive |  |
| material decays exponentially. The number $N$ of nuclei remaining in a sample |  |
| of $N_{0}$ nuclei after time $t$ is |  |

where $\lambda$ is the decay constant for the particular nuclear species in question. The rate of decay is described by the decay constant $\lambda$, the half-life $T_{1 / 2}$, or the lifetime $T_{\text {mean }}$, related by

$$
T_{\text {mean }}=\frac{1}{\lambda}=\frac{T_{1 / 2}}{\ln 2}=\frac{T_{1 / 2}}{0.693} .
$$

Radioactivity can be used beneficially to date artifacts or to diagnose and treat medical conditions. Radioactivity also can be harmful to human tissue, as it can ionize cellular material, but radiation hazards can be reduced through proper precautions.

Nuclear reactions result from the bombardment of a nucleus by a particle. Energy is exchanged in nuclear reactions. A reaction resulting in an excess of kinetic energy is called an exoergic reaction, and a reaction resulting in a deficiency of kinetic energy is called an endoergic reaction.

Fission is the radioactive decay of an unstable nucleus into two or more nuclei. Fission is used to power nuclear reactors. Fusion is a nuclear reaction in which two or more light nuclei combine to form a larger nucleus plus excess energy.

## Conceptual Questions

## 1: Size of nuclei

How much must the mass number of a nucleus increase to double the volume of the nucleus? To double its radius?

## Solution

IDENTIFY, SET UP, AND EXECUTE The radius of a nucleus is proportional to the cube root of the mass number. Volume is proportional to the cube of the radius. To double the volume of a nucleus, the mass number must double.

To double the radius, the mass number must increase by a factor of $2^{3}=8$.
EVALUATE Changes in radius and volume give indications of the structure of nuclei. The relation of the radius to the volume indicates that the nuclear density is roughly constant for all nuclei.

## 2: Causing nuclear fusion

Deuterium nuclei can fuse, liberating energy. Why doesn't a container of deuterium gas begin fusion after being shaken?

## Solution

IDENTIFY, SET UP, AND EXECUTE The deuterium nuclei can fuse, but only when the nuclei contact each other. Deuterium nuclei in deuterium gas are surrounded by electron clouds, and Coulomb repulsion prevents the nuclei from coming into contact with each other. For the fusion reaction to take place, the deuterium nuclei must be given enough energy to overcome the Coulomb repulsion.

EVALUATE The biggest barrier to fusion reactions being used as a source of energy is the electric repulsion between charged nuclei. Even bare nuclei without electron clouds must have enough kinetic energy to overcome that repulsion in order for fusion to occur.

## Problems

## 1: Binding energy of silver and gold

(a) Calculate the total binding energy of ${ }^{107} \mathrm{Ag}$ (atomic mass 106.905097 u ) and ${ }^{197} \mathrm{Au}$ (atomic mass 196.966569 u ). (b) Calculate the binding energy per nucleon for each atom. (c) Which of these nuclei is more tightly bound?

## Solution

IDENTIFY We find the binding energy by converting the mass defect into energy.
SET UP The nucleon number is the number of nucleons, so we will divide the binding energy by 107 and 197 to find the respective binding energies per nucleon.

EXECUTE The mass defect is calculated by subtracting the atomic mass from the combined mass of the protons and neutrons in the nucleus:

$$
\Delta M=Z m_{p}+N m_{n}-M=Z(1.007825 \mathrm{u})+N(1.008665 \mathrm{u})-M
$$

${ }^{107} \mathrm{Ag}$ has 47 protons and 60 neutrons, for a total of 107 nucleons. Its mass defect is

$$
\Delta M=(47)(1.007825 \mathrm{u})+(60)(1.008665 \mathrm{u})-(106.905097 \mathrm{u})=0.9826 \mathrm{u}
$$

Its binding energy is

$$
E_{B}=(0.9826 \mathrm{u})(931.5 \mathrm{MeV} / \mathrm{u})=915.3 \mathrm{MeV}
$$

Its binding energy per nucleon is

$$
E_{B} / A=915.3 \mathrm{MeV} / 107=8.55 \mathrm{MeV} / \text { nucleon } .
$$

${ }^{197} \mathrm{Au}$ has 79 protons and 118 neutrons, for a total of 197 nucleons. Its mass defect is

$$
\Delta M=(79(1.007825 \mathrm{u})+(118)(1.008665 \mathrm{u})-(196.966569 \mathrm{u})=1.6741 \mathrm{u} .
$$

Its binding energy is

$$
E_{B}=(1.6741 \mathrm{u})(931.5 \mathrm{MeV} / \mathrm{u})=1559.4 \mathrm{MeV}
$$

Its binding energy per nucleon is

$$
E_{B} / A=1559.4 \mathrm{MeV} / 197=7.92 \mathrm{MeV} / \text { nucleon. }
$$

The silver atom's nucleons are more tightly bound, since they require more energy per nucleon to become free.

EVALUATE The total binding energy is larger for gold than for silver, but the binding energy per nucleon is less for gold. This state of affairs agrees with our model of the nucleus: Atoms with larger numbers of nucleons have greater radii, causing their nucleons to be less strongly bound than atoms with smaller numbers of nucleons.

## 2: Examining a nuclear reaction

Calculate the energy released or absorbed by the reaction ${ }_{2}^{3} \mathrm{He}+{ }_{1}^{2} \mathrm{H} \rightarrow{ }_{2}^{4} \mathrm{He}+{ }_{1}^{1} \mathrm{H}$.

## Solution

IDENTIFY AND SET UP We find the reaction energy by subtracting the rest masses of the products from the rest masses of the initial particles. Table 43.2 in the text provides the needed rest masses.

EXECUTE The reaction energy is

$$
Q=M_{A}+M_{B}-M_{C}-M_{D},
$$

where $M_{A}$ and $M_{B}$ are the rest masses of ${ }_{2}^{3} \mathrm{He}$ and ${ }_{1}^{2} \mathrm{H}$, respectively, and $M_{C}$ and $M_{D}$ are the rest masses of ${ }_{2}^{4} \mathrm{He}$ and ${ }_{1}^{1} \mathrm{H}$, respectively. Table 30.2 in the text gives the rest masses:

$$
\begin{array}{ll}
{ }_{2}^{3} \mathrm{He}: & 3.016029 \mathrm{u}, \\
{ }_{1}^{2} \mathrm{H}: & 2.014101 \mathrm{u}, \\
{ }_{2}^{4} \mathrm{He}: & 4.002603 \mathrm{u}, \\
{ }_{1}^{1} \mathrm{H}: & 1.007825 \mathrm{u} .
\end{array}
$$

The change in mass is

$$
Q=(3.016029 \mathrm{u})+(2.014101 \mathrm{u})-(4.002603 \mathrm{u})-(1.007825 \mathrm{u})=0.0197 \mathrm{u}
$$

The mass has decreased, indicating that energy is released. The reaction energy is

$$
Q=(0.0197 \mathrm{u})(931.5 \mathrm{MeV} / \mathrm{u})=18.35 \mathrm{MeV}
$$

A total of 18.35 MeV is released in the reaction.
EVALUATE Since the total mass decreases, the kinetic energy of the products is more than the initial energy of the interacting nuclei. This fusion reaction is one candidate being considered for the production of fusion energy.

We also should check that charge is conserved in the reaction. The charge of the incoming nuclei is $+3 e$, and the charge of the product nuclei is $+3 e$, confirming that charge is conserved in the reaction.

## 3: Radioactivity

At a certain time, a sample of radioactive material is measured and is found to decay at a rate of 32 counts per second. Two hours later, the sample is measured to decay at a rate of 13 counts per second. What is the half-life of the sample?

## Solution

IDENTIFY AND SET UP Radioactivity is an exponential decay process. The two decay rates have the same decay constant. We will divide the initial rate by the later rate and substitute the resulting ratio into the radioactive decay formula to find the decay constant. The half-life is then found from the decay constant.

EXECUTE The initial decay rate can be written

$$
\frac{\Delta N}{\Delta t}=32 \text { counts } / \mathrm{s}=-\lambda N_{0}
$$

The later decay rate can be written

$$
\frac{\Delta N}{\Delta t}=13 \text { counts } / \mathrm{s}=-\lambda N_{2}
$$

We divide the second decay rate by the first decay rate to find the ratio $N_{2} / N_{0}$ :

$$
\frac{N_{2}}{N_{0}}=\frac{13 \text { counts } / \mathrm{s}}{32 \text { counts } / \mathrm{s}}=0.4063
$$

The number of counts at 2 hours must be the initial number of counts times $e^{-e \lambda t}$ :

$$
N_{2}=N_{0} e^{-\lambda t}=N_{0} e^{-\lambda(2 \text { hours })} .
$$

We substitute and rearrange terms to solve for $\lambda$ :

$$
\begin{aligned}
& \frac{N_{2}}{N_{0}}=0.4063=e^{-\lambda(7200 \mathrm{~s})} \\
& \ln 0.4063=\ln \left(e^{-\lambda(7200 \mathrm{~s})}=-\lambda(7200 \mathrm{~s})\right. \\
& \quad \lambda=-\frac{\ln 0.4063}{7200 \mathrm{~s}}=1.25 \times 10^{-4} / \mathrm{s}
\end{aligned}
$$

The half-life is then

$$
T_{1 / 2}=\frac{\ln 2}{\lambda}=\frac{\ln 2}{1.25 \times 10^{-4} / \mathrm{s}}=5540 \mathrm{~s}=1.54 \mathrm{hr}
$$

The half-life of the sample is 1.54 hours.
EVALUATE We check our result by noting that the decay rate dropped by more than a factor of 2 in 1 hour, indicating that the half-life was less than 2 hours, in agreement with our result.

## 4: Carbon dating

The half-life of $\mathrm{C}^{14}$ is 5568 years. If $\mathrm{C}^{14}$ dating was done on a piece of 2000-year-old wood, how would the abundance of $\mathrm{C}^{14}$ in the wood compare with that of $\mathrm{C}^{14}$ in freshly cut wood from a similar tree?

## Solution

IDENTIFY AND SET UP $C^{14}$ is a radioactive nucleus, so its mass decreases exponentially with time. We'll use the half-life formula to solve the problem.

EXECUTE The number of nuclei decay according to

$$
N=N_{0} e^{-\lambda t} .
$$

The decay constant for $\mathrm{C}^{14}$ is

$$
\lambda=\frac{0.693}{T_{1 / 2}}=1.24 \times 10^{-4} / \mathrm{yr}
$$

We'll take $N_{0}$ as the number of $\mathrm{C}^{14}$ nuclei in the new wood. $N$ is the number of $\mathrm{C}^{14}$ nuclei in the old wood. Their ratio is

$$
\frac{N}{N_{0}}=e^{-\lambda t}=e^{-\left(1.24 \times 10^{-4} / \mathrm{yr}\right)(2000 \mathrm{yr})}=0.78
$$

The old wood has $78 \%$ of the abundance of $\mathrm{C}^{14}$ in the new wood.
EVALUATE If, instead, we were interested in the abundance of $C^{14}$ in the old wood, we could have found the age of that wood.

## Try It Yourself!

## 1: $\mathrm{N}^{14}$ mass and energy

Find the mass defect, the total binding energy, and the binding energy per nucleon for $\mathrm{N}^{14}$.

## Solution Checkpoints

IDENTIFY AND SET UP Use tables in the text and mass relations to solve.
EXECUTE The mass of the $\mathrm{N}^{14}$ nucleus is

$$
m\left(\mathrm{~N}^{14}\right)=14.00307 \mathrm{u}-7(0.000549 \mathrm{u})=13.99923 \mathrm{u}
$$

The combined mass of seven protons and seven neutrons is 14.11159 u , so the mass defect is 0.124 u .
The mass defect is the binding energy divided by $c^{2}$, so the binding energy is

$$
E_{B}=(0.124 \mathrm{u})(931.5 \mathrm{MeV} / \mathrm{u})=104.7 \mathrm{MeV}
$$

Dividing the binding energy by the 14 nucleons gives a binding energy per nucleon of 7.5 MeV .
EVALUATE If seven protons and seven neutrons were brought together to form an $\mathrm{N}^{14}$ nucleus, would energy be released or used up in the process?

## 2: $A=3$ nuclear decay

(a) Find the mass of the tritium nucleus ${ }_{1} \mathrm{H}^{3}$. (b) Find the mass of the helium nucleus ${ }_{2} \mathrm{He}^{3}$ plus an electron at rest. (c) If a tritium nucleus decays into a helium nucleus, plus an electron, plus a neutrino, how much energy is released as kinetic energy?

## Solution Checkpoints

IDENTIFY AND SET UP Use Table 42.3 in the text and mass relations to solve.
EXECUTE (a) The mass of the tritium nucleus is

$$
m\left({ }_{1} \mathrm{H}^{3}\right)=3.016049 \mathrm{u}-1(0.000549 \mathrm{u})=3.015500 \mathrm{u} .
$$

(b) The mass of the helium nucleus plus an electron at rest is

$$
m\left({ }_{2} \mathrm{He}^{3}\right)+m_{e}=3.016029 \mathrm{u}-2(0.000549 \mathrm{u})+0.000549 \mathrm{u}=3.015480 \mathrm{u} .
$$

(c) In the decay, the kinetic energy is shared between the electron and the neutrino. The mass defect is 0.00002 u , or 18.6 keV .
eVALUATE Tritium is a radioactive hydrogen isotope that emits beta rays. What is the maximum energy of the beta rays emitted?

# Particle Physics and Cosmology 

## Summary

In this chapter, we look at the formation of the universe through an investigation of particle physics. We will examine the four fundamental forces through which particles interact. We investigate the rules that govern these interactions and their influence on the universe as we know it. We will go on to consider the similarities among these forces and how they can be joined, or unified, into fewer forces. We'll then look back and see how the universe has evolved since the Big Bang-the massive explosion that generated the universe.

## Objectives

After studying this chapter, you will understand

- Key fundamental subatomic particles and how they were discovered.
- The four fundamental interactions between particles.
- The rules that govern the four interactions and how to apply them.
- How protons, neutrons, and other particles are built from quarks and gluons.
- How all the particles and interactions fit into the standard model.
- Evidence of the Big Bang and the history of the universe since the Big Bang.

Concepts and Equations

| Term | Description |
| :--- | :--- |
| Fundamental Particles | Protons, neutrons, and other hadrons are made of quarks with fractional <br> charge. Some hadrons, such as protons and neutrons, are made of three <br> quarks. Other hadrons, such as mesons, are made of two quarks. All massive <br> particles (protons, neutrons, electrons, etc.) have antiparticles-particles with <br> the same characteristics, but opposite charge. |
| Particles serve as mediators of the fundamental interactions. The photon is |  |
| the mediator for the electromagnetic force, the pion for the nuclear force. |  |, | Cyccelerators and Detectors |
| :--- |
| Fundamental Interactions |
| cles to high energies so that physicists can study their fundamental interac- |
| tions. Higher energy accelerators allow scientists to probe deeper into the |
| nucleus. Particle detectors are used to identify reaction products and measure |
| their energies and momenta. |

## Conceptual Questions

## 1: Fundamental forces

Which of the four fundamental forces influence electrons, neutrinos, and protons?

## Solution

IDENTIFY, SET UP, AND EXECUTE Electrons are affected by the electromagnetic force, since they have electric charge; the weak force, since they have weak charge; and the gravitational force, since they have mass (i.e. "gravitiational charge").

Neutrinos have no mass or electric charge and so are affected only by the weak force.

Protons are affected by all four forces, since they have mass, electric charge, and are made of quarks and, thus, experience the strong and weak force.

EVALUATE Understanding how particles interact helps us understand the forces between them.

## 2: Charge of quarks

Imagine that you are unaware of the charge of the up and down quarks, but you do know that two up quarks and one down quark make up the proton and one up quark and two down quarks make up the neutron. What are the charges of the up and down quarks?

## Solution

IDENTIFY AND SET UP Protons have charge $+e$ and neutrons are neutral. We will combine the quarks' charges, set them equal to the charges of the proton and neutron, and then solve for the quarks' charges.

EXECUTE The charges of the quarks in the proton add to $+1 e$, giving

$$
2 q_{u}+1 q_{d}=+1 e
$$

The charges of the quarks in the neutron add to 0 , giving

$$
1 q_{u}+2 q_{d}=0
$$

We use the neutron result to substitute into the proton result to solve for the charge on the up quarks, yielding

$$
\begin{aligned}
2 q_{u}-\frac{1}{2} q_{u} & =+1 e, \\
\frac{3}{2} q_{u} & =+1 e, \\
q_{u} & =+\frac{2}{3} e .
\end{aligned}
$$

We then find that the charge on the down quarks is $-e / 3$. So both up and down quarks have fractional charges of $e$ !

EVALUATE A process similar to the one we used in this problem led to the discovery that quarks have fractional charge.

## Problems

## 1: Rho meson decay

A neutral rho meson $\left(\rho^{0}\right.$, mass $\left.770 \mathrm{MeV} / \mathrm{c}^{2}\right)$ decays at rest into a pair of charged pions $\left(\pi^{+}, \pi^{-}\right.$, masses $140 \mathrm{MeV} / \mathrm{c}^{2}$ ). Find the energies and momenta of the pions.

## Solution

IDENTIFY AND SET UP The rest mass of the rho meson will be converted to the rest mass and kinetic energy of the pions. Momentum must be conserved, so each pion carries away equal momentum and energy.

EXECUTE A particle's relativistic energy is

$$
E=\sqrt{m^{2} c^{4}+p^{2} c^{2}}
$$

The rest mass of the rho meson is shared between the two pions. Each pion will have half the total energy of the rho meson:

$$
E_{\text {pion }}=\frac{1}{2} M_{\mathrm{rho}} c^{2}=\frac{1}{2}\left(770 \mathrm{MeV} / \mathrm{c}^{2}\right) c^{2}=385 \mathrm{MeV}
$$

The momentum is found from the relativistic energy:

$$
E_{\text {pion }}=\sqrt{{m_{\text {pion }}}^{2} c^{4}+p_{\text {pion }}^{2} c^{2}}
$$

Solving for the momentum yields

$$
p_{\text {pion }}=\frac{1}{c^{2}} \sqrt{E_{\text {pion }}^{2}-m_{\text {pion }}^{2} c^{4}}=\frac{1}{c} \sqrt{(385 \mathrm{MeV})^{2}-\left(140 \mathrm{MeV} / \mathrm{c}^{2}\right)^{2} c^{4}}=359 \mathrm{MeV} / \mathrm{c}
$$

Each pion has an energy of 385 MeV and a momentum of $359 \mathrm{MeV} / \mathrm{c}$.
EVALUATE This problem illustrates how physicists analyze elementary particle decays and how they must include relativistic kinematics in their work. They also check charge conservation, which holds in this case because the neutral rho meson decays into a positive and negative pair of pions.

## 2: $Z_{0}$ production

A $Z_{0}$ particle, one of the mediators of the weak interaction, has a mass of $91 \mathrm{GeV} / \mathrm{c}^{2}$. It can be produced at rest in a high energy electron-positron collider in which electrons and positrons in counterrotating beams are allowed to collide. Estimate the product of the magnetic field and beam radius $r B$ necessary to create $Z_{0}$ particles.

## Solution

IDENTIFY AND SET UP We will calculate the energy needed by the electrons and positrons to create the $Z_{0}$ particles. Then we will find the design parameter $r B$.

EXECUTE To create $Z_{0}$ particles, electrons and positrons are annihilated in a head-on collision. Each particle must carry half the rest energy of the $Z_{0}$ particle, so each must have 45.5 GeV of energy. The rest energy of the electron and positron is very small compared with 45.5 GeV , so the momentum is simply

$$
p=\frac{E}{c}=45.5 \mathrm{GeV} / \mathrm{c}
$$

The magnetic field necessary to keep a particle in a magnetic field at a fixed radius (recall Chapter 27) is

$$
r=\frac{p}{q B}
$$

This expression holds for these relativistic particles. Solving for $r B$ gives

$$
\begin{aligned}
r B & =\frac{p}{q} \\
& =\frac{45.5 \mathrm{GeV} / \mathrm{c}}{e} \\
& =\frac{(45.5 \mathrm{GeV})\left(1.6 \times 10^{-19} \mathrm{~J} / \mathrm{eV}\right)}{\left(1.6 \times 10^{-19} \mathrm{C}\right)\left(3 \times 10^{8} \mathrm{~m} / \mathrm{s}\right)}=152 \mathrm{~m} \cdot \mathrm{~T}
\end{aligned}
$$

EVALUATE With 0.1-T magnets, how large a ring is necessary? The ring's radius would need to be 1.5 km .

## 3: Possible decays

Consider the following decays of the positively charged pion:
(1) $\pi^{+} \rightarrow \mu^{+} \nu_{\mu}$
(2) $\pi^{+} \rightarrow e^{+} \nu_{e}$
(3) $\pi^{+} \rightarrow \mu^{+} \nu_{\mu} \gamma$
(4) $\pi^{+} \rightarrow e^{+} \nu_{e} \gamma$
(5) $\pi^{+} \rightarrow e^{+} \nu_{e} \pi^{0}$
(6) $\pi^{+} \rightarrow e^{+} \nu_{e} e^{+} e^{-}$
(7) $\pi^{+} \rightarrow e^{+} \nu_{e} \nu \bar{\nu}$
(8) $\pi^{+} \rightarrow \mu^{+} \bar{\nu}_{e}$
(9) $\pi^{+} \rightarrow \mu^{+} \nu_{e}$
(10) $\pi^{+} \rightarrow \mu^{-} e^{+} e^{-} \nu$

Which decays are allowed by conservation of lepton number?

## Solution

IDENTIFY AND SET UP We will use Table 44.2 in the text to determine the lepton numbers of the particles. We will then check whether the lepton numbers are conserved.

EXECUTE For all decays, the positive pion has no lepton number, so we will check whether there is a net lepton number in the decay products. In decay (1), the muon numbers are -1 and 1 , adding to zero. In decay (2), the electron numbers add to zero. In decays (3)-(5), the third particle has no lepton number, so the overall lepton numbers add to zero. In decays (6) and (7), there are four electron numbers, but they add to zero. So the first seven decays conserve lepton number and are allowed.

Decay (8) doesn't conserve lepton number, as the right side has -1 for the positive muon and -1 for the antineutrino. Decays (9) and (10) conserve overall lepton number, but don't conserve electron or muon numbers.

EVALUATE Decays (1) through (6) have been observed in particle accelerators. Decays (7) through (10) have not been observed.

## 4: Size of the universe

According to Hubble's law, matter at a distance $r$ travels away from us at a speed

$$
v=H_{0} r
$$

where $H_{0}=2.3 \times 10^{-18} / \mathrm{s}$ is Hubble's constant. What is the age of the universe according to Hubble's law?

## Solution

IDENTIFY AND SET UP If we assume that Hubble's constant has remained constant throughout the lifetime of the universe, then the age of the universe is the size of the universe divided by its velocity of expansion. We'll use this method to find the age of the universe.

EXECUTE The age of the universe is

$$
T=\frac{r}{v} .
$$

Substituting the velocity from Hubble's law, we find that

$$
T=\frac{r}{H_{0}} r=\frac{1}{H_{0}}=\frac{1}{2.3 \times 10^{-18} / \mathrm{s}}=4.35 \times 10^{17} \mathrm{~s}=13.8 \text { billion years. }
$$

Hubble's law predicts that the age of the universe is roughly 14 billion years old.
EVALUATE This is the currently accepted value for the age of the universe. There is evidence, however, that the universe is expanding at an accelerating rate, meaning that Hubble's constant is not constant. We are finding that the universe is even more interesting than we once thought.

## Try It Yourself!

## 1: Protons in a cyclotron

Protons in a cyclotron spiral out to a radius of 15 cm . The magnetic field has a magnitude of 1.25 T in the cyclotron. (a) Find the frequency of the alternating voltage used to accelerate the protons in the gap. (b) Find the energy of the protons.

## Solution Checkpoints

IDENTIFY AND SET UP Use the cyclotron frequency and energy conservation to solve.
EXECUTE (a) The cyclotron frequency, or angular frequency of rotation, is given by

$$
\omega=\frac{B e}{m}
$$

This corresponds to a voltage frequency of $1.91 \times 10^{7} \mathrm{~Hz}$.
(b) The kinetic energy of the protons is

$$
\frac{1}{2} m(\omega r)^{2}=1.6 \mathrm{MeV}
$$

EVALUATE The properties of the magnetic force in a uniform magnetic field, which we studied in Chapter 27 of the text, is key to building particle accelerators. Introductory physics is used by particle physicists every day.

## 2: Quark flavors

Find the charge and strangeness of all mesons that can be constructed from quark-antiquark pairs with flavors $u, d$, and $s$.

## Solution Checkpoints

IDENTIFY AND SET UP Find all possible combinations of the three quarks and three antiquarks.
EXECUTE There are nine independent combinations of these quarks, given in the following table (the charge and strangeness are found by adding the charge and strangeness of the quarks that make up the meson):

| Combination | Charge | Strangeness |
| :---: | :---: | :---: |
| $\bar{u} \bar{u}, d \bar{d}, s \bar{s}$ | 0 | 0 |
| $u \bar{d}$ | +1 | 0 |
| $d \bar{u}$ | -1 | 0 |
| $u \bar{s}$ | +1 | +1 |
| $s \bar{u}$ | -1 | -1 |
| $s \bar{d}$ | 0 | -1 |
| $d \bar{s}$ | 0 | +1 |

Table 44.1 Try It Yourself 2.
EVALUATE The second and third particles are the positive and negative pions. The last four particles are called kaons and are four varieties of strange mesons.

